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# Estimation of the Location of a 0 -type or $\infty$-type Singularity by Poisson Observations 

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#### Abstract

We consider an inhomogeneous Poisson process $X$ on $[0, T]$. The intensity function of $X$ is supposed to be strictly positive and smooth on $[0, T]$ except at the point $\theta$, in which it has either a 0 -type singularity (tends to 0 like $|x|^{p}, p \in(0,1)$ ), or an $\infty$-type singularity (tends to $\infty$ like $\left.|x|^{p}, p \in(-1,0)\right)$. We suppose that we know the shape of the intensity function, but not the location of the singularity. We consider the problem of estimation of this location (shift) parameter $\theta$ based on $n$ observations of the process $X$. We study the Bayesian estimators and, in the case $p>0$, the maximum likelihood estimator. We show that these estimators are consistent, their rate of convergence is $n^{1 /(p+1)}$, they have different limit distributions, and the Bayesian estimators are asymptotically efficient.


Keywords: inhomogeneous Poisson process, singularity, parameter estimation, Bayesian estimators, maximum likelihood estimator, consistency, limit distribution, convergence of moments, asymptotic efficiency

Mathematics Subject Classification (2000): 62M05

## 1 Introduction

Inhomogeneous Poisson process is one of the simplest point processes (see, for example, Daley and Vere-Jones (2]). However, due to the large choice of intensity functions, this model is reach enough and is widely used in many applied statistical problems, such as optical communications, reliability, biology, medicine, image treatment, and so on (see, for example, Karr [7] and [8], Snyder and Miller (11] and Thompson (12]).

The diversity of applications is also due to the possibility of using the likelihood ratio analysis. In parameter estimation problems the large samples theory is quite close to the one of the classical (i.i.d.) statistics. In particular, let us consider the problem of estimation of the parameter $\theta$ by $n$ independent observations on some fixed interval $[0, T]$ of an inhomogeneous Poisson process $X=\{X(t), 0 \leqslant t \leqslant T\}$ of intensity function $S_{\theta}(t)$. Let us mention that this problem is equivalent to the one of estimation of the parameter by one observation on a growing interval of a periodic inhomogeneous Poisson process. If the problem is regular (the model is locally asymptotically normal), then both the maximum likelihood estimator (MLE) $\widehat{\theta}_{n}$ and the Bayesian estimators (BE) $\widetilde{\theta}_{n}$ are consistent, asymptotically normal:

$$
\sqrt{n}\left(\widehat{\theta}_{n}-\theta\right) \Longrightarrow \mathcal{N}\left(0, I(\theta)^{-1}\right), \quad \sqrt{n}\left(\widetilde{\theta}_{n}-\theta\right) \Longrightarrow \mathcal{N}\left(0, I(\theta)^{-1}\right)
$$

and asymptotically efficient (see, for example, Kutoyants [9] and [10]). Here $I(\theta)$ is the Fisher information given by

$$
I(\theta)=\int_{0}^{T} \frac{\dot{S}_{\theta}^{2}(t)}{S_{\theta}(t)} d t
$$

where $S_{\theta}(t)$ is the intensity function and $\dot{S}_{\theta}(t)=\frac{\partial}{\partial \theta} S_{\theta}(t)$.
If the problem is not regular, then the properties of estimators essentially change. For example, if $S_{\theta}(\cdot)$ is smooth everywhere on $[0, T]$ except at the point $\theta$, in which it has a jump (consider for instance $S_{\theta}(t)=s(t-\theta)$ where $s(\cdot)$ is discontinuous in 0 ), then the MLE and BE are still consistent, but converge at a faster rate:

$$
n\left(\widehat{\theta}_{n}-\theta\right) \Longrightarrow \xi_{1}, \quad n\left(\widetilde{\theta}_{n}-\theta\right) \Longrightarrow \xi_{2}
$$

have different limit distributions ( $\xi_{1}$ and $\xi_{2}$ are different with $\mathbf{E} \xi_{1}^{2}>\mathbf{E} \xi_{2}^{2}$ ), and the BE are asymptotically efficient (see, for example, Kutoyants (9] and [10]).

In this paper we deal with the case where the intensity function $S_{\theta}(\cdot)$ is smooth everywhere on $[0, T]$ except at the point $\theta$, in which it has a singularity
of order $p$. The cusp type singularities were already studied in the preceding paper [1]. Here we consider 0 -type and $\infty$-type singularities. More precisely, we suppose that $S_{\theta}(t)=s(t-\theta)$, where $s(\cdot)$ is some known strictly positive function on $[-T, T] \backslash\{0\}$ and $\theta \in(0, T)$ is some unknown parameter, and that we have the following representation

$$
S_{\theta}(t)=s(t-\theta)=\left\{\begin{array}{ll}
a|t-\theta|^{p}+\psi(t-\theta), & \text { if } t<\theta \\
b|t-\theta|^{p}+\psi(t-\theta), & \text { if } t>\theta
\end{array},\right.
$$

where $a, b>0, p>-1$ (to guarantee the finiteness of intensity measure), and $\psi(\cdot)$ is smooth.

If $\psi(0) \neq 0$ and $p>1 / 2$ then, in spite of the singularity of the intensity function in $\theta$, the Fisher information is finite, and so this case can be treated as the regular one.

If $\psi(0) \neq 0$ and $0<p<1 / 2$ we say that the intensity function has a cusp at $\theta$. This is the case treated in [1] (where instead of $a, b>0$ it was supposed $a^{2}+b^{2}>0$ only). There it was shown that the MLE and the BE are consistent, converge at the rate $n^{1 /(2 p+1)}$ (which is faster than in the regular case but slower than in discontinuous case):

$$
n^{1 /(2 p+1)}\left(\widehat{\theta}_{n}-\theta\right) \Longrightarrow \eta_{1}, \quad n^{1 /(2 p+1)}\left(\widetilde{\theta}_{n}-\theta\right) \Longrightarrow \eta_{2},
$$

have different limit distributions, and the BE are asymptotically efficient. The convergence of moments was equally verified.

If $\psi(0)=0$ and $p>1$ then, as above, the Fisher information is finite and this case can be treated as the regular one.

If $\psi(0)=0$ and $0<p<1$ we say that the intensity function has a 0 type singularity at $\theta$. In this case we study the asymptotic behavior of the MLE and the BE, and we prove that the estimators are consistent, converge at the rate $n^{1 /(p+1)}$ (which is again intermediate between the regular and discontinuous case rates), have different limit distributions, and the BE are asymptotically efficient. We verify also the convergence of moments.

If $-1<p<0$ we say that the intensity function has a $\infty$-type singularity at $\theta$. In this case we study the asymptotic behavior of the BE only (MLE makes no sense in this case). We prove that the estimators are consistent, converge at the rate $n^{1 /(p+1)}$ (which is even faster than in discontinuous case), and are asymptotically efficient. We verify as well the convergence of moments.

Let us note, that the jump can also be considered as a singularity by taking $p=0$ and $a \neq b$, which explains that the rates are slower for $p>0$ and faster for $p<0$.

Let us also mention, that our results are similar to those obtained by Ibragimov and Khasminskii for the problem of estimation of a singularity location of the density for the i.i.d. model of observations. An exhaustive exposition of the results can be found in Chapter 6 of their book [6], but one can also refer to their previous works [4] and [5]. The asymptotic behavior of the MLE and of a wide class of BE obtained for this (i.i.d.) model is similar to the one obtained here for the model of Poisson observations. Particularly, the rate of convergence of the estimators is $n^{1 /(p+1)}$, and the BE are asymptotically efficient.

Finally, let us note that for the study of the asymptotic behavior of the estimators we use the method of Ibragimov and Khasminskii presented in their book [6] (see as well Kutoyants [10], where this method is applied to inhomogeneous Poisson process).

## 2 Main results

Suppose we observe $n$ realizations $\left(X_{1}, \ldots, X_{n}\right)=X^{n}$ of the Poisson process $X=\{X(t), 0 \leqslant t \leqslant T\}$ of intensity function $S_{\theta}(t)=s(t-\theta)$, where $\theta$ is some unknown parameter, $\theta \in \Theta=(\alpha, \beta) \subseteq(0, T)$, and $s(\cdot)$ is some known strictly positive function on $[-T, T] \backslash\{0\}$. We suppose that the function $s(\cdot)$ can be written in the form $s(t)=d(t)|t|^{p}+\psi(t)$, where $p \in(-1,0) \cup(0,1)$,

$$
d(t)=\left\{\begin{array}{ll}
a, & \text { if } t<0 \\
b, & \text { if } t>0
\end{array},\right.
$$

$a, b>0$, and the function $\psi(\cdot)$ is Hölder continuous on $[-T, T]$ of order higher than $(p+1) / 2$, that is $|\psi(x)-\psi(y)| \leqslant L|x-y|^{\kappa}$ for all $x, y \in[-T, T]$ with some fixed constants $L>0$ and $\varkappa>(p+1) / 2$. In the case $p>0$ we suppose equally that $\psi(0)=0$. Our aim is to estimate the parameter $\theta$ and to study the asymptotic behavior of estimators as $n$ goes to infinity.

The likelihood ratio in our problem can be written (see, for example, 10]) as

$$
\begin{aligned}
L\left(\theta, \theta_{1}, X^{n}\right)=\exp \{ & \sum_{i=1}^{n} \int_{0}^{T} \ln \frac{S_{\theta}(t)}{S_{\theta_{1}}(t)} d X_{i}(t) \\
& \left.-n \int_{0}^{T}\left[\frac{S_{\theta}(t)}{S_{\theta_{1}}(t)}-1\right] S_{\theta_{1}}(t) d t\right\}
\end{aligned}
$$

where $\theta_{1}$ is some fixed value of $\theta$.
As usually, introduce the MLE $\widehat{\theta}_{n}$ as one of the solutions of the equation

$$
L\left(\widehat{\theta}_{n}, \theta_{1}, X^{n}\right)=\sup _{\theta \in \Theta} L\left(\theta, \theta_{1}, X^{n}\right)
$$

and the $\mathrm{BE} \widetilde{\theta}_{n}$ for prior density $q$ (supposed to be positive and continuous) and quadratic loss function as

$$
\widetilde{\theta}_{n}=\int_{\alpha}^{\beta} \theta q\left(\theta \mid X^{n}\right) d \theta
$$

where the posterior density

$$
q\left(\theta \mid X^{n}\right)=L\left(\theta, \theta_{1}, X^{n}\right) q(\theta)\left(\int_{\alpha}^{\beta} L\left(\theta, \theta_{1}, X^{n}\right) q(\theta) d \theta\right)^{-1}
$$

Note that the MLE makes no sense in the case $p<0$, since in this case the likelihood equals infinity in any point $\theta$ which is event of one of the Poisson processes $X_{1}, \ldots, X_{n}$.

To describe the properties of these estimators we need to introduce the stochastic process

$$
\begin{aligned}
Z(u)=\exp \{ & p \int_{-\infty}^{+\infty} \ln \left|1-\frac{u}{z}\right| \pi(d z)+\ln \frac{a}{b} \int_{0}^{u} Y(d z) \\
& -\int_{-\infty}^{+\infty}\left[\left|1-\frac{u}{z}\right|^{p}-1-p \ln \left|1-\frac{u}{z}\right|\right] d(z)|z|^{p} d z \\
& \left.-\frac{a-b}{p+1}|u|^{p+1} \operatorname{sign}(u)\right\}, \quad u \in \mathbb{R}
\end{aligned}
$$

Here and in the sequel $Y$ denotes a Poisson process on $\mathbb{R}$ of intensity function $S_{0}(z)=d(z)|z|^{p}$, and $\pi$ is its centered version : $\pi=Y-\mathbf{E} Y$.

We introduce also the random variable $\zeta$, and in the case $p>0$ the random variable $\xi$ by the equations

$$
\zeta=\int_{-\infty}^{+\infty} u Z(u) d u\left(\int_{-\infty}^{+\infty} Z(u) d u\right)^{-1}
$$

and

$$
Z(\xi)=\sup _{u \in \mathbb{R}} Z(u) .
$$

Let us note here, that $\xi$ is well defined in the case $p>0$, since in this case with probability one the process $Z(u)$ attains its maximum in a unique point (see, for example, [3])

Now we can finally state the main results of this paper.

Theorem 1 Under the maid assumptions, the following lower bound on the risks of all estimators holds: for any $\theta_{0} \in \Theta$ we have

$$
\lim _{\delta \rightarrow 0} \underline{\lim _{n \rightarrow \infty}} \inf _{\bar{\theta}_{n}} \sup _{\left|\theta-\theta_{0}\right|<\delta} \mathbf{E}_{\theta}\left(n^{1 /(p+1)}\left(\bar{\theta}_{n}-\theta\right)\right)^{2} \geqslant \mathbf{E} \zeta^{2}
$$

where inf is taken over all possible estimators $\bar{\theta}_{n}$ of $\theta$.

This theorem leads us to introduce the following
Definition 2 We say that the estimator $\bar{\theta}_{n}$ is asymptotically efficient if

$$
\lim _{\delta \rightarrow 0} \lim _{n \rightarrow \infty} \sup _{\left|\theta-\theta_{0}\right|<\delta} \mathbf{E}_{\theta}\left(n^{1 /(p+1)}\left(\bar{\theta}_{n}-\theta\right)\right)^{2}=\mathbf{E} \zeta^{2}
$$

for any $\theta_{0} \in \Theta$.
For the BE we have the following
Theorem 3 The $B E \tilde{\theta}_{n}$ have uniformly in $\theta \in \mathbf{K}($ for any compact $\mathbf{K} \subset \Theta)$ the following properties:

- $\widetilde{\theta}_{n}$ is consistent, that is

$$
\widetilde{\theta}_{n} \xrightarrow{\mathbf{P}_{\theta}} \theta \text { (convergence in probability), }
$$

- the limit distribution of $\widetilde{\theta}_{n}$ is $\zeta$, that is

$$
n^{1 /(p+1)}\left(\widetilde{\theta}_{n}-\theta\right) \Longrightarrow \zeta(\text { convergence in law })
$$

- for any $k>0$ we have

$$
\lim _{n \rightarrow \infty} \mathbf{E}_{\theta}\left|n^{1 /(p+1)}\left(\widetilde{\theta}_{n}-\theta\right)\right|^{k}=\mathbf{E}|\zeta|^{k}
$$

and, moreover, $\tilde{\theta}_{n}$ is asymptotically efficient.

And for the MLE (in the case $p>0$ ) we have the following
Theorem 4 Let $p \in(0,1)$. The MLE $\widehat{\theta}_{n}$ has uniformly in $\theta \in \mathbf{K}$ (for any compact $\mathbf{K} \subset \Theta$ ) the following properties:

- $\widehat{\theta}_{n}$ is consistent, that is

$$
\widehat{\theta}_{n} \xrightarrow{\mathbf{P}_{\theta}} \theta \text { (convergence in probability), }
$$

- the limit distribution of $\widehat{\theta}_{n}$ is $\xi$, that is

$$
n^{1 /(p+1)}\left(\widehat{\theta}_{n}-\theta\right) \Longrightarrow \xi(\text { convergence in law }),
$$

- for any $k>0$ we have

$$
\lim _{n \rightarrow \infty} \mathbf{E}_{\theta}\left|n^{1 /(p+1)}\left(\hat{\theta}_{n}-\theta\right)\right|^{k}=\mathbf{E}|\xi|^{k}
$$

To prove the above stated theorems we apply the method of Ibragimov and Khasminskii (see [6]). For this we denote $\theta_{u}=\theta+u n^{-1 /(p+1)}$ for all $u \in U_{n}=\left(n^{1 /(p+1)}(\alpha-\theta), n^{1 /(p+1)}(\beta-\theta)\right)$, we introduce the normalized likelihood ratio process as

$$
Z_{n}(u)=L\left(\theta_{u}, \theta, X^{n}\right), \quad u \in U_{n},
$$

and we establish (the proofs are in the next section) the following three lemmas.

Lemma 5 The finite-dimensional distributions of $Z_{n}(u)$ converge to those of $Z(u)$ uniformly in $\theta \in \mathbf{K}$ (for any compact $\mathbf{K} \subset \Theta$ ).

Lemma 6 For any compact $\mathbf{K} \subset \Theta$ there exists some positive constant $C$ such that

$$
\mathbf{E}_{\theta}\left|Z_{n}^{1 / 2}\left(u_{1}\right)-Z_{n}^{1 / 2}\left(u_{2}\right)\right|^{2} \leqslant C\left|u_{1}-u_{2}\right|^{p+1}
$$

for all $u_{1}, u_{2} \in U_{n}, \theta \in \mathbf{K}$ and $n$ sufficiently large.
Lemma 7 For any compact $\mathbf{K} \subset \Theta$ there exists some positive constant $c$ such that

$$
\mathbf{E}_{\theta} Z_{n}^{1 / 2}(u) \leqslant \exp \left\{-c|u|^{p+1}\right\}
$$

for all $u \in U_{n}, \theta \in \mathbf{K}$ and $n \in \mathbb{N}$.
Using these lemmas and applying Theorems 1.9.1, 1.10.2 and 1.10.1 of [6],


## 3 Proofs of the Lemmas

For convenience of notation, all throughout this section $C$ and $c$ denote generic positive constants which can differ from formula to formula (and even in the same formula), and we put $\nu=1 /(p+1)$.

First of all let us fix some $\delta>0$ such that $c d(t)|t|^{p} \leqslant s(t) \leqslant C d(t)|t|^{p}$ on $(-\delta, \delta)$, and $s(t) \geqslant c$ on $[-T, T] \backslash(-\delta / 4, \delta / 4)$. To do so, we note that

$$
|\psi(t)| \leqslant|\psi(0)|+C|t|^{\varkappa}=C|t|^{p}\left(|t|^{\varkappa-p}+c|\psi(0)||t|^{-p}\right) \leqslant \min \{a, b\}|t|^{p} / 2
$$

for $t \in(-\delta, \delta)$, since $\varkappa-p>(1-p) / 2>0$ and $\psi(0)=0$ in the case $p>0$. It follows clearly

$$
s(t)=d(t)|t|^{p}+\psi(t) \geqslant(d(t)-\min \{a, b\} / 2)|t|^{p} \geqslant d(t)|t|^{p} / 2
$$

and $s(t) \leqslant 2 d(t)|t|^{p}$. Finally, on the compact set $[-T, T] \backslash(-\delta / 4, \delta / 4)$ we have $s(t) \geqslant c$ since the function $s(\cdot)$ is continuous.

Now, let us fix some sequence $\left(A_{n}\right)$ such that $A_{n} \rightarrow+\infty$ sufficiently slowly. More precisely, we suppose that $A_{n} n^{-\nu} \rightarrow 0$ and we will give some additional conditions below. We split the interval $[0, T]$ in three parts:

$$
\begin{aligned}
& E_{1}=\left\{t:|t-\theta|<A_{n} n^{-\nu}\right\}=\left(\theta-A_{n} n^{-\nu}, \theta+A_{n} n^{-\nu}\right), \\
& E_{2}=\left\{t: A_{n} n^{-\nu}<|t-\theta|<\delta\right\}=\left(\theta-\delta, \theta-A_{n} n^{-\nu}\right) \cup\left(\theta+A_{n} n^{-\nu}, \theta+\delta\right), \\
& E_{3}=\{t: \delta<|t-\theta|\}=(0, \theta-\delta) \cup(\theta+\delta, T) .
\end{aligned}
$$

In order to prove Lemma 5 we will only study the convergence of the one-dimensional (the general case can be considered similarly) distributions of the process

$$
\begin{aligned}
\ln Z_{n}(u) & =\sum_{i=1}^{n} \int_{0}^{T} \ln \frac{S_{\theta_{u}}(t)}{S_{\theta}(t)} d X_{i}(t)-n \int_{0}^{T}\left[\frac{S_{\theta_{u}}(t)}{S_{\theta}(t)}-1\right] S_{\theta}(t) d t \\
& =\sum_{i=1}^{n} \int_{0}^{T} f d X_{i}(t)-n \int_{0}^{T} g S_{\theta}(t) d t,
\end{aligned}
$$

where we denote

$$
f=f(\theta, t, u, n)=\ln \frac{S_{\theta_{u}}(t)}{S_{\theta}(t)} \quad \text { and } \quad g=g(\theta, t, u, n)=\frac{S_{\theta_{u}}(t)}{S_{\theta}(t)}-1 .
$$

The characteristic function of the random variable $\ln Z_{n}(u)$ can be written as (see, for example, Lemma 1.1 of (10])

$$
\phi_{n}(\lambda)=\mathbf{E}_{\theta} \exp \left\{i \lambda \ln Z_{n}(u)\right\}=\exp \left\{n \int_{0}^{T}\left[e^{i \lambda f}-1-i \lambda g\right] S_{\theta}(t) d t\right\}
$$

and hence

$$
\begin{equation*}
\ln \phi_{n}(\lambda)=n \int_{0}^{T}\left[e^{i \lambda f}-1-i \lambda f\right] S_{\theta}(t) d t+i \lambda n \int_{0}^{T}[f-g] S_{\theta}(t) d t \tag{1}
\end{equation*}
$$

To study this expression, let us at first establish the two following properties.
a) For any fixed $u$, we have $\lim _{n \rightarrow \infty} g(\theta, t, u, n)=0$ uniformly in $\theta \in \mathbf{K}$ and $t \in E_{2} \cup E_{3}$.
b) We have

$$
\lim _{n \rightarrow \infty} n \int_{E_{2} \cup E_{3}} g^{2} S_{\theta}(t) d t=0 .
$$

To prove a), we put $y=t-\theta \in\left(E_{2}-\theta\right) \cup\left(E_{3}-\theta\right)$ and we write

$$
\begin{aligned}
|g(\theta, t, u, n)| & =\left|\frac{s\left(t-\theta_{u}\right)}{s(t-\theta)}-1\right| \\
& =\left|\frac{d\left(y-u n^{-\nu}\right)\left|y-u n^{-\nu}\right|^{p}+\psi\left(y-u n^{-\nu}\right)-d(y)|y|^{p}-\psi(y)}{s(y)}\right|= \\
& =\frac{\left|C\left(\left|y-u n^{-\nu}\right|^{p}-|y|^{p}\right)+\psi\left(y-u n^{-\nu}\right)-\psi(y)\right|}{s(y)} \\
& \leqslant C \frac{\left|y-u n^{-\nu}\right|^{p}-|y|^{p} \mid}{s(y)}+\frac{\left|\psi\left(y-u n^{-\nu}\right)-\psi(y)\right|}{s(y)} \\
& =M_{1}+M_{2}
\end{aligned}
$$

with evident notations.
For $y \in E_{2}-\theta$ we have

$$
\left.M_{1} \leqslant C \frac{| | y-\left.u n^{-\nu}\right|^{p}-|y|^{p} \mid}{c|y|^{p}}=C| | 1-\left.\frac{u}{y n^{\nu}}\right|^{p}-1 \right\rvert\, \leqslant \frac{C|u|}{|y| n^{\nu}} \leqslant \frac{C|u|}{A_{n}} \rightarrow 0,
$$

$$
\begin{aligned}
& M_{2} \leqslant \frac{C\left|u n^{-\nu}\right|^{\varkappa}}{c}=C|u|^{\varkappa} n^{-\nu \varkappa} \rightarrow 0 \quad \text { if } p<0, \\
& M_{2} \leqslant \frac{C\left|u n^{-\nu}\right|^{\varkappa}}{c|y|^{p}} \leqslant \frac{C|u|^{\varkappa} n^{-\nu \varkappa}}{\left(A_{n} n^{-\nu}\right)^{p}}=C|u|^{\varkappa} \frac{n^{-\nu(\varkappa-p)}}{A_{n}^{p}} \rightarrow 0 \quad \text { if } p>0 .
\end{aligned}
$$

Finally for $y \in E_{3}-\theta$, using the Hölder continuity of $s(\cdot)$ we have

$$
|g(\theta, t, u, n)| \leqslant \frac{C\left|u n^{-\nu}\right|^{\varkappa}}{c}=C|u|^{\varkappa} n^{-\nu \varkappa} \rightarrow 0 .
$$

So, a) is proved.
To prove b), we first note that

$$
\begin{aligned}
n \int_{E_{3}} g^{2} S_{\theta}(t) d t & =n \int_{E_{3}} \frac{\left(S_{\theta_{u}}(t)-S_{\theta}(t)\right)^{2}}{S_{\theta}(t)} d t=n \int_{E_{3}-\theta} \frac{\left(s\left(y-u n^{-\nu}\right)-s(y)\right)^{2}}{s(y)} d y \\
& \leqslant C n \int_{E_{3}-\theta}\left|u n^{-\nu}\right|^{2 \varkappa} d y=C|u|^{2 \varkappa} n^{-(2 \nu \varkappa-1)} \rightarrow 0
\end{aligned}
$$

since $2 \nu \varkappa-1>0$. To conclude the proof it remains to show that

$$
n \int_{\theta+A_{n} n^{-\nu}}^{\theta+\delta} g^{2} S_{\theta}(t) d t+n \int_{\theta-\delta}^{\theta-A_{n} n^{-\nu}} g^{2} S_{\theta}(t) d t \rightarrow 0
$$

For the first term we have

$$
\begin{aligned}
n \int_{\theta+A_{n} n^{-\nu}}^{\theta+\delta} g^{2} S_{\theta}(t) d t= & n \int_{A_{n} n^{-\nu}}^{\delta} \frac{\left(s\left(y-u n^{-\nu}\right)-s(y)\right)^{2}}{s(y)} d y \\
\leqslant & n \int_{A_{n} n^{-\nu}}^{\delta} \frac{\left(s\left(y-u n^{-\nu}\right)-s(y)\right)^{2}}{c|y|^{p}} d y \\
= & C n \int_{A_{n} n^{-\nu}}^{\delta} \frac{\left(\left|y-u n^{-\nu}\right|^{p}-|y|^{p}\right)^{2}}{|y|^{p}} d y \\
& +C n \int_{A_{n} n^{-\nu}}^{\delta} \frac{\left(\psi\left(y-u n^{-\nu}\right)-\psi(y)\right)^{2}}{|y|^{p}} d y \\
& +C n \int_{A_{n} n^{-\nu}}^{\delta} \frac{\left(\psi\left(y-u n^{-\nu}\right)-\psi(y)\right)\left(\left|y-u n^{-\nu}\right|^{p}-|y|^{p}\right)}{|y|^{p}} d y
\end{aligned}
$$

$$
=n J_{1}+n J_{2}+n J_{3}
$$

with evident notations. Further
$n J_{1}=C n \int_{A_{n}}^{\delta n^{\nu}} \frac{\left(|z-u|^{p}-|z|^{p}\right)^{2} n^{-2 \nu p}}{|z|^{p} n^{-\nu p}} n^{-\nu} d z \leqslant C \int_{A_{n}}^{+\infty} \frac{\left(|z-u|^{p}-|z|^{p}\right)^{2}}{|z|^{p}} d z \rightarrow 0$
since

$$
\frac{\left(|z-u|^{p}-|z|^{p}\right)^{2}}{|z|^{p}}=|z|^{p}\left(\left|1-\frac{u}{z}\right|^{p}-1\right)^{2} \sim|z|^{p}\left(\frac{C}{z}\right)^{2}=C|z|^{p-2}
$$

and $p-2<-1$. Similarly

$$
n J_{2} \leqslant C n \int_{A_{n} n^{-\nu}}^{\delta} \frac{\left|u n^{-\nu}\right|^{2 \varkappa}}{|y|^{p}} d y \leqslant C|u|^{2 \varkappa} n^{-(2 \nu \varkappa-1)} \int_{0}^{\delta}|y|^{-p} d y \rightarrow 0
$$

since $2 \nu \varkappa-1>0$ and $-p>-1$. Finally

$$
\left|n J_{3}\right| \leqslant C \sqrt{\left(n J_{1}\right)\left(n J_{2}\right)} \rightarrow 0
$$

by Cauchy-Schwarz inequality, and so the first term converges to 0 .
The second term can be treated in the same way. So, b) is proved.
Now let us return to the study of the characteristic function $\phi_{n}(\cdot)$. Using (取) we can write

$$
\ln \phi_{n}=\varphi_{1}+\varphi_{2}+\varphi_{3}
$$

where we put

$$
\varphi_{k}=n \int_{E_{k}}\left[e^{i \lambda f}-1-i \lambda f\right] S_{\theta}(t) d t+i \lambda n \int_{E_{k}}[f-g] S_{\theta}(t) d t .
$$

For $\varphi_{3}$ we get

$$
\begin{aligned}
\varphi_{3} & =n \int_{E_{3}}\left[e^{i \lambda f}-1-i \lambda f\right] S_{\theta}(t) d t+i \lambda n \int_{E_{3}}[f-g] S_{\theta}(t) d t \\
& \simeq \frac{1}{2} n \int_{E_{3}}(i \lambda f)^{2} S_{\theta}(t) d t-\frac{1}{2} i \lambda n \int_{E_{3}} g^{2} S_{\theta}(t) d t \\
& \simeq-\frac{1}{2} \lambda^{2} n \int_{E_{3}} g^{2} S_{\theta}(t) d t-\frac{1}{2} i \lambda n \int_{E_{3}} g^{2} S_{\theta}(t) d t \rightarrow 0
\end{aligned}
$$

where the symbol " $\simeq$ " means equality of limits
In the same way we get $\varphi_{2} \rightarrow 0$, and it remains to study $\varphi_{1}$. For this we put $y_{u}=y-u n^{-\nu}, \alpha(y)=d(y)|y|^{p}$,

$$
\beta(y)=\frac{s(y)}{\alpha(y)}=1+\frac{\psi(y)}{d(y)|y|^{p}}
$$

for $y \in[-T, T] \backslash\{0\}$, and $\beta(0)=1$. Note that the function $\beta(\cdot)$ is clearly Hölder continuous of order

$$
\mu=\left\{\begin{array}{ll}
\varkappa-p, & \text { if } p>0 \\
\min \{\varkappa,-p\}, & \text { if } p<0
\end{array} .\right.
$$

We have

$$
\begin{aligned}
\varphi_{1}= & n \int_{E_{1}}\left[e^{i \lambda f}-1-i \lambda f\right] S_{\theta}(t) d t+i \lambda n \int_{E_{1}}[f-g] S_{\theta}(t) d t \\
= & n \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}}\left[\left(\frac{\alpha\left(y_{u}\right) \beta\left(y_{u}\right)}{\alpha(y) \beta(y)}\right)^{i \lambda}-1-i \lambda \ln \frac{\alpha\left(y_{u}\right) \beta\left(y_{u}\right)}{\alpha(y) \beta(y)}\right] \alpha(y) \beta(y) d y \\
& -i \lambda n \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}}\left[\frac{\alpha\left(y_{u}\right) \beta\left(y_{u}\right)}{\alpha(y) \beta(y)}-1-\ln \frac{\alpha\left(y_{u}\right) \beta\left(y_{u}\right)}{\alpha(y) \beta(y)}\right] \alpha(y) \beta(y) d y \\
\simeq & n \int_{A_{n}}^{A_{n} n^{-\nu}}\left[\frac{\alpha^{i \lambda}\left(y_{u}\right)}{\alpha^{i \lambda}(y)}-1-i \lambda \ln \frac{\alpha\left(y_{u}\right)}{\alpha(y)}\right] \alpha(y) d y \\
& +n \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}} \frac{\alpha^{i \lambda}\left(y_{u}\right)}{\alpha^{i \lambda}(y)}\left(\frac{\beta^{i \lambda}\left(y_{u}\right)}{\beta^{i \lambda}(y)}-1\right) \alpha(y) d y \\
& -i \lambda n \int_{-A_{n} n^{n^{-\nu}}}^{A_{n} n^{-\nu}}\left[\frac{\alpha\left(y_{u}\right)}{\alpha(y)}-1-\ln \frac{\alpha\left(y_{u}\right)}{\alpha(y)}\right] \alpha(y) d y \\
& -i \lambda n \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}} \alpha\left(y_{u}\right)\left(\frac{\beta\left(y_{u}\right)}{\beta(y)}-1\right) d y \\
= & I_{1}+n I_{2}-i \lambda n I_{3}-i \lambda n I_{4}
\end{aligned}
$$

with evident notations.

Using the Hölder continuity of $\beta(\cdot)$, we get

$$
\begin{aligned}
\left|n I_{4}\right| & \leqslant n \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}} \alpha\left(y_{u}\right) \frac{\left|\beta\left(y_{u}\right)-\beta(y)\right|}{\beta(y)} d y \leqslant n C\left|u n^{-\nu}\right|^{\mu} \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}} \frac{\alpha\left(y_{u}\right)}{\beta(y)} d y \\
& \simeq C|u|^{\mu} n^{1-\nu \mu} \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}} \alpha\left(y_{u}\right) d y=C|u|^{\mu} n^{1-\nu \mu} \int_{\left(-A_{n}-u\right) n^{-\nu}}^{\left(A_{n}-u\right) n^{-\nu}} d(x)|x|^{p} d x \\
& =C|u|^{\mu} n^{1-\nu \mu}\left[\frac{a}{p+1}\left(A_{n}+u\right)^{p+1}+\frac{b}{p+1}\left(A_{n}-u\right)^{p+1}\right] n^{-\nu(p+1)} \\
& \leqslant C|u|^{\mu}\left(A_{n}+|u|\right)^{p+1} n^{-\nu \mu} \rightarrow 0
\end{aligned}
$$

if $\left(A_{n}\right)$ is chosen so that $A_{n} n^{-\nu^{2} \mu} \rightarrow 0$.
Similarly, noting that $\beta^{i \lambda}(\cdot)$ is also Hölder continuous of order $\mu$, we get

$$
\begin{aligned}
\left|n I_{2}\right| & \leqslant n \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}} \frac{\left|\alpha^{i \lambda}\left(y_{u}\right)\right|}{\left|\alpha^{i \lambda}(y)\right|} \frac{\left|\beta^{i \lambda}\left(y_{u}\right)-\beta^{i \lambda}(y)\right|}{\left|\beta^{i \lambda}(y)\right|} \alpha(y) d y \\
& =n \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}}\left|\beta^{i \lambda}\left(y_{u}\right)-\beta^{i \lambda}(y)\right| \alpha(y) d y \leqslant n C\left|u n^{-\nu}\right|^{\mu} \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}} \alpha(y) d y \\
& =C|u|^{\mu} n^{1-\nu \mu} A_{n}^{p+1} n^{-\nu(p+1)}=C|u|^{\mu} A_{n}^{p+1} n^{-\nu \mu} \rightarrow 0
\end{aligned}
$$

under the same condition on the choice of $\left(A_{n}\right)$.
For $n I_{3}$ we can write

$$
\begin{aligned}
n I_{3} & =n \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}}\left[\frac{d\left(y_{u}\right)\left|y_{u}\right|^{p}}{d(y)|y|^{p}}-1-\ln \frac{d\left(y_{u}\right)\left|y_{u}\right|^{p}}{d(y)|y|^{p}}\right] d(y)|y|^{p} d y \\
& =n \int_{-A_{n}}^{A_{n}}\left[\frac{d(z-u)}{d(z)}\left|1-\frac{u}{z}\right|^{p}-1-\ln \left(\frac{d(z-u)}{d(z)}\left|1-\frac{u}{z}\right|^{p}\right)\right] \frac{d(z)|z|^{p}}{n^{\nu(p+1)}} d z \\
& \rightarrow \int_{-\infty}^{\infty}\left[\frac{d(z-u)}{d(z)}\left|1-\frac{u}{z}\right|^{p}-1-\ln \left(\frac{d(z-u)}{d(z)}\left|1-\frac{u}{z}\right|^{p}\right)\right] d(z)|z|^{p} d z
\end{aligned}
$$

Note that the last integral is finite, since

$$
\frac{d(z-u)}{d(z)}=\left(\frac{a}{b}\right)^{\operatorname{sign}(u) \mathbb{1}_{\left[u^{-}, u^{+}\right]}(z)}
$$

and hence the integrand behaves as $C|z|^{p-2}$ as $z \rightarrow \infty$.
Finally, for $n I_{1}$ we have

$$
\begin{aligned}
n I_{1} & =n \int_{-A_{n} n^{-\nu}}^{A_{n} n^{-\nu}}\left[\frac{d^{i \lambda}\left(y_{u}\right)\left|y_{u}\right|^{i \lambda p}}{d^{i \lambda}(y)|y|^{i \lambda p}}-1-i \lambda \ln \frac{d\left(y_{u}\right)\left|y_{u}\right|^{p}}{d(y)|y|^{p}}\right] d(y)|y|^{p} d y \\
& =n \int_{-A_{n}}^{A_{n}}\left[\left(\frac{d(z-u)}{d(z)}\left|1-\frac{u}{z}\right|^{p}\right)^{i \lambda}-1-i \lambda \ln \left(\frac{d(z-u)}{d(z)}\left|1-\frac{u}{z}\right|^{p}\right)\right] \frac{d(z)|z|^{p}}{n^{\nu(p+1)}} d z \\
& \rightarrow \int_{-\infty}^{\infty}\left[\left(\frac{d(z-u)}{d(z)}\left|1-\frac{u}{z}\right|^{p}\right)^{i \lambda}-1-i \lambda \ln \left(\frac{d(z-u)}{d(z)}\left|1-\frac{u}{z}\right|^{p}\right)\right] d(z)|z|^{p} d z
\end{aligned}
$$

where the last integral is finite as before.
So we get

$$
\begin{aligned}
\ln \phi_{n} \rightarrow \mathcal{L}= & \int_{-\infty}^{\infty}\left[\left(\frac{d(z-u)}{d(z)}\right)^{i \lambda}\left|1-\frac{u}{z}\right|^{i \lambda p}-1-i \lambda p \ln \left|1-\frac{u}{z}\right|\right] d(z)|z|^{p} d z \\
& -i \lambda \int_{-\infty}^{\infty}\left[\frac{d(z-u)}{d(z)}\left|1-\frac{u}{z}\right|^{p}-1-p \ln \left|1-\frac{u}{z}\right|\right] d(z)|z|^{p} d z
\end{aligned}
$$

To terminate the proof of Lemma ${ }^{5}$ it remains to show that $\mathcal{L}=\ln \phi$, where $\phi(\cdot)$ is the characteristic function of $\ln Z(u)$.

Recall that

$$
\begin{aligned}
\ln Z(u)= & p \int_{-\infty}^{+\infty} \ln \left|1-\frac{u}{z}\right| \pi(d z)+\ln \frac{a}{b} \int_{0}^{u} Y(d z) \\
& -\int_{-\infty}^{+\infty}\left[\left|1-\frac{u}{z}\right|^{p}-1-p \ln \left|1-\frac{u}{z}\right|\right] d(z)|z|^{p} d z \\
& -\frac{a-b}{p+1}|u|^{p+1} \operatorname{sign}(u) \\
= & K_{1}+K_{2}-K_{3}-K_{4}
\end{aligned}
$$

with evident notations.
Hence
$\ln \phi(\lambda)=\ln \mathbf{E} \exp \{i \lambda \ln Z(u)\}$

$$
\begin{aligned}
= & \ln \mathbf{E} \exp \left\{i \lambda K_{1}+i \lambda K_{2}\right\}-i \lambda K_{3}-i \lambda K_{4} \\
= & \ln \mathbf{E} \exp \left\{i \lambda \int_{-\infty}^{+\infty}\left[p \ln \left|1-\frac{u}{z}\right|+\ln \left(\frac{a}{b}\right) \operatorname{sign}(u) \mathbb{1}_{\left[u^{-}, u^{+}\right]}(z)\right] Y(d z)\right\} \\
& +i \lambda \ln \frac{a}{b} \int_{0}^{u} d(z)|z|^{p} d z-i \lambda K_{3}-i \lambda K_{4} \\
= & \int_{-\infty}^{+\infty}\left[\exp \left\{i \lambda p \ln \left|1-\frac{u}{z}\right|+i \lambda \ln \left(\frac{a}{b}\right) \operatorname{sign}(u) \mathbb{1}_{\left[u^{-}, u^{+}\right]}(z)\right\}-1\right. \\
& \quad-\left.i \lambda p \ln \left|1-\frac{a}{z} \int_{0}^{u} d(z)\right| z\right|^{p} d z-i \lambda K_{3}-i \lambda K_{4} \\
& \left.+\int_{0}^{+\infty}\left[\left\lvert\, \frac{a}{b}\right.\right) \operatorname{sign}(u) \mathbb{1}_{\left[u^{-}, u^{+}\right]}(z)\right] d(z)|z|^{p} d z \\
= & {\left[1-\left.\frac{u}{z}\right|^{i \lambda p}\left(\frac{d(z-u)}{d(z)}\right)^{i \lambda}-1-i \lambda p \ln \left|1-\frac{u}{z}\right|\right] d(z)|z|^{p} d z } \\
& -i \lambda K_{3}-i \lambda K_{4} \\
= & \mathcal{L}(\lambda)+i \lambda \int_{-\infty}^{+\infty}\left|1-\frac{u}{z}\right|^{p}\left(\frac{d(z-u)}{d(z)}-1\right) d(z)|z|^{p} d z-i \lambda K_{4} \\
= & \mathcal{L}(\lambda)+i \lambda \int_{-\infty}^{+\infty}|z-u|^{p}(a-b) \operatorname{sign}(u) \mathbb{1}_{\left[u^{-}, u u^{+}\right]}(z) d z-i \lambda K_{4} \\
= & \mathcal{L}(\lambda)+i \lambda(a-b) \int_{0}^{u}|z-u|^{p} d z-i \lambda \frac{a-b}{p+1}|u|^{p+1} \operatorname{sign}(u)=\mathcal{L}(\lambda) .
\end{aligned}
$$

So, the convergence of the one-dimensional distributions is proved. The case of higher-dimensional distributions can be treated similarly. The uniformity in $\theta$ on any compact set $\mathbf{K} \subset \Theta$ is evident. Lemma ${ }^{\text {R }}$ is proved.

Now let us prove Lemma 6. For $\left|u_{1}-u_{2}\right| \geqslant 1$ the assertion is evident since for all $\theta$ and $n$ we have

$$
\mathbf{E}_{\theta}\left|Z_{n}^{1 / 2}\left(u_{1}\right)-Z_{n}^{1 / 2}\left(u_{2}\right)\right|^{2} \leqslant 4 \leqslant 4\left|u_{1}-u_{2}\right|^{2 p+1}
$$

Suppose now that $\left|u_{1}-u_{2}\right| \leqslant 1$. Denoting $\Delta=u_{1}-u_{2}$ and using Lemma 1.5 of 10 we can write

$$
\begin{aligned}
\mathbf{E}_{\theta}\left|Z_{n}^{1 / 2}\left(u_{1}\right)-Z_{n}^{1 / 2}\left(u_{2}\right)\right|^{2} & \leqslant n \int_{0}^{T}\left[\sqrt{S_{\theta_{u_{1}}}(t)}-\sqrt{S_{\theta_{u_{2}}}(t)}\right]^{2} d t \\
& =n \int_{0}^{T}\left[\sqrt{s\left(t-\theta-u_{1} n^{-\nu}\right)}-\sqrt{s\left(t-\theta-u_{2} n^{-\nu}\right)}\right]^{2} d t \\
& =n \int_{-\theta-u_{2} n^{-\nu}}^{T-\theta-u_{1} n^{-\nu}}\left[\sqrt{s\left(y-\Delta n^{-\nu}\right)}-\sqrt{s(y)}\right]^{2} d y \\
= & n \int^{T-\theta-u_{1} n^{-\nu}} \frac{\left[s\left(y-\Delta n^{-\nu}\right)-s(y)\right]^{2}}{\left[\sqrt{s\left(y-\Delta n^{-\nu}\right)}+\sqrt{s(y)}\right]^{2}} d y \\
= & n I_{1}+n I_{2}
\end{aligned}
$$

where $I_{1}$ and $I_{2}$ are the integrals of the same function over the interval $(-\delta / 2, \delta / 2)$ and over the set $E=\left(-\theta-u_{2} n^{-\nu}, T-\theta-u_{1} n^{-\nu}\right) \backslash(-\delta / 2, \delta / 2)$ respectively.

On the set $E$ we have $|y| \geqslant \delta / 2$, and hence $\left|y-\Delta n^{-\nu}\right| \geqslant \delta / 4$ for $n$ sufficiently large. Recall that on the set $\{y:|y| \geqslant \delta / 4\}$ the function $s(\cdot)$ is separated from zero and Hölder continuous of order $\varkappa$. So, for $n$ sufficiently large we get

$$
n I_{2} \leqslant n \int_{E} \frac{\left|\Delta n^{-\nu}\right|^{2 \varkappa}}{c} d y \leqslant C n\left|\Delta n^{-\nu}\right|^{p+1}=C\left|u_{1}-u_{2}\right|^{p+1}
$$

Further, for the first integral we have

$$
\begin{aligned}
n I_{1} \leqslant & C n \int_{-\delta / 2}^{\delta / 2} \frac{\left[d\left(y-\Delta n^{-\nu}\right)\left|y-\Delta n^{-\nu}\right|^{p}-d(y)|y|^{p}\right]^{2}}{\left[\sqrt{s\left(y-\Delta n^{-\nu}\right)}+\sqrt{s(y)}\right]^{2}} d y \\
& +C n \int_{-\delta / 2}^{\delta / 2} \frac{\left[\psi\left(y-\Delta n^{-\nu}\right)-\psi(y)\right]^{2}}{s(y)} d y \\
\leqslant & C n \int_{-\delta / 2}^{\delta / 2} \frac{\left[d\left(y-\Delta n^{-\nu}\right)\left|y-\Delta n^{-\nu}\right|^{p}-d(y)|y|^{p}\right]^{2}}{\frac{1}{2}\left[\sqrt{d\left(y-\Delta n^{-\nu}\right)}\left|y-\Delta n^{-\nu}\right|^{p / 2}+\sqrt{d(y)}|y|^{p / 2}\right]^{2}} d y
\end{aligned}
$$

$$
\begin{aligned}
& +C n \int_{-\delta / 2}^{\delta / 2} \frac{\left|\Delta n^{-\nu}\right|^{2 \varkappa}}{\frac{1}{2} d(y)|y|^{p}} d y \\
\leqslant & C n \int_{-\delta / 2}^{\delta / 2}\left[\sqrt{d\left(y-\Delta n^{-\nu}\right)}\left|y-\Delta n^{-\nu}\right|^{p / 2}-\sqrt{d(y)}|y|^{p / 2}\right]^{2} d y \\
& +C n\left|\Delta n^{-\nu}\right|^{p+1} \int_{-\delta / 2}^{\delta / 2} \frac{1}{d(y)}|y|^{-p} d y \\
\leqslant & C n\left|\Delta n^{-\nu}\right|^{p+1} \int_{-\infty}^{\infty}\left[\widetilde{d}(z-1)|z-1|^{p / 2}-\widetilde{d}(z)|z|^{p / 2}\right]^{2} d z+C|\Delta|^{p+1} \\
= & C|\Delta|^{p+1}=C\left|u_{1}-u_{2}\right|^{p+1} .
\end{aligned}
$$

Here in the last integral we have denoted $\widetilde{d}(z)=\sqrt{d(z \Delta)}$ and noticed that the integrand behaves as $C|z|^{p-2}$ as $z \rightarrow \infty$.

So, in the case $\left|u_{1}-u_{2}\right| \leqslant 1$, for all $\theta$ and $n$ sufficiently large we get finally

$$
\mathbf{E}_{\theta}\left|Z_{n}^{1 / 2}\left(u_{1}\right)-Z_{n}^{1 / 2}\left(u_{2}\right)\right|^{2} \leqslant C n I_{1}+C n I_{2} \leqslant C\left|u_{1}-u_{2}\right|^{p+1} .
$$

Lemma 6 is proved.
It remains to verify Lemma 7. Using Lemma 1.5 of [10], for any $n, \theta$ and $u \in U_{n}$ we can write

$$
\mathbf{E}_{\theta} Z_{n}^{1 / 2}(u) \leqslant \exp \left\{-\frac{1}{2} n F\left(u n^{-\nu}\right)\right\},
$$

where for all $u \in(\alpha-\theta, \beta-\theta) \subset(-T, T)$ we denote

$$
F(u)=\int_{0}^{T}\left[\sqrt{S_{\theta+u}(t)}-\sqrt{S_{\theta}(t)}\right]^{2} d t .
$$

First we suppose $|u| \leqslant \delta / 2$. Since $\theta \in \mathbf{K} \subset(0, T)$, we have

$$
F(u)=\int_{0}^{T}[\sqrt{s(t-\theta-u)}-\sqrt{s(t-\theta)}]^{2} d t
$$

$$
=\int_{-\theta}^{T-\theta}[\sqrt{s(y-u)}-\sqrt{s(y)}]^{2} d y \geqslant \int_{-\varepsilon}^{\varepsilon}[\sqrt{s(y-u)}-\sqrt{s(y)}]^{2} d y
$$

where we can take $0<\varepsilon<\delta / 2$. Hence $|y| \leqslant \delta / 2<\delta$ and $|y-u| \leqslant \delta$, and so we get

$$
\begin{aligned}
F(u) \geqslant & \int_{-\varepsilon}^{\varepsilon} \frac{[s(y-u)-s(y)]^{2}}{[\sqrt{s(y-u)}+\sqrt{s(y)}]^{2}} d y \\
\geqslant & \int_{-\varepsilon}^{\varepsilon} \frac{\left[\left(d(y-u)|y-u|^{p}-d(y)|y|^{p}\right)+(\psi(y-u)-\psi(y)]^{2}\right.}{\left[\sqrt{2 d(y-u)}|y-u|^{p / 2}+\sqrt{2 d(y)}|y|^{p / 2}\right]^{2}} d y \\
= & c \int_{-\varepsilon}^{\varepsilon}\left[\sqrt{d(y-u)}|y-u|^{p / 2}-\sqrt{d(y)}|y|^{p / 2}\right]^{2} d y \\
& +c \int_{-\varepsilon}^{\varepsilon} \frac{[\psi(y-u)-\psi(y)]^{2}}{\left[\sqrt{d(y-u)}|y-u|^{p / 2}+\sqrt{d(y)}|y|^{p / 2}\right]^{2}} d y \\
& +c \int_{-\varepsilon}^{\varepsilon} \frac{\left(\sqrt{d(y-u)}|y-u|^{p / 2}-\sqrt{d(y)}|y|^{p / 2}\right)(\psi(y-u)-\psi(y))}{\sqrt{d(y-u)}|y-u|^{p / 2}+\sqrt{d(y)}|y|^{p / 2}} d y \\
= & I_{1}+I_{2} \pm\left|I_{3}\right| \quad
\end{aligned}
$$

with evident notations.
For the first integral we have

$$
I_{1}=C|u|^{p} \int_{-\varepsilon /|u|}^{\varepsilon /|u|}\left[\sqrt{d(u(z-1))}|z-1|^{p / 2}-\sqrt{d(u z)}|z|^{p / 2}\right]^{2} d z,
$$

and so $c|u|^{p} \leqslant I_{1} \leqslant C|u|^{p}$ since the last integral can be bounded from above and from below by the integral of the same function over $\mathbb{R}$ and over $(-\varepsilon / T, \varepsilon / T)$ respectively.

For the second integral we get

$$
I_{2} \leqslant C|u|^{2 \varkappa} \int_{-\varepsilon}^{\varepsilon} \frac{1}{\left[\sqrt{d(y)}|y|^{p / 2}\right]^{2}} d y=C|u|^{2 \varkappa}
$$

Using Cauchy-Schwarz inequality, we obtain $\left|I_{3}\right| \leqslant C \sqrt{I_{1} I_{2}} \leqslant C|u|^{\varkappa+\frac{p+1}{2}}$ for the last integral, and finally

$$
F(u) \geqslant c|u|^{p+1}-C|u|^{\varkappa+\frac{p+1}{2}}=c|u|^{p+1}\left(1-C|u|^{\varkappa-\frac{p+1}{2}}\right) \geqslant c|u|^{p+1}
$$

for $u$ sufficiently small, that is for $|u| \leqslant \sigma$, where $\sigma>0$ is some fixed constant.

On the other hand, we have also

$$
\inf _{|u| \geqslant \sigma} F(u)=c>0,
$$

since otherwise we should have $S_{\theta+u^{*}}(t)=S_{\theta}(t)$ for some fixed $u^{*}$ and almost all $t \in[0, T]$, which is impossible. Hence, for all $|u| \geqslant \sigma$ we can write

$$
F(u) \geqslant c \geqslant c \frac{|u|^{p+1}}{T^{p+1}}=c|u|^{p+1}
$$

So, for all $\theta$ and $u \in(\alpha-\theta, \beta-\theta)$ we have

$$
F(u) \geqslant c|u|^{p+1}
$$

and hence for all $n, \theta$ and $u \in U_{n}$ we can write

$$
\mathbf{E}_{\theta} Z_{n}^{1 / 2}(u) \leqslant \exp \left\{-\frac{1}{2} n F\left(u n^{-\nu}\right)\right\} \leqslant \exp \left\{-c|u|^{p+1}\right\}
$$

Lemma 7 is proved.

## 4 Concluding remarks

1. For simplicity of exposition, in this paper we considered the Bayesian estimators and the notion of asymptotic efficiency in the case of quadratic loss function. In fact, the results hold for a larger class of loss functions (see [6] for more details).
2. Again for simplicity of exposition, we considered the case where the unknown parameter $\theta$ is a shift parameter, that is $S_{\theta}(t)=s(t-\theta)$. In fact, the results hold in a more general situation, for example when the intensity function is strictly positive (except possibly in $\theta$ ) and can be written as

$$
S_{\theta}(t)=d(t-\theta)|t-\theta|^{p}+\Psi(\theta, t),
$$

where $p \in(-1,0) \cup(0,1)$, the function $d(\cdot)$ is as before, and the function $\Psi(\theta, t)$ is continuous, and uniformly in $t$ Hölder continuous (of order higher than $(p+1) / 2)$ with respect to $\theta$. In the case $p>0$ we suppose equally that $\Psi(\theta, \theta)=0$. It is not difficult to obtain for this case the same results as those presented above.
3. Like in Chapter 6 of [6], one can consider a situation when the intensity function has several singularities of the same order. More precisely, we suppose that $t_{1}<\cdots<t_{r}$ with $t_{r}-t_{1}<T$, the unknown parameter $\theta \in \Theta=(\alpha, \beta) \subseteq\left(-t_{1}, T-t_{r}\right)$, and the intensity function is strictly positive and can be written as

$$
S_{\theta}(t)=\sum_{i=1}^{r} d_{i}\left(t-\theta-t_{i}\right)\left|t-\theta-t_{i}\right|^{p}+\Psi(\theta, t)
$$

where $p \in(-1,0) \cup(0,1)$,

$$
d_{i}(x)= \begin{cases}a_{i}, & \text { if } x<0 \\ b_{i}, & \text { if } x>0\end{cases}
$$

$a_{i}, b_{i}>0$, and the function $\Psi(\theta, t)$ is continuous, and uniformly in $t$ Hölder continuous (of order higher than $(p+1) / 2$ ) with respect to $\theta$. In the case $p>0$ we suppose equally that $\Psi\left(\theta, \theta+t_{i}\right)=0$. It is not difficult to obtain for this problem the results similar to those presented above. The difference is that now one needs to introduce the process $Z$ (and hence the random variables $\zeta$ and $\xi$ ) in a slightly different manner. More precisely, for each $i=1, \ldots, n$, one should introduce a process $Z_{i}$ in the same manner (but using constants $a_{i}$ and $b_{i}$ instead of $a$ and $b$ ) as $Z$ was introduced. Further one should consider the process $Z$ defined by

$$
Z(u)=\prod_{i=1}^{r} Z_{i}(u)
$$

where the processes $Z_{i}$ are independent.

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