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Dynkin's isomorphism without symmetry

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1 Introduction:

The purpose of this note is to extend Dynkin's isomorphism involving functionals of the occupation field of a symmetric Markov processes and of the associated Gaussian field to a suitable class of non symmetric Markov processes. This was briefly proposed in [5] using Grassmann variables, extending to the non symmetric case some results of [4]. Here we propose an alternative approach, not relying on Grassmann variables that can be applied to the study of local times, in the spirit of [6]. It works in general on a finite space and on an infinite space under some assumption on the skew symmetric part of the generator which is checked on two examples.

2 The finite case

2.1 Dual processes

Let us first consider the case of an irreducible Markov process on a finite space X , with finite lifetime ζ , generator L and potential $V = (-L)^{-1}$.

Let $m = \mu V$ for any nonnegative probability μ on X . Recall that L can be written in the form $L = q(\Pi - I)$ with q positive and Π submarkovian. Then the m -adjoint \widehat{L} can be expressed similarly with the same q and a possibly different submarkovian matrix $\widehat{\Pi}$. Moreover, $m = \widehat{\mu} \widehat{V}$, with $\widehat{\mu}$ the law of $x_{\zeta-}$ under \mathbb{P}_μ .

Note that for any $z = x + iy \in \mathbb{C}^X$, the "energy" $\langle -Lz, \bar{z} \rangle_m = \sum_x -(Lz)_x \bar{z}_x m_x$ is nonnegative as it can be written $\sum_{x,y} C_{x,y} (z_x - z_y)(\bar{z}_x - \bar{z}_y) + \langle \Pi 1 + \widehat{\Pi} 1 - 2, z \bar{z} \rangle_m$, with $C_{x,y} = C_{y,x} = m_x q_x \Pi_{x,y}$. The highest eigenvector of $\frac{1}{2}(\Pi + \widehat{\Pi})$ is nonnegative by the well known argument which shows that the module contraction lowers the energy. and it follows from the strict submarkovianity that the corresponding eigenvalue is strictly smaller than 1. Hence there is a "mass

gap": For some positive ε , the "energy" $\langle -Lz, \bar{z} \rangle_m$ dominates $\varepsilon \langle z, \bar{z} \rangle_m$ for all z .

2.2 A twisted Gaussian measure

Then, although L is not symmetric, an elementary computation (given in a more general context in the following section and in a footnote p.96 of [4]) shows that for any $\chi \in \mathbb{R}_+^X$, denoting M_χ the diagonal matrix with coefficients given by χ ,

$$\frac{1}{(2\pi)^{|X|}} \int e^{-\langle z\bar{z}, \chi \rangle_m} e^{\frac{1}{2}\langle Lz, \bar{z} \rangle_m} \Pi dx_u dy_u = \det(-L + M_\chi)^{-1}$$

As a consequence, differentiating with respect to χ_x ,

$$\frac{1}{(2\pi)^{|X|}} \int z_x \bar{z}_x e^{-\langle z\bar{z}, \chi \rangle_m} e^{\frac{1}{2}\langle Lz, \bar{z} \rangle_m} \Pi dx_u dy_u = \det(-L + M_\chi)^{-1} (-L + M_\chi)_{xx}^{-1}$$

In a similar way by perturbation using a non diagonal matrix, one obtains

$$\frac{1}{(2\pi)^{|X|}} \int z_x \bar{z}_y e^{-\langle z\bar{z}, \chi \rangle_m} e^{\frac{1}{2}\langle Lz, \bar{z} \rangle_m} \Pi dx_u dy_u = \det(-L + M_\chi)^{-1} (-L + M_\chi)_{xy}^{-1}$$

But with the usual notations for Markov processes, setting $l_t^x = \int_0^{t \wedge \zeta} \frac{1}{m_{x_s}} ds$ and $l_\zeta^x = l^x$, we have

$$(-L + M_\chi)_{xy}^{-1} = \mathbb{E}_x \left(\int_0^\zeta e^{-\langle \chi, l_t \rangle_m} dl_t^y \right)$$

Defining the path measure $\mathbb{E}_{x,y}$ by: $\mathbb{E}_x(\int_0^\zeta G(x_s, s \leq t) dl_t^y) = \mathbb{E}_{x,y}(G)$ the above relation writes

$$(-L + M_\chi)_{xy}^{-1} = \mathbb{E}_{x,y}(e^{-\langle \chi, l \rangle_m})$$

It follows that for any continuous function F on \mathbb{R}_+^X

$$\begin{aligned} (*) \quad & \int z_x \bar{z}_y F(z_u \bar{z}_u, u \in X) e^{\frac{1}{2}\langle Lz, \bar{z} \rangle_m} \Pi dx_u dy_u \\ & = \int \mathbb{E}_{x,y}(F(l^u + z_u \bar{z}_u, u \in X)) e^{\frac{1}{2}\langle Lz, \bar{z} \rangle_m} \Pi dx_u dy_u \end{aligned}$$

2.3 Positivity

It should be noted that setting $\rho_u = \frac{1}{2} z_u \bar{z}_u$ and $z_x = \sqrt{\rho_x} / 2 e^{i\theta_x}$ the image on \mathbb{R}_+^X of the normalized complex measure $\nu_X = \frac{1}{(2\pi)^{|X|}} \det(-L) e^{\frac{1}{2}\langle Lz, \bar{z} \rangle_m} \Pi dx_u dy_u$ by the map $z_u \rightarrow \rho_u$ is an infinitely divisible probability distribution Q on \mathbb{R}_+^X with density $\frac{1}{(2\pi)^{|X|}} \det(-L) \int e^{\langle L\sqrt{\rho} e^{i\theta}, \sqrt{\rho} e^{-i\theta} \rangle_m} \Pi d\theta_u$. Note that the positivity

is not a priori obvious when L is not m -symmetric. This important fact follows easily by considering the moment generating function $\Phi(s) = \frac{\det(-L)}{\det(-L+M_s)} = \det(I + (-L)^{-1}M_s)^{-1}$ defined for all s with non negative coordinates, positive and analytic. The expansion in power series around any s (which appears for exemple in [8]) is explicit:

$$\begin{aligned} \frac{\Phi(s+h)}{\Phi(s)} &= \det(I+(-L+M_s)^{-1}M_h)^{-1} = \exp(-\log(\det(I+(-L+M_s)^{-1}M_h))) \\ &= \exp(\sum \frac{(-1)^k}{k} Tr([(-L+M_s)^{-1}M_h]^k)). \end{aligned}$$

As $(-L+M_s)^{-1}$ is nonnegative, it implies that Φ is completely monotone as in this last expression, all coefficients of h -monomials of order n are of the sign of $(-1)^n$.¹

Note that the same argument works for fractional powers of $\Phi(s)$ which shows the infinite divisibility. Let us incidentally mention it has been known for a long time (cf [8]) that this expansion can be simplified further in terms of permanents.

For $x = y$, formula (*) becomes:

$$(**) \int \rho_x F(\rho_u, u \in X) Q(d\rho) = \int \mathbb{E}_{x,x}(F(l^u + z_u \bar{z}_u, u \in X)) Q(d\rho)$$

Note that this last formula is also obtained in [2] after a direct definition of the measure Q .

REMARK: If $Y \subset X$, it is well known that the trace of the process on Y is a Markov process the potential of which is the restriction of V to $Y \times Y$. The distribution ν_X induces ν_Y . Therefore the formulas (*) and (**) on X and Y are consistent.

EXAMPLE: let us consider, as an example, the case where $X = \{1, 2, \dots, N\}$, $q_i = 1$, $\Pi_{i,j} = \mathbf{1}_{i < N, j=i+1}$, $\mu_i = \mathbf{1}_{i=1}$, $m_i = 1$ and $(-L)_{i,j}^{-1} = \mathbf{1}_{i \leq j}$.

The characteristic polynomial of $\frac{1}{2}(\Pi + \widehat{\Pi})$ is $(-\lambda + \sqrt{\lambda^2 - 1})^N + (-\lambda - \sqrt{\lambda^2 - 1})^N$ hence one gets easily that the mass gap equals $2 \sin^2(\frac{\pi}{2N})$

Under $\mathbb{E}_{x,x}$ all local times vanish except l^x which, follows an exponential distribution. Moreover an easy calculation shows that Q reduces to a product of exponential distribution. The formula (**) reduces to the convolution of two exponentials.

¹Characterisation of completely monotone functions in several variables are related to the moment problem treated in [3]

3 The infinite case

We now explain how in certain situations, the above can be extended to a Markov process on an infinite space X . Of course, a Markov process for which points are not polar can always be viewed elementarily as a consistent system of processes on finite subspaces but we aim at a stronger representation allowing to consider any functional of the occupation field. There are some obvious obstructions to a generalisation. The mass gap property does not always hold: consider for example the case of a constant drift on an interval, analogous to our first exemple. Some conditions have to be assumed in order that the energy controls the antisymmetric part of the generator.

3.1 Some calculations in Gaussian space

Let H be a real Hilbert space with scalar product \langle, \rangle . At first the reader may suppose it finite dimensional and then check that the assumptions we will make allow to extend the results to the infinite dimensional case.

Let ϕ be the canonical Gaussian field indexed by H . Given any ONB e_k of H , $w_k = \phi(e_k)$ are independent normal variables. Recall that for all $f \in H$, $\phi(f) = \sum_k \langle f, e_k \rangle w_k$ and $E(e^{i\phi(f)}) = e^{-\frac{\|f\|^2}{2}}$.

In the following, $\phi(f)$ can be denoted by $\langle \phi, f \rangle$ though of course ϕ does not belong to H .

Let K be any Hilbert-Schmidt operator on H . Note that $K\phi = \sum_k w_k K e_k$ is well defined as a H -valued random variable, and that $E(\|K\phi\|^2) = \text{Tr}(KK^*)$.

Let C be a symmetric non negative trace-class linear operator on H . Recall that the positive integrable random variable $\langle C\phi, \phi \rangle \in \mathbb{L}^1$ can be defined by $\sum_k \langle C e_k, e_k \rangle w_k^2$ for any ONB diagonalizing C , and by $\sum_{k,l} \langle C e_k, e_l \rangle w_k w_l$ for any ONB.

Moreover, $E(e^{-\frac{1}{2}\langle C\phi, \phi \rangle + i\phi(f)}) = \det(I + C)^{-\frac{1}{2}} e^{-\frac{\langle (I+C)^{-1} f, f \rangle}{2}}$ (the determinant can be defined as $\prod(1 + \lambda_i)$, where the λ_i are the eigenvalues of C).

In fact $\det(I + T)$ is well defined for any trace class operator T as $1 + \sum_{n=1}^{\infty} \text{Tr}(T^n)$. (Cf [7] Chapter 3). It extends continuously the determinant defined with finite ranks operators and it verifies the identity:

$\det(I + T_1 + T_2 + T_1 T_2) = \det(I + T_1) \det(I + T_2)$. By Lidskii's theorem, it is also given by the product $\prod(1 + \lambda_i)$ defined by the eigenvalues of the trace class (hence compact) operator T , counted with their algebraic multiplicity.

Let ϕ_1 and ϕ_2 be two independent copies of the canonical Gaussian process indexed by H . Let B be a skew-symmetric Hilbert-Schmidt operator on H . Note that $\langle B\phi_1, \phi_2 \rangle = \sum \langle B e_k, e_l \rangle w_k^1 w_l^2 = -\langle B\phi_2, \phi_1 \rangle$ is well defined in \mathbb{L}^2 and that $E(e^{i\langle B\phi_1, \phi_2 \rangle}) = E(e^{-\frac{1}{2}\|B\phi_1\|^2}) = \det(I + BB^*)^{-\frac{1}{2}}$.

As B is Hilbert-Schmidt, BB^* is trace-class. The renormalized determinant $\det_2(I+B) = \det((I+B)e^{-B})$ is well defined (Cf [7]), and as the eigenvalues of B are purely imaginary and pairwise conjugated, it is strictly positive. Moreover since $B^* = -B$, $\det_2(I+B) = \det_2(I-B) = \det(I+BB^*)^{\frac{1}{2}}$

Finally, it comes that

$$E(e^{i\langle B\phi_1, \phi_2 \rangle}) = \det_2(I+B)^{-1}$$

More generally, setting $\psi = \phi_1 + i\phi_2$,

$$E(e^{-\frac{1}{2}\langle (C-B)\psi, \bar{\psi} \rangle}) = \det_2(I+C+B)^{-1} \exp(-\text{Tr}(C))$$

(Recall that when T is trace class, $\det_2(I+T) \exp(\text{Tr}(T)) = \det(I+T)$).

$$\begin{aligned} \text{Indeed, } E(e^{-\frac{1}{2}\langle (C-B)\psi, \bar{\psi} \rangle}) &= E(e^{-\frac{1}{2}\langle C\phi_1, \phi_1 \rangle - \frac{1}{2}\langle C\phi_2, \phi_2 \rangle + i\langle B\phi_1, \phi_2 \rangle}) \\ &= \det(I+C)^{-\frac{1}{2}} E(e^{-\frac{1}{2}\langle (I+C)^{-1}B\phi_1, B\phi_1 \rangle + \langle C\phi_1, \phi_1 \rangle}) \text{ (by integration in } \phi_2) \\ &= \det(I+C)^{-\frac{1}{2}} \det(I+C - B(I+C)^{-1}B)^{-\frac{1}{2}} \\ &= \det(I - (I+C)^{-\frac{1}{2}}B(I+C)^{-1}B(I+C)^{-\frac{1}{2}})^{-\frac{1}{2}} \det(I+C)^{-1} \\ &= \det((I+(I+C)^{-\frac{1}{2}}B(I+C)^{-\frac{1}{2}})(I - (I+C)^{-\frac{1}{2}}B(I+C)^{-\frac{1}{2}}))^{-\frac{1}{2}} \det(I+C)^{-1} \\ &= \det_2(I+(I+C)^{-\frac{1}{2}}B(I+C)^{-\frac{1}{2}})^{-1} \det(I+C)^{-1} \text{ (since } \det_2(I+(I+C)^{-\frac{1}{2}}B(I+C)^{-\frac{1}{2}}) \\ &= \det_2(I - (I+C)^{-\frac{1}{2}}B(I+C)^{-\frac{1}{2}}) \text{ by skew symmetry as before)} \\ &= \det_2(I+C+B)^{-1} \exp(-\text{Tr}(C)). \end{aligned}$$

Note that $I+C+B$ is always invertible, as $C+B$ is a compact operator and -1 is not an eigenvalue.

Let f_1 and f_2 be two elements of H . Set $D(f) = \langle f, f_1 \rangle f_2$. For small enough ε ,

$$\begin{aligned} E(e^{-\frac{1}{2}\langle (C-B)\psi, \bar{\psi} \rangle + \varepsilon\psi(f_1)\bar{\psi}(f_2)}) &= E(e^{-\frac{1}{2}\langle C\phi_1, \phi_1 \rangle - \frac{1}{2}\langle C\phi_2, \phi_2 \rangle + i\langle B\phi_1, \phi_2 \rangle + \varepsilon(\langle \phi_1(f_1), \phi_1(f_2) \rangle + \langle \phi_2(f_1), \phi_2(f_2) \rangle)}) \\ &= \det_2(I+C+B+\varepsilon D)^{-1} \exp(-\text{Tr}(C+\varepsilon D)) \\ &= \det_2(I+C+B)^{-1} \exp(-\text{Tr}(C)) \det(I+\varepsilon(I+C+B)^{-1}D)^{-1}. \end{aligned}$$

Hence, differentiating both members at $\varepsilon = 0$, $E(\psi(f_1)\bar{\psi}(f_2)e^{-\frac{1}{2}\langle (C-B)\psi, \bar{\psi} \rangle}) = \det_2(I+C+B)^{-1} \exp(-\text{Tr}(C)) \text{Tr}((I+C+B)^{-1}D)$. Therefore

$$\frac{E(\psi(f_1)\bar{\psi}(f_2)e^{-\frac{1}{2}\langle (C-B)\psi, \bar{\psi} \rangle})}{E(e^{-\frac{1}{2}\langle (C-B)\psi, \bar{\psi} \rangle})} = \langle (I+C+B)^{-1}(f_1), f_2 \rangle .$$

If C is only Hilbert-Schmidt, we can consider only the renormalized "Wick product" : $\langle C\phi, \phi \rangle := \sum_k \langle Ce_k, e_k \rangle (w_k^2 - 1)$ for any ONB diagonalizing C and

$$E(e^{-\frac{1}{2}\langle C\phi, \phi \rangle + i\phi(f)}) = \det_2(I+C)^{-\frac{1}{2}} e^{-\frac{\langle (I+C)^{-1}f, f \rangle}{2}}$$

The results given above extend immediatly as follows:

$$E(e^{-\frac{1}{2}\langle (C-B)\psi, \bar{\psi} \rangle}) = \det_2(I+C+B)^{-1}$$

$$\frac{E(\psi(f_1)\bar{\psi}(f_2)e^{-\frac{1}{2}(\langle(C-B)\psi,\bar{\psi}\rangle)})}{E(e^{-\frac{1}{2}(\langle(C-B)\psi,\bar{\psi}\rangle)})} = \langle (I + C + B)^{-1}(f_1), f_2 \rangle .$$

3.2 A class of Markov processes in duality

Let $(V_\alpha, \alpha \geq 0)$ and \widehat{V}_α be two Markovian or submarkovian resolvents in duality in a space $\mathbb{L}^2(X, \mathcal{B}, m)$ with generators L and \widehat{L} , such that, denoting $\mathcal{D} = \mathcal{D}(L) \cap \mathcal{D}(\widehat{L})$:

1) \mathcal{D} and $L(\mathcal{D})$ are dense in $\mathbb{L}^2(m)$

2) $\langle -Lf, f \rangle_m \geq \varepsilon \langle f, f \rangle_m$ for some $\varepsilon > 0$ and any $f \in \mathcal{D}$ (i.e. we assume the existence of a spectral gap: it can always be obtained by adding a negative constant to L).

Let H be the completion of \mathcal{D} with respect to the energy norm. It is a functional space imbedded in $\mathbb{L}^2(X, \mathcal{B}, m)$. Let A be the associated self adjoint generator so that $H = \mathcal{D}(\sqrt{-A})$. On \mathcal{D} , $\frac{1}{2}(L + \widehat{L}) = A$.

The final assumption is crucial to allow the control of the antisymmetric part:

3) $B = (-A)^{-1} \frac{L - \widehat{L}}{2}$ extends in a Hilbert-Schmidt operator on H .

Then, $(-A)^{-\frac{1}{2}} \frac{L - \widehat{L}}{2} (-A)^{-\frac{1}{2}}$ is an antisymmetric Hilbert Schmidt operator on $\mathbb{L}^2(m)$ since for any ONB e_k of H , $(-A)^{\frac{1}{2}} e_k$ is an ONB of $\mathbb{L}^2(m)$. Note that $I - B$ is bounded and invertible on H and that $V_0 = (I - B)^{-1} (-A)^{-1}$ maps $\mathbb{L}^2(m)$ into H . Indeed, one can see first that on \mathcal{D} , $A(I - B) = L$ so that on $L(\mathcal{D})$, $V_0 = (I - B)^{-1} (-A)^{-1}$

EXAMPLES

This applies to the case of the finite space considered above.

Let us mention other examples:

1) Diffusion with drift on the circle: $X = S^1$, $A = \frac{\partial^2}{\partial \theta^2} - \varepsilon$, $L - \widehat{L} = b(\theta) \frac{\partial}{\partial \theta}$, where b is a bounded function on S^1 .

Indeed, considering the orthonormal basis $e^{ik\theta}$ in $\mathbb{L}^2(d\theta)$, $\frac{1}{\sqrt{k^2 + \varepsilon}} e^{ik\theta}$ is an orthonormal basis in $H = H^1$, and

$$\sum_k \left\| (-A)^{-\frac{1}{2}} b(\theta) \frac{\partial}{\partial \theta} \frac{1}{\sqrt{k^2 + \varepsilon}} e^{ik\theta} \right\|_{\mathbb{L}^2(d\theta)}^2 = \sum_{k,l} \frac{k^2}{k^2 + \varepsilon} (\widehat{b}(l - k))^2 \frac{1}{l^2 + \varepsilon} < \infty$$

2) Levy processes on the circle: The Fourier coefficients $a_k + ib_k$ of L should verify $\sum_k \left(\frac{b_k}{a_k}\right)^2 < \infty$

3.3 An extension of Dynkin's isomorphism

Assume X is locally compact and separable, and that functions of H are continuous. By Banach-Steinhaus theorem, given any point $x \in X$, there exists an

element of H , denoted η_x defined by the identity: $f(x) = \langle \eta_x, f \rangle_H$. Note that $\eta_x = \sum e_k(x)e_k$ for all ONB of H .

The resolvent V_λ is Fellerian and induces a strong Markov process. Denote by l_t^x the local time at x of this Markov process. Let x and y be two points of X . Set: $\langle \eta_x, \eta_y \rangle_H = \sum e_k(x)e_k(y) = K(x, y)$ so that $A^{-1}f(x) = \int K(x, y)f(y)m(dy)$ or $A^{-1}f = \int \eta_y f(y)m(dy)$.

Set $V_0(x, y) = \langle (I + B)^{-1}\eta_x, \eta_y \rangle_H$ and note that $V_0(x, y)$ is a kernel for V_0 . Indeed, for any $f, g \in \mathbb{L}^2(m)$, $\langle V_0 f, g \rangle_{\mathbb{L}^2(m)} = \langle (I - B)^{-1}(-A)^{-1}f, g \rangle_{\mathbb{L}^2(m)} = \langle (I - B)^{-1}(-A)^{-1}f, A^{-1}g \rangle_H = \int f(x)g(y)V_0(x, y)m(dx)m(dy)$

As a consequence, $V_0(x, y) = \mathbb{E}_x(l_t^y)$

Applying the construction of the section 3-1, we see the kernel $K(x, y)$ is the covariance of a Gaussian process $(Z_x = \psi(\eta_x) = \sum e_k(x)w_k, x \in X)$, for any ONB e_k of H .

More generally, for any non-negative finitely supported measure $\chi = \sum_1^N p_j \delta_{u_j}$ on X , letting C be the finite rank operator: $C = \sum_1^N p_j \eta_{u_j} \otimes \eta_{u_j}$, set $V_\chi(x, y) = \langle (I - B + C)^{-1}\eta_x, \eta_y \rangle_H$

In a similar way as above for V_0 , we have $V_\chi(x, y) = \mathbb{E}_x(e^{-\int l_t^z \chi(dz)} dl_t^y)$.

Then, from section 3-1

$$\begin{aligned} E(Z_x \bar{Z}_y e^{-\frac{1}{2} \langle (C-B)\psi, \bar{\psi} \rangle_H}) &= E(e^{-\frac{1}{2} \langle (C-B)\psi, \bar{\psi} \rangle_H}) \langle (I - B + C)^{-1}\eta_x, \eta_y \rangle_H \\ &= E(e^{-\frac{1}{2} \langle (C-B)\psi, \bar{\psi} \rangle_H}) V_\chi(x, y) \end{aligned}$$

But $\langle C\psi, \bar{\psi} \rangle_H = \int \psi(\eta_u) \bar{\psi}(\eta_u) \chi(du) = \int Z_u \bar{Z}_u \chi(du)$ and $\langle B\psi, \bar{\psi} \rangle_H = \langle (-A)^{-1} \frac{L-\hat{L}}{2} \psi, \bar{\psi} \rangle_H = \sum \langle (-A)^{-1} \frac{L-\hat{L}}{2} e_k, e_l \rangle_H w_k^1 w_l^2$.

On the other hand, at least formally in general but exactly in the finite dimensional case,

$$\int \frac{L-\hat{L}}{2} Z_u \bar{Z}_u m(du) = \langle (-A)^{-1} \frac{L-\hat{L}}{2} Z, \bar{Z} \rangle_H = \sum w_k^1 w_l^2 \langle (-A)^{-1} \frac{L-\hat{L}}{2} e_k, e_l \rangle_H$$

Hence we can denote: $\langle B\psi, \bar{\psi} \rangle_H$ by $\langle \frac{L-\hat{L}}{2} Z, \bar{Z} \rangle_{\mathbb{L}^2(m)}$. Therefore

$$\begin{aligned} E(Z_x \bar{Z}_y e^{\frac{1}{2} \langle \frac{L-\hat{L}}{2} Z, \bar{Z} \rangle_{\mathbb{L}^2(m)} - \int Z_u \bar{Z}_u \chi(du)}) \\ &= E \otimes \mathbb{E}_x \left(\int_0^\zeta e^{\frac{1}{2} \langle \frac{L-\hat{L}}{2} Z, \bar{Z} \rangle_{\mathbb{L}^2(m)}} e^{-\int l_t^z \chi(dz) - \int Z_u \bar{Z}_u \chi(du)} dl_t^y \right) \\ &= E(\mathbb{E}_{x,y}(e^{-\int l_t^z \chi(dz)}) e^{\frac{1}{2} \langle \frac{L-\hat{L}}{2} Z, \bar{Z} \rangle_{\mathbb{L}^2(m)} - \int Z_u \bar{Z}_u \chi(du)}) \end{aligned}$$

Finally, we get that for any continuous bounded function F of N non negative real coordinates, and any N -uple of points u_j in X ,

$$E(Z_x \bar{Z}_y e^{\frac{1}{2} \langle \frac{L-\hat{L}}{2} Z, \bar{Z} \rangle_{\mathbb{L}^2(m)}} F(Z_{u_j} \bar{Z}_{u_j})) = E \otimes \mathbb{E}_{x,y}(e^{\frac{1}{2} \langle \frac{L-\hat{L}}{2} Z, \bar{Z} \rangle_{\mathbb{L}^2(m)}} F(l_t^{u_j} + Z_{u_j} \bar{Z}_{u_j}))$$

and the formula finally extends to any bounded measurable function of a real field on X :

$$(*bis) E(Z_x \bar{Z}_y e^{\frac{1}{2} \langle \frac{L-\hat{L}}{2} Z, \bar{Z} \rangle_{\mathbb{L}^2(m)}} F(Z \bar{Z})) = E \otimes \mathbb{E}_{x,y}(e^{\frac{1}{2} \langle \frac{L-\hat{L}}{2} Z, \bar{Z} \rangle_{\mathbb{L}^2(m)}} F(l_t + Z \bar{Z}))$$

It induces the formula (*) obtained in the finite case if we consider the trace of the process on any finite subset. Then it follows that the restriction of the twisted Gaussian measure $e^{\frac{1}{2}\langle \frac{L-\tilde{L}}{2} Z, \bar{Z} \rangle_{L^2(m)}} P(dZ)$ to $\sigma(Z_u \bar{Z}_u, u \in X)$ is a probability measure Q under which the distribution of the process $(Z_u \bar{Z}_u, u \in X)$ is infinitely divisible. It is clear that that formula (**) extends in the same way. Moreover, the important fact is that this probability is absolutely continuous with respect to the restriction of P to $\sigma(Z_u \bar{Z}_u, u \in X)$.

Hence it follows for example that the continuity Gaussian field Z implies the continuity of the local time field l under all loop measures P_{xx} .

Note finally that these results, as in the symmetric case, can be extended to some situations where the local time does not exist (like the two dimensional Brownian motion), by considering the occupation field l^x and the "Wick square" $: Z_x \bar{Z}_x :$ (formally given by $Z_x \bar{Z}_x - K(x, x)$) as generalized random fields. This makes sense for the Wick square provided K is a Hilbert-Schmidt operator.

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