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# NONCOMMMUTATIVE BATALIN-VILKOVISKY GEOMETRY AND MATRIX INTEGRALS 

SERGUEI BARANNIKOV


#### Abstract

We describe matrix model associated naturally with any solution to quantum master equation of noncommutative Batalin-Vilkovisky geometry. In the simplest example we get the Kontsevich model of 2-dimensional gravity.


## 1. Introduction.

We describe matrix model associated naturally with solution to quantum master equation of noncommutative Batalin-Vilkovisky geometry. The simplest example is the Kontsevich model of 2-dimensional gravity.

Notations. We work in the tensor category of super vector spaces, over an algebraically closed field $k, \operatorname{char}(k)=0$. Let $V=V^{1 \mid 0} \oplus V^{0 \mid 1}$ be a super vector space. We denote by $\bar{\alpha}$ the parity of an element $\alpha$ and by $\Pi V$ the super vector space with inversed parity. Sometimes we may also assume that $V$ has a differential, we leave to the interested reader to make the appropriate modifications, if they are not mentioned explicitely. For a finite group $G$ acting on a vector space $U$, we denote via $U^{G}$ the space of invariants with respect to the action of $G$.

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## 2. Noncommutative Batalin-Vilkovisky geometry.

2.1. Even inner products. Let $B: V^{\otimes 2} \rightarrow k$ be an even symmetric inner product
on $V$ :

$$
B(x \otimes y)=(-1)^{\bar{x}+\bar{y}} B(y \otimes x)
$$

The inner prodcut $B$ corresponds to an even symplectic structure on $\Pi V$. We introduced in [B1] the space $F=\bigoplus_{n} F_{n}$

$$
\begin{equation*}
F_{n}=\left((\Pi V)^{\otimes n} \otimes k\left[\mathbb{S}_{n}\right]^{\prime}\right)^{\mathbb{S}_{n}} \tag{2.1}
\end{equation*}
$$

where $k\left[\mathbb{S}_{n}\right]^{\prime}$ denotes the super $k$-vector space with the basis indexed by elements $\left(\sigma, \rho_{\sigma}\right)$, where $\sigma \in \mathbb{S}_{n}$ is a permutation with $i_{\sigma}$ cycles $\sigma_{\alpha}$ and $\rho_{\sigma}=\sigma_{1} \wedge \ldots \wedge \sigma_{i_{\sigma}}$, $\rho_{\sigma} \in \operatorname{Det}(\operatorname{Cycle}(\sigma)), \operatorname{Det}(\operatorname{Cycle}(\sigma))=\operatorname{Symm}^{i_{\sigma}}\left(k^{0 \mid i_{\sigma}}\right)$, is one of the generators of

[^0]the one-dimensional determinant of the set of cycles of $\sigma$, i.e. $\rho_{\sigma}$ is an order on the set of cycles defined up to even reordering, and $\left(\sigma,-\rho_{\sigma}\right)=-\left(\sigma, \rho_{\sigma}\right)$. The group $\mathbb{S}_{n}$ acts on $k\left[\mathbb{S}_{n}\right]^{\prime}$ by conjugation. The space $F$ carries the naturally defined operator
$\Delta\left(a_{1} \otimes \ldots \otimes a_{n} \otimes\left(\sigma, \rho_{\sigma}\right)\right)=B\left(a_{n-1} \otimes a_{n}\right)\left(a_{1} \otimes \ldots \otimes a_{n-2} \otimes\left(\pi_{n-1, n}\left(s_{n-1, n} \sigma\right), \rho_{\delta}\right)\right.$
where $s_{n-1, n}$ is the transposition interchanging $n-1$ and $n, \pi_{n-1, n}$ is the operator of erasing of elements $n-1$ and $n$ from the cycles of the representation, and $\rho_{\delta}$ is the natural ordering on the cycles of resulting permutation $\pi_{n-1, n}\left(\sigma s_{n-1, n}\right)$, see section 9 from loc.cit. The space $F$ can be identified with the symmetric power of the parity transposed space of cyclic invariant tensors on $\Pi V$
$$
F=\operatorname{Symm}\left(\oplus_{j} \Pi\left(\Pi V^{\otimes j}\right)^{\mathbb{Z} / j \mathbb{Z}}\right)
$$

If one chooses a basis $\left\{a_{i}\right\}$ in $\Pi V$, in which $B$ has the form $B\left(a_{i}, a_{j}\right)=b_{i j}$ then operator $\Delta$ can be written as

$$
\Delta=\sum_{i j} \gamma b_{i j} \frac{\partial^{2}}{\partial a_{i} \partial a_{j}}
$$

where $\gamma$ denotes the "dissection-gluing" operator acting on cycles, which corresponds to the to the symmetrisation of the operator $\left(\sigma, \rho_{\sigma}\right) \rightarrow\left(\pi_{n-1, n}\left(s_{n-1, n} \sigma\right), \rho_{\delta}\right)$ acting on pairs $\left(\sigma, \rho_{\sigma}\right)$. The operator $\gamma$ is odd since it changes the total number of cycles by one.

The operator $\Delta$ is of the second order with respect to the multiplication on $F$ and by the general formula of Batalin-Vilkovisky geometry it defines the odd Lie bracket on $F$, which we described in loc.cit.,section 5 , and which can be written as

$$
\{x, y\}=\sum_{i j} \gamma b_{i j} \frac{\partial x}{\partial a_{i}} \frac{\partial x}{\partial a_{j}}
$$

Let us denote by $\widetilde{\mathbb{S}}$ the $\mathbb{Z} / 2 \mathbb{Z}$-graded version of the twisted modular $\mathcal{K}$-operad $\tilde{\mathfrak{s}} \Sigma \mathbb{S}[t]$. By (a $\mathbb{Z} / 2 \mathbb{Z}$-graded version) of the theorem 1 from [B1] the solutions of quantum master equation in $F$

$$
d_{V} S+\hbar \Delta S+\frac{1}{2}\{S, S\}=0
$$

$S=\sum_{2(g-1)+n>0} \hbar^{g} S_{g, n}$, are in one-to one correspondence with the structure of algebra over the $\mathbb{Z} / 2 \mathbb{Z}$-graded modular operad $\mathcal{F}_{\mathcal{K}} \widetilde{\mathbb{S}}$ on the vector space $V$ with symetric inner product $B$. Its genus zero part corresponds to the structure of $A_{\infty}$-algebra with the even invariant inner product on $V$.
2.2. Odd inner products. Let $V=V^{1 \mid 0} \oplus V^{0 \mid 1}$ be a super vector space and $B: V^{\otimes 2} \rightarrow \Pi k$ be an odd symmetric inner product, $B(x \otimes y)=(-1)^{\bar{x}+\bar{y}} B(y \otimes x)$. By definition $B$ is an odd symplectic structure on $V$. The space $F$ in this situation has components

$$
\begin{equation*}
F_{n}=\left(V^{\otimes n} \otimes k\left[\mathbb{S}_{n}\right]\right)^{\mathbb{S}_{n}} \tag{2.2}
\end{equation*}
$$

where $k\left[\mathbb{S}_{n}\right]$ is the $k$-vector space with the basis indexed by $\sigma \in \mathbb{S}_{n}$ and $\mathbb{S}_{n}$ acts on $k\left[\mathbb{S}_{n}\right]$ by conjugation. The space $F$ carries the naturally defined operator
$\Delta\left(a_{1} \otimes \ldots \otimes a_{n} \otimes \sigma=(-1)^{\sum_{i=1}^{n-2} \bar{a}_{i}} B\left(a_{n-1} \otimes a_{n}\right)\left(a_{1} \otimes \ldots \otimes a_{n-2} \otimes\left(\pi_{n-1, n}\left(s_{n-1, n} \sigma\right)\right)\right.\right.$
where $s_{n-1, n}$ is the transposition interchanging $n-1$ and $n, \pi_{n-1, n}$ is the operator of erasing of elements $n-1$ and $n$ from the cycles of the representation, see section

9 from loc.cit. The space $F$ can be identified with the symmetric power of the space of cyclic invariant tensors on $V$

$$
F=\operatorname{Symm}\left(\oplus_{i}\left(V^{\otimes i}\right)^{\mathbb{Z} / i \mathbb{Z}}\right)
$$

If one chooses a basis $\left\{a_{i}\right\}$ in $V$, in which $B$ has the form $B\left(a_{i}, a_{j}\right)=b_{i j}$ then operator $\Delta$ can be written as

$$
\Delta=\sum_{i j} \gamma b_{i j} \frac{\partial^{2}}{\partial a_{i} \partial a_{j}}
$$

where $\gamma$ denotes the "dissection-gluing" operator acting on cycles, which corresponds to the symmetrisation of the operator $\sigma \rightarrow \pi_{n-1, n}\left(s_{n-1, n} \sigma\right)$ acting on permutations. The operator $\gamma$ is even in this case since in the definition of $F$ the parity change on the space of cyclic invariant tensors is not involved.

The operator $\Delta$ is of the second order with respect to the multiplication on $F$ and it defines the odd Lie bracket on $F$, which can be written as

$$
\{x, y\}=\sum_{i j} \gamma b_{i j} \frac{\partial x}{\partial a_{i}} \frac{\partial x}{\partial a_{j}}
$$

The odd analogue corresponds to the modular operad, which is natural to denote $\mathbb{S}$, the untwisted analogue of $\mathbb{Z} / 2 \mathbb{Z}$-graded version of $\widetilde{\mathfrak{s}} \Sigma \mathbb{S}[t]$. The spaces $\mathbb{S}((n))$ are simply the group algebras $k\left[\mathbb{S}_{n}\right]$ with the action of $\mathbb{S}_{n}$ by conjugation. The composition maps are defined by operators $\pi_{f, f^{\prime}}\left(s_{f, f^{\prime}} \sigma\right)$ as in section 9 of [B1]. The Feynman transform $\mathcal{F} \mathbb{S}[t]$ is a $\mathcal{K}$-twisted modular operad. The solutions of quantum master equation in $F$ are in one-to one correspondence with the structure of algebra over the modular operad $\mathcal{F S}[t]$ on the vector space $V$ with odd symmetric inner product $B$. The genus zero part of $\mathcal{F} \mathbb{S}[t]$ action on $V$ corresponds to the structure of $A_{\infty}$-algebra with the odd invariant inner product on $V$. It is straightforward generalisation of the theorem 2 from loc.cit. that the complexes $\mathcal{F} \mathbb{S}[t]((0, \gamma, \nu)))$ compute the cohomology of the compactified moduli spaces with coefficients in the local system $H^{i}\left(\overline{\mathcal{M}}_{\gamma, \nu}^{\prime} / \mathbb{S}_{\nu}, \mathcal{L}\right), \mathcal{L}=\operatorname{Det}\left(P_{\Sigma}\right)$, where $P_{\Sigma}$ is the set of marked points on the Riemann surface $\Sigma$.

## 3. Matrix integrals.

We showed in [B1] that solutions to quantum master equation corresponding to a twisted modular operad $\mathcal{P}$ define natural cocycles in the graph complex which is the Feynmann transform of $\mathcal{P}$. In this section we show that for $\widetilde{\mathbb{S}}$ and $\mathbb{S}$ these cocycles can be naturally interpreted as the terms of Feynman diagram expansions of matrix integrals of $\exp (S / \hbar)$.
3.1. Odd inner product. Let us assume that

$$
S\left(a_{i}, \hbar\right)=\sum_{2(g-1)+n>0} \hbar^{g} S_{n, g}\left(a_{i}\right)
$$

$S_{n, g} \in F_{n}$, be a solution to quantum master equation in the space (2.2). Let us consider the super vector space

$$
M=\operatorname{Hom}\left(V, \operatorname{End}_{k}(U)\right)
$$

where $U$ is a super vector space $\operatorname{dim} U<\infty$. If $V=k^{(d \mid f)}$, then $M=\left(E n d_{k}(U)\right)^{\times d} \times$ $\left(\Pi E n d_{k}(U)\right)^{\times f}$. The supertrace gives a natural even symmetric inner product on the space $\operatorname{End}_{k}(U)$ :

$$
\operatorname{Tr}(X \otimes Y)=\Sigma_{\alpha \beta}(-1)^{\bar{\alpha}} X_{\beta}^{\alpha} Y_{\alpha}^{\beta}
$$

$X, Y \in E n d_{k}(U)$, it is invariant with respect to the action of the group $G L(U)$. It gives an extension to $M$ of the odd symmetric inner product on $\operatorname{Hom}(V, k)$ dual to $B$. In coordinates it is given by

$$
B^{-1}(X, Y)=\sum_{i, j, \alpha, \beta}\left(B^{-1}\right)^{i j}(-1)^{\bar{\alpha}} X_{i, \beta}^{\alpha} Y_{j, \alpha}^{\beta}
$$

Let us extend $S$ to a function $S_{U}$ on $M$ in the obvious way, so that instead of each cyclically symmetric tensor we put the supertrace of the product of the corresponding matrices from $E n d_{k}(U)$

$$
\sum a_{i_{1}} \otimes \ldots \otimes a_{i_{k}} \rightarrow \operatorname{Tr}\left(X_{i_{1}} \cdot \ldots \cdot X_{i_{k}}\right)
$$

and the product of cyclic words goes to the product of traces.
We claim that the tensors which we associated with $S\left(a_{i}, \hbar\right)$ via the statistical sum model construction to a graph $G$, equipped with the marking corresponding to basis elements from $\mathbb{M}^{\text {dual }}((G))$, are obtained naturally from the Feynman diagrams corresponding to the asympthotic expansion of the following integral

$$
\log \int \exp \frac{1}{\hbar}\left(B^{-1}(X, \Xi)-\frac{1}{2} B^{-1}(X, X)+S_{U}(X)\right) d X /(2 \pi \hbar)^{\frac{N}{2}}
$$

(compare with the formula (0.1) from [GK]). This is immediate from the standard rules of the Feynman diagrammatics, in particular the combinatorics of the terms in $S_{U}(X)$ reproduce exactly the data associated with vertices in the complex $\mathbb{M S}^{\text {dual }}((G))$, i.e. symmetric product of cyclic permutations and an integer number.
3.2. Even inner product. Let now $S\left(a_{i}, \hbar\right)=\sum_{2(g-1)+n>0} \hbar^{g} S_{n, g}\left(a_{i}\right)$ is a solution to quantum master equation in the Batalin-Vilkovisky algebra 2.1. In this case the space $M$ is the super vector space

$$
M=\operatorname{Hom}_{k}\left(\Pi V, \Pi \operatorname{End}_{k}(U) \oplus \operatorname{End}_{k}(U)\right)
$$

where $U$ is a purely even vector space. The analogue of trace defines on the space $\Pi E n d_{k}(U) \oplus \operatorname{End}_{k}(U)$ an odd symplectic structure, which after choosing a basis $\left\{e_{\alpha}\right\}$ in $U$ can be written as

$$
\operatorname{Tr}(X \otimes Y)=\sum_{\alpha \beta} X_{\beta}^{\alpha} Y_{\alpha}^{\beta}
$$

$X \in \Pi \operatorname{End}_{k}(U), Y \in \operatorname{End}_{k}(U)$. Notice that there are natural isomorphisms

$$
\Pi E n d_{k}(U)=\operatorname{Hom}_{k}(\Pi U, U)=\operatorname{Hom}_{k}(U, \Pi U)
$$

and also

$$
\operatorname{Hom}_{k}(\Pi U, \Pi U)=\operatorname{End}_{k}(U)
$$

Therefore elements of the space $\Pi E n d_{k}(U) \oplus E n d_{k}(U)$ can be naturally composed so that this space is in fact an associative algebra with odd invariant inner product.

We have an extension to $M$ of the even symmetric inner product on $\operatorname{Hom}(V, k)$ dual to $B$ :

$$
B^{-1}(X, Y)=\sum_{i, j, \alpha, \beta}\left(B^{-1}\right)^{i j}(-1)^{\bar{\alpha}} X_{i, \beta}^{\alpha} Y_{j, \alpha}^{\beta}
$$

$X, Y \in \Pi_{E n d}^{k}(U) \oplus \operatorname{End}_{k}(U)$.
We extend as above the solution $S$ to a function $S_{U}$ on $M$ in the obvious way, so that instead of each cyclically symmetric tensor in $(\Pi V)^{\otimes k}$ we put the trace of the product of the corresponding elements from $\left(\Pi E n d_{k}(U) \oplus \operatorname{End}_{k}(U)\right)$

$$
\sum a_{i_{1}} \otimes \ldots \otimes a_{i_{k}} \rightarrow \operatorname{Tr}\left(X_{i_{1}} \cdot \ldots \cdot X_{i_{k}}\right)
$$

and the product of cyclic words goes to the product of traces.
Then the standard rules of the Feynman diagrammatics imply immediately that the asymptotic expansion of the integral

$$
\log \int \exp \frac{1}{\hbar}\left(B^{-1}(X, \Xi)-\frac{1}{2} B^{-1}(X, X)+S_{U}(X)\right) d X /(2 \pi \hbar)^{\frac{N}{2}}
$$

$X, \Xi \in M$, give the statistical sum tensors on graphs which were associated with the solution $S$ in [B1].
Example 1. Let us consider the case of even one-dimensional space $V$ with even symmetric inner product. It is easy to see that in this case any linear combination of cyclic words $X^{3}, X^{5}, \ldots, X^{2 n+1}$ is a solution of quantum master equation. The corresponding matrix integral reproduces the matrix integrals of $[\mathrm{K}]$.

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