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Rearrangement inequalities and applications to isoperimetric problems for eigenvalues

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Abstract

Let Ω be a bounded C^2 domain in \mathbb{R}^n , where n is any positive integer, and let Ω^* be the Euclidean ball centered at 0 and having the same Lebesgue measure as Ω . Consider the operator $L = -\operatorname{div}(A\nabla) + v \cdot \nabla + V$ on Ω with Dirichlet boundary condition, where the symmetric matrix field A is in $W^{1,\infty}(\Omega)$, the vector field v is in $L^\infty(\Omega, \mathbb{R}^n)$ and V is a continuous function in $\overline{\Omega}$. We prove that minimizing the principal eigenvalue of L when the Lebesgue measure of Ω is fixed and when A , v and V vary under some constraints is the same as minimizing the principal eigenvalue of some operators L^* in the ball Ω^* with smooth and radially symmetric coefficients. The constraints which are satisfied by the original coefficients in Ω and the new ones in Ω^* are expressed in terms of some distribution functions or some integral, pointwise or geometric quantities. Some strict comparisons are also established when Ω is not a ball. To these purposes, we associate to the principal eigenfunction φ of L a new symmetric rearrangement defined on Ω^* , which is different from the classical Schwarz symmetrization, and which preserves the integral of $\operatorname{div}(A\nabla\varphi)$ on suitable equi-measurable sets. A substantial part of the paper is devoted to the proofs of pointwise and integral inequalities of independent interest which are satisfied by this rearrangement. The comparisons for the eigenvalues hold for general operators of the type L and they are new even for symmetric operators. Furthermore they generalize, in particular, and provide an alternative proof of the well-known Rayleigh-Faber-Krahn isoperimetric inequality about the principal eigenvalue of the Laplacian under Dirichlet boundary condition on a domain with fixed Lebesgue measure.

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1 Introduction

Throughout all the paper, we fix an integer $n \geq 1$ and denote by $\alpha_n = \pi^{n/2}/\Gamma(n/2 + 1)$ the Lebesgue measure of the Euclidean unit ball in \mathbb{R}^n . By “domain”, we mean a non-empty open connected subset of \mathbb{R}^n , and we denote by \mathcal{C} the set of all bounded domains of \mathbb{R}^n which are of class C^2 . Throughout all the paper, unless otherwise specified, Ω will always be in the class \mathcal{C} . For any measurable subset $A \subset \mathbb{R}^n$, $|A|$ stands for the Lebesgue measure of A . If $\Omega \in \mathcal{C}$, Ω^* will denote the Euclidean ball centered at 0 such that

$$|\Omega^*| = |\Omega|.$$

Define also $C(\overline{\Omega})$ (resp. $C(\overline{\Omega}, \mathbb{R}^n)$) the space of real-valued (resp. \mathbb{R}^n -valued) continuous functions on $\overline{\Omega}$. For all $x \in \mathbb{R}^n \setminus \{0\}$, set

$$e_r(x) = \frac{x}{|x|}, \tag{1.1}$$

where $|x|$ denotes the Euclidean norm of x . Finally, if $\Omega \in \mathcal{C}$, if $v : \Omega \rightarrow \mathbb{R}^n$ is measurable and if $1 \leq p \leq +\infty$, we say that $v \in L^p(\Omega, \mathbb{R}^n)$ if $|v| \in L^p(\Omega)$, and write (somewhat abusively) $\|v\|_p$ or $\|v\|_{L^p(\Omega, \mathbb{R}^n)}$ instead of $\||v|\|_{L^p(\Omega, \mathbb{R})}$.

Various rearrangement techniques for functions defined on Ω were considered in the literature. The most famous one is the Schwarz symmetrization. Let us briefly recall what the idea of this symmetrization is. For any function $u \in L^1(\Omega)$, denote by μ_u the distribution function of u , given by

$$\mu_u(t) = |\{x \in \Omega; u(x) > t\}|$$

for all $t \in \mathbb{R}$. Note that μ is right-continuous, non-increasing and $\mu_u(t) \rightarrow 0$ (resp. $\mu_u(t) \rightarrow |\Omega|$) as $t \rightarrow +\infty$ (resp. $t \rightarrow -\infty$). For all $x \in \Omega^* \setminus \{0\}$, define

$$u^*(x) = \sup \{t \in \mathbb{R}; \mu_u(t) \geq \alpha_n |x|^n\}.$$

The function u^* is clearly radially symmetric, non-increasing in the variable $|x|$ and it satisfies

$$|\{x \in \Omega, u(x) > \zeta\}| = |\{x \in \Omega^*, u^*(x) > \zeta\}|$$

for all $\zeta \in \mathbb{R}$. An essential property of the Schwarz symmetrization is the following one: if $u \in H_0^1(\Omega)$, then $|u|^* \in H_0^1(\Omega^*)$ and (see [39])

$$\||u|^*\|_{L^2(\Omega^*)} = \|u\|_{L^2(\Omega)} \quad \text{and} \quad \|\nabla |u|^*\|_{L^2(\Omega^*)} \leq \|\nabla u\|_{L^2(\Omega)}. \quad (1.2)$$

One of the main applications of this rearrangement technique is the resolution of optimization problems for the eigenvalues of some second-order elliptic operators on Ω . Let us briefly recall some of these problems. If $\lambda_1(\Omega)$ denotes the first eigenvalue of the Laplace operator in Ω with Dirichlet boundary condition, it is well-known that $\lambda_1(\Omega) \geq \lambda_1(\Omega^*)$ and that the inequality is strict unless Ω is a ball (remember that Ω is always assumed to be in the class \mathcal{C}). Since $\lambda_1(\Omega^*)$ can be explicitly computed, this result provides the classical Rayleigh-Faber-Krahn inequality, which states that

$$\lambda_1(\Omega) \geq \lambda_1(\Omega^*) = \left(\frac{1}{|\Omega|}\right)^{2/n} \alpha_n^{2/n} (j_{n/2-1,1})^2, \quad (1.3)$$

where $j_{m,1}$ the first positive zero of the Bessel function J_m . Moreover, equality in (1.3) is attained if and only if Ω is a ball. This result was first conjectured by Rayleigh for $n = 2$ ([40] pp. 339-340), and proved independently by Faber ([19]) and Krahn ([28]) for $n = 2$, and by Krahn for all n in [29] (see [30] for the English translation). The proof of the inequality $\lambda_1(\Omega) \geq \lambda_1(\Omega^*)$ is an immediate consequence of the following variational formula for $\lambda_1(\Omega)$:

$$\lambda_1(\Omega) = \min_{v \in H_0^1(\Omega) \setminus \{0\}} \frac{\int_{\Omega} |\nabla v(x)|^2 dx}{\int_{\Omega} |v(x)|^2 dx}, \quad (1.4)$$

and of the properties (1.2) of the Schwarz symmetrization.

Lots of optimization results involving other eigenvalues of the Laplacian (or more general elliptic symmetric operators of the form $-\text{div}(A\nabla)$) on Ω under Dirichlet boundary condition

have also been established. For instance, the minimum of $\lambda_2(\Omega)$ (the second eigenvalue of the Laplace operator in Ω under Dirichlet boundary condition) among bounded open sets of \mathbb{R}^n with given Lebesgue measure is achieved by the union of two identical balls (this result is attributed to Szegő, see [37]). Very few things seem to be known about optimization problems for the other eigenvalues, see [17, 23, 37, 38, 47]. Various optimization results are also known for functions of the eigenvalues. For instance, it is proved in [5] that $\lambda_2(\Omega)/\lambda_1(\Omega) \leq \lambda_2(\Omega^*)/\lambda_1(\Omega^*)$, and the equality is attained if and only if Ω is a ball. The same result was also extended in [5] to elliptic operators in divergence form with definite weight. We also refer to [6, 7, 9, 12, 18, 26, 27, 31, 32, 35, 36, 38] for further bounds or other optimization results for some eigenvalues or some functions of the eigenvalues in fixed or varying domains of \mathbb{R}^n (or of manifolds).

Other boundary conditions may also be considered. For instance, if $\mu_2(\Omega)$ is the first non-trivial eigenvalue of $-\Delta$ under the Neumann boundary condition, then $\mu_2(\Omega) \leq \mu_2(\Omega^*)$ and the equality is attained if, and only if, Ω is a ball (see [42] in dimension $n = 2$, and [46] in any dimension). Bounds or optimization results for other eigenvalues of the Laplacian under Neumann boundary condition ([38, 42, 46], see also [10] for inhomogeneous problems), for Robin boundary condition ([15]) or for the Stekloff eigenvalue problem ([16]) have also been established. We also mention another Rayleigh conjecture for the lowest eigenvalue of the clamped plate. If $\Omega \subset \mathbb{R}^2$, denote by $\Lambda_1(\Omega)$ the lowest eigenvalue of the operator Δ^2 , so that $\Delta^2 u_1 = \Lambda_1(\Omega)u_1$ in Ω with $u_1 = \nu \cdot \nabla u_1 = 0$ on $\partial\Omega$, where u_1 denotes the principal eigenfunction and ν denotes the outward unit normal on $\partial\Omega$. The second author proved in [33] that $\Lambda_1(\Omega) \geq \Lambda_1(\Omega^*)$ and that equality holds if and only if Ω is a ball, that is a disk in dimension $n = 2$. The analogous result was also established in \mathbb{R}^3 in [8], while the problem is still open in higher dimensions. Much more complete surveys of all these topics can be found in [11, 23, 24].

It is important to observe that the variational formula (1.4) relies heavily on the fact that $-\Delta$ is symmetric on $L^2(\Omega)$. More generally, all the optimization problems considered hitherto concern symmetric operators, and their resolution relies on a ‘‘Rayleigh’’ quotient (that is, a variational formula similar to (1.4)) and the Schwarz symmetrization. Before going further, let us recall that other rearrangement techniques than the Schwarz symmetrization can be found in the literature. For instance, even if this kind of problem is quite different from the ones we are interested in for the present paper, the Steiner symmetrization is the key tool to show that, among all triangles with fixed area, the principal eigenvalue of the Laplacian with Dirichlet boundary condition is minimal for the equilateral triangle (see [39]). Steiner symmetrization is indeed relevant to take into account the polygonal geometry of the domain.

A natural question then arises: can inequalities on eigenvalues of non-symmetric operators be obtained? In view of what we have just explained, such problems require different rearrangement techniques.

Actually, even for symmetric operators, some optimization problems cannot be solved by means of the Schwarz symmetrization, and other rearrangements have to be used. For instance, consider an operator $L = -\operatorname{div}(A\nabla)$ on a domain Ω under Dirichlet boundary condition. Assume that $A(x) \geq \Lambda(x)\operatorname{Id}$ on Ω in the sense of quadratic forms (see below for precise definitions; Id denotes the $n \times n$ identity matrix) for some positive function Λ , and that the L^1 norm of Λ^{-1} is given. Then, what can be said about the infimum of the principal eigenvalue of L under this constraint? In particular, is this infimum greater than the corresponding one on

Ω^* , which is a natural conjecture in view of all the previous results ? Solving such a problem, which is one of our results in the present paper, does not seem to be possible by means of a variational formula for λ_1 (although the operator L is symmetric in $L^2(\Omega)$) and the Schwarz or Steiner symmetrizations.

More general constraints (given distribution functions; integral, pointwise or geometric constraints) on the coefficients A , v and V of non-symmetric operators L of the type $L = -\operatorname{div}(A\nabla) + v \cdot \nabla + V$ under Dirichlet boundary condition will also be investigated. In general, the operator L is non-symmetric, and there is no simple variational formulation of its first eigenvalue such as (1.4) –min-max formulations of the pointwise type (see [14]) or of the integral type (see [25]) certainly hold, but they do not help in our context.

The purpose of the present paper is twofold. First, we present a new rearrangement technique and we show some properties of the rearranged function. The inequalities we obtain between the function in Ω and its symmetrization in Ω^* are of independent interest. Then, we show how this technique can be used to cope with new comparisons between the principal eigenvalues of general non-symmetric elliptic operators of the type $-\operatorname{div}(A\nabla) + v \cdot \nabla + V$ in Ω and of some symmetrized operators in Ω^* . Actually, the comparisons we establish are new even when the operators are symmetric or one-dimensional.

2 Main results

Let us now give precise statements. We are interested in operators of the form

$$L = -\operatorname{div}(A\nabla) + v \cdot \nabla + V$$

in $\Omega \in \mathcal{C}$ under Dirichlet boundary condition.

Throughout the paper, we denote by $\mathcal{S}_n(\mathbb{R})$ the set of $n \times n$ symmetric matrices with real entries. We always assume that $A : \Omega \rightarrow \mathcal{S}_n(\mathbb{R})$ is in $W^{1,\infty}(\Omega)$. This assumption will be denoted by $A = (a_{i,j})_{1 \leq i,j \leq n} \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$: all the components $a_{i,j}$ are in $W^{1,\infty}(\Omega)$ and they can therefore be assumed to be continuous in $\overline{\Omega}$ up to a modification on a zero-measure set. We set

$$\|A\|_{W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))} = \max_{1 \leq i,j \leq n} \|a_{i,j}\|_{W^{1,\infty}(\Omega)},$$

where

$$\|a_{i,j}\|_{W^{1,\infty}(\Omega)} = \|a_{i,j}\|_{L^\infty(\Omega)} + \sum_{1 \leq k \leq n} \left\| \frac{\partial a_{i,j}}{\partial x_k} \right\|_{L^\infty(\Omega)}.$$

We always assume that A is uniformly elliptic on $\overline{\Omega}$, which means that there exists $\delta > 0$ such that, for all $x \in \overline{\Omega}$ and for all $\xi \in \mathbb{R}^n$,

$$A(x)\xi \cdot \xi \geq \delta |\xi|^2.$$

For $B = (b_{i,j})_{1 \leq i,j \leq n} \in \mathcal{S}_n(\mathbb{R})$, $\xi = (\xi_1, \dots, \xi_n) \in \mathbb{R}^n$ and $\xi' = (\xi'_1, \dots, \xi'_n) \in \mathbb{R}^n$, we denote $B\xi \cdot \xi' = \sum_{1 \leq i,j \leq n} b_{i,j} \xi_j \xi'_i$. Actually, in some statements we compare the matrix field A with a matrix field of the type $x \mapsto \Lambda(x)\operatorname{Id}$. We call

$$L_+^\infty(\Omega) = \{\Lambda \in L^\infty(\Omega), \operatorname{ess\,inf}_\Omega \Lambda > 0\},$$

and, for $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$ and $\Lambda \in L_+^\infty(\Omega)$, we say that $A \geq \Lambda \text{ Id}$ almost everywhere (a.e.) in Ω if, for almost every $x \in \Omega$,

$$\forall \xi \in \mathbb{R}^n, A(x)\xi \cdot \xi \geq \Lambda(x)|\xi|^2.$$

For instance, if, for each $x \in \overline{\Omega}$, $\Lambda[A](x)$ denotes the smallest eigenvalue of the matrix $A(x)$, then $\Lambda[A] \in L_+^\infty(\Omega)$ and there holds $A(x) \geq \Lambda[A](x)\text{Id}$ (this inequality is actually satisfied for all $x \in \overline{\Omega}$).

We also always assume that the vector field v is in $L^\infty(\Omega, \mathbb{R}^n)$ and that the potential V is in $L^\infty(\Omega)$. In some statements, V will be in the space $C(\overline{\Omega})$ of continuous functions on $\overline{\Omega}$.

Denote by $\lambda_1(\Omega, A, v, V)$ the principal eigenvalue of $L = -\text{div}(A\nabla) + v \cdot \nabla + V$ with Dirichlet boundary condition on Ω , and $\varphi_{\Omega, A, v, V}$ the corresponding (unique) nonnegative eigenfunction with L^∞ -norm equal to 1. Recall that the following properties hold for $\varphi_{\Omega, A, v, V}$ (see [14]):

$$\begin{cases} -\text{div}(A\nabla\varphi_{\Omega, A, v, V}) + v \cdot \nabla\varphi_{\Omega, A, v, V} + V\varphi_{\Omega, A, v, V} = \lambda_1(\Omega, A, v, V)\varphi_{\Omega, A, v, V} \text{ in } \Omega, \\ \varphi_{\Omega, A, v, V} > 0 \text{ in } \Omega, \varphi_{\Omega, A, v, V} = 0 \text{ on } \partial\Omega, \|\varphi_{\Omega, A, v, V}\|_{L^\infty(\Omega)} = 1. \end{cases} \quad (2.1)$$

and $\varphi_{\Omega, A, v, V} \in W^{2,p}(\Omega)$ for all $1 \leq p < +\infty$ by standard elliptic estimates, whence $\varphi_{\Omega, A, v, V} \in C^{1,\alpha}(\overline{\Omega})$ for all $0 \leq \alpha < 1$. Recall also that $\lambda_1(\Omega, A, v, V) > 0$ if and only if the operator L satisfies the maximum principle in Ω , and that the inequality

$$\lambda_1(\Omega, A, v, V) > \text{ess inf}_\Omega V$$

always holds (see [14] for details and further results).

We are interested in optimization problems for $\lambda_1(\Omega, A, v, V)$ when Ω, A, v and V vary and satisfy some constraints. Our goal is to compare $\lambda_1(\Omega, A, v, V)$ with the principal eigenvalue $\lambda_1(\Omega^*, A^*, v^*, V^*)$ for some fields A^*, v^* and V^* which are defined in the ball Ω^* and satisfy the same constraints as A, v and V . The constraints may be of different types: integral type, L^∞ type, given distribution function of V^- , or bounds on the determinant of A and on another symmetric function of the eigenvalues of A . Throughout the paper, we denote

$$s^- = \max(-s, 0) \text{ and } s^+ = \max(s, 0) \text{ for all } s \in \mathbb{R}.$$

2.1 Constraints on the distribution function of V^- and on some integrals involving Λ and v

We fix here the L^1 norms of Λ^{-1} and $|v|^2\Lambda^{-1}$, some L^∞ bounds on Λ and v , as well as the distribution function of the negative part of V , under the condition that $\lambda_1(\Omega, A, v, V) \geq 0$. Then we can associate some fields A^*, v^* and V^* satisfying the same constraints in Ω^* , and for which the principal eigenvalue is not too much larger than $\lambda_1(\Omega, A, v, V)$, with the extra property that $A^*, |v^*|$ and V^* are smooth and radially symmetric.

Theorem 2.1 *Let $\Omega \in \mathcal{C}$, $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$, $\Lambda \in L_+^\infty(\Omega)$, $v \in L^\infty(\Omega, \mathbb{R}^n)$ and $V \in C(\overline{\Omega})$. Assume that $A \geq \Lambda \text{ Id}$ a.e. in Ω , and that $\lambda_1(\Omega, A, v, V) \geq 0$. Then, for all $\varepsilon > 0$, there exist*

three radially symmetric $C^\infty(\overline{\Omega^*})$ fields $\Lambda^* > 0$, $\omega^* \geq 0$ and $\overline{V^*} \leq 0$ such that, for $v^* = \omega^* e_r$ in $\overline{\Omega^*} \setminus \{0\}$,

$$\begin{cases} \operatorname{ess\,inf}_\Omega \Lambda \leq \min_{\overline{\Omega^*}} \Lambda^* \leq \max_{\overline{\Omega^*}} \Lambda^* \leq \operatorname{ess\,sup}_\Omega \Lambda, & \|(\Lambda^*)^{-1}\|_{L^1(\Omega^*)} = \|\Lambda^{-1}\|_{L^1(\Omega)}, \\ \|v^*\|_{L^\infty(\Omega^*, \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}, & \| |v^*|^2 (\Lambda^*)^{-1} \|_{L^1(\Omega^*)} = \| |v|^2 \Lambda^{-1} \|_{L^1(\Omega)}, \\ \mu_{|\overline{V^*}|} = \mu_{(\overline{V^*})^-} \leq \mu_{V^-}, \end{cases} \quad (2.2)$$

and

$$\lambda_1(\Omega^*, \Lambda^* \operatorname{Id}, v^*, \overline{V^*}) \leq \lambda_1(\Omega, A, v, V) + \varepsilon. \quad (2.3)$$

There also exists a nonpositive radially symmetric $L^\infty(\Omega^*)$ field V^* such that $\mu_{V^*} = \mu_{-V^-}$, $V^* \leq \overline{V^*} \leq 0$ in Ω^* and $\lambda_1(\Omega^*, \Lambda^* \operatorname{Id}, v^*, V^*) \leq \lambda_1(\Omega^*, \Lambda^* \operatorname{Id}, v^*, \overline{V^*}) \leq \lambda_1(\Omega, A, v, V) + \varepsilon$.

If one further assumes that Λ is equal to a constant $\gamma > 0$ in Ω , then there exist two radially symmetric bounded functions $\omega_0^* \geq 0$ and $V_0^* \leq 0$ in Ω^* such that, for $v_0^* = \omega_0^* e_r$,

$$\begin{cases} \|v_0^*\|_{L^\infty(\Omega^*, \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}, & \|v_0^*\|_{L^2(\Omega^*, \mathbb{R}^n)} \leq \|v\|_{L^2(\Omega, \mathbb{R}^n)}, \\ -\max_{\overline{\Omega}} V^- \leq V_0^* \leq 0 \text{ a.e. in } \Omega^*, & \|V_0^*\|_{L^p(\Omega^*)} \leq \|V^-\|_{L^p(\Omega)} \text{ for all } 1 \leq p \leq +\infty, \end{cases} \quad (2.4)$$

and

$$\lambda_1(\Omega^*, \gamma \operatorname{Id}, v_0^*, V_0^*) \leq \lambda_1(\Omega, A, v, V). \quad (2.5)$$

Remember (see [14]) that the inequality $\lambda_1(\Omega, A, v, V) \geq \lambda_1(\Omega, A, v, -V^-)$ always holds. This is the reason why, in order to decrease $\lambda_1(\Omega, A, v, V)$, the rearranged potentials had better be nonpositive in Ω^* , and only the negative part of V plays a role. Notice that the quantities such as the integral of Λ^{-1} , which are preserved here after symmetrization, also appear in other contexts, like in homogenization of elliptic or parabolic equations.

In the case when Λ is a constant, then the number ε can be dropped in (2.3). The price to pay is that the new fields in Ω^* may not be smooth anymore and the distribution function of the new potential V_0^* in Ω^* is no longer equal to that of $-V^-$.

However, in the general case, neither Λ is constant in Ω nor Λ^* is constant in Ω^* (see Remark 5.5 for details). For instance, as already underlined, an admissible Λ is the continuous positive function $\Lambda[A]$, which is not constant in general. Actually, even in the case of operators L which are written in a self-adjoint form (that is, with $v = 0$), the comparison result stated in Theorem 2.1 is new.

An optimization result follows immediately from Theorem 2.1. To state it, we need a few notations. Given

$$m > 0, \quad \overline{M}_\Lambda \geq \underline{m}_\Lambda > 0, \quad \alpha \in \left[\frac{m}{\overline{M}_\Lambda}, \frac{m}{\underline{m}_\Lambda} \right], \quad \overline{M}_v \geq 0, \quad \tau \in [0, \alpha \overline{M}_v^2], \quad \overline{M}_V \geq 0 \quad (2.6)$$

and

$$\mu \in \mathcal{F}_{0, \overline{M}_V}(m) := \{ \rho : \mathbb{R} \rightarrow [0, m], \rho \text{ is right-continuous, non-increasing, } \rho = m \text{ on } (-\infty, 0), \rho = 0 \text{ on } [\overline{M}_V, +\infty) \},$$

we set, for all open set $\Omega \in \mathcal{C}$ such that $|\Omega| = m$,

$$\begin{aligned} \mathcal{G}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) = & \left\{ (A, v, V) \in W^{1, \infty}(\Omega, \mathcal{S}_n(\mathbb{R})) \times L^\infty(\Omega, \mathbb{R}^n) \times C(\overline{\Omega}); \right. \\ & \exists \Lambda \in L_+^\infty(\Omega), A \geq \Lambda \text{ Id a.e. in } \Omega, \\ & \underline{m}_\Lambda \leq \text{ess inf}_\Omega \Lambda \leq \text{ess sup}_\Omega \Lambda \leq \overline{M}_\Lambda, \|\Lambda^{-1}\|_{L^1(\Omega)} = \alpha, \\ & \left. \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \overline{M}_v, \| |v|^2 \Lambda^{-1} \|_{L^1(\Omega)} = \tau \text{ and } \mu_{V^-} \leq \mu \right\} \end{aligned}$$

and

$$\underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) = \inf_{(A, v, V) \in \mathcal{G}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega)} \lambda_1(\Omega, A, v, V). \quad (2.7)$$

Notice that, given $\mu \in \mathcal{F}_{0, \overline{M}_V}(m)$ and $\Omega \in \mathcal{C}$ such that $|\Omega| = m$, there exists $V \in C(\overline{\Omega})$ such that $\mu_{V^-} \leq \mu$ (for instance, $V = 0$ is admissible; furthermore, there is $V \in L^\infty(\Omega)$ such that $\mu_{V^-} = \mu$, see Appendix 7.2), and, necessarily, $V \geq -\overline{M}_V$ in $\overline{\Omega}$. It is immediate to see that, under the conditions (2.6), each set $\mathcal{G}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega)$ is not empty.

Corollary 2.2 *Let $m, \overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V$ be as in (2.6), $\mu \in \mathcal{F}_{0, \overline{M}_V}(m)$ and Ω^* be the Euclidean ball centered at the origin such that $|\Omega^*| = m$. If $\underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) \geq 0$ for all $\Omega \in \mathcal{C}$ such that $|\Omega| = m$, then*

$$\min_{\Omega \in \mathcal{C}, |\Omega|=m} \underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) = \underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega^*). \quad (2.8)$$

Furthermore, in the definition of $\underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega^*)$ in (2.7), the data A, v and V can be assumed to be such that $A = \Lambda \text{ Id}$, $v = \omega e_r = |v| e_r$ and $V \leq 0$ in Ω^* , where Λ, ω and V are $C^\infty(\overline{\Omega}^*)$ and radially symmetric.

Let us now discuss about the non-negativity condition $\lambda_1(\Omega, A, v, V) \geq 0$ in Theorem 2.1, as well as that of Corollary 2.2. We recall (see [14]) that

$$\lambda_1(\Omega, A, v, V) > \min_{\overline{\Omega}} V.$$

Therefore, the condition $\lambda_1(\Omega, A, v, V) \geq 0$ is satisfied in particular if $V \geq 0$ in $\overline{\Omega}$, and the condition $\underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) \geq 0$ in Corollary 2.2 is satisfied if $\overline{M}_V = 0$. Another more complex condition which also involves A and v can be derived. To do so, assume $A \geq \Lambda \text{ Id}$ a.e. in Ω with $m_\Lambda := \text{ess inf}_\Omega \Lambda > 0$, and call $M_v = \|v\|_\infty$ and $m_V = \min_{\overline{\Omega}} V$. Multiply by $\varphi = \varphi_{\Omega, A, v, V}$ the equation (2.1) and integrate by parts over Ω . It follows that, for all $\beta \in (0, 1]$,

$$\begin{aligned} \lambda_1(\Omega, A, v, V) \int_\Omega \varphi^2 & \geq \int_\Omega \Lambda |\nabla \varphi|^2 - \int_\Omega |v| |\nabla \varphi| \varphi + m_V \int_\Omega \varphi^2 \\ & \geq (1 - \beta) \int_\Omega \Lambda |\nabla \varphi|^2 + m_V \int_\Omega \varphi^2 - \frac{1}{4\beta} \int_\Omega |v|^2 \Lambda^{-1} \varphi^2 \\ & \geq [(1 - \beta) m_\Lambda \lambda_1(\Omega) + m_V - (4\beta m_\Lambda)^{-1} M_v^2] \int_\Omega \varphi^2, \end{aligned}$$

where $\lambda_1(\Omega) = \lambda_1(\Omega, \text{Id}, 0, 0) = \min_{\phi \in H_0^1(\Omega), \|\phi\|_2=1} \int_\Omega |\nabla \phi|^2$. If $M_v > 0$ and $m_\Lambda \sqrt{\lambda_1(\Omega)} \geq M_v$, then the value $\beta = M_v / (2m_\Lambda \sqrt{\lambda_1(\Omega)}) \in (0, 1]$ gives the best inequality, that is $\lambda_1(\Omega, A, v, V) \geq$

$m_V + \sqrt{\lambda_1(\Omega)}(m_\Lambda \sqrt{\lambda_1(\Omega)} - M_v)$. The same inequality also holds from the previous calculations if $M_v = 0$. Therefore, the following inequality always holds:

$$\lambda_1(\Omega, A, v, V) \geq m_V + \sqrt{\lambda_1(\Omega)} \times \max(0, m_\Lambda \sqrt{\lambda_1(\Omega)} - M_v).$$

As a consequence, under the notations of Corollary 2.2, it follows from (1.3) that

$$\underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) \geq -\overline{M}_V + m^{-1/n} \alpha_n^{1/n} j_{n/2-1,1} \times \max(0, \underline{m}_\Lambda m^{-1/n} \alpha_n^{1/n} j_{n/2-1,1} - \overline{M}_v)$$

for all $\Omega \in \mathcal{C}$ such that $|\Omega| = m$. The conclusion of Corollary 2.2 is then true if the right-hand side of the above inequality is nonnegative. In particular, for given $n \in \mathbb{N} \setminus \{0\}$, $\underline{m}_\Lambda > 0$, $\overline{M}_v \geq 0$ and $\overline{M}_V \geq 0$, this holds if $m > 0$ is small enough.

To complete this section, we now give a more precise version of Theorem 2.1 when Ω is not a ball.

Theorem 2.3 *Under the notation of Theorem 2.1, assume that $\Omega \in \mathcal{C}$ is not a ball and let $\overline{M}_A > 0$, $\underline{m}_\Lambda > 0$, $\overline{M}_v \geq 0$ and $\overline{M}_V \geq 0$ be such that*

$$\|A\|_{W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))} \leq \overline{M}_A, \operatorname{ess\,inf}_\Omega \Lambda \geq \underline{m}_\Lambda, \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \overline{M}_v \text{ and } \|V\|_{L^\infty(\Omega, \mathbb{R})} \leq \overline{M}_V. \quad (2.9)$$

Then there exists a positive constant $\theta = \theta(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V) > 0$ depending only on Ω , n , \overline{M}_A , \underline{m}_Λ , \overline{M}_v and \overline{M}_V , such that if $\lambda_1(\Omega, A, v, V) > 0$, then there exist three radially symmetric $C^\infty(\overline{\Omega}^)$ fields $\Lambda^* > 0$, $\omega^* \geq 0$, $\overline{V}^* \leq 0$ and a nonpositive radially symmetric $L^\infty(\overline{\Omega}^*)$ field V^* , which satisfy (2.2), $\mu_{V^*} = \mu_{-V^*}$, $V^* \leq \overline{V}^* \leq 0$ and are such that*

$$\lambda_1(\Omega^*, \Lambda^* \operatorname{Id}, v^*, V^*) \leq \lambda_1(\Omega^*, \Lambda^* \operatorname{Id}, v^*, \overline{V}^*) \leq \frac{\lambda_1(\Omega, A, v, V)}{1 + \theta},$$

where $v^* = \omega^* e_r$ in $\overline{\Omega}^* \setminus \{0\}$.

Notice that the assumption $A \geq \Lambda \operatorname{Id}$ a.e. in Ω and the bounds (2.9) imply necessarily that $\overline{M}_A \geq \underline{m}_\Lambda$.

2.2 Constraints on the determinant and another symmetric function of the eigenvalues of A

For our second type of comparison result, we keep the same constraints on v and V as in Theorem 2.1 but we modify the one on A : we now prescribe some conditions on the determinant and another symmetric function of the eigenvalues of A . **We assume in this subsection that $n \geq 2$.** If $A \in \mathcal{S}_n(\mathbb{R})$, if $p \in \{1, \dots, n-1\}$ and if $\lambda_1[A] \leq \dots \leq \lambda_n[A]$ denote the eigenvalues of A , then we call

$$\sigma_p(A) = \sum_{I \subset \{1, \dots, n\}, \operatorname{card}(I)=p} \left(\prod_{i \in I} \lambda_i[A] \right).$$

Throughout the paper, the notation $\operatorname{card}(I)$ means the cardinal of a finite set I . If A is nonnegative, it follows from the arithmetico-geometrical inequality that $C_n^p \times (\det(A))^{p/n} \leq \sigma_p(A)$, where C_n^p is the binomial coefficient $C_n^p = n!/(p! \times (n-p)!)$.

Our third result is as follows:

Theorem 2.4 Assume $n \geq 2$. Let $\Omega \in \mathcal{C}$, $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$, $v \in L^\infty(\Omega, \mathbb{R}^n)$, $V \in C(\overline{\Omega})$ and let $p \in \{1, \dots, n-1\}$, $\omega > 0$ and $\sigma > 0$ be given. Assume that $A \geq \gamma \text{Id}$ in $\overline{\Omega}$ for some constant $\gamma > 0$, that

$$\det(A(x)) \geq \omega, \quad \sigma_p(A(x)) \leq \sigma \text{ for all } x \in \overline{\Omega}, \quad (2.10)$$

and that $\lambda_1(\Omega, A, v, V) \geq 0$. Then, there are two positive numbers $0 < a_1 \leq a_2$ which only depend on n, p, ω and σ , such that, for all $\varepsilon > 0$, there exist a matrix field $A^* \in C^\infty(\overline{\Omega^*} \setminus \{0\}, \mathcal{S}_n(\mathbb{R}))$, two radially symmetric $C^\infty(\overline{\Omega^*})$ fields $\omega^* \geq 0$ and $\overline{V^*} \leq 0$, and a nonpositive radially symmetric $L^\infty(\Omega^*)$ field V^* , such that, for $v^* = \omega^* e_r$ in $\overline{\Omega^*} \setminus \{0\}$,

$$\begin{cases} A \geq a_1 \text{Id in } \Omega, \quad A^* \geq a_1 \text{Id in } \Omega^*, \\ \det(A^*(x)) = \omega, \quad \sigma_p(A^*(x)) = \sigma \text{ for all } x \in \overline{\Omega^*} \setminus \{0\}, \\ \|v^*\|_{L^\infty(\Omega^*, \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}, \quad \|v^*\|_{L^2(\Omega^*, \mathbb{R}^n)} = \|v\|_{L^2(\Omega, \mathbb{R}^n)}, \\ \mu_{|\overline{V^*}|} \leq \mu_{V^-}, \quad \mu_{V^*} = \mu_{-V^-}, \quad V^* \leq \overline{V^*} \leq 0 \text{ in } \Omega^* \end{cases} \quad (2.11)$$

and

$$\lambda_1(\Omega^*, A^*, v^*, V^*) \leq \lambda_1(\Omega^*, A^*, v^*, \overline{V^*}) \leq \lambda_1(\Omega, A, v, V) + \varepsilon.$$

Furthermore, the matrix field A^* is defined, for all $x \in \overline{\Omega^*} \setminus \{0\}$, by:

$$A^*(x)x \cdot x = a_1|x|^2 \text{ and } A^*(x)y \cdot y = a_2|y|^2 \text{ for all } y \perp x.$$

Lastly, there exist two radially symmetric bounded functions $\omega_0^* \geq 0$ and $V_0^* \leq 0$ in Ω^* satisfying (2.4) and $\lambda_1(\Omega^*, A^*, v_0^*, V_0^*) \leq \lambda_1(\Omega, A, v, V)$, where $v_0^* = \omega_0^* e_r$ in Ω^* .

Remark 2.5 Notice that the assumptions of Theorem 2.4 imply necessarily that $C_n^p \omega^{p/n} \leq \sigma$. Actually, the matrix field A^* cannot be extended by continuity at 0, unless $a_1 = a_2$, namely $C_n^p \omega^{p/n} = \sigma$. As a consequence, A^* is not in $W^{1,\infty}(\Omega^*, \mathcal{S}_n(\mathbb{R}))$ if $C_n^p \omega^{p/n} \neq \sigma$, but we can still define $\lambda_1(\Omega^*, A^*, v^*, V^*)$. Indeed, for $\tilde{A}^* = a_1 \text{Id}$ in $\overline{\Omega^*}$, the principal eigenfunction $\tilde{\varphi}^*$ (resp. φ^*) of the operator $-\text{div}(\tilde{A}^* \nabla) + v^* \cdot \nabla + \overline{V^*}$ (resp. $-\text{div}(A^* \nabla) + v^* \cdot \nabla + V^*$) is radially symmetric and belongs to all $W^{2,p}(\Omega^*)$ spaces for all $1 \leq p < +\infty$. Hence,

$$A^* \nabla \tilde{\varphi}^* = \tilde{A}^* \nabla \tilde{\varphi}^* = a_1 \nabla \tilde{\varphi}^*$$

(resp. $A^* \nabla \varphi^* = \tilde{A}^* \nabla \varphi^* = a_1 \nabla \varphi^*$). With a slight abuse of notation, we say that $\tilde{\varphi}^*$ (resp. φ^*) is the principal eigenfunction of $-\text{div}(A^* \nabla) + v^* \cdot \nabla + \overline{V^*}$ (resp. $-\text{div}(A^* \nabla) + v^* \cdot \nabla + V^*$) and we call

$$\lambda_1(\Omega^*, A^*, v^*, \overline{V^*}) = \lambda_1(\Omega^*, \tilde{A}^*, v^*, \overline{V^*})$$

(resp. $\lambda_1(\Omega^*, A^*, v^*, V^*) = \lambda_1(\Omega^*, \tilde{A}^*, v^*, V^*)$).

An interpretation of the conditions (2.10) is that they provide some bounds for the local deformations induced by the matrices $A(x)$, uniformly with respect to $x \in \overline{\Omega}$. Notice that these constraints are saturated for the matrix field A^* in the ball Ω^* .

As for Theorem 2.1, an optimization result follows immediately from Theorem 2.4.

Corollary 2.6 *Assume $n \geq 2$. Given $m > 0$, $p \in \{1, \dots, n-1\}$, $\omega > 0$, $\sigma \geq C_n^p \omega^{p/n}$, $\overline{M}_v \geq 0$, $\tau \in [0, \sqrt{m} \times \overline{M}_v]$, $\overline{M}_V \geq 0$ and $\mu \in \mathcal{F}_{0, \overline{M}_V}(m)$, we set, for all $\Omega \in \mathcal{C}$ such that $|\Omega| = m$,*

$$\mathcal{G}'_{p, \omega, \sigma, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) = \left\{ (A, v, V) \in W^{1, \infty}(\Omega, \mathcal{S}_n(\mathbb{R})) \times L^\infty(\Omega, \mathbb{R}^n) \times C(\overline{\Omega}); \right. \\ \left. \begin{aligned} &\exists \gamma > 0, A(x) \geq \gamma \text{Id for all } x \in \overline{\Omega}, \\ &\det(A(x)) \geq \omega, \sigma_p(A(x)) \leq \sigma \text{ for all } x \in \overline{\Omega}, \\ &\|v\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \overline{M}_v, \|v\|_{L^2(\Omega, \mathbb{R}^n)} = \tau \text{ and } \mu_{V^-} \leq \mu \end{aligned} \right\}$$

and

$$\underline{\lambda}'_{p, \omega, \sigma, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) = \inf_{(A, v, V) \in \mathcal{G}'_{p, \omega, \sigma, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega)} \lambda_1(\Omega, A, v, V).$$

If $\underline{\lambda}'_{p, \omega, \sigma, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) \geq 0$ for all $\Omega \in \mathcal{C}$ such that $|\Omega| = m$, then

$$\inf_{\Omega \in \mathcal{C}, |\Omega|=m} \underline{\lambda}'_{p, \omega, \sigma, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega) = \inf_{(v^*, V^*) \in \mathcal{G}^*_{\overline{M}_v, \tau, \overline{M}_V, \mu}} \lambda_1(\Omega^*, A^*, v^*, V^*),$$

where Ω^* is the ball centered at the origin such that $|\Omega^*| = m$, A^* is given as in Theorem 2.4 and

$$\mathcal{G}^*_{\overline{M}_v, \tau, \overline{M}_V, \mu} = \left\{ (v^*, V^*) \in L^\infty(\Omega^*, \mathbb{R}^n) \times C(\overline{\Omega}^*), v^* = |v^*|e_r, V^* \leq 0, \right. \\ \left. |v^*| \text{ and } V^* \text{ are radially symmetric and } C^\infty(\overline{\Omega}^*), \right. \\ \left. \|v^*\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \overline{M}_v, \|v^*\|_{L^2(\Omega, \mathbb{R}^n)} = \tau \text{ and } \mu_{(V^*)^-} \leq \mu \right\}.$$

Notice also that a sufficient condition for $\underline{\lambda}'_{p, \omega, \sigma, \overline{M}_v, \tau, \overline{M}_V, \mu}(\Omega)$ to be nonnegative for all $\Omega \in \mathcal{C}$ such that $|\Omega| = m$ is:

$$-\overline{M}_V + m^{-1/n} \alpha_n^{1/n} j_{n/2-1,1} \times \max(0, a_1 m^{-1/n} \alpha_n^{1/n} j_{n/2-1,1} - \overline{M}_v) \geq 0,$$

where $a_1 > 0$ is the same as in Theorem 2.4 and only depends on n, p, ω and σ (see Lemma 5.7 for its definition). When $n, p, \omega, \sigma, \overline{M}_v$ and \overline{M}_V are given, the above inequality is satisfied in particular if $m > 0$ is small enough.

When $\Omega \in \mathcal{C}$ is not a ball, we can make Theorem 2.4 more precise: under the same notations as in Theorem 2.4, if $\overline{M}_A > 0$, $\overline{M}_v \geq 0$ and $\overline{M}_V \geq 0$ are such that $\|A\|_{W^{1, \infty}(\Omega, \mathcal{S}_n(\mathbb{R}))} \leq \overline{M}_A$, $\|v\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \overline{M}_v$ and $\|V\|_{L^\infty(\Omega, \mathbb{R})} \leq \overline{M}_V$, then there exists a positive constant

$$\theta' = \theta'(\Omega, n, p, \omega, \sigma, \overline{M}_A, \overline{M}_v, \overline{M}_V) > 0$$

depending only on $\Omega, n, p, \omega, \sigma, \overline{M}_A, \overline{M}_v$ and \overline{M}_V , such that if $\lambda_1(\Omega, A, v, V) > 0$, then there exist a matrix field $A^* \in C^\infty(\overline{\Omega}^* \setminus \{0\}, \mathcal{S}_n(\mathbb{R}))$ (the same as in Theorem 2.4), two radially symmetric $C^\infty(\overline{\Omega}^*)$ fields $\omega^* \geq 0$, $\overline{V}^* \leq 0$ and a nonpositive radially symmetric $L^\infty(\Omega^*)$ field V^* , which satisfy (2.11), $\mu_{V^*} = \mu_{V^-}$, $V^* \leq \overline{V}^* \leq 0$ and are such that

$$\lambda_1(\Omega^*, A^*, v^*, V^*) \leq \lambda_1(\Omega^*, A^*, v^*, \overline{V}^*) \leq \frac{\lambda_1(\Omega, A, v, V)}{1 + \theta'},$$

where $v^* = \omega^* e_r$ in $\overline{\Omega}^* \setminus \{0\}$. It is immediate to see that this fact is a consequence of Theorems 2.3 and 2.4 (notice in particular that the eigenvalues of $A(x)$ are between two positive constants which only depend on n, p, ω and σ).

2.3 Faber-Krahn inequalities for non-symmetric operators

An immediate corollary of Theorem 2.1 is an optimization result, slightly different from Corollary 2.2, where the constraint over the potential V is stated in terms of L^p norms. Namely, given $m > 0$, $\overline{M}_\Lambda \geq \underline{m}_\Lambda > 0$, $\alpha \in \left[\frac{m}{\overline{M}_\Lambda}, \frac{m}{\underline{m}_\Lambda}\right]$, $\overline{M}_v \geq 0$, $\tau \in [0, \alpha \overline{M}_v^2]$, $\tau_V \geq 0$, $1 \leq p \leq +\infty$ and $\Omega \in \mathcal{C}$ such that $|\Omega| = m$, set

$$\begin{aligned} \mathcal{H}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \tau_V, p}(\Omega) = & \left\{ (A, v, V) \in W^{1, \infty}(\Omega, \mathcal{S}_n(\mathbb{R})) \times L^\infty(\Omega, \mathbb{R}^n) \times C(\overline{\Omega}); \right. \\ & \exists \Lambda \in L_+^\infty(\Omega), A \geq \Lambda \text{ Id a.e. in } \Omega, \\ & \underline{m}_\Lambda \leq \text{ess inf}_\Omega \Lambda \leq \text{ess sup}_\Omega \Lambda \leq \overline{M}_\Lambda, \|\Lambda^{-1}\|_{L^1(\Omega)} = \alpha, \\ & \left. \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \overline{M}_v, \| |v|^2 \Lambda^{-1} \|_{L^1(\Omega)} = \tau \text{ and } \|V^-\|_{L^p(\Omega)} \leq \tau_V \right\} \end{aligned}$$

and

$$\underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \tau_V, p}(\Omega) = \inf_{(A, v, V) \in \mathcal{H}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \tau_V, p}(\Omega)} \lambda_1(\Omega, A, v, V).$$

Since, in Theorem 2.1, the L^p norm of \overline{V}^* is smaller than the one of V^- (because the distribution functions of their absolute values are ordered this way), it follows from Theorem 2.1 that

$$\min_{\Omega \in \mathcal{C}, |\Omega|=m} \underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \tau_V, p}(\Omega) = \underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \tau_V, p}(\Omega^*),$$

assuming that $\underline{\lambda}_{\overline{M}_\Lambda, \underline{m}_\Lambda, \alpha, \overline{M}_v, \tau, \tau_V, p}(\Omega) \geq 0$ for all $\Omega \in \mathcal{C}$ such that $|\Omega| = m$. In other words, the infimum of $\lambda_1(\Omega, A, v, V)$ over all the previous constraints when Ω varies but still satisfies $|\Omega| = m$ is the same as the infimum in the ball Ω^* . Observe that we do not know in general if this infimum is actually a minimum. However, specializing to the case of L^∞ constraints for v and V , we can solve a slightly different optimization problem and establish, as an application of Theorems 2.4 and 6.8 (see Section 6 below), a generalization of the classical Rayleigh-Faber-Krahn inequality for the principal eigenvalue of the Laplace operator.

Theorem 2.7 *Let $\Omega \in \mathcal{C}$, $\overline{M}_A > 0$, $\underline{m}_\Lambda > 0$, $\tau_1 \geq 0$ and $\tau_2 \geq 0$ be given. Assume that Ω is not a ball. Consider $A \in W^{1, \infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$, $\Lambda \in L_+^\infty(\Omega)$, $v \in L^\infty(\Omega, \mathbb{R}^n)$ and $V \in L^\infty(\Omega)$ satisfying*

$$\begin{cases} A \geq \Lambda \text{ Id a.e. in } \Omega, \|A\|_{W^{1, \infty}(\Omega, \mathcal{S}_n(\mathbb{R}))} \leq \overline{M}_A, \text{ess inf}_\Omega \Lambda \geq \underline{m}_\Lambda, \\ \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \tau_1 \text{ and } \|V\|_{L^\infty(\Omega)} \leq \tau_2. \end{cases}$$

Then there exists a positive constant $\eta = \eta(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \tau_1) > 0$ depending only on Ω , n , \overline{M}_A , \underline{m}_Λ and τ_1 , and there exists a radially symmetric $C^\infty(\overline{\Omega}^)$ field $\Lambda^* > 0$ such that*

$$\text{ess inf}_\Omega \Lambda \leq \min_{\overline{\Omega}^*} \Lambda^* \leq \max_{\overline{\Omega}^*} \Lambda^* \leq \text{ess sup}_\Omega \Lambda, \|(\Lambda^*)^{-1}\|_{L^1(\Omega^*)} = \|\Lambda^{-1}\|_{L^1(\Omega)}, \quad (2.12)$$

and

$$\lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, -\tau_2) \leq \lambda_1(\Omega, A, v, V) - \eta. \quad (2.13)$$

Notice that, as in Theorem 2.3, the assumptions of Theorem 2.7 imply necessarily that $\overline{M}_A \geq \underline{m}_A$. Notice also that, in Theorem 2.7, contrary to our other results, we do not assume that $\lambda_1(\Omega, A, v, V) \geq 0$. In the previous results, we imposed a constraint on the distribution function of the negative part of the potential and we needed the nonnegativity of $\lambda_1(\Omega, A, v, V)$. Here, we first write

$$\lambda_1(\Omega, A, v, V) \geq \lambda_1(\Omega, A, v, -\tau_2) = -\tau_2 + \lambda_1(\Omega, A, v, 0)$$

and we apply Theorem 2.3 to $\lambda_1(\Omega, A, v, 0)$, which is positive. We complete the proof with further results which are established in Section 6.

Observe also that, in the inequality (2.13), the constraints τ_1 and τ_2 on the L^∞ norms of the drift and the potential are saturated in the ball Ω^* .

Actually, in Theorem 2.7, if we replace the assumption $\|V\|_{L^\infty(\Omega)} \leq \tau_2$ by $\text{ess inf}_\Omega V \geq \tau_3$ (where $\tau_3 \in \mathbb{R}$), then inequality (2.13) is changed into

$$\lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, \tau_3) \leq \lambda_1(\Omega, A, v, V) - \eta.$$

Since $\lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, \tau) = \lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, 0) + \tau$ for all $\tau \in \mathbb{R}$, the previous inequality is better than (2.13). In the following corollary, we choose to compare directly V with $\text{ess inf}_\Omega V$.

Corollary 2.8 *Let $\Omega \in \mathcal{C}$, $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$, $v \in L^\infty(\Omega, \mathbb{R}^n)$ and $V \in L^\infty(\Omega)$. Call $\Lambda[A](x)$ the smallest eigenvalue of the matrix $A(x)$ at each point $x \in \overline{\Omega}$ and assume that $\gamma_A = \min_{\overline{\Omega}} \Lambda[A] > 0$. Then*

$$\lambda_1(\Omega, A, v, V) \geq F_n(|\Omega|, \min_{\overline{\Omega}} \Lambda[A], \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}, \text{ess inf}_\Omega V), \quad (2.14)$$

where $F_n : (0, +\infty) \times (0, +\infty) \times [0, +\infty) \times \mathbb{R} \rightarrow \mathbb{R}$ is defined by

$$F_n(m, \gamma, \alpha, \beta) = \lambda_1(B_{(m/\alpha_n)^{1/n}}^n, \gamma \text{Id}, \alpha e_r, \beta)$$

for all $(m, \gamma, \alpha, \beta) \in (0, +\infty) \times (0, +\infty) \times [0, +\infty) \times \mathbb{R}$, and $B_{(m/\alpha_n)^{1/n}}^n$ denotes the Euclidean ball of \mathbb{R}^n with center 0 and radius $(m/\alpha_n)^{1/n}$. Furthermore, the inequality (2.14) is strict if Ω is not a ball.

In Corollary 2.8, formula (2.14) reduces to (1.3) when $A = \text{Id}$ and $v = 0$, $V = 0$. Theorem 2.7 can then be viewed as a natural extension of the first Rayleigh conjecture to more general elliptic operators with potential, drift and general diffusion. We refer to Remark 6.9 for further comments on these results.

2.4 Some comparisons with results in the literature

If in Theorem 2.1, the function Λ is identically equal to a constant $\gamma > 0$ in Ω , and if $V \geq 0$, then inequality (2.5) could also be derived implicitly from Theorem 1 by Talenti [43]. In [43], Talenti's argument relies on the Schwarz symmetrization and one of the key inequalities which is used in [43] is

$$\int_{\Omega} -\text{div}(A \nabla \varphi) \times \varphi = \int_{\Omega} A \nabla \varphi \cdot \nabla \varphi \geq \gamma \int_{\Omega} |\nabla \varphi|^2.$$

This kind of inequality cannot be used directly for our purpose since it does not take into account the fact that $A \geq \Lambda \text{Id}$ a.e. in Ω , where the function Λ may not be constant. The proofs of the present paper use a completely different rearrangement technique which has its own interest, and which allows us to take into account any non-constant function $\Lambda \in L_+^\infty(\Omega)$. Actually, paper [43] was not concerned with eigenvalue problems, but with various comparison results for solutions of elliptic problems (see also [2, 3, 4, 44]). Even in the case when Λ is constant and $V \geq 0$, proving the inequality (2.5) between the principal eigenvalues of the initial and rearranged operators by means of Talenti's results requires several extra arguments, some of them using results contained in Section 6 of the present paper. We also refer to Section 6.2 for additional comments in the case when Λ is constant.

But, once again, besides the own interest and the novelty of the tools we use in the present paper, one of the main features in Theorem 2.1 (and in Theorems 2.3 and 2.7) is that the ellipticity function Λ and its symmetrization Λ^* are not constant in general (see Remark 5.5). Optimizing with non-constant coefficients in the second-order terms creates additional and substantial difficulties. In particular, the conclusion of Theorem 2.1 does not follow from previous works, even implicitly and even if the lower-order terms are zero. More generally speaking, all the comparison results of the present paper are new even when $v = 0$, namely when the operator L is symmetric. Moreover, all the results are new also when the operators are one-dimensional (except Theorem 2.4 the statement of which does make sense only when $n \geq 2$).

The improved version of Theorem 2.1 when Ω is not a ball, namely Theorem 2.3, is also new and does not follow from earlier results.

As far as Theorem 2.4 is concerned, optimization problems for eigenvalues when the constraint on A is expressed in terms of the determinant and the trace, or more general symmetric functions of the eigenvalues of A , have not been considered hitherto.

Let us now focus on Theorem 2.7 and Corollary 2.8. In a previous work ([21, 22]), we proved a somewhat more complete version of this Faber-Krahn inequality in the case of the Laplace operator with a drift term. Namely, let Ω be a $C^{2,\alpha}$ non empty bounded domain of \mathbb{R}^n for some $0 < \alpha < 1$. For any vector field $v \in L^\infty(\Omega, \mathbb{R}^n)$, denote by

$$\lambda_1(\Omega, v) = \lambda_1(\Omega, \text{Id}, v, 0) \tag{2.15}$$

the principal eigenvalue of $-\Delta + v \cdot \nabla$ in Ω under Dirichlet boundary condition. Then, the following Faber-Krahn type inequality holds:

Theorem 2.9 [21, 22] *Let Ω be a $C^{2,\alpha}$ non-empty bounded connected open subset of \mathbb{R}^n for some $0 < \alpha < 1$, let $\tau \geq 0$ and $v \in L^\infty(\Omega, \mathbb{R}^n)$ be such that $\|v\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \tau$. Then*

$$\lambda_1(\Omega, v) \geq \lambda_1(\Omega^*, \tau e_r), \tag{2.16}$$

and the equality holds if and only if, up to translation, $\Omega = \Omega^$ and $v = \tau e_r$.*

Remark 2.10 Here we quote exactly the statement of [21, 22], but actually it is enough to assume that Ω is of class C^2 .

Notice that we can recover Theorem 2.9 from the results of the present paper. Indeed, when Ω is not ball, the strict inequality in (2.16) follows at once from Theorem 2.7, and when Ω is a ball (say, with center 0) and $v \neq \tau e_r$, this strict inequality will follow from Theorem 6.8 (see Section 6 below). Strictly speaking, the inequality (2.16) could also be derived from Theorem 2 in [43] (see also [2, 3]) and from extra arguments similar to the ones used in Section 6.1. But the case of equality is new, while Theorem 2.7 is entirely new. Indeed, an important feature in Theorem 2.7 is the fact that the diffusion A is assumed to be bounded from below by ΛId where Λ is a possibly non-constant function, and that $\lambda_1(\Omega, A, v, V)$ is compared with $\lambda_1(\Omega^*, \Lambda^* \text{Id}, \|v\|_\infty e_r, -\|V\|_\infty)$, where Λ^* is also possibly nonconstant (in other words, the operator $\text{div}(\Lambda^* \nabla)$ is not necessarily equal to a constant times the Laplace operator). Furthermore, another novelty in Theorem 2.7 is that, when Ω is not a ball, the difference $\lambda_1(\Omega, A, v, V) - \lambda_1(\Omega^*, \Lambda^* \text{Id}, \|v\|_\infty e_r, -\|V\|_\infty)$ is estimated from below by a positive quantity depending only on Ω , n and on some structural constants of the operator. All these observations imply that Theorem 2.7 is definitely more general than Theorem 2.9 and is not implicit in [43], or even in more recent works in the same spirit (like [4], for instance).

When the vector field v is divergence free (in the sense of distributions), then $\lambda_1(\Omega, v) \geq \lambda_1(\Omega)$ (multiply $-\Delta \varphi_{\Omega, \text{Id}, v, 0} + v \cdot \nabla \varphi_{\Omega, \text{Id}, v, 0} = \lambda_1(\Omega, v) \varphi_{\Omega, \text{Id}, v, 0}$ by $\varphi_{\Omega, \text{Id}, v, 0}$ and integrate by parts over Ω).¹ Thus, minimizing $\lambda_1(\Omega, v)$ when $|\Omega| = m$ and v is divergence free and satisfies $\|v\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \tau$ (with given $m > 0$ and $\tau \geq 0$), is the same as minimizing $\lambda_1(\Omega)$ in the Rayleigh conjecture. We also refer to [21] and [22] for further optimization results for $\lambda_1(\Omega, v)$ with L^∞ constraints on the drifts.

Remark 2.11 For non-empty connected and possibly unbounded open sets Ω with finite measure, the principal eigenvalue $\lambda_1(\Omega, A, v, V)$ of the operator $L = -\text{div}(A \nabla) + v \cdot \nabla + V$ can be defined as

$$\lambda_1(\Omega, A, v, V) = \sup \{ \lambda \in \mathbb{R}, \exists \phi \in C^2(\Omega), \phi > 0 \text{ in } \Omega, (-L + \lambda)\phi \leq 0 \text{ in } \Omega \}.$$

When Ω is bounded, this definition is taken from [14] (see also [1, 34]), and it coincides with the characterization (2.1) when $\Omega \in \mathcal{C}$. It follows from the arguments of Chapter 2 of [14] that

$$\lambda_1(\Omega, A, v, V) = \inf_{\Omega' \subset \subset \Omega, \Omega' \in \mathcal{C}} \lambda_1(\Omega', A|_{\Omega'}, v|_{\Omega'}, V|_{\Omega'}), \quad (2.17)$$

where $A|_{\Omega'}$, $v|_{\Omega'}$, $V|_{\Omega'}$ denote the restrictions of the fields A , v and V to Ω' . When Ω is a general non-empty open set with finite measure, we then define

$$\lambda_1(\Omega, A, v, V) = \inf_{j \in J} \lambda_1(\Omega_j, A|_{\Omega_j}, v|_{\Omega_j}, V|_{\Omega_j}), \quad (2.18)$$

where the Ω_j 's are the connected components of Ω .

Some of the comparison results which are stated in the previous subsections can then be extended to the class of general open sets Ω with finite measure (see Remarks 5.6, 5.8 and 6.10).

¹We refer to [13] for a detailed analysis of the behavior of $\lambda_1(\Omega, A, Bv, V)$ when $B \rightarrow +\infty$ and v is a fixed divergence free vector field in $L^\infty(\Omega)$.

2.5 Main tools: a new type of symmetrization

As already underlined, the proofs of Theorems 2.1, 2.3, 2.4 and 2.7 do not use the usual Schwarz symmetrization. The key tool in the proofs is a new (up to our knowledge) rearrangement technique for some functions on Ω , which can take into account non-constant ellipticity functions Λ . Roughly speaking, given Ω , A , v and V such that $A \geq \Lambda \text{Id}$, if $\varphi = \varphi_{\Omega, A, v, V}$ denotes the principal eigenfunction of the operator $-\text{div}(A\nabla) + v \cdot \nabla + V$ in Ω under Dirichlet boundary condition (that is, φ solves (2.1)), we associate to φ , Λ , v and V some rearranged functions or vector fields, which are called $\tilde{\varphi}$, $\hat{\Lambda}$, \hat{v} and \hat{V} . They are defined on Ω^* and are built so that some quantities are preserved. The precise definitions will be given in Section 3, but let us quickly explain how the function $\tilde{\varphi}$ is defined. Denote by R the radius of Ω^* . For all $0 \leq a < 1$, define

$$\Omega_a = \{x \in \Omega, a < \varphi(x) \leq 1\}$$

and define $\rho(a) \in (0, R]$ such that $|\Omega_a| = |B_{\rho(a)}|$, where B_s denotes the open Euclidean ball of radius $s > 0$ and centre 0. Define also $\rho(1) = 0$. The function $\rho : [0, 1] \rightarrow [0, R]$ is decreasing, continuous, one-to-one and onto. Then, the rearrangement of φ is the radially symmetric decreasing function $\tilde{\varphi} : \overline{\Omega^*} \rightarrow \mathbb{R}$ vanishing on $\partial\Omega^*$ such that, for all $0 \leq a < 1$,

$$\int_{\Omega_a} \text{div}(A\nabla\varphi)(x)dx = \int_{B_{\rho(a)}} \text{div}(\hat{\Lambda}\nabla\tilde{\varphi})(x)dx$$

(we do not wish to give the explicit expression of the function $\hat{\Lambda}$ right now). The fundamental inequality satisfied by $\tilde{\varphi}$ is the fact that, for all $x \in \overline{\Omega^*}$,

$$\tilde{\varphi}(x) \geq \rho^{-1}(|x|) \tag{2.19}$$

(see Corollary 3.6 below, and Lemma 4.3 for strict inequalities when Ω is not a ball).

This symmetrization is definitely different from the Schwarz symmetrization since the distribution functions of φ and $\tilde{\varphi}$ are not the same in general. Moreover, the L^1 norm of the gradient of $\tilde{\varphi}$ on Ω^* is larger than or equal to that of φ on Ω , and, when $A = \gamma\text{Id}$ (for a positive constant γ), the L^2 norm of the gradient of $\tilde{\varphi}$ on Ω^* is larger than or equal to that of φ on Ω (see Remark 3.13 below).

Actually, the function φ is not regular enough for this construction to be correct, and we have to deal with suitable approximations of φ . We refer to Section 3 and the following ones for exact and complete statements and proofs. Let us just mention that the proof of (2.19) relies, apart from the definition of $\tilde{\varphi}$, on the usual isoperimetric inequality on \mathbb{R}^n .

Notice that the tools which are developed in this paper not only give new comparison results for symmetric and non-symmetric second-order operators with non-constant coefficients, but they also provide an alternative proof of the Rayleigh-Faber-Krahn isoperimetric inequality (1.3) for the Dirichlet Laplacian.

Finally, the new rearrangement we introduce in this paper is likely to be used in other problems involving elliptic partial differential equations.

Let us give a few open problems related to our results. In all our results, several minimization problems for the principal eigenvalue of a second-order elliptic operator in a

domain Ω under some constraints have been reduced to the same problems on the ball Ω^* centered at 0 with the same Lebesgue measure and for operators with radially symmetric coefficients. However, even in the case of the ball and for operators with radially symmetric coefficients, some of these optimization problems remain open. For instance, in Corollary 2.2, is it possible to compute explicitly the right-hand side of (2.8) ? An analogous question may be asked for the other theorems, corresponding to different constraints (even for Theorem 2.7).

When we combine Theorems 2.1 and 2.3, it follows that the inequality (2.5) is strict when Ω is not a ball and $\lambda_1(\Omega, A, v, V) > 0$. But in Theorem 2.1, when Ω is a ball, for which A , v and V does the case of equality occur in (2.5) ? Does this require that the initial data should be all radially symmetric ? The same question can be asked in Theorem 2.4 as well. An answer to these questions would provide a complete analogue of Theorem 2.9 for general second-order elliptic operators in divergence form. Furthermore, in Theorem 2.1, in the general case when Λ is not constant and even if Ω is a ball, can one state a result without ε but with still keeping the constraints (2.2) ?

When $\Omega = \Omega^*$, Λ^* is fixed and v and V vary with some constraints on their L^∞ norms, we prove in Section 6 that there exist a unique v and a unique V minimizing $\lambda_1(\Omega^*, \Lambda^* \text{Id}, v, V)$. In particular, if Λ^* is radially symmetric, then we show that v and V are given by inequality (2.13) of Theorem 2.7. Many other optimization results in the ball can be asked if some of the fields Λ^* , v^* and V^* are fixed while the others vary under some constraints. We intend to come back to all these issues in a forthcoming paper.

Here are some other open problems. In Theorem 2.4, can one replace the determinant of A by more general functions of the eigenvalues of A , namely $\sigma_q(A)$ with $p < q \leq n - 1$?

It would also be very interesting to obtain results similar to ours for general second-order elliptic operators of the form

$$-\sum_{i,j} a_{i,j} \partial_{i,j} + \sum_i b_i \partial_i + c,$$

where the $a_{i,j}$'s are continuous in $\overline{\Omega}$ (but do not necessarily belong to $W^{1,\infty}(\Omega)$), and the b_i 's and c are bounded in Ω (recall that such operators still have a real principal eigenvalue, see [14]), and to consider other boundary conditions (Neumann, Robin, Stekloff problems...)

Outline of the paper. The paper is organized as follows. Section 3 is devoted to the precise definitions of the rearranged function and the proof of the inequalities satisfied by this rearrangement, whereas improved inequalities are obtained in Section 4 when Ω is not a ball. The proofs of Theorems 2.1, 2.3 and 2.4 are given in Section 5, while the Faber-Krahn inequalities (Theorem 2.7 and Corollary 2.8) are established in Section 6. Some optimization results in a fixed domain, which are interesting in their own right and are also required for the proof of Theorem 2.7, are also proved in Section 6. Finally, the appendix contains the proof of a technical approximation result (which is used in the proofs of Section 5), a short remark about distribution functions and some useful asymptotics of $\lambda_1(\Omega^*, \tau e_r) = \lambda_1(\Omega^*, \text{Id}, \tau e_r, 0)$ when $\tau \rightarrow +\infty$.

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3 Inequalities for the rearranged functions

In this section, we present a new spherical rearrangement of functions and we prove some pointwise and integral inequalities for the rearranged data. The results are of independent interest and this is the reason why we present them in a separate section.

3.1 General framework, definitions of the rearrangements and basic properties

In this subsection, we give some assumptions which will remain valid throughout all Section 3. Fix $\Omega \in \mathcal{C}$, $A_\Omega \in C^1(\overline{\Omega}, \mathcal{S}_n(\mathbb{R}))$, $\Lambda_\Omega \in C^1(\overline{\Omega})$, $\omega \in C(\overline{\Omega})$ and $V \in C(\overline{\Omega})$. Assume that

$$A_\Omega(x) \geq \Lambda_\Omega(x)\text{Id for all } x \in \overline{\Omega}, \quad (3.1)$$

and that there exists $\gamma > 0$ such that

$$\Lambda_\Omega(x) \geq \gamma \text{ for all } x \in \overline{\Omega}.$$

Let ψ be a $C^1(\overline{\Omega})$ function, analytic and positive in Ω , such that $\psi = 0$ on $\partial\Omega$ and

$$\nabla\psi(x) \neq 0 \text{ for all } x \in \partial\Omega,$$

so that $\nu \cdot \nabla\psi < 0$ on $\partial\Omega$, where ν denotes the outward unit normal to $\partial\Omega$. We always assume throughout this section that

$$f := -\text{div}(A_\Omega \nabla\psi) \text{ in } \Omega$$

is a non-zero polynomial, so that $\psi \in W^{2,p}(\Omega)$ for all $1 \leq p < +\infty$ and $\psi \in C^{1,\alpha}(\overline{\Omega})$ for all $0 \leq \alpha < 1$.

Set

$$M = \max_{x \in \overline{\Omega}} \psi(x).$$

For all $a \in [0, M)$, define

$$\Omega_a = \{x \in \Omega, \psi(x) > a\}$$

and, for all $a \in [0, M]$,

$$\Sigma_a = \{x \in \overline{\Omega}, \psi(x) = a\}.$$

The set $\{x \in \overline{\Omega}, \nabla\psi(x) = 0\}$ is included in some compact set $K \subset \Omega$, which implies that the set

$$Z = \{a \in [0, M], \exists x \in \Sigma_a, \nabla\psi(x) = 0\}$$

of the critical values of ψ is finite ([41]) and can then be written as

$$Z = \{a_1, \dots, a_m\}$$

for some $m \in \mathbb{N}^* = \mathbb{N} \setminus \{0\}$. Observe also that $M \in Z$ and that $0 \notin Z$. One can then assume without loss of generality that

$$0 < a_1 < \dots < a_m = M.$$

The set $Y = [0, M] \setminus Z$ of the non critical values of ψ is open relatively to $[0, M]$ and can be written as

$$Y = [0, M] \setminus Z = [0, a_1) \cup (a_1, a_2) \cup \cdots \cup (a_{m-1}, M).$$

For all $a \in Y$, the hypersurface Σ_a is of class C^2 (notice also that $\Sigma_0 = \partial\Omega$ is of class C^2 by assumption) and $|\nabla\psi|$ does not vanish on Σ_a . Therefore, the functions defined on Y by

$$\begin{cases} g : Y \ni a \mapsto \int_{\Sigma_a} |\nabla\psi(y)|^{-1} d\sigma_a(y) \\ h : Y \ni a \mapsto \int_{\Sigma_a} f(y) |\nabla\psi(y)|^{-1} d\sigma_a(y) \\ i : Y \ni a \mapsto \int_{\Sigma_a} d\sigma_a(y) \end{cases} \quad (3.2)$$

are (at least) continuous in Y and C^1 in $Y \setminus \{0\}$, where $d\sigma_a$ denotes the surface measure on Σ_a for $a \in Y$.

Denote by R the radius of Ω^* (the open Euclidean ball centered at the origin and such that $|\Omega^*| = |\Omega|$, that is $\Omega^* = B_R$). For all $a \in [0, M)$, let $\rho(a) \in (0, R]$ be defined so that

$$|\Omega_a| = |B_{\rho(a)}| = \alpha_n \rho(a)^n.$$

Recall that α_n is the volume of the unit ball B_1 . The function ρ is extended at M by

$$\rho(M) = 0.$$

Lemma 3.1 *The function ρ is a continuous decreasing map from $[0, M]$ onto $[0, R]$.*

Proof. The function $\rho : [0, M] \rightarrow [0, R]$ is clearly decreasing since

$$|\{x \in \Omega, a < \psi(x) \leq b\}| > 0$$

for all $0 \leq a < b \leq M$. Fix now any $a \in (0, M]$. Since $\psi \in W^{2,p}(\Omega)$ (actually, for all $1 \leq p < +\infty$), one has

$$\frac{\partial^2 \psi}{\partial x_i \partial x_j} \times \mathbf{1}_{\{\psi=a\}} = \frac{\partial \psi}{\partial x_i} \times \mathbf{1}_{\{\psi=a\}} = 0 \text{ almost everywhere in } \Omega$$

for all $1 \leq i, j \leq n$, where $\mathbf{1}_E$ denotes the characteristic function of a set E . Therefore, $f \times \mathbf{1}_{\{\psi=a\}} = 0$ almost everywhere in Ω . Since f is a nonzero polynomial, one gets that

$$|\Sigma_a| = 0 \text{ for all } a \in (0, M].$$

Notice that $|\Sigma_0| = |\partial\Omega| = 0$ as well. Lastly, $\rho(0) = R$ and $\rho(M) = 0$. As a conclusion, the function ρ is continuous on $[0, M]$ and is a one-to-one and onto map from $[0, M]$ to $[0, R]$. \square

Lemma 3.2 *The function ρ is of class C^1 in Y and*

$$\forall a \in Y, \quad \rho'(a) = -(n\alpha_n \rho(a)^{n-1})^{-1} g(a) = -(n\alpha_n \rho(a)^{n-1})^{-1} \int_{\Sigma_a} |\nabla\psi(y)|^{-1} d\sigma_a(y).$$

Proof. Fix $a \in Y$. Let $\eta > 0$ be such that $[a, a + \eta] \subset Y$. For $t \in (0, \eta)$,

$$\begin{aligned} \alpha_n[\rho(a+t)^n - \rho(a)^n] &= |\Omega_{a+t}| - |\Omega_a| = - \int_{\{a < \psi(x) \leq a+t\}} dx \\ &= - \int_a^{a+t} \left(\int_{\Sigma_b} |\nabla \psi(y)|^{-1} d\sigma_b(y) \right) db \end{aligned}$$

from the co-area formula. Hence,

$$\frac{\alpha_n[\rho(a+t)^n - \rho(a)^n]}{t} \rightarrow -g(a) \quad \text{as } t \rightarrow 0^+$$

for all $a \in Y$, due to the continuity of g on Y . Similarly, one has that

$$\frac{\alpha_n[\rho(a+t)^n - \rho(a)^n]}{t} \rightarrow -g(a) \quad \text{as } t \rightarrow 0^-$$

for all $a \in Y \setminus \{0\}$. The conclusion of the lemma follows since $Y \subset [0, M)$, whence $\rho(a) \neq 0$ for all $a \in Y$. \square

We now define the function $\tilde{\psi}$ in $\overline{\Omega^*}$, which is a spherical rearrangement of ψ by means of a new type of symmetrization. The definition of $\tilde{\psi}$ involves the rearrangement of the datum Λ_Ω .

First, call

$$E = \{x \in \overline{\Omega^*}, |x| \in \rho(Y)\}.$$

The set E is a finite union of spherical shells and, from Lemma 3.1, it is open relatively to $\overline{\Omega^*}$ and can be written as

$$E = \{x \in \mathbb{R}^n, |x| \in (0, \rho(a_{m-1})) \cup \dots \cup (\rho(a_2), \rho(a_1)) \cup (\rho(a_1), R)\}.$$

with

$$0 = \rho(a_m) = \rho(M) < \rho(a_{m-1}) < \dots < \rho(a_1) < R.$$

Notice that $0 \notin E$.

Next, for all $r \in \rho(Y)$, set

$$G(r) = \frac{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}}{\int_{\Sigma_{\rho^{-1}(r)}} \Lambda_\Omega(y)^{-1} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}} > 0, \quad (3.3)$$

where $\rho^{-1} : [0, R] \rightarrow [0, M]$ denotes the reciprocal of the function ρ . For all $x \in E$, define

$$\widehat{\Lambda}(x) = G(|x|). \quad (3.4)$$

The function $\widehat{\Lambda}$ is then defined almost everywhere in $\overline{\Omega^*}$. By the observations above and since Λ_Ω is positive and $C^1(\overline{\Omega})$, the function $\widehat{\Lambda}$ is continuous on E and C^1 on $E \cap \Omega^*$. Furthermore, $\widehat{\Lambda} \in L^\infty(\Omega^*)$ and

$$0 < \min_{\overline{\Omega}} \Lambda_\Omega \leq \operatorname{ess\,inf}_{\Omega^*} \widehat{\Lambda} \leq \operatorname{ess\,sup}_{\Omega^*} \widehat{\Lambda} \leq \max_{\overline{\Omega}} \Lambda_\Omega. \quad (3.5)$$

For any two real numbers $a < b$ such that $[a, b] \subset Y$, the co-area formula gives

$$\begin{aligned} \int_{\Omega_a \setminus \Omega_b} \Lambda_\Omega(y)^{-1} dy &= \int_a^b \left(\int_{\Sigma_s} \Lambda_\Omega(y)^{-1} |\nabla \psi(y)|^{-1} d\sigma_s(y) \right) ds \\ &= \int_{\rho(b)}^{\rho(a)} \left(\frac{\int_{\Sigma_{\rho^{-1}(t)}} \Lambda_\Omega(y)^{-1} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(t)}(y)}{\int_{\Sigma_{\rho^{-1}(t)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(t)}(y)} \right) n\alpha_n t^{n-1} dt. \end{aligned}$$

The last equality is obtained from Lemma 3.2 after the change of variables $s = \rho^{-1}(t)$. Since $\widehat{\Lambda}$ is radially symmetric, it follows by (3.3-3.4) that

$$\int_{\Omega_a \setminus \Omega_b} \Lambda_\Omega(y)^{-1} dy = \int_{S_{\rho(b), \rho(a)}} \widehat{\Lambda}(x)^{-1} dx,$$

where, for any $0 \leq s < s'$, $S_{s, s'}$ denotes

$$S_{s, s'} = \{x \in \mathbb{R}^n, s < |x| < s'\}.$$

Lebesgue's dominated convergence theorem then implies that

$$\int_{\Omega} \Lambda_\Omega(y)^{-1} dy = \int_{\Omega^*} \widehat{\Lambda}(x)^{-1} dx. \quad (3.6)$$

Lastly, set $F(0) = 0$ and, for all $r \in \rho(Y)$, set

$$F(r) = \frac{1}{n\alpha_n r^{n-1} G(r)} \int_{\Omega_{\rho^{-1}(r)}} \operatorname{div}(A_\Omega \nabla \psi)(x) dx. \quad (3.7)$$

The function F is then defined almost everywhere in $[0, R]$.

Lemma 3.3 *The function F belongs to $L^\infty([0, R])$ and is continuous on $\rho(Y) \cup \{0\}$. Moreover, $F < 0$ on $\rho(Y)$.*

Proof. The continuity of F on $\rho(Y)$ is a consequence of Lemma 3.1, of the continuity of $\widehat{\Lambda}$ on E and of the fact that $\operatorname{div}(A_\Omega \nabla \psi) = f$ in Ω , with f continuous and thus bounded in $\overline{\Omega}$.

Observe that, since $\Lambda_\Omega(x) \geq \gamma > 0$ for all $x \in \overline{\Omega}$, one has $\widehat{\Lambda}(x) \geq \gamma$ for all $x \in E$. For $0 < r \leq R$ with $r \in \rho(Y)$ ($\supset (0, \rho(a_{m-1}))$), one has

$$|F(r)| \leq (n\alpha_n r^{n-1} \gamma)^{-1} \|f\|_{L^\infty(\Omega)} \alpha_n r^n = (n\gamma)^{-1} \|f\|_{L^\infty(\Omega)} r,$$

thus F is continuous at 0 as well and belongs to $L^\infty([0, R])$. Finally, for all $r \in \rho(Y)$, since $\psi(y) = \rho^{-1}(r)$ for all $y \in \Sigma_{\rho^{-1}(r)}$ and since $\psi > \rho^{-1}(r)$ in $\Omega_{\rho^{-1}(r)}$ and $|\nabla \psi(y)| \neq 0$ for all $y \in \Sigma_{\rho^{-1}(r)}$, one has

$$\nu_{\rho^{-1}(r)} \cdot \nabla \psi < 0 \text{ on } \Sigma_{\rho^{-1}(r)},$$

where, for any $a \in Y$, ν_a denotes the outward unit normal on $\partial\Omega_a$. Therefore

$$\nabla\psi(y) = -|\nabla\psi(y)|\nu_{\rho^{-1}(r)}(y) \text{ for all } r \in \rho(Y) \text{ and } y \in \Sigma_{\rho^{-1}(r)}.$$

As a consequence, the Green-Riemann formula yields that, for all $r \in \rho(Y)$,

$$\begin{aligned} \int_{\Omega_{\rho^{-1}(r)}} \operatorname{div}(A_\Omega \nabla\psi)(y) dy &= \int_{\Sigma_{\rho^{-1}(r)}} A_\Omega(y) \nabla\psi(y) \cdot \nu_{\rho^{-1}(r)}(y) d\sigma_{\rho^{-1}(r)}(y) \\ &= - \int_{\Sigma_{\rho^{-1}(r)}} A_\Omega(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y) |\nabla\psi(y)| d\sigma_{\rho^{-1}(r)}(y) < 0, \end{aligned}$$

which ends the proof. \square

For all $x \in \overline{\Omega^*}$, set

$$\tilde{\psi}(x) = - \int_{|x|}^R F(r) dr. \quad (3.8)$$

The function $\tilde{\psi}$ is then radially symmetric and it vanishes on $\partial\Omega^* = \partial B_R$. From Lemma 3.3,

$$\tilde{\psi} > 0 \text{ in } \Omega^*,$$

$\tilde{\psi}$ is continuous on $\overline{\Omega^*}$, decreasing with respect to $|x|$ in $\overline{\Omega^*}$, and C^1 on $E \cup \{0\}$ (remember that $F(0) = 0$). Note that

$$\tilde{\psi} \in H_0^1(\Omega^*) \cap W^{1,\infty}(\Omega^*).$$

Moreover, the following statement holds true:

Lemma 3.4 *The function $\tilde{\psi}$ is of class C^2 in $E \cap \Omega^*$.*

Proof. By definition of $\tilde{\psi}$ and since $\hat{\Lambda}$ is C^1 in $E \cap \Omega^*$, it is enough to prove that the function

$$z : r \mapsto \int_{\Omega_{\rho^{-1}(r)}} \operatorname{div}(A_\Omega \nabla\psi)(x) dx = - \int_{\Omega_{\rho^{-1}(r)}} f(x) dx$$

is of class C^1 on $\rho(Y)$. It would actually be enough to prove that z is C^1 on $\rho(Y) \setminus \{R\}$.

Let r be fixed in $\rho(Y) = (0, \rho(a_{m-1})) \cup \dots \cup (\rho(a_2), \rho(a_1)) \cup (\rho(a_1), R]$ and let $\eta > 0$ be such that $[r - \eta, r] \subset \rho(Y)$. For $t \in (0, \eta)$, one has

$$\begin{aligned} z(r-t) - z(r) &= \int_{\{\rho^{-1}(r) < \psi(x) \leq \rho^{-1}(r-t)\}} f(x) dx \\ &= \int_{\rho^{-1}(r)}^{\rho^{-1}(r-t)} \left(\int_{\Sigma_a} f(y) |\nabla\psi(y)|^{-1} d\sigma_a(y) \right) da = \int_{\rho^{-1}(r)}^{\rho^{-1}(r-t)} h(a) da, \end{aligned}$$

where h is defined in (3.2). Since ρ^{-1} is of class C^1 on $\rho(Y)$ from Lemma 3.2 and since h is continuous on Y , it follows that

$$\frac{z(r-t) - z(r)}{-t} \rightarrow h(\rho^{-1}(r))(\rho^{-1})'(r) = -\frac{n\alpha_n r^{n-1} h(\rho^{-1}(r))}{g(\rho^{-1}(r))} \text{ as } t \rightarrow 0^+.$$

The same limit holds as $t \rightarrow 0^-$ for all $r \in \rho(Y) \setminus \{R\}$. Therefore, the function z is differentiable on $\rho(Y)$ and

$$z'(r) = -\frac{n\alpha_n r^{n-1} h(\rho^{-1}(r))}{g(\rho^{-1}(r))} \text{ for all } r \in \rho(Y).$$

Since ρ^{-1} is continuous on $[0, R]$, and g and h are continuous on Y , the function z is of class C^1 on $\rho(Y)$. That completes the proof of Lemma 3.4. \square

We now define a rearranged drift \widehat{v} and a rearranged potential \widehat{V} . For all $x \in E$, define

$$\widehat{v}(x) = \left(\frac{\int_{\Sigma_{\rho^{-1}(|x|)}} \omega(y)^2 \Lambda_\Omega(y)^{-1} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(|x|)}(y)}{\int_{\Sigma_{\rho^{-1}(|x|)}} \Lambda_\Omega(y)^{-1} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(|x|)}(y)} \right)^{1/2} e_r(x), \quad (3.9)$$

(remember that e_r is defined by (1.1)). The vector field \widehat{v} is then defined almost everywhere in $\overline{\Omega^*}$. Notice also that $|\widehat{v}|$ is radially symmetric, that $\widehat{v}(x)$ points in the direction $e_r(x)$ at each point $x \in E$, that \widehat{v} belongs to $L^\infty(\Omega^*)$ and that

$$\operatorname{ess\,inf}_\Omega |\omega| \leq \operatorname{ess\,inf}_{\Omega^*} |\widehat{v}| \leq \operatorname{ess\,sup}_{\Omega^*} |\widehat{v}| \leq \operatorname{ess\,sup}_\Omega |\omega| = \|\omega\|_{L^\infty(\Omega)}. \quad (3.10)$$

Furthermore, since Λ_Ω^{-1} and ω are continuous in $\overline{\Omega}$, the vector field \widehat{v} is continuous in E , and, as it was done for (3.6), it is easy to check that

$$\int_\Omega \omega(y)^2 \Lambda_\Omega(y)^{-1} dy = \int_{\Omega^*} |\widehat{v}(x)|^2 \widehat{\Lambda}(x)^{-1} dx. \quad (3.11)$$

Lastly, for all $x \in E$, define

$$\widehat{V}(x) = \frac{-\int_{\Sigma_{\rho^{-1}(|x|)}} V^-(y) |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(|x|)}(y)}{\int_{\Sigma_{\rho^{-1}(|x|)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(|x|)}(y)}, \quad (3.12)$$

where $V^-(y)$ denotes the negative part of $V(y)$, that is $V^-(y) = \max(0, -V(y))$. The function \widehat{V} is then defined almost everywhere in $\overline{\Omega^*}$. Observe that \widehat{V} is radially symmetric, nonpositive, belongs to $L^\infty(\Omega^*)$, is continuous in E , and that

$$-\|V\|_{L^\infty(\Omega)} \leq \min_{\overline{\Omega}}(-V^-) \leq \operatorname{ess\,inf}_{\Omega^*} \widehat{V} \leq \operatorname{ess\,sup}_{\Omega^*} \widehat{V} \leq 0. \quad (3.13)$$

3.2 Pointwise comparison between ψ and $\widetilde{\psi}$

The first interest of the spherical rearrangement which was defined in the previous subsection is that the functions ψ and $\widetilde{\psi}$ can be compared on the sets Σ_a and $\partial B_{\rho(a)}$. Namely, the function $\widetilde{\psi}$ satisfies the following key inequalities, which are summarized in Proposition 3.5 and Corollary 3.6:

Proposition 3.5 For any unit vector e of \mathbb{R}^n , the function

$$\begin{aligned}\tilde{\Psi} &: [0, M] \rightarrow \mathbb{R}_+ \\ a &\mapsto \tilde{\psi}(\rho(a)e)\end{aligned}$$

is continuous on $[0, M]$, differentiable on Y , and

$$\forall a \in Y, \quad \tilde{\Psi}'(a) \geq 1. \quad (3.14)$$

Before giving the proof of Proposition 3.5, let us first establish the following important corollary.

Corollary 3.6 For all $x \in \overline{\Omega^*}$,

$$\tilde{\psi}(x) \geq \rho^{-1}(|x|).$$

Proof. Since $\tilde{\Psi}$ is continuous on $[0, M]$ and differentiable on $[0, M]$ except on a finite set of points and since $\tilde{\Psi}(0) = 0$, the mean-value theorem and (3.14) show that $\tilde{\Psi}(a) \geq a$ for all $a \in [0, M]$, which means that $\tilde{\psi}(\rho(a)e) \geq a$ for all $a \in [0, M]$ and all unit vector e . Since $\tilde{\psi}$ is radially symmetric, Corollary 3.6 follows from Lemma 3.1. \square

Proof of Proposition 3.5. Let us first observe that the function $\tilde{\Psi}$ is differentiable on Y , from Lemma 3.2 and the fact that $\tilde{\psi}$ is C^1 in E (and even in $E \cup \{0\}$). Furthermore, since $\tilde{\psi}$ is radially symmetric, and decreasing with respect to the variable $|x|$ and since ρ is itself decreasing, it is enough to prove that

$$\forall x \in E, \quad |\rho'(\rho^{-1}(|x|))| \times |\nabla \tilde{\psi}(x)| \geq 1. \quad (3.15)$$

We will make use of the following inequality:

$$\forall x \in E, \quad \frac{\int_{\Sigma_{\rho^{-1}(|x|)}} A_{\Omega}(y) \nu_{\rho^{-1}(|x|)}(y) \cdot \nu_{\rho^{-1}(|x|)}(y) |\nabla \psi(y)| d\sigma_{\rho^{-1}(|x|)}(y)}{\int_{\Sigma_{\rho^{-1}(|x|)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(|x|)}(y)} \leq \widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2, \quad (3.16)$$

where one recalls that $\nu_{\rho^{-1}(|x|)}$ denotes the outward unit normal on $\partial\Omega_{\rho^{-1}(|x|)}$. We postpone the proof of (3.16) to the end of this subsection and go on in the proof of Proposition 3.5.

Fix $x \in E$ and set $r = |x|$. Since $\rho^{-1}(r) \in Y$, there exists $\eta > 0$ such that $\rho^{-1}(r-t) \in Y$ for all $t \in [0, \eta]$. For $t \in (0, \eta]$, the Cauchy-Schwarz inequality gives

$$\begin{aligned}\left(\frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} |\nabla \psi(y)| dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \right)^2 &\leq \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} \Lambda_{\Omega}(y)^{-1} dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \\ &\times \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} \Lambda_{\Omega}(y) |\nabla \psi(y)|^2 dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|}.\end{aligned} \quad (3.17)$$

The left-hand side of (3.17) is equal to

$$\left(\frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} |\nabla \psi(y)| dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \right)^2 = \left(\frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} |\nabla \psi(y)| dy}{\rho^{-1}(r-t) - \rho^{-1}(r)} \times \frac{\rho^{-1}(r-t) - \rho^{-1}(r)}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \right)^2.$$

By the co-area formula,

$$\lim_{t \rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} |\nabla \psi(y)| dy}{\rho^{-1}(r-t) - \rho^{-1}(r)} = \int_{\Sigma_{\rho^{-1}(r)}} d\sigma_{\rho^{-1}(r)}(y) = i(\rho^{-1}(r)),$$

and

$$\lim_{t \rightarrow 0^+} \frac{\rho^{-1}(r-t) - \rho^{-1}(r)}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} = \frac{1}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y)} = \frac{1}{n\alpha_n r^{n-1} |\rho'(\rho^{-1}(r))|}$$

from Lemma 3.2. By the isoperimetric inequality applied to $\Sigma_{\rho^{-1}(r)} = \partial\Omega_{\rho^{-1}(r)}$ and ∂B_r , one has

$$0 < n\alpha_n r^{n-1} \leq i(\rho^{-1}(r)) = \int_{\Sigma_{\rho^{-1}(r)}} d\sigma_{\rho^{-1}(r)}(y), \quad (3.18)$$

Therefore, one obtains

$$\begin{aligned} \lim_{t \rightarrow 0^+} \left(\frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} |\nabla \psi(y)| dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \right)^2 &\geq \left(\frac{i(\rho^{-1}(r))}{n\alpha_n r^{n-1}} \right)^2 \times \frac{1}{|\rho'(\rho^{-1}(r))|^2} \\ &\geq \frac{1}{|\rho'(\rho^{-1}(r))|^2}. \end{aligned} \quad (3.19)$$

The first factor of the right-hand side of (3.17) is equal to

$$\frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} \Lambda_{\Omega}(y)^{-1} dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} = \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} \Lambda_{\Omega}(y)^{-1} dy}{\rho^{-1}(r-t) - \rho^{-1}(r)} \times \frac{\rho^{-1}(r-t) - \rho^{-1}(r)}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|},$$

and the co-area formula therefore shows that

$$\lim_{t \rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} \Lambda_{\Omega}(y)^{-1} dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} = \widehat{\Lambda}(x)^{-1} \quad (3.20)$$

from (3.3) and (3.4). Finally, the coarea formula again implies that

$$\begin{aligned}
\lim_{t \rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} \Lambda_{\Omega}(y) |\nabla \psi(y)|^2 dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} &= \frac{\int_{\Sigma_{\rho^{-1}(r)}} \Lambda_{\Omega}(y) |\nabla \psi(y)| d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y)} \\
&\leq \frac{\int_{\Sigma_{\rho^{-1}(r)}} A_{\Omega}(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y) |\nabla \psi(y)| d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y)} \\
&\leq \widehat{\Lambda}(x) |\nabla \widetilde{\psi}(x)|^2
\end{aligned} \tag{3.21}$$

by (3.16).

Finally, (3.17), (3.19), (3.20) and (3.21) imply that

$$\frac{1}{|\rho'(\rho^{-1}(r))|^2} \leq \left(\frac{i(\rho^{-1}(r))}{n\alpha_n r^{n-1}} \right)^2 \times \frac{1}{|\rho'(\rho^{-1}(r))|^2} \leq |\nabla \widetilde{\psi}(x)|^2. \tag{3.22}$$

Therefore, inequality (3.15) holds and so does inequality (3.14). \square

Remark 3.7 Observe that (3.17), (3.20) and (3.21) together with the co-area formula yield

$$\frac{\int_{\Sigma_{\rho^{-1}(|x|)}} d\sigma_{\rho^{-1}(|x|)}(y)}{\int_{\Sigma_{\rho^{-1}(|x|)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(|x|)}(y)} = \lim_{t \rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(|x|)} \setminus \Omega_{\rho^{-1}(|x|-t)}} |\nabla \psi(y)| dy}{|\Omega_{\rho^{-1}(|x|)} \setminus \Omega_{\rho^{-1}(|x|-t)}|} \leq |\nabla \widetilde{\psi}(x)| \tag{3.23}$$

for all $x \in E$.

We now give the

Proof of (3.16). Fix $x \in E$ and call $r = |x|$. Notice first that, as was already observed, for all $y \in \partial\Omega_{\rho^{-1}(r)}$,

$$\nabla \psi(y) = -|\nabla \psi(y)| \nu_{\rho^{-1}(r)}(y).$$

The Green-Riemann formula and the choice of $\widetilde{\psi}$ therefore yield

$$\begin{aligned}
\int_{\Sigma_{\rho^{-1}(r)}} A_{\Omega}(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y) |\nabla \psi(y)| d\sigma_{\rho^{-1}(r)}(y) &= - \int_{\Omega_{\rho^{-1}(r)}} \operatorname{div}(A_{\Omega} \nabla \psi)(y) dy \\
&= -n\alpha_n r^{n-1} \widehat{\Lambda}(x) F(r) \\
&= n\alpha_n r^{n-1} \widehat{\Lambda}(x) |\nabla \widetilde{\psi}(x)|.
\end{aligned} \tag{3.24}$$

By Cauchy-Schwarz,

$$\begin{aligned}
i(\rho^{-1}(r))^2 &= \left(\int_{\Sigma_{\rho^{-1}(r)}} d\sigma_{\rho^{-1}(r)}(y) \right)^2 \\
&\leq \int_{\Sigma_{\rho^{-1}(r)}} A_{\Omega}(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y) |\nabla\psi(y)| d\sigma_{\rho^{-1}(r)}(y) \\
&\quad \times \int_{\Sigma_{\rho^{-1}(r)}} (A_{\Omega}(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y))^{-1} |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y) \\
&\leq \int_{\Sigma_{\rho^{-1}(r)}} A_{\Omega}(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y) |\nabla\psi(y)| d\sigma_{\rho^{-1}(r)}(y) \\
&\quad \times \int_{\Sigma_{\rho^{-1}(r)}} \Lambda_{\Omega}(y)^{-1} |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y) \\
&= \widehat{\Lambda}(x)^{-1} \times \int_{\Sigma_{\rho^{-1}(r)}} A_{\Omega}(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y) |\nabla\psi(y)| d\sigma_{\rho^{-1}(r)}(y) \\
&\quad \times \int_{\Sigma_{\rho^{-1}(r)}} |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y).
\end{aligned}$$

In other words,

$$\begin{aligned}
&\frac{\int_{\Sigma_{\rho^{-1}(r)}} A_{\Omega}(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y) |\nabla\psi(y)| d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y)} \\
&\leq \widehat{\Lambda}(x)^{-1} \times \left(\frac{\int_{\Sigma_{\rho^{-1}(r)}} A_{\Omega}(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y) |\nabla\psi(y)| d\sigma_{\rho^{-1}(r)}(y)}{i(\rho^{-1}(r))} \right)^2 \\
&= \left(\frac{n\alpha_n r^{n-1}}{i(\rho^{-1}(r))} \right)^2 \widehat{\Lambda}(x) |\nabla\tilde{\psi}(x)|^2
\end{aligned}$$

by (3.24). The isoperimetric inequality (3.18) ends the proof of (3.16). \square

3.3 A pointwise differential inequality for the rearranged data

In the previous subsection, we could compare the values of ψ and of its symmetrized function $\tilde{\psi}$. Here, we prove a partial differential inequality involving ψ and $\tilde{\psi}$, as well as the rearranged data $\widehat{\Lambda}$, \widehat{v} and \widehat{V} .

Proposition 3.8 *Let $\omega_0 \in \mathbb{R}_+$ and $x \in E \cap \Omega^*$. Then, there exists $y \in \Omega$ such that $\psi(y) = \rho^{-1}(|x|)$, that is $y \in \Sigma_{\rho^{-1}(|x|)}$, and*

$$\begin{aligned}
&-\operatorname{div}(\widehat{\Lambda}\nabla\tilde{\psi})(x) + \widehat{v}(x) \cdot \nabla\tilde{\psi}(x) - \omega_0|\nabla\tilde{\psi}(x)| + \widehat{V}(x)\tilde{\psi}(x) \\
&\leq -\operatorname{div}(A_{\Omega}\nabla\psi)(y) - |\omega(y)| \times |\nabla\psi(y)| - \omega_0|\nabla\psi(y)| + V(y)\psi(y).
\end{aligned}$$

Notice that $\widehat{v}(x) \cdot \nabla\tilde{\psi}(x) = -|\widehat{v}(x)| \times |\nabla\tilde{\psi}(x)|$.

Proof. Let $x \in E \cap \Omega^*$, $r = |x|$ and $\eta > 0$ such that $\overline{S_{r-\eta, r}} \subset E \cap \Omega^*$. As done in the proof of Proposition 3.5, the co-area formula and Cauchy-Schwarz inequality yield

$$\begin{aligned}
& \left(\lim_{t \rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} |\omega(y)| \times |\nabla \psi(y)| \, dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \right)^2 \\
&= \left(\frac{\int_{\Sigma_{\rho^{-1}(r)}} |\omega(y)| \, d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla \psi(y)|^{-1} \, d\sigma_{\rho^{-1}(r)}(y)} \right)^2 \\
&\leq \frac{\int_{\Sigma_{\rho^{-1}(r)}} \Lambda_{\Omega}(y)^{-1} \omega(y)^2 |\nabla \psi(y)|^{-1} \, d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla \psi(y)|^{-1} \, d\sigma_{\rho^{-1}(r)}(y)} \times \frac{\int_{\Sigma_{\rho^{-1}(r)}} \Lambda_{\Omega}(y) |\nabla \psi(y)| \, d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla \psi(y)|^{-1} \, d\sigma_{\rho^{-1}(r)}(y)}.
\end{aligned} \tag{3.25}$$

Using (3.16), one obtains

$$\begin{aligned}
& \frac{\int_{\Sigma_{\rho^{-1}(r)}} \Lambda_{\Omega}(y) |\nabla \psi(y)| \, d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla \psi(y)|^{-1} \, d\sigma_{\rho^{-1}(r)}(y)} \\
&\leq \frac{\int_{\Sigma_{\rho^{-1}(r)}} A_{\Omega}(y) \nu_{\rho^{-1}(r)}(y) \cdot \nu_{\rho^{-1}(r)}(y) |\nabla \psi(y)| \, d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla \psi(y)|^{-1} \, d\sigma_{\rho^{-1}(r)}(y)} \\
&\leq \widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2.
\end{aligned} \tag{3.26}$$

Finally, (3.25) and (3.26), together with definitions (3.3-3.4) and (3.9), give that

$$\lim_{t \rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} |\omega(y)| \times |\nabla \psi(y)| \, dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \leq |\widehat{v}(x)| \times |\nabla \tilde{\psi}(x)| = -\widehat{v}(x) \cdot \nabla \tilde{\psi}(x). \tag{3.27}$$

The last equality follows also from (3.8) and Lemma 3.3.

Remember also from (3.23) that

$$\lim_{t \rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} |\nabla \psi(y)| \, dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \leq |\nabla \tilde{\psi}(x)| = -e_r(x) \cdot \nabla \tilde{\psi}(x). \tag{3.28}$$

As far as V is concerned, for any fixed unit vector e in \mathbb{R}^n and for any $t \in (0, \eta)$, it follows from (3.12) and Lemma 3.2 that

$$\begin{aligned}
& \int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} V(y)\psi(y)dy \\
& \geq \int_{\rho^{-1}(r)}^{\rho^{-1}(r-t)} \left(\int_{\Sigma_a} (-V^-(y))\psi(y) |\nabla\psi(y)|^{-1} d\sigma_a(y) \right) da \\
& = - \int_{\rho^{-1}(r)}^{\rho^{-1}(r-t)} a \left(\int_{\Sigma_a} V^-(y) |\nabla\psi(y)|^{-1} d\sigma_a(y) \right) da \\
& = - \int_{r-t}^r \left(\int_{\Sigma_{\rho^{-1}(s)}} V^-(y) |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(s)}(y) \right) \frac{\rho^{-1}(s)ds}{|\rho'(\rho^{-1}(s))|} \\
& = -n\alpha_n \int_{r-t}^r s^{n-1}\rho^{-1}(s) \left(\frac{\int_{\Sigma_{\rho^{-1}(s)}} V^-(y) |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(s)}(y)}{\int_{\Sigma_{\rho^{-1}(s)}} |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(s)}(y)} \right) ds \\
& = n\alpha_n \int_{r-t}^r s^{n-1}\rho^{-1}(s) \widehat{V}(se) ds.
\end{aligned}$$

Moreover, the radial symmetry of \widehat{V} and $\widetilde{\psi}$ yields

$$\int_{S_{r-t,r}} \widehat{V}(y)\widetilde{\psi}(y)dy = n\alpha_n \int_{r-t}^r s^{n-1}\widehat{V}(se)\widetilde{\psi}(se)ds.$$

Corollary 3.6 and the facts that $|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}| = |S_{r-t,r}|$ and that $\widehat{V} \leq 0$ therefore show that

$$\frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} V(y)\psi(y)dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \geq \frac{\int_{S_{r-t,r}} \widehat{V}(y)\widetilde{\psi}(y)dy}{|S_{r-t,r}|}.$$

Since \widehat{V} and $\widetilde{\psi}$ are continuous in E and radially symmetric, one therefore obtains, together with the co-area formula,

$$\begin{aligned}
\frac{\int_{\Sigma_{\rho^{-1}(r)}} V(y)\psi(y)|\nabla\psi(y)|^{-1}d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla\psi(y)|^{-1}d\sigma_{\rho^{-1}(r)}(y)} &= \lim_{t \rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}} V(y)\psi(y)dy}{|\Omega_{\rho^{-1}(r)} \setminus \Omega_{\rho^{-1}(r-t)}|} \\
&\geq \widehat{V}(x)\widetilde{\psi}(x).
\end{aligned} \tag{3.29}$$

Let now t be any real number in $(0, \eta)$. Since $\widetilde{\psi}$ (resp. $\widehat{\Lambda}$) is radially symmetric, and C^2 (resp. C^1) on $\overline{S_{r-t,r}} \subset E \cap \Omega^*$, the Green Riemann formula gives

$$\begin{aligned}
\int_{S_{r-t,r}} \operatorname{div}(\widehat{\Lambda}\nabla\widetilde{\psi})(y)dy &= \int_{\partial S_{r-t,r}} \widehat{\Lambda}(y)\nabla\widetilde{\psi}(y) \cdot \nu(y)d\sigma(y) \\
&= n\alpha_n[r^{n-1}G(r)F(r) - (r-t)^{n-1}G(r-t)F(r-t)],
\end{aligned} \tag{3.30}$$

where $d\sigma$ and ν here denote the superficial measure on $\partial S_{r-t,r}$ and the outward unit normal to $S_{r-t,r}$, and G and F were defined in (3.3) and (3.7). By definition of F , one gets that

$$\int_{S_{r-t,r}} \operatorname{div}(\widehat{\Lambda}\nabla\tilde{\psi})(y)dy = \int_{\Omega_{\rho^{-1}(r)}\setminus\Omega_{\rho^{-1}(r-t)}} \operatorname{div}(A_\Omega\nabla\psi)(y)dy,$$

whence

$$\begin{aligned} & \frac{\int_{\Sigma_{\rho^{-1}(r)}} \operatorname{div}(A_\Omega\nabla\psi)(y)|\nabla\psi(y)|^{-1}d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla\psi(y)|^{-1}d\sigma_{\rho^{-1}(r)}(y)} \\ &= \lim_{t\rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(r)}\setminus\Omega_{\rho^{-1}(r-t)}} \operatorname{div}(A_\Omega\nabla\psi)(y)dy}{|\Omega_{\rho^{-1}(r)}\setminus\Omega_{\rho^{-1}(r-t)}|} = \operatorname{div}(\widehat{\Lambda}\nabla\tilde{\psi})(x) \end{aligned} \quad (3.31)$$

since $|S_{r-t,r}| = |\Omega_{\rho^{-1}(r)}\setminus\Omega_{\rho^{-1}(r-t)}|$.

It follows from (3.27), (3.28), (3.29) and (3.31) that

$$\begin{aligned} & \lim_{t\rightarrow 0^+} \frac{\int_{\Omega_{\rho^{-1}(r)}\setminus\Omega_{\rho^{-1}(r-t)}} [\operatorname{div}(A_\Omega\nabla\psi)(y) + |\omega(y)| \times |\nabla\psi(y)| + \omega_0|\nabla\psi(y)| - V(y)\psi(y)] dy}{|\Omega_{\rho^{-1}(r)}\setminus\Omega_{\rho^{-1}(r-t)}|} \\ & \leq \operatorname{div}(\widehat{\Lambda}\nabla\tilde{\psi})(x) - \widehat{v}(x) \cdot \nabla\tilde{\psi}(x) + \omega_0|\nabla\tilde{\psi}(x)| - \widehat{V}(x)\tilde{\psi}(x). \end{aligned}$$

To finish the proof, pick any sequence of positive numbers $(\varepsilon_l)_{l\in\mathbb{N}}$ such that $\varepsilon_l \rightarrow 0$ as $l \rightarrow +\infty$. Since ψ is C^2 in Ω , since A_Ω is C^1 in Ω (and even in $\overline{\Omega}$) and ω and V are continuous in Ω (and even in $\overline{\Omega}$), the previous inequality provides the existence of a sequence of positive numbers $(t_l)_{l\in\mathbb{N}} \in (0, \eta)$ such that $t_l \rightarrow 0$ as $l \rightarrow +\infty$, and a sequence of points

$$y_l \in \overline{\Omega_{\rho^{-1}(r)}\setminus\Omega_{\rho^{-1}(r-t_l)}} \subset \overline{\Omega_{\rho^{-1}(r)}} \subset \Omega$$

such that

$$\begin{aligned} & \operatorname{div}(A_\Omega\nabla\psi)(y_l) + |\omega(y_l)| \times |\nabla\psi(y_l)| + \omega_0|\nabla\psi(y_l)| - V(y_l)\psi(y_l) \\ & \leq \operatorname{div}(\widehat{\Lambda}\nabla\tilde{\psi})(x) - \widehat{v}(x) \cdot \nabla\tilde{\psi}(x) + \omega_0|\nabla\tilde{\psi}(x)| - \widehat{V}(x)\tilde{\psi}(x) + \varepsilon_l. \end{aligned}$$

Since $\rho^{-1}(r) \leq \psi(y_l) \leq \rho^{-1}(r - t_l)$ and ρ^{-1} is continuous, the points y_l converge, up to the extraction of some subsequence, to a point $y \in \Sigma_{\rho^{-1}(r)}$ such that

$$\begin{aligned} & \operatorname{div}(A_\Omega\nabla\psi)(y) + |\omega(y)| \times |\nabla\psi(y)| + \omega_0|\nabla\psi(y)| - V(y)\psi(y) \\ & \leq \operatorname{div}(\widehat{\Lambda}\nabla\tilde{\psi})(x) - \widehat{v}(x) \cdot \nabla\tilde{\psi}(x) + \omega_0|\nabla\tilde{\psi}(x)| - \widehat{V}(x)\tilde{\psi}(x), \end{aligned}$$

which is the conclusion of Proposition 3.8. \square

Corollary 3.9 *If there are $\omega_0 \geq 0$ and $\mu \geq 0$ such that*

$$-\operatorname{div}(A_\Omega\nabla\psi)(y) - |\omega(y)| \times |\nabla\psi(y)| - \omega_0|\nabla\psi(y)| + V(y)\psi(y) \leq \mu\psi(y) \quad \text{for all } y \in \Omega,$$

then

$$-\operatorname{div}(\widehat{\Lambda}\nabla\tilde{\psi})(x) + \widehat{v}(x) \cdot \nabla\tilde{\psi}(x) - \omega_0|\nabla\tilde{\psi}(x)| + \widehat{V}(x)\tilde{\psi}(x) \leq \mu\tilde{\psi}(x) \quad \text{for all } x \in E \cap \Omega^*.$$

Proof. It follows immediately from Corollary 3.6 and Proposition 3.8. \square

3.4 An integral inequality for the rearranged data

A consequence of the pointwise comparisons which were established in the previous subsections is the following integral comparison result:

Proposition 3.10 *With the previous notations, assume that, for some $(\omega_0, \mu) \in \mathbb{R}^2$,*

$$-\operatorname{div}(\widehat{\Lambda}\nabla\tilde{\psi})(x) + \widehat{v}(x) \cdot \nabla\tilde{\psi}(x) - \omega_0|\nabla\tilde{\psi}(x)| + \widehat{V}(x)\tilde{\psi}(x) \leq \mu\tilde{\psi}(x) \quad \text{for all } x \in E \cap \Omega^*. \quad (3.32)$$

Fix a unit vector $e \in \mathbb{R}^n$. For all $r \in [0, R]$, define

$$H(r) = \int_0^r |\widehat{v}(re)| \widehat{\Lambda}(re)^{-1} dr \quad (3.33)$$

and, for all $x \in \overline{\Omega^*}$, let

$$U(x) = H(|x|). \quad (3.34)$$

Then, the following integral inequality is valid:

$$\int_{\Omega^*} \left[\widehat{\Lambda}(x)|\nabla\tilde{\psi}(x)|^2 - \omega_0|\nabla\tilde{\psi}(x)|\tilde{\psi}(x) + \widehat{V}(x)\tilde{\psi}(x)^2 \right] e^{-U(x)} dx \leq \mu \int_{\Omega^*} \tilde{\psi}(x)^2 e^{-U(x)} dx. \quad (3.35)$$

Proof. Note first that, since $|\widehat{v}|$ and $\widehat{\Lambda}$ are radially symmetric and since $|\widehat{v}| \in L^\infty(\Omega^*)$ and $\widehat{\Lambda}$ satisfies (3.5), the function H is well-defined and continuous in $[0, R]$. Furthermore, its definition is independent from the choice of e . The radially symmetric function U is then continuous in $\overline{\Omega^*}$ and, since the radially symmetric functions $\widehat{v} = |\widehat{v}|e_r$ and $1/\widehat{\Lambda}$ are (at least) continuous in E , the function U is of class C^1 in E and

$$\nabla U(x) = \widehat{\Lambda}(x)^{-1}\widehat{v}(x) \quad \text{for all } x \in E. \quad (3.36)$$

Observe also that the integrals in (3.35) are all well-defined since $\tilde{\psi} \in H_0^1(\Omega^*)$ and $\widehat{\Lambda}, \widehat{V}, U \in L^\infty(\Omega^*)$ (even, $U \in C(\overline{\Omega^*})$).

Now, recall that the set of critical values of ψ is $Z = \{a_1, \dots, a_m\}$ with

$$0 < a_1 < \dots < a_m = M$$

and remember that the function ρ defined in Subsection 3.1 is continuous and decreasing from $[0, M]$ onto $[0, R]$, from Lemma 3.1. Fix $j \in \{1, \dots, m-1\}$ and r, r' such that

$$0 \leq \rho(a_{j+1}) < r < r' < \rho(a_j) < R.$$

Multiplying (3.32) by the nonnegative function $\tilde{\psi}e^{-U}$ and integrating over $S_{r,r'}$ yields

$$\begin{aligned} & \int_{S_{r,r'}} \left[-\operatorname{div}(\widehat{\Lambda}\nabla\tilde{\psi})(x) + \widehat{v}(x) \cdot \nabla\tilde{\psi}(x) - \omega_0|\nabla\tilde{\psi}(x)| + \widehat{V}(x)\tilde{\psi}(x) \right] \tilde{\psi}(x)e^{-U(x)} dx \\ & \leq \mu \int_{S_{r,r'}} \tilde{\psi}(x)^2 e^{-U(x)} dx. \end{aligned} \quad (3.37)$$

Notice that all integrals above are well-defined since $\tilde{\psi}$ in C^2 in $E \cap \Omega^*$, $\widehat{\Lambda}$ is C^1 in $E \cap \Omega^*$, \widehat{v} , \widehat{V} are continuous in E , U is continuous in $\overline{\Omega^*}$ and $\overline{S_{r,r'}} \subset E \cap \Omega^*$. Furthermore, as in (3.30), the Green-Riemann formula yields

$$\begin{aligned} & \int_{S_{r,r'}} -\operatorname{div}(\widehat{\Lambda} \nabla \tilde{\psi})(x) \tilde{\psi}(x) e^{-U(x)} dx \\ &= \int_{S_{r,r'}} \widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 e^{-U(x)} dx - \int_{S_{r,r'}} \widehat{\Lambda}(x) \tilde{\psi}(x) \nabla \tilde{\psi}(x) \cdot \nabla U(x) e^{-U(x)} dx \\ & \quad - n\alpha_n (r')^{n-1} G(r') F(r') \tilde{\psi}(r'e) e^{-H(r')} + n\alpha_n r^{n-1} G(r) F(r) \tilde{\psi}(re) e^{-H(r)}. \end{aligned}$$

By (3.36), it follows then that

$$\begin{aligned} & \int_{S_{r,r'}} \left[-\operatorname{div}(\widehat{\Lambda} \nabla \tilde{\psi})(x) + \widehat{v}(x) \cdot \nabla \tilde{\psi}(x) - \omega_0 |\nabla \tilde{\psi}(x)| + \widehat{V}(x) \tilde{\psi}(x) \right] \tilde{\psi}(x) e^{-U(x)} dx \\ &= \int_{S_{r,r'}} \left[\widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 - \omega_0 |\nabla \tilde{\psi}(x)| \tilde{\psi}(x) + \widehat{V}(x) \tilde{\psi}(x)^2 \right] e^{-U(x)} dx \\ & \quad - n\alpha_n (r')^{n-1} G(r') F(r') \tilde{\psi}(r'e) e^{-H(r')} + n\alpha_n r^{n-1} G(r) F(r) \tilde{\psi}(re) e^{-H(r)}. \end{aligned} \quad (3.38)$$

On the other hand, for all $s \in \rho(Y)$,

$$n\alpha_n s^{n-1} F(s) G(s) = \int_{\Omega_{\rho^{-1}(s)}} \operatorname{div}(A_\Omega \nabla \psi)(x) dx,$$

by (3.7). The function

$$s \mapsto I(s) = n\alpha_n s^{n-1} F(s) G(s),$$

which was a priori defined only in $\rho(Y)$, can then be extended continuously in $[0, R]$ from the results in Lemma 3.1 and since $\operatorname{div}(A_\Omega \nabla \psi) = -f$ is bounded in Ω . The continuous extension of I in $[0, R]$ is still called I . Passing to the limit as $r \rightarrow \rho(a_{j+1})^+$ and $r' \rightarrow \rho(a_j)^-$ in (3.37) and (3.38) yields, for each $j \in \{1, \dots, m-1\}$,

$$\begin{aligned} & \int_{S_{\rho(a_{j+1}), \rho(a_j)}} \left[\widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 - \omega_0 |\nabla \tilde{\psi}(x)| \tilde{\psi}(x) + \widehat{V}(x) \tilde{\psi}(x)^2 \right] e^{-U(x)} dx \\ & - I(\rho(a_j)) \tilde{\psi}(\rho(a_j)e) e^{-H(\rho(a_j))} + I(\rho(a_{j+1})) \tilde{\psi}(\rho(a_{j+1})e) e^{-H(\rho(a_{j+1}))} \\ & \leq \mu \int_{S_{\rho(a_{j+1}), \rho(a_j)}} \tilde{\psi}(x)^2 e^{-U(x)} dx. \end{aligned} \quad (3.39)$$

Once again, all integrals above are well-defined. Arguing similarly in the spherical shell $S_{\rho(a_1), R}$ and since $\psi(Re) = 0$, one obtains

$$\begin{aligned} & \int_{S_{\rho(a_1), R}} \left[\widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 - \omega_0 |\nabla \tilde{\psi}(x)| \tilde{\psi}(x) + \widehat{V}(x) \tilde{\psi}(x)^2 \right] e^{-U(x)} dx \\ & + I(\rho(a_1)) \tilde{\psi}(\rho(a_1)e) e^{-H(\rho(a_1))} \\ & \leq \mu \int_{S_{\rho(a_1), R}} \tilde{\psi}(x)^2 e^{-U(x)} dx. \end{aligned} \quad (3.40)$$

Summing up (3.39) for all $1 \leq j \leq m-1$ and (3.40) and using the fact that $I(\rho(a_m)) = I(0) = 0$ yield (3.35). \square

Corollary 3.11 *If there are $\omega_0 \geq 0$ and $\mu \geq 0$ such that*

$$-\operatorname{div}(A_\Omega \nabla \psi)(y) - |\omega(y)| \times |\nabla \psi(y)| - \omega_0 |\nabla \psi(y)| + V(y)\psi(y) \leq \mu \psi(y) \quad \text{for all } y \in \Omega,$$

then, under the notation of Proposition 3.10,

$$\int_{\Omega^*} \left[\widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 - \omega_0 |\nabla \tilde{\psi}(x)| \tilde{\psi}(x) + \widehat{V}(x) \tilde{\psi}(x)^2 \right] e^{-U(x)} dx \leq \mu \int_{\Omega^*} \tilde{\psi}(x)^2 e^{-U(x)} dx.$$

Proof. It follows immediately from Corollary 3.9 and Proposition 3.10. \square

We complete this section by two remarks which proceed from the previous results and provide comparisons between some norms of the function ψ and its symmetrization $\tilde{\psi}$.

Remark 3.12 The calculations of the previous subsections (see in particular the proof of Proposition 3.8) and Corollary 3.6 imply that, for any nondecreasing function $\Theta : [0, +\infty) \rightarrow [0, +\infty)$,

$$\int_{\Omega_{\rho^{-1}(s)} \setminus \Omega_{\rho^{-1}(r)}} \Theta(\psi(y)) dy = \int_{S_{r,s}} \Theta(\rho^{-1}(|x|)) dx \leq \int_{S_{r,s}} \Theta(\tilde{\psi}(x)) dx$$

for all $0 < r < s \leq R$ such that $[r, s] \subset \rho(Y)$, and then for all $0 \leq r < s \leq R$ from Lebesgue's dominated convergence theorem. In particular,

$$\int_{\Omega} (\psi(y))^p dy \leq \int_{\Omega^*} (\tilde{\psi}(x))^p dx$$

for all $0 \leq p < +\infty$. Remember also (as an immediate consequence of Corollary 3.6) that

$$\max_{\Omega^*} \tilde{\psi} = \tilde{\psi}(0) \geq \rho^{-1}(0) = \max_{\Omega} \psi.$$

Remark 3.13 For any $0 < r < s \leq R$ such that $[r, s] \subset \rho(Y)$, it follows from the co-area formula and a change of variables that

$$\begin{aligned} & \int_{\Omega_{\rho^{-1}(s)} \setminus \Omega_{\rho^{-1}(r)}} A_\Omega(y) \nabla \psi(y) \cdot \nabla \psi(y) dy \\ &= n \alpha_n \int_r^s t^{n-1} \times \left(\frac{\int_{\Sigma_{\rho^{-1}(t)}} A_\Omega(y) \nu_{\rho^{-1}(t)}(y) \cdot \nu_{\rho^{-1}(t)}(y) |\nabla \psi(y)| d\sigma_{\rho^{-1}(t)}(y)}{\int_{\Sigma_{\rho^{-1}(t)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(t)}(y)} \right) dt \\ &\leq \int_{S_{r,s}} \widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 dx, \end{aligned}$$

where the last inequality is due to (3.16). As usual, Lebesgue's dominated convergence theorem then yields

$$\int_{\Omega_{\rho^{-1}(s)} \setminus \Omega_{\rho^{-1}(r)}} A_\Omega(y) \nabla \psi(y) \cdot \nabla \psi(y) dy \leq \int_{S_{r,s}} \widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 dx$$

for all $0 \leq r < s \leq R$, whence

$$\int_{\Omega} A_{\Omega}(y) \nabla \psi(y) \cdot \nabla \psi(y) dy \leq \int_{\Omega^*} \widehat{\Lambda}(x) |\nabla \widetilde{\psi}(x)|^2 dx.$$

As a consequence,

$$\|\nabla \psi\|_{L^2(\Omega, \mathbb{R}^n)} \leq \sqrt{\frac{M_{\Lambda}}{m_{\Lambda}}} \times \|\nabla \widetilde{\psi}\|_{L^2(\Omega^*, \mathbb{R}^n)}$$

from (3.1) and (3.5), where $M_{\Lambda} = \max_{\overline{\Omega}} \Lambda_{\Omega}$ and $m_{\Lambda} = \min_{\overline{\Omega}} \Lambda_{\Omega}$. In particular, $\|\nabla \psi\|_{L^2(\Omega, \mathbb{R}^n)} \leq \|\nabla \widetilde{\psi}\|_{L^2(\Omega^*, \mathbb{R}^n)}$ if Λ is constant in Ω .

With the same notations as above, it follows from (3.23) that

$$\begin{aligned} \int_{\Omega_{\rho^{-1}(s)} \setminus \Omega_{\rho^{-1}(r)}} |\nabla \psi(y)| dy &= n \alpha_n \int_r^s t^{n-1} \times \left(\frac{\int_{\Sigma_{\rho^{-1}(t)}} d\sigma_{\rho^{-1}(t)}(y)}{\int_{\Sigma_{\rho^{-1}(t)}} |\nabla \psi(y)|^{-1} d\sigma_{\rho^{-1}(t)}(y)} \right) dt \\ &\leq \int_{S_{r,s}} |\nabla \widetilde{\psi}(x)| dx. \end{aligned}$$

The inequality then holds for all $0 \leq r < s \leq R$, whence

$$\|\nabla \psi\|_{L^1(\Omega, \mathbb{R}^n)} \leq \|\nabla \widetilde{\psi}\|_{L^1(\Omega^*, \mathbb{R}^n)}.$$

4 Improved inequalities when Ω is not a ball

Throughout this section, we assume that $\Omega \in \mathcal{C}$ is not a ball. Fix real numbers $\alpha \in (0, 1)$, $N > 0$ and $\delta > 0$. Denote by $E_{\alpha, N, \delta}(\Omega)$ the set of all functions $\psi \in C^{1, \alpha}(\overline{\Omega})$, positive in Ω , vanishing on $\partial\Omega$, such that

$$\|\psi\|_{C^{1, \alpha}(\overline{\Omega})} \leq N \text{ and } \psi(x) \geq \delta \times d(x, \partial\Omega) \text{ for all } x \in \overline{\Omega}, \quad (4.1)$$

where $d(x, \partial\Omega) = \min_{y \in \partial\Omega} |x - y|$ denotes the distance between x and $\partial\Omega$, and we set

$$\|\psi\|_{C^{1, \alpha}(\overline{\Omega})} = \|\psi\|_{L^{\infty}(\Omega)} + \|\nabla \psi\|_{L^{\infty}(\Omega, \mathbb{R}^n)} + \sup_{z \neq z' \in \Omega} \frac{|\nabla \psi(z) - \nabla \psi(z')|}{|z - z'|^{\alpha}}.$$

Notice that, for each $\psi \in E_{\alpha, N, \delta}(\Omega)$, one has

$$\nabla \psi(y) \cdot \nu(y) = -|\nabla \psi(y)| \leq -\delta \text{ for all } y \in \partial\Omega,$$

where ν denotes the outward unit normal on $\partial\Omega$.

Our goal here is to prove stronger versions of Corollary 3.6 and Corollary 3.11, using the fact that Ω is not a ball. In the sequel, unless explicitly mentioned, all the constants only depend on some of the data Ω , n , α , N and δ .

Denote again by R the radius of Ω^* , so that $\Omega^* = B_R$. First, the isoperimetric inequality yields the existence of $\beta = \beta(\Omega, n) > 0$ such that

$$\text{area}(\partial\Omega) = \int_{\partial\Omega} d\sigma_{\partial\Omega}(y) \geq (1 + \beta)n\alpha_n R^{n-1}, \quad (4.2)$$

where the left-hand side is the $(n - 1)$ -dimensional measure of $\partial\Omega$.

For all $\gamma > 0$, define

$$U_\gamma = \{x \in \overline{\Omega}, d(x, \partial\Omega) \leq \gamma\}.$$

Since $\partial\Omega$ is of class C^2 , there exists $\gamma_1 = \gamma_1(\Omega) > 0$ such that $\overline{\Omega} \setminus \overline{U_{\gamma_1}} \neq \emptyset$ and the segments $[y, y - \gamma_1\nu(y)]$ are included in $\overline{\Omega}$ and pairwise disjoint when y describes $\partial\Omega$. Thus, for all $\gamma \in (0, \gamma_1]$, the segments $[y, y - \gamma\nu(y)]$ (resp. $(y, y - \gamma\nu(y))$) describe the set U_γ (resp. $\{x \in \Omega, d(x, \partial\Omega) \leq \gamma\}$) as y describes $\partial\Omega$.

Lemma 4.1 *There exists a constant*

$$\gamma_2 = \gamma_2(\Omega, \alpha, N, \delta) \in (0, \gamma_1]$$

such that, for all $\psi \in E_{\alpha, N, \delta}(\Omega)$, one has:

$$|\nabla\psi| \geq \frac{\delta}{2} \text{ in } U_{\gamma_2},$$

$$\nabla\psi(y - r\nu(y)) \cdot \nu(y) \leq -\frac{\delta}{2} \text{ for all } y \in \partial\Omega \text{ and } r \in [0, \gamma_2],$$

and

$$\psi \geq \frac{\gamma\delta}{2} \text{ in } \overline{\Omega} \setminus \overline{U_\gamma} \text{ for all } \gamma \in [0, \gamma_2].$$

Proof. Assume that the conclusion of the lemma does not hold. Then there exists a sequence of positive numbers $(\gamma^l)_{l \in \mathbb{N}} \rightarrow 0$ and a sequence of functions $(\psi^l)_{l \in \mathbb{N}} \in E_{\alpha, N, \delta}(\Omega)$ such that one of the three following cases occur:

- 1) either for each $l \in \mathbb{N}$, there is a point $x^l \in U_{\gamma^l}$ such that $|\nabla\psi^l(x^l)| < \delta/2$;
- 2) or for each $l \in \mathbb{N}$, there are a point $y^l \in \partial\Omega$ and a number $r^l \in [0, \gamma^l]$ such that $\nabla\psi^l(y^l - r^l\nu(y^l)) \cdot \nu(y^l) > -\delta/2$;
- 3) or for each $l \in \mathbb{N}$, there is a point $x^l \in \overline{\Omega} \setminus \overline{U_{\gamma^l}}$ such that $\psi^l(x^l) < \gamma^l\delta/2$.

Observe first that, by Ascoli theorem, up to the extraction of a subsequence, there exists a function $\psi \in E_{\alpha, N, \delta}(\Omega)$ such that

$$\psi^l \rightarrow \psi \text{ in } C^1(\overline{\Omega}) \text{ as } l \rightarrow +\infty.$$

If case 1) occurs, then, up to some subsequence, one can assume without loss of generality that $x^l \rightarrow x \in \partial\Omega$ as $l \rightarrow +\infty$, and one obtains $|\nabla\psi(x)| \leq \delta/2$, which is impossible since $\psi \in E_{\alpha, N, \delta}(\Omega)$ and $\delta > 0$. Similarly, if case 2) occurs, $y^l \rightarrow y \in \partial\Omega$ as $l \rightarrow +\infty$ up to a subsequence, and one has $-|\nabla\psi(y)| = \nabla\psi(y) \cdot \nu(y) \geq -\delta/2$, which is also impossible.

Therefore, only case 3) can occur. For each $l \in \mathbb{N}$, let $y^l \in \partial\Omega$ be such that

$$d^l := |x^l - y^l| = d(x^l, \partial\Omega) \geq \gamma^l > 0.$$

Up to extraction of some subsequence, one has

$$x^l \rightarrow x \in \overline{\Omega} \text{ as } l \rightarrow +\infty \text{ and } \psi(x) \leq 0$$

by passing to the limit as $l \rightarrow +\infty$ in the inequality $\psi^l(x^l) < \gamma^l \delta / 2$. Since

$$\psi(x) \geq \delta d(x, \partial\Omega) > 0 \text{ for all } x \in \Omega,$$

it follows that $x \in \partial\Omega$, whence $|x^l - y^l| \rightarrow 0$ and $y^l \rightarrow x$ as $l \rightarrow +\infty$. On the one hand, the mean value theorem implies that

$$\frac{\psi^l(x^l) - \psi^l(y^l)}{|x^l - y^l|} = \nabla\psi^l(z^l) \cdot \frac{x^l - y^l}{|x^l - y^l|} \rightarrow -\nabla\psi(x) \cdot \nu(x) = |\nabla\psi(x)| \text{ as } l \rightarrow +\infty,$$

where z^l is a point lying on the segment between x^l and y^l (whence, $z^l \rightarrow x$ as $l \rightarrow +\infty$). On the other hand, since $\psi^l = 0$ on $\partial\Omega$,

$$\frac{\psi^l(x^l) - \psi^l(y^l)}{|x^l - y^l|} = \frac{\psi^l(x^l)}{d^l} < \frac{\gamma^l \delta}{2\gamma^l} = \frac{\delta}{2}.$$

Hence, $|\nabla\psi(x)| \leq \delta/2$ at the limit as $l \rightarrow +\infty$, which contradicts the positivity of δ and the fact that $\psi \in E_{\alpha, N, \delta}(\Omega)$.

Therefore, case 3) is ruled out too and the proof of Lemma 4.1 is complete. \square

In the sequel, we assume that $\psi \in E_{\alpha, N, \delta}(\Omega)$ and is analytic in Ω . The data A_Ω , Λ_Ω , ω and V are as in Section 3. We assume that

$$\operatorname{div}(A_\Omega \nabla \psi) = -f \text{ in } \Omega,$$

where f is a non zero polynomial, and we use the same sets Z , Y , E , Ω_a , Σ_a and the same functions ρ , $\tilde{\psi}$, $\hat{\Lambda}$, \hat{v} , \hat{V} and U as in Section 3.

Lemma 4.2 *Assume that $\Omega \in \mathcal{C}$ is not a ball. Then there exists a constant $a_0 = a_0(\Omega, n, \alpha, N, \delta) > 0$ only depending on Ω , n , α , N and δ such that $[0, a_0] \subset Y$ and*

$$i(a) = \int_{\Sigma_a} d\sigma_a(y) = \operatorname{area}(\Sigma_a) \geq \left(1 + \frac{\beta}{2}\right) n\alpha_n R^{n-1}$$

for all $a \in [0, a_0]$, where $\beta = \beta(\Omega, n) > 0$ was given in (4.2).

Proof. Let $\gamma_2 = \gamma_2(\Omega, \alpha, N, \delta) > 0$ be as in Lemma 4.1. Since $0 < \gamma_2 \leq \gamma_1$, it follows that $\overline{\Omega \setminus U_{\gamma_2}} \neq \emptyset$, and

$$\psi \geq \frac{\gamma_2 \delta}{2} > 0 \text{ in } \overline{\Omega \setminus U_{\gamma_2}}$$

from Lemma 4.1. Therefore,

$$M = \max_{\overline{\Omega}} \psi \geq \frac{\gamma_2 \delta}{2}$$

and

$$\Sigma_a = \{x \in \overline{\Omega}, \psi(x) = a\} \subset U_{\gamma_2} \text{ for all } a \in \left[0, \frac{\gamma_2 \delta}{4}\right]. \quad (4.3)$$

Call

$$a'_0 = a'_0(\Omega, \alpha, N, \delta) = \frac{\gamma_2 \delta}{4} > 0.$$

From Lemma 4.1 and the assumptions made on ψ and $\partial\Omega$, one has that $|\nabla\psi| \neq 0$ everywhere on the C^2 hypersurface Σ_a for all $a \in [0, a'_0]$. Thus,

$$[0, a'_0] \subset Y. \quad (4.4)$$

On the other hand, for all $y \in \partial\Omega$, the segment $[y, y - \gamma_2\nu(y)]$ is included in $\overline{\Omega}$ and there exists $\theta \in [0, 1]$ such that

$$\psi(y - \gamma_2\nu(y)) = \underbrace{\psi(y)}_{=0} - \gamma_2\nu(y) \cdot \nabla\psi(y - \theta\gamma_2\nu(y)) \geq \frac{\gamma_2\delta}{2},$$

again from Lemma 4.1. Actually, more precisely, for each $y \in \partial\Omega$, the function

$$\kappa : [0, \gamma_2] \rightarrow \mathbb{R}, \quad s \mapsto \psi(y - s\nu(y))$$

is differentiable and $\kappa'(s) \geq \delta/2$ for all $s \in [0, \gamma_2]$. It follows that, for all $a \in [0, a'_0]$ and $y \in \partial\Omega$, there exists a unique point

$$\phi_a(y) \in [y, y - \gamma_2\nu(y)] \cap \Sigma_a.$$

Moreover, for such a choice of a , the map ϕ_a is one-to-one since the segments $[y, y - \gamma_2\nu(y)]$ are pairwise disjoint (and describe U_{γ_2}) when y describes $\partial\Omega$ (because $0 < \gamma_2 \leq \gamma_1$). Lastly,

$$\Sigma_a = \{\phi_a(y), y \in \partial\Omega\} \quad (4.5)$$

from (4.3).

Let us now prove that the area of Σ_a is close to that of $\partial\Omega$ for $a \geq 0$ small enough. To do so, call

$$\mathcal{B} = \{x' = (x_1, \dots, x_{n-1}), |x'| < 1\}$$

and represent $\partial\Omega$ by a finite number of C^2 maps y^1, \dots, y^p (for some $p = p(\Omega) \in \mathbb{N}^*$) defined in $\overline{\mathcal{B}}$, depending only on Ω , and for which

$$\partial_1 y^j(x') \times \dots \times \partial_{n-1} y^j(x') \neq 0 \text{ for all } 1 \leq j \leq p \text{ and } x' \in \overline{\mathcal{B}}.$$

Here,

$$\partial_i y^j(x') = (\partial_{x_i} y_1^j(x'), \dots, \partial_{x_i} y_n^j(x')) \in \mathbb{R}^n$$

for each $1 \leq i \leq n-1$, $1 \leq j \leq p$ and $x' \in \overline{\mathcal{B}}$, where $y^j(x') = (y_1^j(x'), \dots, y_n^j(x'))$. The maps y^j are chosen so that

$$\partial\Omega = \{y^j(x'), 1 \leq j \leq p, x' \in \overline{\mathcal{B}}\}.$$

For each $a \in [0, a'_0]$ and for each $1 \leq j \leq p$, there exists then a map $t_a^j : \overline{\mathcal{B}} \rightarrow [0, \gamma_2]$ such that

$$\psi(y^j(x') - t_a^j(x')\nu(y^j(x'))) = a \text{ for all } x' \in \overline{\mathcal{B}}, \quad (4.6)$$

and

$$\Sigma_a = \{y^j(x') - t_a^j(x')\nu(y^j(x')), 1 \leq j \leq p, x' \in \overline{\mathcal{B}}\}.$$

Namely,

$$y^j(x') - t_a^j(x')\nu(y^j(x')) = \phi_a(y^j(x')) \text{ for all } 1 \leq j \leq p \text{ and } x' \in \overline{\mathcal{B}}.$$

From the arguments above, each real number $t_a^j(x')$ is then uniquely determined, and $t_0^j(x') = 0$. Since the functions ψ , y^j and $\nu \circ y^j$ (for all $1 \leq j \leq p$) are (at least) of class C^1 (respectively in $\overline{\Omega}$, $\overline{\mathcal{B}}$ and $\overline{\mathcal{B}}$), it follows from the implicit function theorem and Lemma 4.1 that the functions t_a^j (for all $a \in [0, a'_0]$ and $1 \leq j \leq p$) are of class $C^1(\overline{\mathcal{B}})$ and that the functions

$$h_{x'}^j : [0, a_1] \rightarrow [0, \gamma_2], \quad a \mapsto t_a^j(x')$$

(for all $1 \leq j \leq p$ and $x' \in \overline{\mathcal{B}}$) are of class $C^1([0, a'_0])$. From the chain rule applied to (4.6), it is straightforward to check that, for all $a \in [0, a'_0]$, $1 \leq j \leq p$ and $x' \in \overline{\mathcal{B}}$,

$$(h_{x'}^j)'(a) = \frac{-1}{\nu(y^j(x')) \cdot \nabla \psi(y^j(x') - t_a^j(x')\nu(y^j(x')))} \in (0, 2\delta^{-1}] \quad (\text{from Lemma 4.1}),$$

whence

$$0 \leq t_a^j(x') = h_{x'}^j(a) \leq 2\delta^{-1}a \quad (4.7)$$

because $h_{x'}^j(0) = t_0^j(x') = 0$. Similarly,

$$\partial_{x_i} t_a^j(x') = \frac{[\partial_i y^j(x') - t_a^j(x')\partial_i(\nu \circ y^j)(x')] \cdot \nabla \psi(y^j(x') - t_a^j(x')\nu(y^j(x')))}{\nu(y^j(x')) \cdot \nabla \psi(y^j(x') - t_a^j(x')\nu(y^j(x')))} \quad (4.8)$$

for all $a \in [0, a'_0]$, $1 \leq i \leq n-1$, $1 \leq j \leq p$ and $x' \in \overline{\mathcal{B}}$. For all $1 \leq j \leq p$ and $x' \in \overline{\mathcal{B}}$, one has $\psi(y^j(x')) = 0$, whence $\partial_i y^j(x') \cdot \nabla \psi(y^j(x')) = 0$ (for all $1 \leq i \leq n-1$). As a consequence,

$$|\partial_i y^j(x') \cdot \nabla \psi(y^j(x') - t_a^j(x')\nu(y^j(x')))| \leq C_1 \times (t_a^j(x'))^\alpha$$

for all $a \in [0, a'_0]$, $1 \leq i \leq n-1$, $1 \leq j \leq p$, $x' \in \overline{\mathcal{B}}$, and for some constant C_1 defined by

$$C_1 = \max_{1 \leq i' \leq n-1, 1 \leq j' \leq p, \xi \in \overline{\mathcal{B}}} |\partial_{i'} y^{j'}(\xi)| \times \sup_{z \neq z' \in \Omega} \frac{|\nabla \psi(z) - \nabla \psi(z')|}{|z - z'|^\alpha} < +\infty.$$

Observe that $C_1 = C_1(\Omega, N)$ only depends on Ω and N (remember that $\|\psi\|_{C^{1,\alpha}(\overline{\Omega})} \leq N$). Call now

$$C_2 = C_2(\Omega, N) = \max_{1 \leq i' \leq n-1, 1 \leq j' \leq p, \xi \in \overline{\mathcal{B}}} |\partial_{i'}(\nu \circ y^{j'}) (\xi)| \times \sup_{z \in \overline{\Omega}} |\nabla \psi(z)| < +\infty,$$

which also depends on Ω and N only. Together with (4.8) and Lemma 4.1, the above arguments imply that

$$\begin{aligned} |\partial_{x_i} t_a^j(x')| &\leq 2\delta^{-1}[C_1 \times (t_a^j(x'))^\alpha + C_2 \times t_a^j(x')] \\ &\leq 2\delta^{-1}[C_1 \times (2\delta^{-1}a)^\alpha + 2C_2 \times \delta^{-1}a] \quad \text{from (4.7)} \end{aligned}$$

for all $a \in [0, a'_0]$, $1 \leq i \leq n-1$, $1 \leq j \leq p$ and $x' \in \overline{\mathcal{B}}$.

It follows that, for all $\eta > 0$, there exists $a''_0 = a''_0(\Omega, \alpha, N, \delta, \eta) \in (0, a'_0]$ such that, for all $a \in [0, a''_0]$,

$$\sup_{1 \leq i \leq n-1, 1 \leq j \leq p, x' \in \overline{\mathcal{B}}} |\partial_{x_i}(t_a^j \nu \circ y^j)(x')| \leq \eta. \quad (4.9)$$

Finally, there are some open sets $U^1, \dots, U^p \subset \mathcal{B}$ such that

$$\text{area}(\partial\Omega) = \sum_{j=1}^p \int_{U^j} |\partial_1 y^j(x') \times \dots \times \partial_{n-1} y^j(x')| dx',$$

where the sets $\{y^j(x'), x' \in U^j\}$ for $j = 1, \dots, p$ are pairwise disjoint and, for any $\varepsilon > 0$, there are some measurable sets $V^1, \dots, V^p \subset \overline{\mathcal{B}}$ such that

$$\partial\Omega = \{y^j(x'), 1 \leq j \leq p, x' \in V^j\},$$

and $V^j \supset U^j$, $\int_{\mathcal{B}} \mathbf{1}_{V^j \setminus U^j}(x') dx' \leq \varepsilon$ for all $1 \leq j \leq p$. Since all functions y^j and $t_a^j \nu \circ y^j$ (for all $a \in [0, a'_0]$ and $1 \leq j \leq p$) are of class $C^1(\overline{\mathcal{B}})$, since each function ϕ_a is one-to-one and since (4.5) holds, it follows that

$$\text{area}(\Sigma_a) = \sum_{j=1}^p \int_{U^j} |\partial_1(y^j - t_a^j \nu \circ y^j)(x') \times \dots \times \partial_{n-1}(y^j - t_a^j \nu \circ y^j)(x')| dx'$$

for all $a \in [0, a'_0]$.

One concludes from (4.9) that, since $\beta = \beta(\Omega, n)$ in (4.2) is positive, there exists a positive constant $a_0 = a_0(\Omega, n, \alpha, N, \delta) \in (0, a'_0]$ which only depends on Ω , n , α , N and δ , and which is such that

$$|\text{area}(\Sigma_a) - \text{area}(\partial\Omega)| \leq \frac{\beta}{2} n \alpha_n R^{n-1} \text{ for all } a \in [0, a_0].$$

As a consequence, one has

$$i(a) = \text{area}(\Sigma_a) \geq \left(1 + \frac{\beta}{2}\right) n \alpha_n R^{n-1}$$

for all $a \in [0, a_0]$. The area of Σ_a is well-defined for all $a \in [0, a_0]$ since $[0, a_0] \subset [0, a'_0] \subset Y$ because of (4.4).

That completes the proof of Lemma 4.2. \square

Lemma 4.3 *Assume that $\Omega \in \mathcal{C}$ is not a ball. Then, with the notations of Section 3 and Lemma 4.2, one has*

$$\tilde{\psi}(x) \geq \left(1 + \frac{\beta}{2}\right) \rho^{-1}(|x|)$$

for all $x \in \overline{\Omega^*}$ such that $\rho(a_0) \leq |x| \leq R$.

Proof. From Lemma 4.2, one knows that

$$\overline{S_{\rho(a_0), R}} \subset E.$$

Notice that $0 < \rho(a_0) < R$. Fix any $x \in \overline{\Omega^*}$ such that $r = |x| \in [\rho(a_0), R]$ (that is $\rho^{-1}(r) \in [0, a_0] \subset Y$). Formula (3.22) of Section 3 implies that

$$\frac{i(\rho^{-1}(r))}{n\alpha_n r^{n-1}} \leq |\rho'(\rho^{-1}(r))| \times |\nabla \tilde{\psi}(x)|.$$

But

$$\begin{aligned} i(\rho^{-1}(r)) &= \int_{\Sigma_{\rho^{-1}(r)}} d\sigma_{\rho^{-1}(r)}(y) = \text{area}(\Sigma_{\rho^{-1}(r)}) \geq \left(1 + \frac{\beta}{2}\right) n\alpha_n R^{n-1} \\ &\geq \left(1 + \frac{\beta}{2}\right) n\alpha_n r^{n-1} \end{aligned}$$

from Lemma 4.2. Thus,

$$1 + \frac{\beta}{2} \leq |\rho'(\rho^{-1}(r))| \times |\nabla \tilde{\psi}(x)|.$$

The conclusion of Lemma 4.3 follows from the above inequality, as in the proof of Corollary 3.6. \square

The improved version of Corollary 3.6 is the following:

Corollary 4.4 *Assume that $\Omega \in \mathcal{C}$ is not a ball. Then there exists a positive constant $\eta = \eta(\Omega, n, \alpha, N, \delta) > 0$ depending only on Ω, n, α, N and δ , such that*

$$\tilde{\psi}(x) \geq (1 + \eta) \rho^{-1}(|x|)$$

for all $x \in \overline{\Omega^*}$.

Proof. Let e be any unit vector in \mathbb{R}^n . Let $\tilde{\Phi}$ be the function defined in $[0, R]$ by

$$\tilde{\Phi}(r) = \tilde{\psi}(re) \text{ for all } r \in [0, R].$$

This function is continuous in $[0, R]$, differentiable (except at a finite set of points) and decreasing in $[0, R]$. Furthermore, Proposition 3.5 and the fact that ρ^{-1} is decreasing in $[0, R]$ imply that

$$-\tilde{\Phi}'(r) \geq -\frac{d}{dr}(\rho^{-1}(r)) \geq 0$$

for all $r \in \rho(Y) = (0, \rho(a_{m-1})) \cup \dots \cup (\rho(a_2), \rho(a_1)) \cup (\rho(a_1), R]$. As in the proof of Corollary 3.6, the mean value theorem yields

$$\tilde{\psi}(re) - \tilde{\psi}(\rho(a_0)e) \geq \rho^{-1}(r) - a_0$$

for all $r \in [0, \rho(a_0)]$. For each such a r in $[0, \rho(a_0)]$, one has $\rho^{-1}(r) \in [a_0, M] \subset (0, M]$, whence

$$\frac{\tilde{\psi}(re)}{\rho^{-1}(r)} \geq 1 + \frac{\tilde{\psi}(\rho(a_0)e) - a_0}{\rho^{-1}(r)} \geq 1 + \frac{\beta a_0}{2\rho^{-1}(r)}$$

from Lemma 4.3. Remind that M denotes the maximum of ψ , so that $\rho^{-1}(r) \leq M \leq N$. Hence, one obtains

$$\tilde{\psi}(re) \geq \left(1 + \frac{\beta a_0}{2N}\right) \rho^{-1}(r) \text{ for all } r \in [0, \rho(a_0)].$$

As in the proof of Corollary 3.6, the conclusion of Corollary 4.4 follows from the above inequality and from Lemma 4.3, with the choice

$$\eta = \eta(\Omega, n, \alpha, N, \delta) = \min\left(\frac{\beta}{2}, \frac{\beta a_0}{2N}\right) = \frac{\beta a_0}{2N} > 0$$

for instance (notice that $a_0 \leq M \leq N$). □

Lastly, the following corollary is an improved version of Corollary 3.11.

Corollary 4.5 *Assume that $\Omega \in \mathcal{C}$ is not a ball. If there are $\omega_0 \geq 0$ and $\mu \geq 0$ such that*

$$-\operatorname{div}(A_\Omega \nabla \psi)(y) - |\omega(y)| \times |\nabla \psi(y)| - \omega_0 |\nabla \psi(y)| + V(y)\psi(y) \leq \mu \psi(y) \text{ for all } y \in \Omega,$$

then, under the notation of Section 3,

$$\int_{\Omega^*} \left[\widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 - \omega_0 |\nabla \tilde{\psi}(x)| \tilde{\psi}(x) + \widehat{V}(x) \tilde{\psi}(x)^2 \right] e^{-U(x)} dx \leq \frac{\mu}{1 + \eta} \int_{\Omega^*} \tilde{\psi}(x)^2 e^{-U(x)} dx,$$

where the positive constant $\eta = \eta(\Omega, n, \alpha, N, \delta) > 0$ is given in Corollary 4.4 and depends only on Ω, n, α, N and δ .

Proof. Under the assumptions of Corollary 4.5, it follows from Proposition 3.8 and Corollary 4.4 that

$$-\operatorname{div}(\widehat{\Lambda} \nabla \tilde{\psi})(x) + \widehat{v}(x) \cdot \nabla \tilde{\psi}(x) - \omega_0 |\nabla \tilde{\psi}(x)| + \widehat{V}(x) \tilde{\psi}(x) \leq \mu \rho^{-1}(|x|) \leq \frac{\mu}{1 + \eta} \tilde{\psi}(x)$$

for all $x \in E \cap \Omega^*$. The conclusion of Corollary 4.5 then follows from Proposition 3.10. □

5 Application to eigenvalue problems

The present section is devoted to the proofs of some of the main theorems which were stated in Section 2. We apply the rearrangement inequalities of the previous two sections to get some comparison results for the principal eigenvalues of operators which are defined in Ω and in Ω^* . Here, the data have given averages or given distribution functions, or satisfy other types of pointwise constraints.

We shall use a triple approximation process. First, we approximate the diffusion and the drift coefficients in Ω by smooth functions. Second, we approximate the principal eigenfunctions in Ω by analytic functions. Lastly, we approximate the symmetrized data in Ω^* by coefficients having the same distribution functions or satisfying the same constraints as the

original data in Ω . Section 5.1 is concerned with the latter approximation process. In Section 5.2, we deal with the case of general non-symmetric operators for which the inverse of the lower bound Λ of the diffusion matrix field A has a given L^1 norm, the drift v has a given L^2 norm with weight Λ^{-1} and the negative part of the potential V has a given distribution function. Lastly, in Section 5.3, we consider diffusion matrix fields A whose trace and determinant satisfy some pointwise constraints.

5.1 Approximation of symmetrized fields by fields having given distribution functions

In this subsection, Ω denotes an open connected bounded non-empty C^1 subset of \mathbb{R}^n and Ω^* is the open Euclidean ball which is centered at the origin and such that $|\Omega^*| = |\Omega|$. In this subsection we do not require Ω to be of class C^2 . Call $R > 0$ the radius of Ω^* , that is $\Omega^* = B_R$. For $0 \leq s < s' \leq R$, one recalls that $S_{s,s'} = \{z \in \mathbb{R}^n, s < |z| < s'\}$.

Let $\psi : \overline{\Omega} \rightarrow \mathbb{R}$ be a continuous function. Call

$$m = \min_{\overline{\Omega}} \psi \quad \text{and} \quad M = \max_{\overline{\Omega}} \psi.$$

For all $a < b \in \mathbb{R}$, denote

$$\Omega_a = \{x \in \Omega, a < \psi(x)\}, \quad \Omega_{a,b} = \{x \in \Omega, a < \psi(x) \leq b\} \quad \text{and} \quad \Sigma_a = \{x \in \Omega, \psi(x) = a\}.$$

One assumes that

$$|\Sigma_a| = 0 \quad \text{for all } a \in \mathbb{R}. \quad (5.1)$$

It follows then that $m < M$. For each $a \in [m, M]$, set

$$\rho(a) = \left(\frac{|\Omega_a|}{\alpha_n} \right)^{1/n}.$$

The function $\rho : [m, M] \rightarrow \mathbb{R}$ is then continuous, decreasing and it ranges onto $[0, R]$.

Lemma 5.1 *Under assumption (5.1), let g be in $L^\infty(\Omega, \mathbb{R})$ and h in $L^\infty(\Omega^*, \mathbb{R})$ and radially symmetric, and assume that*

$$\forall m \leq a < b \leq M, \quad \int_{\Omega_{a,b}} g = \int_{S_{\rho(b), \rho(a)}} h. \quad (5.2)$$

Then there exist a sequence of radially symmetric functions $(g_k)_{k \in \mathbb{N}}$ in $L^\infty(\Omega^, \mathbb{R})$ and two sequences of radially symmetric $C^\infty(\overline{\Omega^*}, \mathbb{R})$ functions $(\underline{g}_k)_{k \in \mathbb{N}}$ and $(\overline{g}_k)_{k \in \mathbb{N}}$ such that*

$$g_k \rightharpoonup h, \quad \underline{g}_k \rightharpoonup h \quad \text{and} \quad \overline{g}_k \rightharpoonup h \quad \text{as } k \rightarrow +\infty$$

in $\sigma(L^p(\Omega^*), L^{p'}(\Omega^*))$ for all $1 < p \leq +\infty$ (weak convergence, with $1/p + 1/p' = 1$ for $1 < p < +\infty$, and weak- * convergence for $p = +\infty$ and $p' = 1$),² and

$$\left\{ \begin{array}{l} |\{x \in \Omega, g(x) > t\}| = |\{x \in \Omega^*, g_k(x) > t\}| \quad \text{for all } t \in \mathbb{R} \text{ and } k \in \mathbb{N}, \\ |\{x \in \Omega, g(x) \geq t\}| = |\{x \in \Omega^*, g_k(x) \geq t\}| \quad \text{for all } t \in \mathbb{R} \text{ and } k \in \mathbb{N}, \\ \operatorname{ess\,inf}_{\Omega} g \leq \underline{g}_k \leq g_k \leq \bar{g}_k \leq \operatorname{ess\,sup}_{\Omega} g \quad \text{a.e. in } \Omega^* \text{ for all } k \in \mathbb{N}. \end{array} \right.$$

Remark 5.2 The functions g_k and \underline{g}_k are actually constructed so that

$$g_k - \underline{g}_k \rightarrow 0 \quad \text{and} \quad g_k - \bar{g}_k \rightarrow 0 \quad \text{as } k \rightarrow +\infty \text{ in } L^q(\Omega^*)$$

for all $q \in [1, +\infty)$ (see the proof of Lemma 5.1).

Remark 5.3 The functions g_k then have the same distribution function as g and the functions \underline{g}_k (resp. \bar{g}_k) then have distribution functions which are less than (resp. larger than) or equal to that of g . Moreover, if one further assumes that g and h are nonnegative almost everywhere in Ω and Ω^* respectively, then the functions g_k (resp. $\underline{g}_k, \bar{g}_k$) are nonnegative almost everywhere in Ω^* (resp. everywhere in $\bar{\Omega}^*$). The same property holds good with nonpositivity instead of nonnegativity.

In order not to lengthen the reading of the paper, the proof of this lemma is postponed in the Appendix. To finish this subsection, we just point out an immediate corollary of Lemma 5.1.

Corollary 5.4 *In addition to (5.1), assume that ψ is in $C^1(\bar{\Omega})$, $\psi = m$ on $\partial\Omega$, $\psi > m$ in Ω , and ψ has a finite number p of critical values a_i with $m < a_1 < \dots < a_p = M$. Let $g \in L^\infty(\Omega, \mathbb{R})$ and $\hat{g} \in L^\infty(\Omega^*, \mathbb{R})$ be defined by*

$$\hat{g}(x) = \frac{\int_{\Sigma_{\rho^{-1}(|x|)}} g(y) |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(|x|)}(y)}{\int_{\Sigma_{\rho^{-1}(|x|)}} |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(|x|)}(y)} \quad (5.3)$$

for almost every $x \in \Omega^*$ such that $|x| \neq \rho(a_p), \dots, \rho(a_1)$, where $d\sigma_a$ denotes the surface measure on Σ_a for $a \in [m, M]$ which is not a critical value of ψ . Then the conclusion of Lemma 5.1 holds for $h = \hat{g}$.

Proof. It is enough to prove that the function \hat{g} defined by (5.3) is indeed well-defined, bounded and radially symmetric and that property (5.2) is satisfied with $h = \hat{g}$. To do so, choose first any two real numbers $a < b$ in $[m, M]$ such that $a_i \notin [a, b]$ for all $i = 1, \dots, p$. From the co-area formula and Fubini's theorem, one has

$$\int_{\Omega_{a,b}} |g(y)| dy = \int_a^b \left(\int_{\Sigma_s} |g(y)| \times |\nabla\psi(y)|^{-1} d\sigma_s(y) \right) ds < +\infty \quad (5.4)$$

²This convention is used throughout the paper.

and the quantity

$$\int_{\Sigma_s} |g(y)| \times |\nabla\psi(y)|^{-1} d\sigma_s(y)$$

is therefore finite for almost every $s \in [a, b]$. The quantity

$$\frac{\int_{\Sigma_{\rho^{-1}(r)}} g(y) |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y)}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y)}$$

is then finite for almost every r in $[\rho^{-1}(b), \rho^{-1}(a)]$ and for all $m \leq a < b \leq M$ such that $\{a_1, \dots, a_p\} \cap [a, b] = \emptyset$. In other words, \widehat{g} is well-defined for almost every $x \in \Omega^*$ such that $|x| \neq \rho(a_1), \dots, \rho(a_p)$. Moreover, the function \widehat{g} is in $L^\infty(\Omega^*, \mathbb{R})$,

$$\operatorname{ess\,inf}_{\Omega^*} g \leq \operatorname{ess\,inf}_{\Omega^*} \widehat{g} \leq \operatorname{ess\,sup}_{\Omega^*} \widehat{g} \leq \operatorname{ess\,sup}_{\Omega^*} g$$

and \widehat{g} is clearly radially symmetric.

On the other hand, the same calculations as the ones which were done in Lemma 3.2 in Section 3 imply that the function ρ is actually differentiable at each value $a \notin \{a_1, \dots, a_p\}$, with

$$\rho'(a) = -(n\alpha_n \rho(a)^{n-1})^{-1} \int_{\Sigma_a} |\nabla\psi(y)|^{-1} d\sigma_a(y).$$

Coming back to (5.4), the change of variables $s = \rho^{-1}(r)$ then yields

$$\int_{\Omega_{a,b}} g(y) dy = \int_{\rho(b)}^{\rho(a)} \left(\int_{\Sigma_{\rho^{-1}(r)}} g(y) |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y) \right) \frac{n\alpha_n r^{n-1}}{\int_{\Sigma_{\rho^{-1}(r)}} |\nabla\psi(y)|^{-1} d\sigma_{\rho^{-1}(r)}(y)} dr,$$

that is

$$\int_{\Omega_{a,b}} g(y) dy = \int_{S_{\rho(b), \rho(a)}} \widehat{g}(x) dx$$

for all $a < b$ in $[m, M]$ such that $\{a_1, \dots, a_p\} \cap [a, b] = \emptyset$. Since $|\Sigma_c| = 0$ for all $c \in [m, M]$ (and then ρ is continuous), one gets from Lebesgue's dominated convergence theorem that

$$\int_{\Omega_{a,b}} g = \int_{S_{\rho(b), \rho(a)}} \widehat{g}$$

for all $a < b$ in $[m, M]$ and then the conclusion of Corollary 5.4 follows from Lemma 5.1. \square

5.2 Operators whose coefficients have given averages or given distribution functions

In this subsection, we consider operators for which $A \geq \Lambda \operatorname{Id}$, and Λ and v satisfy some integral constraints while the negative part of the potential V will be fixed.

Proof of Theorem 2.1. It will be divided into several steps. Throughout the proof, we fix

$$\Omega \in \mathcal{C}, A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R})), \Lambda \in L^{\infty}_+(\Omega), v \in L^{\infty}(\Omega, \mathbb{R}^n) \text{ and } V \in C(\overline{\Omega}).$$

Call

$$0 < m_{\Lambda} = \operatorname{ess\,inf}_{\Omega} \Lambda \leq \operatorname{ess\,sup}_{\Omega} \Lambda = M_{\Lambda} < +\infty \quad (5.5)$$

and assume that

$$A \geq \Lambda \operatorname{Id} \text{ a.e. in } \Omega. \quad (5.6)$$

and that

$$\lambda_1(\Omega, A, v, V) \geq 0.$$

Step 1: Approximation of Λ in Ω . Write first

$$A(x) = (a_{i,j}(x))_{1 \leq i,j \leq n}.$$

Each function $a_{i,j}$ is in $W^{1,\infty}(\Omega)$ and is therefore continuous in $\overline{\Omega}$ (up to a modification on a zero-measure set). For each $x \in \overline{\Omega}$, call $\Lambda[A](x)$ the lowest eigenvalue of $A(x)$, that is:

$$\forall x \in \overline{\Omega}, \quad \Lambda[A](x) = \min_{\xi \in \mathbb{R}^n, |\xi|=1} A(x)\xi \cdot \xi.$$

The function $\Lambda[A]$ is then continuous in $\overline{\Omega}$ and

$$\Lambda[A](x) \geq \Lambda(x) \text{ a.e. in } \Omega$$

because of (5.6). In particular, $\Lambda[A](x) \geq m_{\Lambda}$ for all $x \in \overline{\Omega}$, where $m_{\Lambda} > 0$ has been defined in (5.5). There exists then a continuous function $\overline{\Lambda}$ in \mathbb{R}^n such that

$$\overline{\Lambda}(x) = \Lambda[A](x) \text{ for all } x \in \overline{\Omega}, \text{ and } m_{\Lambda} \leq \overline{\Lambda}(x) \leq \|\Lambda[A]\|_{L^{\infty}(\Omega)} \text{ for all } x \in \mathbb{R}^n.$$

We first consider the case when $m_{\Lambda} < M_{\Lambda}$. Thus,

$$\int_{\Omega} \Lambda(y)^{-1}(y) dy > M_{\Lambda}^{-1} |\Omega|,$$

whence

$$\overline{\varepsilon} := M_{\Lambda} - \|\Lambda^{-1}\|_{L^1(\Omega)}^{-1} |\Omega| \in (0, M_{\Lambda} - m_{\Lambda}].$$

Pick any ε in $(0, \overline{\varepsilon})$. Let J_{ε} be the function defined in $[0, M_{\Lambda} - m_{\Lambda} - \varepsilon]$ by

$$\forall \tau \in [0, M_{\Lambda} - m_{\Lambda} - \varepsilon], \quad J_{\varepsilon}(\tau) = \int_{\Omega} \max(\min(\Lambda(y), M_{\Lambda} - \varepsilon), m_{\Lambda} + \tau)^{-1} dy.$$

This means that the function Λ is truncated between $m_{\Lambda} + \tau$ and $M_{\Lambda} - \varepsilon$. The function J_{ε} is Lipschitz-continuous and nonincreasing in $[0, M_{\Lambda} - m_{\Lambda} - \varepsilon]$. Furthermore,

$$J_{\varepsilon}(0) = \int_{\Omega} \min(\Lambda(y), M_{\Lambda} - \varepsilon)^{-1} dy > \int_{\Omega} \Lambda(y)^{-1} dy$$

since $M_\Lambda - \varepsilon < M_\Lambda = \text{ess sup}_\Omega \Lambda$, while

$$J_\varepsilon(M_\Lambda - m_\Lambda - \varepsilon) = (M_\Lambda - \varepsilon)^{-1}|\Omega| < (M_\Lambda - \bar{\varepsilon})^{-1}|\Omega| = \int_\Omega \Lambda(y)^{-1} dy$$

owing to the definition of $\bar{\varepsilon}$. Therefore,

$$\tau(\varepsilon) := \min \left\{ \tau \in [0, M_\Lambda - m_\Lambda - \varepsilon], J_\varepsilon(\tau) = \int_\Omega \Lambda(y)^{-1} dy \right\}$$

is well-defined and $0 < \tau(\varepsilon) < M_\Lambda - m_\Lambda - \varepsilon$.

Moreover, we claim that

$$\tau(\varepsilon) \rightarrow 0^+ \text{ as } \varepsilon \rightarrow 0^+. \quad (5.7)$$

If not, there exist $\tau_\infty \in (0, M_\Lambda - m_\Lambda]$ and a sequence $(\varepsilon_p)_{p \in \mathbb{N}} \in (0, \bar{\varepsilon})$ such that $\varepsilon_p \rightarrow 0$ and $\tau(\varepsilon_p) \rightarrow \tau_\infty$ as $p \rightarrow +\infty$. Then,

$$\begin{aligned} \int_\Omega \Lambda(y)^{-1} dy = J_{\varepsilon_p}(\tau(\varepsilon_p)) &= \int_\Omega \max(\min(\Lambda(y), M_\Lambda - \varepsilon_p), m_\Lambda + \tau(\varepsilon_p))^{-1} dy \\ &\xrightarrow{p \rightarrow +\infty} \int_\Omega \max(\min(\Lambda(y), M_\Lambda), m_\Lambda + \tau_\infty)^{-1} dy, \end{aligned}$$

whence

$$\int_\Omega \Lambda(y)^{-1} dy = \int_\Omega \max(\Lambda(y), m_\Lambda + \tau_\infty)^{-1} dy$$

and $\Lambda \geq m_\Lambda + \tau_\infty > m_\Lambda$ a.e. in Ω , which is impossible.

Choose a sequence $(\varepsilon_k)_{k \in \mathbb{N}}$ of real numbers such that

$$0 < \varepsilon_k < \bar{\varepsilon} \text{ for all } k \in \mathbb{N}, \text{ and } \varepsilon_k \rightarrow 0 \text{ as } k \rightarrow +\infty.$$

For each $k \in \mathbb{N}$, call

$$\tau_k = \tau(\varepsilon_k) \in (0, M_\Lambda - m_\Lambda - \varepsilon_k).$$

It then follows from (5.7) that

$$\tau_k \rightarrow 0 \text{ as } k \rightarrow +\infty.$$

Then, for each $k \in \mathbb{N}$, denote

$$\Lambda_{\Omega,k} = \max(\min(\Lambda, M_\Lambda - \varepsilon_k), m_\Lambda + \tau_k) \text{ a.e. in } \Omega \quad (5.8)$$

and define the function $\underline{\Lambda}_k$ in \mathbb{R}^n (almost everywhere) by

$$\underline{\Lambda}_k(y) = \begin{cases} \Lambda_{\Omega,k}(y) & \text{if } y \in \Omega, \\ m_\Lambda + \tau_k & \text{if } y \in \mathbb{R}^n \setminus \Omega, \end{cases}$$

and the continuous function $\bar{\Lambda}_k$ by

$$\bar{\Lambda}_k(y) = \max(\bar{\Lambda}(y), m_\Lambda + \tau_k) \text{ for all } y \in \mathbb{R}^n.$$

Notice that

$$\int_{\Omega} \underline{\Delta}_k(y)^{-1} dy = \int_{\Omega} \Lambda_{\Omega,k}(y)^{-1} dy = J_{\varepsilon_k}(\tau_k) = \int_{\Omega} \Lambda(y)^{-1} dy \quad (5.9)$$

for each $k \in \mathbb{N}$. Observe also that

$$0 < m_{\Lambda} < m_{\Lambda} + \tau_k \leq \underline{\Delta}_k \leq M_{\Lambda} - \varepsilon_k < M_{\Lambda}, \quad \underline{\Delta}_k \leq \bar{\Lambda}_k \text{ a.e. in } \mathbb{R}^n,$$

and that

$$\bar{\Lambda}_k \leq \max(\|\Lambda[A]\|_{L^{\infty}(\Omega)}, m_{\Lambda} + \tau_k) = \|\Lambda[A]\|_{L^{\infty}(\Omega)} \text{ in } \mathbb{R}^n,$$

because $\|\Lambda[A]\|_{L^{\infty}(\Omega)} \geq \text{ess sup}_{\Omega} \Lambda = M_{\Lambda} \geq m_{\Lambda} + \tau_k$.

Let $(\rho_{k'})_{k' \in \mathbb{N}}$ be a sequence of mollifiers in \mathbb{R}^n . For each $(k, k') \in \mathbb{N}^2$, call

$$\bar{\Lambda}_{k,k'} = \left(\rho_{k'} * \bar{\Lambda}_k^{-1} \right)^{-1} \text{ and } \underline{\Delta}_{k,k'} = \left(\rho_{k'} * \underline{\Delta}_k^{-1} \right)^{-1}$$

that is

$$\bar{\Lambda}_{k,k'}(y)^{-1} = \int_{\mathbb{R}^n} \rho_{k'}(z) \bar{\Lambda}_k(y-z)^{-1} dz \text{ and } \underline{\Delta}_{k,k'}(y)^{-1} = \int_{\mathbb{R}^n} \rho_{k'}(z) \underline{\Delta}_k(y-y)^{-1} dz$$

for all $y \in \mathbb{R}^n$. The functions $\bar{\Lambda}_{k,k'}$ and $\underline{\Delta}_{k,k'}$ are of class $C^{\infty}(\mathbb{R}^n)$ and they satisfy

$$0 < m_{\Lambda} < m_{\Lambda} + \tau_k \leq \underline{\Delta}_{k,k'}(y) \leq M_{\Lambda} - \varepsilon_k < M_{\Lambda}$$

and

$$\underline{\Delta}_{k,k'}(y) \leq \bar{\Lambda}_{k,k'}(y) \leq \|\Lambda[A]\|_{L^{\infty}(\Omega)} \text{ for all } y \in \mathbb{R}^n.$$

Furthermore, for each $k \in \mathbb{N}$,

$$\underline{\Delta}_{k,k'}^{-1} \rightarrow \underline{\Delta}_k^{-1} \text{ in } L^p_{loc}(\mathbb{R}^n) \text{ for all } 1 \leq p < +\infty \text{ as } k' \rightarrow +\infty \quad (5.10)$$

and

$$\bar{\Lambda}_{k,k'} \rightarrow \bar{\Lambda}_k = \max(\Lambda[A], m_{\Lambda} + \tau_k) \text{ uniformly in } \bar{\Omega} \text{ as } k' \rightarrow +\infty.$$

Actually, since $\bar{\Lambda}_k^{-1} \rightarrow \bar{\Lambda}^{-1}$ uniformly in \mathbb{R}^n as $k \rightarrow +\infty$ (because $\tau_k \rightarrow 0$ as $k \rightarrow +\infty$ and $\bar{\Lambda} \geq m_{\Lambda} > 0$ in \mathbb{R}^n), one even has that

$$\bar{\Lambda}_{k,k'} - \bar{\Lambda}_k \rightarrow 0 \text{ uniformly in } \bar{\Omega} \text{ as } (k, k') \rightarrow (+\infty, +\infty). \quad (5.11)$$

Call, for each $(k, k') \in \mathbb{N}^2$,

$$\alpha_{k,k'} = \frac{\|\underline{\Delta}_{k,k'}^{-1}\|_{L^1(\Omega)}}{\|\Lambda^{-1}\|_{L^1(\Omega)}} > 0.$$

Define the function $\underline{\Delta}$ almost everywhere in \mathbb{R}^n by:

$$\underline{\Delta} = \begin{cases} \Lambda & \text{in } \Omega, \\ m_{\Lambda} & \text{in } \mathbb{R}^n \setminus \Omega. \end{cases}$$

Since $0 < m_{\Lambda} \leq \underline{\Delta} \leq M_{\Lambda}$ a.e. in \mathbb{R}^n and $\varepsilon_k, \tau_k \rightarrow 0$ as $k \rightarrow +\infty$, it follows that

$$\|\underline{\Delta}_k - \underline{\Delta}\|_{L^{\infty}(\mathbb{R}^n)} \rightarrow 0 \text{ and } \|\underline{\Delta}_k^{-1} - \underline{\Delta}^{-1}\|_{L^{\infty}(\mathbb{R}^n)} \rightarrow 0 \text{ as } k \rightarrow +\infty.$$

Thus,

$$\|\underline{\Lambda}_{k,k'}^{-1} - \Lambda^{-1}\|_{L^1(\Omega)} \leq \|\rho_{k'} * (\underline{\Lambda}_k^{-1} - \underline{\Lambda}^{-1})\|_{L^1(\Omega)} + \|\rho_{k'} * \underline{\Lambda}^{-1} - \underline{\Lambda}^{-1}\|_{L^1(\Omega)} \rightarrow 0 \text{ as } (k, k') \rightarrow (+\infty, +\infty),$$

whence

$$\alpha_{k,k'} - 1 = \frac{\|\underline{\Lambda}_{k,k'}^{-1}\|_{L^1(\Omega)} - \|\Lambda^{-1}\|_{L^1(\Omega)}}{\|\Lambda^{-1}\|_{L^1(\Omega)}} \rightarrow 0 \text{ as } (k, k') \rightarrow (+\infty, +\infty). \quad (5.12)$$

Furthermore, because of (5.9) and (5.10), there holds

$$\alpha_{k,k'} - 1 = \frac{\|\underline{\Lambda}_{k,k'}^{-1}\|_{L^1(\Omega)} - \|\underline{\Lambda}_k^{-1}\|_{L^1(\Omega)}}{\|\Lambda^{-1}\|_{L^1(\Omega)}} \rightarrow 0 \text{ as } k' \rightarrow +\infty, \text{ for each } k \in \mathbb{N}. \quad (5.13)$$

Define now

$$\Lambda_{k,k'}(y) = \alpha_{k,k'} \underline{\Lambda}_{k,k'}(y) \text{ for all } y \in \overline{\Omega} \text{ and } (k, k') \in \mathbb{N}^2.$$

The functions $\Lambda_{k,k'}$ are of class $C^\infty(\overline{\Omega})$ and they satisfy

$$\int_{\Omega} \Lambda_{k,k'}(y)^{-1} dy = \int_{\Omega} \Lambda(y)^{-1} dy \text{ for all } (k, k') \in \mathbb{N}^2, \quad (5.14)$$

and

$$\begin{cases} 0 < \alpha_{k,k'} \times (m_\Lambda + \tau_k) \leq \Lambda_{k,k'} \leq \alpha_{k,k'} \times (M_\Lambda - \varepsilon_k) \text{ in } \overline{\Omega} \text{ for all } (k, k') \in \mathbb{N}^2, \\ \alpha_{k,k'} \times (m_\Lambda + \tau_k) \xrightarrow[k' \rightarrow +\infty]{} m_\Lambda + \tau_k \text{ for all } k \in \mathbb{N}, \\ \alpha_{k,k'} \times (M_\Lambda - \varepsilon_k) \xrightarrow[k' \rightarrow +\infty]{} M_\Lambda - \varepsilon_k \text{ for all } k \in \mathbb{N}, \end{cases} \quad (5.15)$$

together with

$$\|\Lambda_{k,k'}^{-1} - \Lambda^{-1}\|_{L^1(\Omega)} \rightarrow 0 \text{ as } (k, k') \rightarrow (+\infty, +\infty). \quad (5.16)$$

Lastly,

$$\Lambda_{k,k'} = \alpha_{k,k'} \underline{\Lambda}_{k,k'} \leq \alpha_{k,k'} \overline{\Lambda}_{k,k'} \text{ in } \overline{\Omega} \text{ for all } (k, k') \in \mathbb{N}^2 \quad (5.17)$$

and

$$\alpha_{k,k'} \overline{\Lambda}_{k,k'} - \overline{\Lambda}_k \rightarrow 0 \text{ uniformly in } \overline{\Omega} \text{ as } (k, k') \rightarrow (+\infty, +\infty) \quad (5.18)$$

from (5.11) and (5.12).

In the case when $m_\Lambda = M_\Lambda$, namely when Λ is equal to a constant (up to modification on a zero-measure set), then one sets $\Lambda_{k,k'} = \underline{\Lambda}_{k,k'} = \underline{\Lambda}_k = \Lambda_{\Omega,k} = \Lambda$, $\alpha_{k,k'} = 1$, $\overline{\Lambda}_k = \overline{\Lambda}$, $\varepsilon_k = \tau_k = 0$ and properties (5.14), (5.15), (5.16), (5.17) and (5.18) hold immediately.

Step 2: Approximation of A in Ω . Let us now approximate the $W^{1,\infty}(\Omega)$ matrix field $A = (a_{i,j})_{1 \leq i,j \leq n}$. First, each function $a_{i,j}$ can be extended to a $W^{1,\infty}(\mathbb{R}^n)$ function $\overline{a}_{i,j}$ such that

$$\|\overline{a}_{i,j}\|_{L^\infty(\mathbb{R}^n)} = \|a_{i,j}\|_{L^\infty(\Omega)} \text{ and } \|\nabla \overline{a}_{i,j}\|_{L^\infty(\mathbb{R}^n)} = \|\nabla a_{i,j}\|_{L^\infty(\Omega)},$$

whence

$$\|\overline{a}_{i,j}\|_{W^{1,\infty}(\mathbb{R}^n)} = \|a_{i,j}\|_{W^{1,\infty}(\Omega)} \leq \|A\|_{W^{1,\infty}(\Omega)},$$

where we recall that $\|A\|_{W^{1,\infty}(\Omega)} = \max_{1 \leq i,j \leq n} \|a_{i,j}\|_{W^{1,\infty}(\Omega)}$. Since the matrix field $A = (a_{i,j})_{1 \leq i,j \leq n}$ is symmetric, the matrix field $(\bar{a}_{i,j})_{1 \leq i,j \leq n}$ can be assumed to be symmetric. For each $1 \leq i, j \leq n$, the functions $\rho_{k'} * \bar{a}_{i,j}$ are of class $C^\infty(\mathbb{R}^n)$ and converge uniformly to $a_{i,j}$ in $\bar{\Omega}$ as $k' \rightarrow +\infty$. Furthermore,

$$\|\nabla(\rho_{k'} * \bar{a}_{i,j})\|_{L^\infty(\mathbb{R}^n)} \leq \|\nabla \bar{a}_{i,j}\|_{L^\infty(\mathbb{R}^n)} = \|\nabla a_{i,j}\|_{L^\infty(\Omega)} \leq \|a_{i,j}\|_{W^{1,\infty}(\Omega)} \leq \|A\|_{W^{1,\infty}(\Omega)}$$

for all $k' \in \mathbb{N}$ and $1 \leq i, j \leq n$. For each $k' \in \mathbb{N}$, the matrix field $(\rho_{k'} * \bar{a}_{i,j})_{1 \leq i,j \leq n}$ can be approximated in $C^1(\bar{\Omega})$ norm by symmetric matrix fields with polynomial entries in $\bar{\Omega}$. Therefore, there exists a sequence of symmetric matrix fields $(A'_{k'})_{k' \in \mathbb{N}} = ((a'_{k',i,j})_{1 \leq i,j \leq n})_{k' \in \mathbb{N}}$ in $\bar{\Omega}$ with polynomial entries $a'_{k',i,j}$ such that, for all $1 \leq i, j \leq n$,

$$a'_{k',i,j} \rightarrow a_{i,j} \text{ uniformly in } \bar{\Omega} \text{ as } k' \rightarrow +\infty \quad (5.19)$$

and

$$\limsup_{k' \rightarrow +\infty} \|\nabla a'_{k',i,j}\|_{L^\infty(\Omega)} \leq \|\nabla a_{i,j}\|_{L^\infty(\Omega)} \leq \|A\|_{W^{1,\infty}(\Omega)}.$$

Call

$$\eta_{k,k'} = n^2 \times \max_{1 \leq i,j \leq n} \|a'_{k',i,j} - a_{i,j}\|_{L^\infty(\Omega)} + \|\bar{\Lambda}_k - \alpha_{k,k'} \bar{\Lambda}_{k,k'}\|_{L^\infty(\Omega)} + \tau_k.$$

Because of (5.18), (5.19), and since $\tau_k \rightarrow 0$ as $k \rightarrow +\infty$, there holds

$$\eta_{k,k'} \rightarrow 0 \text{ as } (k, k') \rightarrow (+\infty, +\infty).$$

The symmetric matrix fields

$$A_{k,k'} = A'_{k'} + \eta_{k,k'} \text{Id} = (a_{k,k',i,j})_{1 \leq i,j \leq n}$$

with polynomial entries $a_{k,k',i,j}$ are such that

$$a_{k,k',i,j} \rightarrow a_{i,j} \text{ uniformly in } \bar{\Omega} \text{ as } (k, k') \rightarrow (+\infty, +\infty) \quad (5.20)$$

and

$$\sup_{k \in \mathbb{N}} \limsup_{k' \rightarrow +\infty} \|\nabla a_{k,k',i,j}\|_{L^\infty(\Omega)} \leq \|\nabla a_{i,j}\|_{L^\infty(\Omega)} \leq \|A\|_{W^{1,\infty}(\Omega)}. \quad (5.21)$$

Furthermore, for each $(k, k') \in \mathbb{N}^2$, $x \in \bar{\Omega}$ and $\xi = (\xi_1, \dots, \xi_n) \in \mathbb{R}^n$ with $|\xi| = 1$, there holds

$$\begin{aligned} & A_{k,k'}(x) \xi \cdot \xi - \Lambda_{k,k'}(x) \\ &= \sum_{1 \leq i,j \leq n} (a'_{k',i,j}(x) - a_{i,j}(x)) \xi_i \xi_j + n^2 \times \max_{1 \leq i,j \leq n} \|a'_{k',i,j} - a_{i,j}\|_{L^\infty(\Omega)} \\ & \quad + A(x) \xi \cdot \xi - \Lambda_{k,k'}(x) + \|\bar{\Lambda}_k - \alpha_{k,k'} \bar{\Lambda}_{k,k'}\|_{L^\infty(\Omega)} + \tau_k \\ & \geq \Lambda[A](x) + \tau_k - \alpha_{k,k'} \bar{\Lambda}_{k,k'}(x) + \|\bar{\Lambda}_k - \alpha_{k,k'} \bar{\Lambda}_{k,k'}\|_{L^\infty(\Omega)} \text{ (because of (5.17))} \\ & \geq 0 \end{aligned}$$

because $\Lambda[A] + \tau_k \geq \max(\Lambda[A], m_\Lambda + \tau_k) = \bar{\Lambda}_k$ in $\bar{\Omega}$. In other words,

$$A_{k,k'}(x) \geq \Lambda_{k,k'}(x) \text{Id for all } (k, k') \in \mathbb{N}^2 \text{ and } x \in \bar{\Omega}, \quad (5.22)$$

in the sense of symmetric matrices.

Step 3: Approximation of v in Ω . Call

$$0 \leq m_v = \operatorname{ess\,inf}_{\Omega} |v| \leq \operatorname{ess\,sup}_{\Omega} |v| = \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} = \|v\|_\infty = M_v < +\infty.$$

We first consider the case when $m_v < M_v$ and $m_\Lambda < M_\Lambda$. In particular, it follows that $M_v = \|v\|_\infty > 0$ and that there exists

$$\bar{\varepsilon}' \in (0, M_v - m_v) \text{ such that } (M_v - \bar{\varepsilon}')^2 \times \int_{\Omega} \Lambda(y)^{-1} dy > \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy. \quad (5.23)$$

Call K the function defined for all $\varepsilon' \in [0, \bar{\varepsilon}']$ by

$$K(\varepsilon') = \int_{\Omega} [|v(y)|^2 - \min(|v(y)|^2, (M_v - \varepsilon')^2)] \times \Lambda(y)^{-1} dy.$$

The function K is continuous and nondecreasing in $[0, \bar{\varepsilon}']$, vanishes at 0 and is positive in $(0, \bar{\varepsilon}']$ due to the definition of M_v . Let the sequences of positive numbers $(\varepsilon_k)_{k \in \mathbb{N}}$ and $(\tau_k)_{k \in \mathbb{N}}$ be as in Step 1. Since $\max(\varepsilon_k, \tau_k) \rightarrow 0$ as $k \rightarrow +\infty$, one can assume without loss of generality that

$$\frac{|\Omega| M_v^2}{m_\Lambda^2} \times \max(\varepsilon_k, \tau_k) < K(\bar{\varepsilon}') \text{ for all } k \in \mathbb{N}.$$

For each $k \in \mathbb{N}$, call

$$\varepsilon'_k = \min \left\{ \varepsilon' \in [0, \bar{\varepsilon}'], \frac{|\Omega| M_v^2}{m_\Lambda^2} \times \max(\varepsilon_k, \tau_k) = K(\varepsilon') \right\}.$$

From the above remarks, ε'_k is well-defined and $0 < \varepsilon'_k < \bar{\varepsilon}'$. Furthermore, $K(\varepsilon'_k) = \max(\varepsilon_k, \tau_k) \times |\Omega| M_v^2 m_\Lambda^{-2} \rightarrow 0$ as $k \rightarrow +\infty$, whence

$$\varepsilon'_k \rightarrow 0^+ \text{ as } k \rightarrow +\infty.$$

Fix now any unit vector $e \in \mathbb{R}^n$. For each $k \in \mathbb{N}$, let L_k be the function defined for all $\tau' \in [0, M_v - m_v - \varepsilon'_k]$ by

$$L_k(\tau') = \int_{\Omega} |v_{\Omega, k, \tau'}(y)|^2 \Lambda_{\Omega, k}(y)^{-1} dy, \quad (5.24)$$

where

$$v_{\Omega, k, \tau'}(y) = \begin{cases} (M_v - \varepsilon'_k) |v(y)|^{-1} v(y) & \text{if } |v(y)| > M_v - \varepsilon'_k, \\ v(y) & \text{if } m_v + \tau' \leq |v(y)| \leq M_v - \varepsilon'_k, \\ (m_v + \tau') |v(y)|^{-1} v(y) & \text{if } m_v < |v(y)| < m_v + \tau', \\ (m_v + \tau') |v(y)|^{-1} v(y) & \text{if } |v(y)| = m_v \text{ and } m_v > 0, \\ (m_v + \tau') e & \text{if } |v(y)| = 0 \text{ and } m_v = 0. \end{cases} \quad (5.25)$$

Each function L_k is Lipschitz-continuous in $[0, M_v - m_v - \varepsilon'_k]$ and

$$L_k(0) = \int_{\Omega} \min(|v(y)|^2, (M_v - \varepsilon'_k)^2) \Lambda_{\Omega, k}(y)^{-1} dy,$$

whence

$$\begin{aligned}
L_k(0) - \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy &= \int_{\Omega} \min(|v(y)|^2, (M_v - \varepsilon'_k)^2) \cdot (\Lambda_{\Omega,k}(y)^{-1} - \Lambda(y)^{-1}) dy - K(\varepsilon'_k) \\
&\leq \frac{|\Omega| M_v^2}{m_{\Lambda}^2} \times \max(\varepsilon_k, \tau_k) - K(\varepsilon'_k) \\
&= 0
\end{aligned}$$

due to the definitions of $\Lambda_{\Omega,k}$ (see (5.8)) and m_{Λ} , M_{Λ} , M_v and ε'_k . Furthermore,

$$L_k(M_v - m_v - \varepsilon'_k) = (M_v - \varepsilon'_k)^2 \int_{\Omega} \Lambda_{\Omega,k}(y)^{-1} dy = (M_v - \varepsilon'_k)^2 \int_{\Omega} \Lambda(y)^{-1} dy > \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy$$

from (5.9) and (5.23) (remember that $0 < \varepsilon'_k < \bar{\varepsilon}' < M_v - m_v \leq M_v$). Therefore, by continuity of L_k , the real number

$$\tau'_k = \min \left\{ \tau' \in [0, M_v - m_v - \varepsilon'_k), L_k(\tau') = \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy \right\} \quad (5.26)$$

is well-defined and $0 \leq \tau'_k < M_v - m_v - \varepsilon'_k$. Moreover, $\tau'_k \rightarrow 0$ as $k \rightarrow +\infty$. Otherwise, up to extraction, there exists $\tau > 0$ such that $\tau'_k \rightarrow \tau$ as $k \rightarrow +\infty$, and

$$L_k(\tau'_k) \xrightarrow{k \rightarrow +\infty} \int_{\Omega} \max(|v(y)|^2, (m_v + \tau)^2) \Lambda(y)^{-1} dy > \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy$$

by definition of m_v . This is impossible by definition of τ'_k .

Call now

$$v_{\Omega,k} = v_{\Omega,k,\tau'_k}.$$

Notice that, for each $k \in \mathbb{N}$, the vector field $v_{\Omega,k}$ is in $L^{\infty}(\Omega, \mathbb{R}^n)$ and it satisfies

$$m_v + \tau'_k \leq \operatorname{ess\,inf}_{\Omega} |v_{\Omega,k}| \leq \operatorname{ess\,sup}_{\Omega} |v_{\Omega,k}| \leq M_v - \varepsilon'_k$$

and

$$\int_{\Omega} |v_{\Omega,k}(y)|^2 \Lambda_{\Omega,k}(y)^{-1} dy = L_k(\tau'_k) = \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy. \quad (5.27)$$

Write now the vector fields v and $v_{\Omega,k}$ as $v = (v_1, \dots, v_n)$ and $v_{\Omega,k} = (v_{k,1}, \dots, v_{k,n})$, extend all functions v_i and $v_{k,i}$ by 0 in $\mathbb{R}^n \setminus \Omega$, call \underline{v}_i and $\underline{v}_{k,i}$ these extensions. Set

$$\underline{v} = (\underline{v}_1, \dots, \underline{v}_n) \text{ and } \underline{v}_k = (\underline{v}_{k,1}, \dots, \underline{v}_{k,n}).$$

One then has that

$$\|\underline{v}_k - \underline{v}\|_{L^{\infty}(\mathbb{R}^n, \mathbb{R}^n)} \rightarrow 0 \text{ as } k \rightarrow +\infty \quad (5.28)$$

since $(\varepsilon'_k, \tau'_k) \rightarrow (0, 0)$ as $k \rightarrow +\infty$. For each $(k, k') \in \mathbb{N}^2$, denote

$$\underline{v}_{k,k'} = (\rho_{k'} * \underline{v}_{k,1}, \dots, \rho_{k'} * \underline{v}_{k,n}),$$

where $(\rho_{k'})_{k' \in \mathbb{N}}$ is the same sequence of mollifiers as in Step 1. The vector fields $\underline{v}_{k,k'}$ are then of class $C^\infty(\mathbb{R}^n, \mathbb{R}^n)$ and they satisfy

$$\|\underline{v}_{k,k'}\|_{L^\infty(\mathbb{R}^n, \mathbb{R}^n)} \leq \|\underline{v}_k\|_{L^\infty(\mathbb{R}^n, \mathbb{R}^n)} \leq M_v - \varepsilon'_k < M_v = \|v\|_\infty \text{ for all } (k, k') \in \mathbb{N}^2.$$

For each fixed $k \in \mathbb{N}$, the fields $\underline{v}_{k,k'}$ converge to $\underline{v}_k|_\Omega = v_{\Omega,k}$ as $k' \rightarrow +\infty$ in all spaces $L^p(\Omega)$ for $1 \leq p < +\infty$. Actually, one also has that

$$\|\underline{v}_{k,k'} - v_{\Omega,k}\|_{L^p(\Omega, \mathbb{R}^n)} + \|\underline{v}_{k,k'} - v\|_{L^p(\Omega, \mathbb{R}^n)} \rightarrow 0 \text{ as } (k, k') \rightarrow (+\infty, +\infty)$$

because of (5.28). Together with (5.12), (5.15) and (5.16), it follows that

$$\int_\Omega |\underline{v}_{k,k'}(y)|^2 \Lambda_{k,k'}(y)^{-1} dy \rightarrow \int_\Omega |v(y)|^2 \Lambda(y)^{-1} dy \text{ as } (k, k') \rightarrow (+\infty, +\infty).$$

Remember that the limit in the right-hand side is positive because $M_v > 0$ here. Then the positive real numbers

$$\beta_{k,k'} = \left(\frac{\int_\Omega |v(y)|^2 \Lambda(y)^{-1} dy}{\int_\Omega |\underline{v}_{k,k'}(y)|^2 \Lambda_{k,k'}(y)^{-1} dy} \right)^{1/2}$$

are well-defined for k and k' large enough (one can then assume for all $(k, k') \in \mathbb{N}^2$ without loss of generality) and $\beta_{k,k'} \rightarrow 1$ as $(k, k') \rightarrow (+\infty, +\infty)$. Moreover, for each $k \in \mathbb{N}$,

$$\beta_{k,k'}^{-2} = \frac{\int_\Omega |\underline{v}_{k,k'}(y)|^2 \Lambda_{k,k'}(y)^{-1} dy}{\int_\Omega |v(y)|^2 \Lambda(y)^{-1} dy} \xrightarrow{k' \rightarrow +\infty} \frac{\int_\Omega |\underline{v}_k(y)|^2 \Lambda_k(y)^{-1} dy}{\int_\Omega |v(y)|^2 \Lambda(y)^{-1} dy} = \frac{\int_\Omega |v_{\Omega,k}(y)|^2 \Lambda_{\Omega,k}(y)^{-1} dy}{\int_\Omega |v(y)|^2 \Lambda(y)^{-1} dy} = 1$$

because of (5.10), (5.13) and (5.27).

Set

$$v_{k,k'}(y) = \beta_{k,k'} \underline{v}_{k,k'}(y) \text{ for all } y \in \overline{\Omega} \text{ and } (k, k') \in \mathbb{N}^2.$$

The vector fields $v_{k,k'}$ are in $C^\infty(\overline{\Omega}, \mathbb{R}^n)$ and they satisfy

$$\int_\Omega |v_{k,k'}(y)|^2 \Lambda_{k,k'}(y)^{-1} dy = \int_\Omega |v(y)|^2 \Lambda(y)^{-1} dy \text{ for all } (k, k') \in \mathbb{N}^2 \quad (5.29)$$

and

$$\begin{cases} \|v_{k,k'}\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \beta_{k,k'} \|\underline{v}_{k,k'}\|_{L^\infty(\mathbb{R}^n, \mathbb{R}^n)} \leq \beta_{k,k'} \times (\|v\|_\infty - \varepsilon'_k) \text{ for all } (k, k') \in \mathbb{N}^2, \\ \beta_{k,k'} \times (\|v\|_\infty - \varepsilon'_k) \rightarrow \|v\|_\infty - \varepsilon'_k \text{ as } k' \rightarrow +\infty \text{ for all } k \in \mathbb{N}, \\ \|v_{k,k'} - v\|_{L^p(\Omega, \mathbb{R}^n)} \rightarrow 0 \text{ as } (k, k') \rightarrow (+\infty, +\infty) \text{ for all } 1 \leq p < +\infty. \end{cases} \quad (5.30)$$

Consider now the case when $m_v < M_v$ and $m_\Lambda = M_\Lambda$. Namely, up to modification on a zero-measure set, Λ is constant. Choose $\bar{\varepsilon}' \in (0, M_v - m_v)$ such that (5.23) holds, namely

$$(M_v - \bar{\varepsilon}')^2 |\Omega| > \int_\Omega |v(y)|^2 dy. \quad (5.31)$$

Take any sequence $(\varepsilon'_k)_{k \in \mathbb{N}}$ in $(0, \bar{\varepsilon}')$ such that $\varepsilon'_k \rightarrow 0$ as $k \rightarrow +\infty$. For each $k \in \mathbb{N}$ and $\tau' \in [0, M_v - m_v - \varepsilon'_k]$, define $L_k(\tau')$ as in (5.24) with $\Lambda_{\Omega, k} = \Lambda$. Each function L_k is Lipschitz-continuous. Moreover,

$$L_k(0) = \Lambda^{-1} \int_{\Omega} \min(|v(y)|^2, (M_v - \varepsilon'_k)^2) dy < \Lambda^{-1} \int_{\Omega} |v(y)|^2 dy = \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy$$

since $0 \leq M_v - \varepsilon'_k < M_v$, and

$$\begin{aligned} L_k(M_v - m_v - \varepsilon'_k) &= (M_v - \varepsilon'_k)^2 \Lambda^{-1} |\Omega| > (M_v - \bar{\varepsilon}')^2 \Lambda^{-1} |\Omega| \\ &> \Lambda^{-1} \int_{\Omega} |v(y)|^2 dy = \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy \end{aligned}$$

from (5.31). Therefore, the real numbers τ'_k given in (5.26) are well-defined and are such that $0 < \tau'_k < M_v - m_v - \varepsilon'_k$ for each $k \in \mathbb{N}$. We then keep the same definitions of $v_{\Omega, k}$, \underline{v} , \underline{v}_k , $\underline{v}_{k, k'}$, $\beta_{k, k'}$ and $v_{k, k'}$ as above and properties (5.29) and (5.30) hold.

Lastly, in the case when $m_v = M_v = \|v\|_{\infty}$, namely when $|v|$ is equal to the constant $\|v\|_{\infty}$ almost everywhere, then v is kept unchanged. We set $\underline{v}_{k, k'} = v_{k, k'} = v$, $\beta_{k, k'} = 1$, $\varepsilon'_k = 0$, and properties (5.29) and (5.30) hold.

Step 4: Approximation of the eigenvalue $\lambda_1(\Omega, A, v, V)$. Consider first the case when $m_v < M_v = \|v\|_{\infty}$. For each $k \in \mathbb{N}$, it follows from (5.14), (5.15), (5.22), (5.29) and (5.30) that there exists an integer $k'(k) \geq k$ such that the C^{∞} fields

$$A_k = (a_{k, i, j})_{1 \leq i, j \leq n} = A_{k, k'(k)}, \quad \Lambda_k = \Lambda_{k, k'(k)} \quad \text{and} \quad v_k = v_{k, k'(k)}$$

satisfy, for all $k \in \mathbb{N}$,

$$\left\{ \begin{array}{l} \int_{\Omega} \Lambda_k(y)^{-1} dy = \int_{\Omega} \Lambda(y)^{-1} dy, \\ 0 < m_{\Lambda} \leq m_{\Lambda} + \frac{\varepsilon_k}{2} \leq \Lambda_k \leq M_{\Lambda} - \frac{\varepsilon_k}{2} \leq M_{\Lambda} \text{ in } \bar{\Omega}, \\ A_k(y) \geq \Lambda_k(y) \text{Id for all } y \in \bar{\Omega}, \\ \int_{\Omega} |v_k(y)|^2 \Lambda_k(y)^{-1} dy = \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy, \\ \|v_k\|_{L^{\infty}(\Omega, \mathbb{R}^n)} \leq \|v\|_{L^{\infty}(\Omega, \mathbb{R}^n)} - \frac{\varepsilon'_k}{2} \leq \|v\|_{L^{\infty}(\Omega, \mathbb{R}^n)}. \end{array} \right. \quad (5.32)$$

Notice that this is possible in both cases $m_{\Lambda} < M_{\Lambda}$ (then, ε_k and τ_k are all positive) or $m_{\Lambda} = M_{\Lambda}$ (then, $\varepsilon_k = \tau_k = 0$ for all $k \in \mathbb{N}$). Notice also that $\varepsilon'_k > 0$ for all $k \in \mathbb{N}$ in this case when $m_v < M_v$. Furthermore, the matrix field A_k is symmetric with polynomial entries $a_{k, i, j}$ in $\bar{\Omega}$, and, by (5.20), (5.21) and (5.30),

$$\left\{ \begin{array}{l} a_{k, i, j} \rightarrow a_{i, j} \text{ uniformly in } \bar{\Omega} \text{ as } k \rightarrow +\infty, \\ \limsup_{k \rightarrow +\infty} \|\nabla a_{k, i, j}\|_{L^{\infty}(\Omega)} \leq \|\nabla a_{i, j}\|_{L^{\infty}(\Omega)} \leq \|A\|_{W^{1, \infty}(\Omega)} \text{ for all } 1 \leq i, j \leq n, \\ \|v_k - v\|_{L^p(\Omega, \mathbb{R}^n)} \rightarrow 0 \text{ as } k \rightarrow +\infty \text{ for all } 1 \leq p < +\infty. \end{array} \right. \quad (5.33)$$

In the case when $m_v = M_v = \|v\|_{\infty}$, then one sets $v_k = v$ for all $k \in \mathbb{N}$, and properties (5.32) and (5.33) still hold (with $\varepsilon'_k = 0$ in that case).

Let us now prove that

$$\lambda_1(\Omega, A_k, v_k, V) \rightarrow \lambda_1(\Omega, A, v, V) \text{ as } k \rightarrow +\infty. \quad (5.34)$$

Notice first that each operator $-\operatorname{div}(A_k \nabla) + v_k \cdot \nabla + V$ is elliptic because of (5.32). Fix an open non-empty ball B such that $B \subset \Omega$. It follows from [14] that

$$\min_{\overline{\Omega}} V < \lambda_1(\Omega, A_k, v_k, V) \leq \lambda_1(B, A_k, v_k, V) \text{ for all } k \in \mathbb{N}. \quad (5.35)$$

Furthermore, properties (5.32) and (5.33) imply that the sequences of matrix fields $(A_k)_{k \in \mathbb{N}}$ and $(A_k^{-1})_{k \in \mathbb{N}}$ are bounded in $W^{1,\infty}(\Omega)$ (and then in $W^{1,\infty}(B)$), and that the sequence of vector fields $(v_k)_{k \in \mathbb{N}}$ is bounded in $L^\infty(\Omega)$, and then in $L^\infty(B)$. From Lemma 1.1 in [14], there exists then a constant C independent from k such that

$$\lambda_1(B, A_k, v_k, V) \leq C \text{ for all } k \in \mathbb{N}.$$

Together with (5.35), it resorts that the sequence $(\lambda_1(\Omega, A_k, v_k, V))_{k \in \mathbb{N}}$ is bounded. Thus, for a sequence of integers $n(k) \rightarrow +\infty$ as $k \rightarrow +\infty$, one has that

$$\lambda_1(\Omega, A_{n(k)}, v_{n(k)}, V) \rightarrow \lambda \in \mathbb{R} \text{ as } k \rightarrow +\infty.$$

For each $k \in \mathbb{N}$, call φ_k the principal eigenfunction of the operator $-\operatorname{div}(A_k \nabla) + v_k \cdot \nabla + V$ in Ω with Dirichlet boundary condition, such that $\max_{\overline{\Omega}} \varphi_k = 1$. Namely, each function φ_k satisfies

$$\left\{ \begin{array}{ll} -\operatorname{div}(A_k \nabla \varphi_k) + v_k \cdot \nabla \varphi_k + V \varphi_k = \lambda_1(\Omega, A_k, v_k, V) \varphi_k & \text{in } \Omega, \\ \varphi_k > 0 & \text{in } \Omega, \\ \|\varphi_k\|_{L^\infty(\Omega)} = 1, & \\ \varphi_k = 0 & \text{on } \partial\Omega. \end{array} \right. \quad (5.36)$$

From standard elliptic estimates, each function φ_k is in $W^{2,p}(\Omega)$ for all $1 \leq p < +\infty$ and in $C^{1,\alpha}(\overline{\Omega})$ for all $0 \leq \alpha < 1$. Furthermore, since the eigenvalues $\lambda_1(\Omega, A_k, v_k, V)$ are bounded and $\|\varphi_k\|_{L^\infty(\Omega)} = 1$, it follows from (5.32) and (5.33) that the sequence $(\varphi_k)_{k \in \mathbb{N}}$ is bounded in all $W^{2,p}(\Omega)$ and $C^{1,\alpha}(\overline{\Omega})$, for all $1 \leq p < +\infty$ and $0 \leq \alpha < 1$. Up to extraction of another subsequence, one can assume that there exists $\varphi_\infty \in \bigcap_{1 \leq p < +\infty} W^{2,p}(\Omega)$ such that

$$\varphi_{n(k)} \rightarrow \varphi_\infty \text{ as } k \rightarrow +\infty, \text{ weakly in } W^{2,p}(\Omega) \text{ and strongly in } C^{1,\alpha}(\overline{\Omega})$$

for all $1 \leq p < +\infty$ and $0 \leq \alpha < 1$. Notice that (5.33) implies that $\partial_{x_{i'}} a_{k,i,j} \rightharpoonup \partial_{x_{i'}} a_{i,j}$ in $\sigma(L^\infty(\Omega), L^1(\Omega))$ as $k \rightarrow +\infty$, for all $1 \leq i', i, j \leq n$. Multiplying (5.36) for $n(k)$ by any test function in $\mathcal{D}(\Omega)$, integrating over Ω and passing to the limit as $k \rightarrow +\infty$ leads to, because of (5.32) and (5.33),

$$-\operatorname{div}(A \nabla \varphi_\infty) + v \cdot \nabla \varphi_\infty + V \varphi_\infty = \lambda \varphi_\infty \text{ in } \Omega,$$

together with

$$\varphi_\infty = 0 \text{ on } \partial\Omega, \quad \varphi_\infty \geq 0 \text{ in } \overline{\Omega} \text{ and } \max_{\overline{\Omega}} \varphi_\infty = 1.$$

The strong maximum principle and the characterization of the principal eigenvalue and eigenfunction thanks to Krein-Rutman theory imply that $\lambda = \lambda_1(\Omega, A, v, V)$ and φ_∞ is the principal eigenfunction of the operator $-\operatorname{div}(A\nabla) + v \cdot \nabla + V$ in Ω with Dirichlet boundary condition. The limiting function φ_∞ is uniquely determined because of the normalization $\max_{\overline{\Omega}} \varphi_\infty = 1$. Since the limits $\lambda = \lambda_1(\Omega, A, v, V)$ and φ_∞ do not depend on any subsequence, one concludes that (5.34) holds and that the whole sequence $(\varphi_k)_{k \in \mathbb{N}}$ converges to φ_∞ as $k \rightarrow +\infty$, weakly in $W^{2,p}(\Omega)$ and strongly in $C^{1,\alpha}(\overline{\Omega})$ for all $1 \leq p < +\infty$ and $0 \leq \alpha < 1$.

Step 5: Approximation of the principal eigenfunction of φ_K for a large K . Choose any arbitrary $\varepsilon > 0$. Let $\varepsilon' \in (0, 1)$ be such that

$$\frac{\lambda_1(\Omega, A, v, V) + 3\varepsilon' + \varepsilon' m_\Lambda^{-1} + \varepsilon' \|V^-\|_\infty}{1 - \varepsilon'} \leq \lambda_1(\Omega, A, v, V) + \varepsilon. \quad (5.37)$$

Thanks to (5.34), there exists then $K \in \mathbb{N}$ such that

$$\lambda_1(\Omega, A_K, v_K, V) \leq \lambda_1(\Omega, A, v, V) + \varepsilon'. \quad (5.38)$$

Remember that (5.32) holds with $k = K$. Let φ_K be the (unique) solution (5.36) with $k = K$. Notice that

$$v \cdot \nabla \varphi_K = -|\nabla \varphi_K| < 0 \text{ on } \partial\Omega$$

from Hopf lemma.

Call \mathcal{F} , $\underline{\mathcal{F}}$, $\overline{\mathcal{F}}$, f and \overline{f} the functions defined in $\overline{\Omega}$ by

$$\begin{cases} \mathcal{F} = -V\varphi_K + \lambda_1(\Omega, A_K, v_K, V)\varphi_K, \\ \underline{\mathcal{F}} = -v_K \cdot \nabla \varphi_K, \quad \overline{\mathcal{F}} = |v_K| \times |\nabla \varphi_K|, \\ f = \underline{\mathcal{F}} + \mathcal{F} = -v_K \cdot \nabla \varphi_K - V\varphi_K + \lambda_1(\Omega, A_K, v_K, V)\varphi_K, \\ \overline{f} = \overline{\mathcal{F}} + \mathcal{F} = |v_K| \times |\nabla \varphi_K| - V\varphi_K + \lambda_1(\Omega, A_K, v_K, V)\varphi_K. \end{cases}$$

The function \mathcal{F} is continuous in $\overline{\Omega}$. There exists then a sequence $(\mathcal{F}_l)_{l \in \mathbb{N}}$ of polynomials such that

$$\mathcal{F}_l \rightarrow \mathcal{F} \text{ uniformly in } \overline{\Omega} \text{ as } l \rightarrow +\infty.$$

Observe also that the function $\underline{\mathcal{F}}$ is in $L^\infty(\Omega)$, and that the function $\overline{\mathcal{F}}$ is nonnegative and continuous in $\overline{\Omega}$: this is true if $m_v < M_v$ because v_K is then actually of class $C^\infty(\overline{\Omega})$ and $\varphi_K \in C^1(\overline{\Omega})$; this is also true if $m_v = M_v$ because v_K is then equal to v and $|v_K| = M_v$ in $\overline{\Omega}$ up to a modification on a zero-measure set. Let $R_0 > 0$ be such that

$$\overline{\Omega} \subset B_{R_0},$$

where B_{R_0} is the open Euclidean ball of radius R_0 and center 0. Denote by $\overline{\mathcal{F}}$ (with a slight abuse of notation) a continuous extension of $\overline{\mathcal{F}}$ in \mathbb{R}^n such that $\overline{\mathcal{F}} \geq 0$ in \mathbb{R}^n and $\overline{\mathcal{F}} = 0$ in $\mathbb{R}^n \setminus B_{R_0}$. Extend by 0 the function $\underline{\mathcal{F}}$ in $\mathbb{R}^n \setminus \overline{\Omega}$ and still call $\underline{\mathcal{F}}$ this extension. Notice that

$$\underline{\mathcal{F}} \leq \overline{\mathcal{F}} \text{ in } \mathbb{R}^n. \quad (5.39)$$

For each $l \in \mathbb{N}$, call ζ_l the function defined in $\overline{B_{2R_0}}$ by

$$\zeta_l(z) = \frac{[(2R_0)^2 - |z|^2]^l}{\int_{B_{2R_0}} [(2R_0)^2 - |z'|^2]^l dz'}.$$

Extend the functions ζ_l by 0 outside B_{2R_0} and still call ζ_l these extensions. In $\overline{\Omega}$, define the functions

$$\underline{\mathcal{F}}_l = \zeta_l * \underline{\mathcal{F}} \text{ and } \overline{\mathcal{F}}_l = \zeta_l * \overline{\mathcal{F}} \text{ for all } l \in \mathbb{N}.$$

Owing to the choices of R_0 and ζ_l , the functions $\underline{\mathcal{F}}_l$ and $\overline{\mathcal{F}}_l$ are polynomials in $\overline{\Omega}$. Furthermore, since $\underline{\mathcal{F}} \in L^\infty(\mathbb{R}^n)$ and $\overline{\mathcal{F}} \in C(\mathbb{R}^n)$, there holds

$$\|\underline{\mathcal{F}}_l - \underline{\mathcal{F}}\|_{L^p(\Omega)} \xrightarrow{l \rightarrow +\infty} 0 \text{ for all } 1 \leq p < +\infty \text{ and } \overline{\mathcal{F}}_l \xrightarrow{l \rightarrow +\infty} \overline{\mathcal{F}} \text{ uniformly in } \overline{\Omega}.$$

Furthermore,

$$\underline{\mathcal{F}}_l = \zeta_l * \underline{\mathcal{F}} \leq \zeta_l * \overline{\mathcal{F}} = \overline{\mathcal{F}}_l \text{ in } \overline{\Omega} \text{ for all } l \in \mathbb{N}$$

because of (5.39). It follows that

$$\underline{f}_l = \underline{\mathcal{F}}_l + \mathcal{F}_l \xrightarrow{l \rightarrow +\infty} \underline{f} \text{ in } L^p(\Omega) \text{ for all } 1 \leq p < +\infty, \quad \overline{f}_l = \overline{\mathcal{F}}_l + \mathcal{F}_l \xrightarrow{l \rightarrow +\infty} \overline{f} \text{ uniformly in } \overline{\Omega} \quad (5.40)$$

together with

$$\underline{f}_l \leq \overline{f}_l \text{ in } \overline{\Omega} \text{ for all } l \in \mathbb{N}. \quad (5.41)$$

Remember that the function φ_K satisfies

$$-\operatorname{div}(A_K \nabla \varphi_K) = \underline{f} \text{ in } \Omega,$$

with $\varphi_K = 0$ on $\partial\Omega$. For each $l \in \mathbb{N}$, call ψ_l the solution of

$$\begin{cases} -\operatorname{div}(A_K \nabla \psi_l) &= \underline{f}_l \text{ in } \Omega, \\ \psi_l &= 0 \text{ on } \partial\Omega. \end{cases} \quad (5.42)$$

Each function ψ_l is then analytic in Ω (remember that A_K is a field of symmetric positive definite matrices with polynomial entries, and that each \underline{f}_l is a polynomial in $\overline{\Omega}$). From standard elliptic estimates, the functions ψ_l converge to the function φ_K as $l \rightarrow +\infty$ in $W^{2,p}(\Omega)$ and $C^{1,\alpha}(\overline{\Omega})$ for all $1 \leq p < +\infty$ and $0 \leq \alpha < 1$. Since $\varphi_K > 0$ in Ω and $|\nabla \varphi_K| > 0 = \varphi_K$ on $\partial\Omega$, one then has

$$\psi_l > 0 \text{ in } \Omega \text{ and } |\nabla \psi_l| > 0 \text{ on } \partial\Omega \text{ for } l \text{ large enough.} \quad (5.43)$$

Furthermore, from (5.41), there holds

$$\begin{aligned} -\operatorname{div}(A_K \nabla \psi_l) - |v_K| \times |\nabla \psi_l| - \varepsilon' |\nabla \psi_l| + [V - \lambda_1(\Omega, A, v, V) - 2\varepsilon'] \psi_l \\ \leq \overline{f}_l - |v_K| \times |\nabla \psi_l| - \varepsilon' |\nabla \psi_l| + [V - \lambda_1(\Omega, A, v, V) - 2\varepsilon'] \psi_l \quad \text{in } \Omega. \end{aligned}$$

From (5.40) and the definition of \overline{f} , it follows that

$$\begin{aligned} \overline{f}_l - |v_K| \times |\nabla \psi_l| - \varepsilon' |\nabla \psi_l| + [V - \lambda_1(\Omega, A, v, V) - 2\varepsilon'] \psi_l \\ \rightarrow [\lambda_1(\Omega, A_K, v_K, V) - \lambda_1(\Omega, A, v, V) - \varepsilon'] \varphi_K - \varepsilon' (|\nabla \varphi_K| + \varphi_K) \end{aligned}$$

as $l \rightarrow +\infty$, uniformly in $\overline{\Omega}$. Using (5.38) and the properties of φ_K , it follows that the continuous function

$$[\lambda_1(\Omega, A_K, v_K, V) - \lambda_1(\Omega, A, v, V) - \varepsilon'] \varphi_K - \varepsilon' (|\nabla \varphi_K| + \varphi_K)$$

is negative in $\overline{\Omega}$. Therefore, there is $L \in \mathbb{N}$ large enough so that (5.43) holds with $l = L$ and

$$-\operatorname{div}(A_K \nabla \psi_L) - |v_K| \times |\nabla \psi_L| - \varepsilon' |\nabla \psi_L| + [V - \lambda_1(\Omega, A, v, V) - 2\varepsilon'] \psi_L < 0 \text{ in } \Omega. \quad (5.44)$$

Step 6: An inequality for the rearranged fields in the ball Ω^* . Apply then the results of Section 3 to the function

$$\psi = \psi_L$$

and to the data

$$(A_\Omega, \Lambda_\Omega, \omega, V) = (A_K, \Lambda_K, |v_K|, V).$$

From the previous steps, these fields satisfy all assumptions of Section 3. Given ψ , one can then define the sets Z , Y , E and the function ρ as in Section 3.1. Given ψ and the data $A_\Omega = A_K$, $\Lambda_\Omega = \Lambda_K$, $\omega = |v_K|$ and V , one can also define the corresponding rearranged fields $\tilde{\psi}$, $\widehat{\Lambda}$, \widehat{v} , \widehat{V} and U given by (3.3), (3.4), (3.7), (3.8), (3.9), (3.12), (3.33) and (3.34).

One recalls that $\lambda_1(\Omega, A, v, V) \geq 0$ by assumption. From (5.44) and Corollary 3.11 applied with $\omega_0 = \varepsilon' \geq 0$ and

$$\mu = \lambda_1(\Omega, A, v, V) + 2\varepsilon' \geq 0,$$

it follows that

$$\int_{\Omega^*} \left[\widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 - \varepsilon' |\nabla \tilde{\psi}(x)| \tilde{\psi}(x) + \widehat{V}(x) \tilde{\psi}(x)^2 \right] e^{-U(x)} dx \leq \mu \int_{\Omega^*} \tilde{\psi}(x)^2 e^{-U(x)} dx. \quad (5.45)$$

Remember that the function $\tilde{\psi}$ is radially symmetric, continuous and decreasing with respect to $|x|$ in $\overline{\Omega^*}$, and that $\tilde{\psi} \in H_0^1(\Omega^*)$. The field $\widehat{\Lambda}$ is radially symmetric and belongs in $L^\infty(\Omega^*)$. Furthermore,

$$\int_{\Omega^*} \widehat{\Lambda}(x)^{-1} dx = \int_{\Omega} \Lambda_K(y)^{-1} dy = \int_{\Omega} \Lambda(y)^{-1} dy \quad (5.46)$$

from (3.6) and (5.32). On the other hand,

$$0 < m_\Lambda \leq m_\Lambda + \frac{\tau_K}{2} \leq \min_{\overline{\Omega}} \Lambda_K \leq \operatorname{ess\,inf}_{\Omega^*} \widehat{\Lambda} \leq \operatorname{ess\,sup}_{\Omega^*} \widehat{\Lambda} \leq \max_{\overline{\Omega}} \Lambda_K \leq M_\Lambda - \frac{\varepsilon_K}{2} \leq M_\Lambda \quad (5.47)$$

from (3.5) and (5.32).

From (5.45) and (5.47), it follows that

$$(1 - \varepsilon') \int_{\Omega^*} \widehat{\Lambda} |\nabla \tilde{\psi}|^2 e^{-U} + \int_{\Omega^*} \widehat{V} \tilde{\psi}^2 e^{-U} \leq \mu \int_{\Omega^*} \tilde{\psi}^2 e^{-U} + \frac{\varepsilon'}{4} \int_{\Omega^*} \widehat{\Lambda}^{-1} \tilde{\psi}^2 e^{-U} \leq (\mu + \varepsilon' m_\Lambda^{-1}) \int_{\Omega^*} \tilde{\psi}^2 e^{-U}.$$

On the other hand, the field \widehat{V} is radially symmetric, it belongs to $L^\infty(\Omega^*)$, and

$$-\|V\|_\infty \leq \min_{\overline{\Omega}}(-V^-) \leq \operatorname{ess\,inf}_{\Omega^*} \widehat{V} \leq \operatorname{ess\,sup}_{\Omega^*} \widehat{V} \leq 0$$

(see (3.13)). Therefore,

$$(1 - \varepsilon') \int_{\Omega^*} \left(\widehat{\Lambda} |\nabla \tilde{\psi}|^2 + \widehat{V} \tilde{\psi}^2 \right) e^{-U} \leq (\mu + \varepsilon' m_\Lambda^{-1} + \varepsilon' \|V^-\|_\infty) \int_{\Omega^*} \tilde{\psi}^2 e^{-U},$$

that is

$$\int_{\Omega^*} \left(\widehat{\Lambda} |\nabla \tilde{\psi}|^2 + \widehat{V} \tilde{\psi}^2 \right) e^{-U} \leq \frac{\lambda_1(\Omega, A, v, V) + 2\varepsilon' + \varepsilon' m_\Lambda^{-1} + \varepsilon' \|V^-\|_\infty}{1 - \varepsilon'} \times \int_{\Omega^*} \tilde{\psi}^2 e^{-U}. \quad (5.48)$$

Step 7: Approximation of the rearranged fields in Ω^* . First, define the function $\underline{\widehat{\Lambda}}$ almost everywhere in \mathbb{R}^n by:

$$\underline{\widehat{\Lambda}}(x) = \begin{cases} \widehat{\Lambda}(x) & \text{if } x \in \Omega^*, \\ \min_{\overline{\Omega}} \Lambda_K & \text{if } x \in \mathbb{R}^n \setminus \Omega^*, \end{cases}$$

and then, for each $m \in \mathbb{N}$,

$$\underline{\Lambda}_m^* = \left(\rho_m * \underline{\widehat{\Lambda}}^{-1} \right)^{-1} \text{ in } \mathbb{R}^n,$$

where the mollifiers ρ_m can be assumed to be radially symmetric for all $m \in \mathbb{N}$. Next, for every $m \in \mathbb{N}$, call

$$\gamma_m = \frac{\|(\underline{\Lambda}_m^*)^{-1}\|_{L^1(\Omega^*)}}{\|\Lambda^{-1}\|_{L^1(\Omega)}} > 0$$

and

$$\Lambda_m^*(x) = \gamma_m \underline{\Lambda}_m^*(x) \text{ for all } x \in \overline{\Omega^*}.$$

As in Step 1, it follows from the above definitions and from (5.46) and (5.47) that each function Λ_m^* is radially symmetric and of class $C^\infty(\overline{\Omega^*})$, that $\lim_{m \rightarrow +\infty} \gamma_m = 1$, that

$$\int_{\Omega^*} \Lambda_m^*(x)^{-1} dx = \int_{\Omega} \Lambda(y)^{-1} dy \text{ for all } m \in \mathbb{N} \quad (5.49)$$

and that

$$\begin{cases} 0 < \gamma_m \times \left(m_\Lambda + \frac{\tau_K}{2} \right) \leq \min_{\overline{\Omega^*}} \Lambda_m^* \leq \max_{\overline{\Omega^*}} \Lambda_m^* \leq \gamma_m \times \left(M_\Lambda - \frac{\varepsilon_K}{2} \right) \text{ for all } m \in \mathbb{N}, \\ \gamma_m \times \left(m_\Lambda + \frac{\tau_K}{2} \right) \xrightarrow{m \rightarrow +\infty} m_\Lambda + \frac{\tau_K}{2}, \quad \gamma_m \times \left(M_\Lambda - \frac{\varepsilon_K}{2} \right) \xrightarrow{m \rightarrow +\infty} M_\Lambda - \frac{\varepsilon_K}{2}. \end{cases} \quad (5.50)$$

From reciprocal of Lebesgue's theorem, one can also assume without loss of generality (even if it means extracting a subsequence) that

$$\Lambda_m^*(x) \rightarrow \underline{\widehat{\Lambda}}(x) = \widehat{\Lambda}(x) \text{ as } m \rightarrow +\infty \text{ for a.e. } x \in \Omega^*.$$

Remember also that if $m_\Lambda < M_\Lambda$ then $\varepsilon_K > 0$, $\tau_K > 0$ and notice that if $m_\Lambda = M_\Lambda$ then $\varepsilon_K = \tau_K = 0$, $\gamma_m = 1$, $\Lambda_m^* = m_\Lambda$ in $\overline{\Omega^*}$ for all $m \in \mathbb{N}$ and properties (5.49) and (5.50) hold immediately.

Next, owing to the definition (3.9), the vector field \widehat{v} can be written as

$$\widehat{v}(x) = |\widehat{v}(x)| e_r(x) \text{ in } \Omega^*,$$

where $|\widehat{v}|$ is radially symmetric. Furthermore, as in Step 3 and since $\omega = |v_K|$, it follows from (3.10), (3.11) and (5.32) that

$$\int_{\Omega^*} |\widehat{v}(x)|^2 \widehat{\Lambda}(x)^{-1} dx = \int_{\Omega} |v_K(y)|^2 \Lambda_K(y)^{-1} dy = \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy \quad (5.51)$$

and

$$\|\widehat{v}\|_{L^\infty(\Omega^*, \mathbb{R}^n)} \leq \|\omega\|_{L^\infty(\Omega)} = \|v_K\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} - \frac{\varepsilon'_K}{2} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}. \quad (5.52)$$

Call

$$\underline{\omega}(x) = \begin{cases} |\widehat{v}(x)| & \text{if } x \in \Omega^*, \\ 0 & \text{if } x \in \mathbb{R}^n \setminus \Omega^* \end{cases}$$

and, for each $m \in \mathbb{N}$,

$$\underline{v}_m^*(x) = (\rho_m * \underline{\omega})(x) e_r(x) \text{ for all } x \in \mathbb{R}^n \setminus \{0\}.$$

One has $\|\underline{v}_m^*\|_{L^\infty(\mathbb{R}^n, \mathbb{R}^n)} \leq \|\underline{\omega}\|_{L^\infty(\mathbb{R}^n, \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}$ and one can assume without loss of generality that

$$\underline{v}_m^*(x) \rightarrow \widehat{v}(x) \text{ as } m \rightarrow +\infty \text{ for a.e. } x \in \Omega^*.$$

Consider first the case when $\|v\|_{L^\infty(\Omega, \mathbb{R}^n)} > 0$. Therefore, for m large enough (one can then assume that this holds for all m without loss of generality), the real numbers

$$\delta_m = \left(\frac{\int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy}{\int_{\Omega^*} |\underline{v}_m^*(x)|^2 \Lambda_m^*(x)^{-1} dx} \right)^{1/2} = \left(\frac{\int_{\Omega^*} |\widehat{v}(x)|^2 \widehat{\Lambda}(x)^{-1} dx}{\int_{\Omega^*} |\underline{v}_m^*(x)|^2 \Lambda_m^*(x)^{-1} dx} \right)^{1/2} \quad (5.53)$$

are well-defined, positive, and they are such that $\delta_m \rightarrow 1$ as $m \rightarrow +\infty$. Therefore, the vector fields defined by

$$v_m^*(x) = \delta_m \underline{v}_m^*(x) \text{ for all } x \in \overline{\Omega^*} \setminus \{0\}$$

are of class $C^\infty(\overline{\Omega^*} \setminus \{0\})$ and converge to $\widehat{v}(x)$ as $m \rightarrow +\infty$ for almost every $x \in \Omega^*$. These fields can be written as

$$v_m^*(x) = |v_m^*(x)| e_r(x) \text{ for all } x \in \overline{\Omega^*} \setminus \{0\}$$

and $|v_m^*|$ is radially symmetric, of class $C^\infty(\overline{\Omega^*} \setminus \{0\})$ and can actually be extended at 0 to a $C^\infty(\overline{\Omega^*})$ function. Furthermore, it follows from (5.52) and (5.53) that

$$\int_{\Omega^*} |v_m^*(x)|^2 \Lambda_m^*(x)^{-1} dx = \int_{\Omega} |v(y)|^2 \Lambda(y)^{-1} dy \quad (5.54)$$

and

$$\|v_m^*\|_{L^\infty(\Omega^*, \mathbb{R}^n)} \leq \delta_m \|\widehat{v}\|_{L^\infty(\Omega^*, \mathbb{R}^n)} \leq \delta_m \times \left(\|v\|_{L^\infty(\Omega, \mathbb{R}^n)} - \frac{\varepsilon'_K}{2} \right) \xrightarrow{m \rightarrow +\infty} \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} - \frac{\varepsilon'_K}{2}. \quad (5.55)$$

Lastly, remember that if $m_v < M_v$, then $\varepsilon'_K > 0$, and notice that if $m_v = M_v$ (this is the case if $\|v\|_{L^\infty(\Omega, \mathbb{R}^n)} = 0$), then $\varepsilon'_K = 0$, $\omega = m_v$ in Ω , $\widehat{v} = v_m^* = m_v e_r$ in Ω^* and properties (5.54) and (5.55) hold immediately with $\delta_m = 1$.

Fix now an arbitrary unit vector e in \mathbb{R}^n and define, for each $m \in \mathbb{N}$,

$$\forall x \in \overline{\Omega^*}, \quad U_m^*(x) = \int_0^{|x|} |v_m^*(re)| \Lambda_m^*(re)^{-1} dr.$$

As in Proposition 3.10, the definition of U_m^* does not depend on the choice of e . Furthermore, each function U_m^* is continuous in $\overline{\Omega^*}$, radially symmetric, of class $C^\infty(\overline{\Omega^*} \setminus \{0\})$ and it satisfies

$$\nabla U_m^*(x) = \Lambda_m^*(x)^{-1} |v_m^*(x)| e_r(x) = \Lambda_m^*(x)^{-1} v_m^*(x) \text{ for all } x \in \overline{\Omega^*} \setminus \{0\}. \quad (5.56)$$

On the other hand, each function U_m^* is nonnegative in $\overline{\Omega^*}$ and it follows from (5.50) and (5.55) that

$$\|U_m^*\|_{L^\infty(\Omega^*)} \leq \delta_m m_\Lambda^{-1} R \times \left(\|v\|_{L^\infty(\Omega, \mathbb{R}^n)} - \frac{\varepsilon'_K}{2} \right) \xrightarrow{m \rightarrow +\infty} m_\Lambda^{-1} R \times \left(\|v\|_{L^\infty(\Omega, \mathbb{R}^n)} - \frac{\varepsilon'_K}{2} \right). \quad (5.57)$$

Moreover, since all fields $|v_m^*|$, $|\widehat{v}|$, Λ_m^* and $\widehat{\Lambda}$ are radially symmetric, it follows from the above estimates and Lebesgue's dominated convergence theorem that

$$U_m^*(x) \xrightarrow{m \rightarrow +\infty} U(x) \text{ for all } x \in \overline{\Omega^*},$$

where U is given in (3.34).

Lastly, Corollary 5.4 applied with $g = -V^-$ and $\widehat{g} = \widehat{V}$ provides the existence of two sequences of radially symmetric fields $(V_m^*)_{m \in \mathbb{N}}$ and $(\overline{V}_m^*)_{m \in \mathbb{N}}$ in Ω^* such that, for each $m \in \mathbb{N}$, $V_m^* \in L^\infty(\Omega^*)$, $\overline{V}_m^* \in C^\infty(\overline{\Omega^*})$,

$$-\|V\|_{L^\infty(\Omega)} \leq \min_{\overline{\Omega}}(-V^-) \leq V_m^* \leq \overline{V}_m^* \leq 0 \text{ in } \Omega^* \quad (5.58)$$

and the distribution functions of V_m^* and \overline{V}_m^* satisfy

$$\mu_{V_m^*} = \mu_{-V^-} \text{ and } \mu_{|\overline{V}_m^*|} \leq \mu_{V^-} \leq \mu_{|V|}.$$

Furthermore, the fields V_m^* and \overline{V}_m^* are constructed so that

$$V_m^*, \overline{V}_m^* \xrightarrow{m \rightarrow +\infty} \widehat{V} \text{ in } \sigma(L^p(\Omega^*), L^{p'}(\Omega^*))$$

for all $1 < p \leq +\infty$, with $1/p + 1/(p') = 1$.

Step 8: An inequality for the eigenvalue $\lambda_1(\Omega^*, \Lambda_m^* \text{Id}, v_m^*, \bar{V}_m^*)$ for large m . Remember first that the function $\tilde{\psi}$ is continuous nonnegative in $\bar{\Omega}^*$ and that $\tilde{\psi}(0) \geq \max_{\bar{\Omega}} \psi_L > 0$ because of Corollary 3.6 and (5.43). It also follows from Lebesgue's dominated convergence theorem and all estimates of Step 7 that

$$\begin{cases} \int_{\Omega} \left(\Lambda_m^* |\nabla \tilde{\psi}|^2 + \bar{V}_m^* \tilde{\psi}^2 \right) e^{-U_m^*} \rightarrow \int_{\Omega} \left(\hat{\Lambda} |\nabla \tilde{\psi}|^2 + \hat{V} \tilde{\psi}^2 \right) e^{-U}, \\ \int_{\Omega} \tilde{\psi}^2 e^{-U_m^*} \rightarrow \int_{\Omega} \tilde{\psi}^2 e^{-U} > 0, \end{cases}$$

as $m \rightarrow +\infty$. Therefore, from (5.48), there exists $M \in \mathbb{N}$ such that

$$\begin{aligned} \int_{\Omega} \left(\Lambda_M^* |\nabla \tilde{\psi}|^2 + \bar{V}_M^* \tilde{\psi}^2 \right) e^{-U_M^*} &\leq \frac{\lambda_1(\Omega, A, v, V) + 3\varepsilon' + \varepsilon' m_{\Lambda}^{-1} + \varepsilon' \|V^-\|_{\infty}}{1 - \varepsilon'} \\ &\quad \times \int_{\Omega^*} \tilde{\psi}^2 e^{-U_M^*}. \end{aligned} \quad (5.59)$$

Remember that (5.49) and (5.54) hold with $m = M$. Furthermore, because of (5.50) and (5.55), one can choose M large enough so that

$$m_{\Lambda} \leq \Lambda_M^* \leq M_{\Lambda} \text{ in } \bar{\Omega}^* \text{ and } \|v_M^*\|_{L^{\infty}(\Omega^*, \mathbb{R}^n)} \leq \|v\|_{L^{\infty}(\Omega, \mathbb{R}^n)}.$$

Notice that this holds also when $m_{\Lambda} = M_{\Lambda}$ (in this case, $\Lambda_M^* = m_{\Lambda}$ in $\bar{\Omega}^*$) or when $m_v = M_v$ (in this case, $v_M^* = m_v e_r$ in $\bar{\Omega}^* \setminus \{0\}$).

Call now

$$I = \inf_{\phi \in H_0^1(\Omega^*) \setminus \{0\}} \frac{\int_{\Omega^*} \left(\Lambda_M^* |\nabla \phi|^2 + \bar{V}_M^* \phi^2 \right) e^{-U_M^*}}{\int_{\Omega^*} \phi^2 e^{-U_M^*}}.$$

It follows from (5.59) that

$$I \leq \frac{\lambda_1(\Omega, A, v, V) + 3\varepsilon' + \varepsilon' m_{\Lambda}^{-1} + \varepsilon' \|V^-\|_{\infty}}{1 - \varepsilon'}. \quad (5.60)$$

Furthermore, I is clearly finite and $I \geq \min_{\bar{\Omega}^*} \bar{V}_M^* \geq -\|V^-\|_{\infty}$. It is classical to check that I is actually a minimum, which is reached at a function $\varphi_M^* \in H_0^1(\Omega^*) \setminus \{0\}$ such that $\varphi_M^* \geq 0$ a.e. in Ω^* and

$$\int_{\Omega^*} \left(\Lambda_M^* \nabla \varphi_M^* \cdot \nabla \phi + \bar{V}_M^* \varphi_M^* \phi \right) e^{-U_M^*} = I \int_{\Omega^*} \varphi_M^* \phi e^{-U_M^*}$$

for all $\phi \in H_0^1(\Omega^*)$. Because of (5.56) with $m = M$, the change of functions $\phi = \Phi e^{U_M^*}$ leads to

$$\int_{\Omega^*} \Lambda_M^* \nabla \varphi_M^* \cdot \nabla \Phi + v_M^* \cdot \nabla \varphi_M^* \Phi + \bar{V}_M^* \varphi_M^* \Phi = I \int_{\Omega^*} \varphi_M^* \Phi$$

for all $\Phi \in H_0^1(\Omega^*)$. From H^2 regularity and $W^{2,p}$ estimates, it then follows that φ_M^* is actually in all $W^{2,p}(\Omega^*)$ and $C^{1,\alpha}(\bar{\Omega}^*)$ for all $1 \leq p < +\infty$ and $0 \leq \alpha < 1$ and that

$$\begin{cases} -\text{div}(\Lambda_M^* \nabla \varphi_M^*) + v_M^* \cdot \nabla \varphi_M^* + \bar{V}_M^* \varphi_M^* = I \varphi_M^* & \text{in } \Omega^*, \\ \varphi_M^* = 0 & \text{on } \partial\Omega^*, \\ \varphi_M^* \geq 0 & \text{in } \Omega^*. \end{cases}$$

Since $\varphi_M^* \neq 0$, one concludes from Krein-Rutman theory that φ_M^* is –up to multiplication by a positive constant– the principal eigenfunction of the operator $-\operatorname{div}(\Lambda_M^* \nabla) + v_M^* \cdot \nabla + \overline{V}_M^*$ with Dirichlet boundary condition, and that I is the principal eigenvalue $I = \lambda_1(\Omega^*, \Lambda_M^* \operatorname{Id}, v_M^*, \overline{V}_M^*)$. Together with (5.37) and (5.60), it resorts that

$$\lambda_1(\Omega^*, \Lambda_M^* \operatorname{Id}, v_M^*, \overline{V}_M^*) \leq \frac{\lambda_1(\Omega, A, v, V) + 3\varepsilon' + \varepsilon' m_\Lambda^{-1} + \varepsilon' \|V^-\|_\infty}{1 - \varepsilon'} \leq \lambda_1(\Omega, A, v, V) + \varepsilon.$$

Lastly, since $V_M^* \leq \overline{V}_M^*$ in Ω (see (5.58) with $m = M$), it follows from [14] that

$$\lambda_1(\Omega^*, \Lambda_M^* \operatorname{Id}, v_M^*, V_M^*) \leq \lambda_1(\Omega^*, \Lambda_M^* \operatorname{Id}, v_M^*, \overline{V}_M^*) \leq \lambda_1(\Omega, A, v, V) + \varepsilon.$$

Step 9: Conclusion. Since $\varepsilon \in (0, 1)$ was arbitrary, the proof of the first part of Theorem 2.1 is complete with the choice

$$(\Lambda^*, \omega^*, v^*, V^*, \overline{V}^*) = (\Lambda_M^*, |v_M^*|, v_M^*, V_M^*, \overline{V}_M^*).$$

Step 10: The case when Λ is equal to a constant $\gamma > 0$ in Ω . It follows from the previous steps that Λ^* is equal to the same constant γ in Ω^* . Furthermore, there exists a family $(v_\varepsilon^*, V_\varepsilon^*)_{\varepsilon > 0}$ of fields satisfying the same bounds (2.2) as v^* and V^* , together with

$$\lambda_1(\Omega^*, \gamma \operatorname{Id}, v_\varepsilon^*, V_\varepsilon^*) \leq \lambda_1(\Omega, A, v, V) + \varepsilon$$

for all $\varepsilon > 0$. Furthermore, $v_\varepsilon^* = |v_\varepsilon^*| e_r$, and $|v_\varepsilon^*|$ and V_ε^* are $C^\infty(\overline{\Omega^*})$ and radially symmetric.

Take any sequence $(\varepsilon_k)_{k \in \mathbb{N}}$ of positive numbers such that $\varepsilon_k \rightarrow 0$ as $k \rightarrow +\infty$. Up to extraction of a subsequence, there exist two radially symmetric functions $\omega_0^* \geq 0$ and $V_0^* \leq 0$ such that $|v_{\varepsilon_k}^*| \rightharpoonup \omega_0^*$ and $V_{\varepsilon_k}^* \rightharpoonup V_0^*$ weakly in $L^p(\Omega^*)$ for all $1 \leq p < +\infty$ and weak- $*$ in $L^\infty(\Omega^*)$. Furthermore, the fields $v_0^* = \omega_0^* e_r$ and V_0^* satisfy the bounds (2.4). Since $-\max_{\overline{\Omega}} V^- < \lambda_1(\Omega^*, \gamma \operatorname{Id}, v_{\varepsilon_k}^*, V_{\varepsilon_k}^*) \leq \lambda_1(\Omega, A, v, V) + \varepsilon_k$ for all $k \in \mathbb{N}$, the sequence $(\lambda_1(\Omega^*, \gamma \operatorname{Id}, v_{\varepsilon_k}^*, V_{\varepsilon_k}^*))_{k \in \mathbb{N}}$ converges, up to extraction of a subsequence, to a real number $\lambda_0^* \in [-\max_{\overline{\Omega}} V^-, \lambda_1(\Omega, A, v, V)]$. From standard elliptic estimates, the principal eigenfunctions $\varphi_k = \varphi_{\Omega^*, \gamma \operatorname{Id}, v_{\varepsilon_k}^*, V_{\varepsilon_k}^*}$ are bounded independently of k in all $W^{2,r}(\Omega^*)$ for $1 \leq r < \infty$. Up to extraction of a subsequence, they converge weakly in $W^{2,r}(\Omega^*)$ for all $1 \leq r < \infty$ and strongly in $C^{1,\alpha}(\overline{\Omega^*})$ for all $0 \leq \alpha < 1$ to a solution φ_0^* of

$$-\gamma \Delta \varphi_0^* + v_0^* \cdot \nabla \varphi_0^* + V_0^* \varphi_0^* = \lambda_0^* \varphi_0^* \text{ in } \Omega^*$$

such that $\varphi_0^* \geq 0$ in Ω^* , $\varphi_0^* = 0$ on $\partial\Omega^*$ and $\|\varphi_0^*\|_{L^\infty(\Omega^*)} = 1$. By uniqueness, it resorts that $\lambda_0^* = \lambda_1(\Omega^*, \gamma \operatorname{Id}, v_0^*, V_0^*)$ and $\varphi_0^* = \varphi_{\Omega^*, \gamma \operatorname{Id}, v_0^*, V_0^*}$. Thus,

$$\lambda_1(\Omega^*, \gamma \operatorname{Id}, v_0^*, V_0^*) \leq \lambda_1(\Omega, A, v, V).$$

Notice that the radially symmetric fields $|v_0^*|$ and V_0^* satisfy the bounds (2.4), but may not be smooth anymore. \square

Remark 5.5 Assume now that $0 < m_\Lambda < M_\Lambda$ and that $\Lambda = m_\Lambda$ in a neighbourhood of $\partial\Omega$, that is $\Lambda = m_\Lambda$ in the set $\{x \in \Omega, d(x, \partial\Omega) < \gamma\}$ for some $\gamma > 0$. Then, under the notations of the previous proof, besides the aforementioned conditions of Steps 1 to 5, one can choose $K \in \mathbb{N}$ large enough so that $\|\Lambda^{-1}\|_{L^1(\Omega)} < |\Omega| \times (m_\Lambda + 4\tau_K)^{-1}$ and

$$\Lambda_K = \min_{\overline{\Omega}} \Lambda_K \in \left[m_\Lambda + \frac{\tau_K}{2}, m_\Lambda + 2\tau_K \right]$$

in a (smaller) neighbourhood of $\partial\Omega$. Then, owing to the definition of $\widehat{\Lambda}$ in Step 6 (from Section 3), it follows that $\widehat{\Lambda} = \min_{\overline{\Omega}} \Lambda_K = \text{ess inf}_{\Omega^*} \widehat{\Lambda}$ in a neighbourhood of $\partial\Omega^*$. Finally, besides the conditions of Steps 7 and 8, one can choose $M \in \mathbb{N}$ large enough so that

$$\Lambda_M^* = \min_{\overline{\Omega^*}} \Lambda_M^* \in \left[m_\Lambda + \frac{\tau_K}{4}, m_\Lambda + 4\tau_K \right]$$

in a neighbourhood of $\partial\Omega^*$. Since $\|\Lambda^{-1}\|_{L^1(\Omega)} < |\Omega| \times (m_\Lambda + 4\tau_K)^{-1}$ and $\|(\Lambda_M^*)^{-1}\|_{L^1(\Omega^*)} = \|\Lambda^{-1}\|_{L^1(\Omega)}$, one concludes that Λ_M^* ($= \Lambda^*$) is not constant in $\overline{\Omega^*}$.

It then follows that, in Theorem 2.1, the functions Λ and Λ^* are not constant in general. Actually, for the same reason, the same observation is true for Theorems 2.3 and 2.7.

Remark 5.6 Consider now the case when Ω is a general open subset of \mathbb{R}^n with finite measure, and let $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$, $v \in L^\infty(\Omega, \mathbb{R}^n)$, $V \in L^\infty(\Omega) \cap C(\overline{\Omega})$ be such that $A \geq \gamma \text{Id}$ in $\overline{\Omega}$ for some constant $\gamma > 0$. Assume that $\lambda_1(\Omega, A, v, V) \geq 0$, where $\lambda_1(\Omega, A, v, V)$ has been characterized in Remark 2.11. We claim that we get a similar conclusion as in Theorem 2.1 with $\Lambda = \gamma$. Indeed, for each $\varepsilon > 0$, it follows from (2.17) and (2.18) that there exists $\Omega' \in \mathcal{C}$ such that $\Omega' \subset\subset \Omega$ and

$$0 \leq \lambda_1(\Omega, A, v, V) \leq \lambda_1(\Omega', A, v, V) \leq \lambda_1(\Omega, A, v, V) + \varepsilon.$$

To make notations simpler, we use the same symbols for the fields in Ω and their restrictions in Ω' . Apply Theorem 2.1 to Ω' ($\subset \Omega$) and to the fields A, v, V with $\Lambda = \gamma$, and call Ω_ε^* the ball centered at the origin with the same measure as Ω' . There exist two radially symmetric bounded functions $\omega_\varepsilon^* \geq 0$ and $V_\varepsilon^* \leq 0$ in Ω_ε^* such that, for $v_\varepsilon^* = \omega_\varepsilon^* e_r$ in Ω_ε^* ,

$$\lambda_1(\Omega_\varepsilon^*, \gamma \text{Id}, v_\varepsilon^*, V_\varepsilon^*) \leq \lambda_1(\Omega', A, v, V) \leq \lambda_1(\Omega, A, v, V) + \varepsilon \quad (5.61)$$

and

$$\begin{cases} \|v_\varepsilon^*\|_{L^\infty(\Omega_\varepsilon^*, \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega', \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}, & \|v_\varepsilon^*\|_{L^2(\Omega_\varepsilon^*, \mathbb{R}^n)} \leq \|v\|_{L^2(\Omega', \mathbb{R}^n)} \leq \|v\|_{L^2(\Omega, \mathbb{R}^n)}, \\ -\text{ess sup } V^- \leq -\max_{\Omega'} V^- \leq V_\varepsilon^* \leq 0 \text{ a.e. in } \Omega_\varepsilon^*, \\ \|V_\varepsilon^*\|_{L^p(\Omega_\varepsilon^*)} \leq \|V^-\|_{L^p(\Omega')} \leq \|V^-\|_{L^p(\Omega)} \text{ for all } 1 \leq p \leq +\infty. \end{cases}$$

Extend v_ε^* by the vector 0 in $\Omega^* \setminus \Omega_\varepsilon^*$ and extend V_ε^* by 0 in $\Omega^* \setminus \Omega_\varepsilon^*$ too, and still call v_ε^* and V_ε^* these extensions, which are now defined in Ω^* ($\supset \Omega_\varepsilon^*$) and satisfy the same bounds as above in Ω^* (remember that Ω^* denotes the ball centered at the origin and having the same Lebesgue measure as Ω). Furthermore,

$$\lambda_1(\Omega^*, \gamma \text{Id}, v_\varepsilon^*, V_\varepsilon^*) \leq \lambda_1(\Omega_\varepsilon^*, \gamma \text{Id}, v_\varepsilon^*, V_\varepsilon^*) \leq \lambda_1(\Omega, A, v, V) + \varepsilon$$

because $\Omega_\varepsilon^* \subset \Omega^*$ and because of (5.61). One can then argue as in Step 10 of the proof of Theorem 2.1 and one can then pass to the limit for a sequence $\varepsilon_k \rightarrow 0^+$ as $k \rightarrow +\infty$. There exist then two radially symmetric bounded functions $\omega_0^* \geq 0$ and $V_0^* \leq 0$ in Ω^* such that, for $v_0^* = \omega_0^* e_r$ in Ω^* ,

$$\lambda_1(\Omega^*, \gamma \text{Id}, v_0^*, V_0^*) \leq \lambda_1(\Omega, A, v, V)$$

and the bounds (2.4) are satisfied (with $-\text{ess sup}_\Omega V^-$ instead of $-\max_{\overline{\Omega}} V^-$ in (2.4)).

Let us now turn to the proof of Theorem 2.3. It shall use the results of Section 4 and it follows the same scheme as the one of Theorem 2.1.

Proof of Theorem 2.3. Assume here that $\Omega \in \mathcal{C}$ is not a ball. Let $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$, $\Lambda \in L_+^\infty(\Omega)$, $v \in L^\infty(\Omega, \mathbb{R}^n)$, $V \in C(\overline{\Omega})$, $\overline{M}_A > 0$, $\underline{m}_\Lambda > 0$, $\overline{M}_v \geq 0$ and $\overline{M}_V \geq 0$ be such that $A \geq \Lambda \text{Id}$ a.e. in Ω and

$$\|A\|_{W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))} \leq \overline{M}_A, \text{ess inf}_\Omega \Lambda \geq \underline{m}_\Lambda, \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} \leq \overline{M}_v \text{ and } \|V\|_{L^\infty(\Omega)} \leq \overline{M}_V.$$

Throughout the proof, the notation $\rho = \rho(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V)$ denotes a constant ρ which only depends on Ω , n , \overline{M}_A , \underline{m}_Λ , \overline{M}_v and \overline{M}_V .

Assume that $\lambda_1(\Omega, A, v, V) > 0$ and call φ the unique principal eigenfunction of the operator $-\text{div}(A\nabla) + v \cdot \nabla + V$ in Ω with Dirichlet boundary condition, such that $\max_{\overline{\Omega}} \varphi = 1$. Namely, the function φ satisfies

$$\begin{cases} -\text{div}(A\nabla\varphi) + v \cdot \nabla\varphi + V\varphi = \lambda_1(\Omega, A, v, V)\varphi & \text{in } \Omega, \\ \varphi > 0 & \text{in } \Omega, \\ \|\varphi\|_{L^\infty(\Omega)} = 1, & \\ \varphi = 0 & \text{on } \partial\Omega, \end{cases} \quad (5.62)$$

and it is in $W^{2,p}(\Omega)$ and in $C^{1,\alpha}(\overline{\Omega})$ for all $1 \leq p < +\infty$ and $0 \leq \alpha < 1$.

First, remember that

$$-\overline{M}_V \leq \min_{\overline{\Omega}} V < \lambda_1(\Omega, A, v, V).$$

Then, let B be an open ball included in Ω . As observed in Step 4 of the proof of Theorem 2.1, it follows from Lemma 1.1 in [14] that there exists a constant $C > 0$ only depending on B , n , \overline{M}_A , \underline{m}_Λ , \overline{M}_v and \overline{M}_V , such that

$$\lambda_1(\Omega, A, v, V) \leq \lambda_1(B, A, v, V) \leq C.$$

From standard elliptic estimates, there exists then a constant

$$N' = N'(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V) > 0$$

such that

$$\|\varphi\|_{C^{1,1/2}(\overline{\Omega})} \leq N'.$$

We now claim the existence of a positive constant

$$\delta' = \delta'(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V) > 0$$

such that

$$\varphi(x) \geq \delta' \times d(x, \partial\Omega) \text{ for all } x \in \overline{\Omega}.$$

Assume not. Then there is a sequence of fields $(A_p, \Lambda_p, v_p, V_p)_{p \in \mathbb{N}}$ in $W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R})) \times L^{\infty}_+(\Omega) \times L^{\infty}(\Omega, \mathbb{R}^n) \times C(\overline{\Omega})$ such that $A_p \geq \Lambda_p \text{ Id}$ a.e. in Ω ,

$$\|A_p\|_{W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))} \leq \overline{M}_A, \text{ ess inf } \Lambda_p \geq \underline{m}_\Lambda, \|v_p\|_{L^{\infty}(\Omega, \mathbb{R}^n)} \leq \overline{M}_v, \|V_p\|_{L^{\infty}(\Omega)} \leq \overline{M}_V$$

for all $p \in \mathbb{N}$, and a sequence of points $(x_p)_{p \in \mathbb{N}}$ in $\overline{\Omega}$ such that

$$0 \leq \varphi_p(x_p) < \frac{d(x_p, \partial\Omega)}{p+1} \text{ for all } p \in \mathbb{N}, \quad (5.63)$$

where one calls $(\lambda_1(\Omega, A_p, v_p, V_p), \varphi_p)$ the unique solution of

$$\begin{cases} -\text{div}(A_p \nabla \varphi_p) + v_p \cdot \nabla \varphi_p + V_p \varphi_p = \lambda_1(\Omega, A_p, v_p, V_p) \varphi_p & \text{in } \Omega, \\ \varphi_p > 0 & \text{in } \Omega, \\ \|\varphi_p\|_{L^{\infty}(\Omega)} = 1, & \\ \varphi_p = 0 & \text{on } \partial\Omega \end{cases}$$

for each $p \in \mathbb{N}$. We have already noticed that the sequence $(\lambda_1(\Omega, A_p, v_p, V_p))_{p \in \mathbb{N}}$ is bounded. From standard elliptic estimates, the sequence $(\varphi_p)_{p \in \mathbb{N}}$ is also bounded in $W^{2,q}(\Omega)$ and $C^{1,\alpha}(\overline{\Omega})$ for each $1 \leq q < +\infty$ and $0 \leq \alpha < 1$. Up to a subsequence, one can assume without loss of generality that $A_p \rightarrow A_\infty$ (componentwise) uniformly in $\overline{\Omega}$, $\nabla A_p \rightarrow \nabla A_\infty$ (componentwise) in $\sigma(L^\infty, L^1)$, $v_p \rightarrow v_\infty$ (componentwise) in $\sigma(L^\infty, L^1)$, $V_p \rightarrow V_\infty$ in $\sigma(L^\infty, L^1)$, $\varphi_p \rightarrow \varphi_\infty$ weakly in $W^{2,q}(\Omega)$ for all $1 \leq q < +\infty$ and strongly in $C^{1,\alpha}(\overline{\Omega})$ for all $0 \leq \alpha < 1$, $\lambda_1(\Omega, A_p, v_p, V_p) \rightarrow \lambda$ and $x_p \rightarrow x_\infty \in \overline{\Omega}$ as $p \rightarrow +\infty$. It follows that

$$\begin{cases} -\text{div}(A_\infty \nabla \varphi_\infty) + v_\infty \cdot \nabla \varphi_\infty + V_\infty \varphi_\infty = \lambda \varphi_\infty & \text{in } \Omega, \\ \varphi_\infty \geq 0 & \text{in } \Omega, \\ \|\varphi_\infty\|_{L^{\infty}(\Omega)} = 1, & \\ \varphi_\infty = 0 & \text{on } \partial\Omega. \end{cases}$$

and $\varphi_\infty(x_\infty) = 0$. Since $A_\infty \geq \underline{m}_\Lambda \text{ Id}$ in Ω , the strong maximum principle yields that $\varphi_\infty > 0$ in Ω , whence $x_\infty \in \partial\Omega$. On the other hand, (5.63) implies that $|\nabla \varphi_\infty(x_\infty)| = 0$, which is impossible from Hopf lemma. One has then reached a contradiction.

Therefore, coming back to the fields (A, Λ, v, V) and to the function φ solving (5.62), we get the existence of $\delta' = \delta'(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V) > 0$ such that

$$\varphi(x) \geq \delta' \times d(x, \partial\Omega) \text{ for all } x \in \overline{\Omega}.$$

In other words, the function φ is in the set $E_{1/2, N', \delta'}(\Omega)$.

Call

$$N = 3N'(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V) > 0$$

and

$$\delta = \frac{\delta'(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V)}{3} > 0.$$

Then, call $\eta > 0$ the positive constant which is given in Corollary 4.5, with the choice $\alpha = 1/2$. It only depends on Ω , N and δ , and therefore,

$$\eta = \eta(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V).$$

Call now

$$\theta = \theta(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V) = \frac{\eta(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \overline{M}_v, \overline{M}_V)}{2} > 0.$$

Then, choose any $\varepsilon' \in (0, 1)$ such that

$$\frac{\frac{\lambda_1(\Omega, A, v, V) + 2\varepsilon'}{1 + \eta} + \varepsilon' + \varepsilon' m_\Lambda^{-1} + \varepsilon' \|V^-\|_\infty}{1 - \varepsilon'} \leq \frac{\lambda_1(\Omega, A, v, V)}{1 + \theta}. \quad (5.64)$$

It is indeed possible to choose such a ε' since $\lambda_1(\Omega, A, v, V) > 0$ and $0 < \theta < \eta$.

Under the notation of Step 4 of the proof of Theorem 2.1, there exists a sequence of C^∞ fields $(A_k, \Lambda_k, v_k)_{k \in \mathbb{N}}$ satisfying (5.32) and (5.33), and such that the solutions φ_k of (5.36) converge to φ as $k \rightarrow +\infty$ weakly in $W^{2,p}(\Omega)$ and strongly in $C^{1,\alpha}(\overline{\Omega})$ for all $1 \leq p < +\infty$ and $0 \leq \alpha < 1$. Furthermore, it has been proved that $\lambda_1(\Omega, A_k, v_k, V) \rightarrow \lambda_1(\Omega, A, v, V)$ as $k \rightarrow +\infty$. Then there exists $K \in \mathbb{N}$ such that (5.38) holds and

$$\varphi_K \in E_{1/2, 2N', \delta/2}(\Omega).$$

Under the same notation as in Step 5 of the proof of Theorem 2.1, the functions ψ_l converge to φ_K as $l \rightarrow +\infty$ in $W^{2,p}(\Omega)$ and $C^{1,\alpha}(\overline{\Omega})$ for all $1 \leq p < +\infty$ and $0 \leq \alpha < 1$. Therefore, as in Step 5, there exists $L \in \mathbb{N}$ such that (5.44) holds and

$$\psi_L \in E_{1/2, 3N', \delta/3}(\Omega) = E_{1/2, N, \delta}(\Omega).$$

Therefore, all assumptions of Corollary 4.5 are satisfied with

$$(A_\Omega, \Lambda_\Omega, \omega, V, \psi, \omega_0, \mu) = (A_K, \Lambda_K, |v_K|, V, \psi_L, \varepsilon', \lambda_1(\Omega, A, v, V) + 2\varepsilon').$$

Notice especially that ω_0 and μ are nonnegative. With the same notations as in Sections 3 and 4, it then follows from Corollary 4.5 that

$$\int_{\Omega^*} \left[\widehat{\Lambda}(x) |\nabla \tilde{\psi}(x)|^2 - \varepsilon' |\nabla \tilde{\psi}(x)| \tilde{\psi}(x) + \widehat{V}(x) \tilde{\psi}(x)^2 \right] e^{-U(x)} dx \leq \frac{\mu}{1 + \eta} \int_{\Omega^*} \tilde{\psi}(x)^2 e^{-U(x)} dx.$$

The same calculations as in Steps 6, 7 and 8 of the proof of Theorem 2.1 can be carried out, where $\mu = \lambda_1(\Omega, A, v, V) + 2\varepsilon'$ is replaced by $(\lambda_1(\Omega, A, v, V) + 2\varepsilon')/(1 + \eta)$ in (5.45). One then gets the existence of three radially symmetric $C^\infty(\overline{\Omega}^*)$ fields $\Lambda^* > 0$, $\omega^* \geq 0$, $\overline{V}^* \leq 0$ and a nonpositive radially symmetric $L^\infty(\Omega^*)$ field V^* , which satisfy (2.2) and are such that $\mu_{V^*} = \mu_{-V^-}$, $V^* \leq \overline{V}^*$ and

$$\begin{aligned} \lambda_1(\Omega^*, \Lambda^* \text{Id}, v^*, V^*) &\leq \lambda_1(\Omega^*, \Lambda^* \text{Id}, v^*, \overline{V}^*) \\ &\leq \frac{\frac{\lambda_1(\Omega, A, v, V) + 2\varepsilon'}{1 + \eta} + \varepsilon' + \varepsilon' m_\Lambda^{-1} + \varepsilon' \|V^-\|_\infty}{1 - \varepsilon'}, \end{aligned}$$

where $v^* = \omega^* e_r$ in Ω^* . As a consequence, there holds

$$\lambda_1(\Omega^*, \Lambda^* \text{Id}, v^*, V^*) \leq \lambda_1(\Omega^*, \Lambda^* \text{Id}, v^*, \bar{V}^*) \leq \frac{\lambda_1(\Omega, A, v, V)}{1 + \theta}$$

from the choice of ε' in (5.64). The proof of Theorem 2.3 is now complete. \square

5.3 Constraints on the eigenvalues of the matrix field A

We now give the proof of Theorem 2.4. The following elementary lemma will be needed:

Lemma 5.7 *Let $n \geq 2$, $p \in \{1, \dots, n-1\}$, $\omega > 0$, $\sigma > 0$ and $A \in \mathcal{S}_n(\mathbb{R})$ be positive definite such that $\det(A) \geq \omega$ and $\sigma_p(A) \leq \sigma$. Then, there exist two positive numbers a_1, a_2 which only depend on n, p, ω and σ such that $\det(D) = \omega$, $\sigma_p(D) = \sigma$, $A \geq a_1 \text{Id}$ and $D \geq a_1 \text{Id}$, where D is the diagonal matrix $D = \text{diag}(a_1, a_2, \dots, a_2) \in \mathcal{S}_n(\mathbb{R})$.*

Proof. Notice first that, as already underlined in the introduction, the assumptions of the lemma imply that $C_n^p \omega^{p/n} \leq \sigma$. Let $f(s)$ be defined for all $s > 0$ as

$$f(s) = \omega C_{n-1}^{p-1} s^{p-n} + C_{n-1}^p s^p.$$

The function s is continuous and strictly convex in $(0, +\infty)$. Furthermore, $f(0^+) = f(+\infty) = +\infty$ and elementary calculations give:

$$\min_{s \in (0, +\infty)} f(s) = f(\omega^{1/n}) = C_n^p \omega^{p/n} \leq \sigma.$$

Call

$$a_2 = \max \{s > 0, f(s) \leq \sigma\}, \quad a_1 = \frac{\omega}{(a_2)^{n-1}}, \quad D = \text{diag}(a_1, a_2, \dots, a_2) \in \mathcal{S}_n(\mathbb{R}).$$

The real numbers a_1 and a_2 are well-defined and positive. They only depend on n, p, ω and σ . Furthermore, $a_1 \leq \omega^{1/n} \leq a_2$, $D \in \mathcal{S}_n(\mathbb{R})$ and

$$\det(D) = \omega, \quad \sigma_p(D) = C_{n-1}^{p-1} a_1 (a_2)^{p-1} + C_{n-1}^p (a_2)^p = f(a_2) = \sigma.$$

Denote $0 < \lambda_1 \leq \dots \leq \lambda_n$ the eigenvalues of A . Call

$$0 < \tilde{a}_2 = \left(\prod_{2 \leq i \leq n} \lambda_i \right)^{1/(n-1)}$$

and

$$0 < \tilde{a}_1 = \frac{\omega}{(\tilde{a}_2)^{n-1}} \leq \frac{\det(A)}{\prod_{2 \leq i \leq n} \lambda_i} = \lambda_1. \quad (5.65)$$

Since $\omega = \tilde{a}_1(\tilde{a}_2)^{n-1} \leq \lambda_1(\tilde{a}_2)^{n-1}$, one has

$$f(\tilde{a}_2) = \omega C_{n-1}^{p-1}(\tilde{a}_2)^{p-n} + C_{n-1}^p(\tilde{a}_2)^p \leq C_{n-1}^{p-1}\lambda_1 \times \left(\prod_{2 \leq i \leq n} \lambda_i \right)^{\frac{p-1}{n-1}} + C_{n-1}^p \times \left(\prod_{2 \leq i \leq n} \lambda_i \right)^{\frac{p}{n-1}}.$$

For $q = p - 1$ or $q = p$, call

$$\mathcal{J}_q = \{J \subset \{2, \dots, n\}, \text{card}(J) = q\}$$

and, for all $I \subset \{1, \dots, n\}$,

$$\pi_I = \prod_{i \in I} \lambda_i.$$

Observe that $\text{card}(\mathcal{J}_q) = C_{n-1}^q$ and that

$$\prod_{2 \leq i \leq n} \lambda_i = \left(\prod_{J \in \mathcal{J}_{p-1}} \pi_J \right)^{1/C_{n-2}^{p-2}} = \left(\prod_{J \in \mathcal{J}_p} \pi_J \right)^{1/C_{n-2}^{p-1}}.$$

Thus,

$$\begin{aligned} f(\tilde{a}_2) &\leq C_{n-1}^{p-1}\lambda_1 \times \left(\prod_{J \in \mathcal{J}_{p-1}} \pi_J \right)^{\frac{p-1}{(n-1)C_{n-2}^{p-2}}} + C_{n-1}^p \times \left(\prod_{J \in \mathcal{J}_p} \pi_J \right)^{\frac{p}{(n-1)C_{n-2}^{p-1}}} \\ &= C_{n-1}^{p-1}\lambda_1 \times \left(\prod_{J \in \mathcal{J}_{p-1}} \pi_J \right)^{1/C_{n-1}^{p-1}} + C_{n-1}^p \times \left(\prod_{J \in \mathcal{J}_p} \pi_J \right)^{1/C_{n-1}^p} \\ &\leq \lambda_1 \times \sum_{J \in \mathcal{J}_{p-1}} \pi_J + \sum_{J \in \mathcal{J}_p} \pi_J \\ &= \sum_{I \subset \{1, \dots, n\}, \text{card}(I)=p} \pi_I \\ &= \sigma_p(A) \\ &\leq \sigma. \end{aligned}$$

Hence, $\tilde{a}_2 \leq a_2$ and

$$\tilde{a}_1 = \frac{\omega}{(\tilde{a}_2)^{n-1}} \geq \frac{\omega}{(a_2)^{n-1}} = a_1.$$

Together with (5.65), it follows that $\lambda_1 \geq a_1$, whence $A \geq a_1 \text{Id}$. \square

Proof of Theorem 2.4. Let $\Omega, A, v, V, p, \omega, \sigma$ be as in Theorem 2.4, and let $a_1 > 0$ and $a_2 > 0$ be given by Lemma 5.7. Therefore,

$$A(x) \geq a_1 \text{Id for all } x \in \overline{\Omega}.$$

For all $x \in \overline{\Omega^*} \setminus \{0\}$, define now $A^*(x)$ as in Theorem 2.4. Thus, $A^*(x) \geq a_1 \text{Id}$ and there is an invertible matrix $P(x)$ of size $n \times n$ such that

$$A^*(x) = P(x)DP(x)^{-1},$$

where $D = \text{diag}(a_1, a_2, \dots, a_n) \in \mathcal{S}_n(\mathbb{R})$. Hence,

$$\det(A^*(x)) = \det(D) = \omega \text{ and } \sigma_p(A^*(x)) = \sigma_p(D) = \sigma \text{ for all } x \in \overline{\Omega^*} \setminus \{0\}.$$

Let $\varepsilon > 0$ be given. From Theorem 2.1 applied with Λ equal to the positive constant a_1 , it follows that there exist two radially symmetric $C^\infty(\overline{\Omega^*})$ fields $\omega^* \geq 0$ and $\overline{V}^* \leq 0$, and a nonpositive radially symmetric $L^\infty(\Omega^*)$ field V^* , such that, for $v^* = \omega^* e_r$ in $\overline{\Omega^*} \setminus \{0\}$,

$$\begin{cases} \|v^*\|_{L^\infty(\Omega^*, \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}, & \|v^*\|_{L^2(\Omega^*)} = \|v\|_{L^2(\Omega)}, \\ \mu_{|\overline{V}^*|} \leq \mu_{V^-}, & \mu_{V^*} = \mu_{V^-}, \quad V^* \leq \overline{V}^* \end{cases}$$

and

$$\lambda_1(\Omega^*, a_1 \text{Id}, v^*, V^*) \leq \lambda_1(\Omega^*, a_1 \text{Id}, v^*, \overline{V}^*) \leq \lambda_1(\Omega, A, v, V) + \varepsilon.$$

From Remark 2.5, one concludes that

$$\lambda_1(\Omega^*, A^*, v^*, V^*) \leq \lambda_1(\Omega^*, A^*, v^*, \overline{V}^*) \leq \lambda_1(\Omega, A, v, V) + \varepsilon.$$

Lastly, since $A \geq a_1 \text{Id}$ in Ω and $A^* \nabla \phi = a_1 \nabla \phi$ for any radially symmetric function ϕ in Ω^* , the existence of two radially symmetric bounded functions $\omega_0^* \geq 0$ and $V_0^* \leq 0$ in Ω^* satisfying (2.4) and

$$\lambda_1(\Omega^*, A^*, v_0^*, V_0^*) \leq \lambda_1(\Omega, A, v, V),$$

where $v_0^* = \omega_0^* e_r$ in Ω^* , can be done as in Step 10 of the proof of Theorem 2.1

This completes the proof of Theorem 2.4. \square

Remark 5.8 Consider now the case when Ω is a general open subset of \mathbb{R}^n with finite measure, and let (A, v, V) be as in Theorem 2.4, with the extra assumption $V \in L^\infty(\Omega)$. Since $A \geq a_1 \text{Id}$ in Ω , it follows, as in Remark 5.6, that there exist two radially symmetric bounded functions $\omega_0^* \geq 0$ and $V_0^* \leq 0$ in Ω^* satisfying (2.4) (with $-\text{ess sup}_\Omega V^-$ instead of $-\max_{\overline{\Omega}} V^-$) and $\lambda_1(\Omega^*, A^*, v_0^*, V_0^*) \leq \lambda_1(\Omega, A, v, V)$, where $v_0^* = \omega_0^* e_r$ in Ω^* and A^* is the same as in Theorem 2.4.

6 The cases of L^p constraints

In this section, we focus on optimization and comparison results for $\lambda_1(\Omega, A, v, V)$ when Ω has fixed Lebesgue measure, A satisfies the same constraints as in Theorem 2.1 and v and V satisfy L^p constraints. We first give in Section 6.1 some optimization results when Ω is fixed. Then, relying on Theorem 2.1, we derive in Section 6.2 some Faber-Krahn type inequalities.

6.1 Optimization in fixed domains

In Theorem 2.1, we were able to compare the eigenvalue $\lambda_1(\Omega, A, v, V)$ to $\lambda_1(\Omega^*, \Lambda^* \text{Id}, v^*, V^*)$ in Ω^* , where the L^1 norms of Λ^{-1} (with $A \geq \Lambda \text{Id}$) and $|v|^2 \Lambda^{-1}$ and the distribution function of V^- were the same as for the new fields in Ω^* (in particular, all L^p norms of V^- were preserved). In this section, we fix the domain Ω and we minimize and maximize $\lambda_1(\Omega, A, v, V)$ when $A \geq \Lambda \text{Id}$ and v and V satisfy some $L^p \cap L^\infty$ constraints with some given weights. Furthermore, we prove the uniqueness of the optimal fields when A is fixed and v and V satisfy given L^∞ constraints.

6.1.1 The case of L^p constraints, $1 < p \leq +\infty$

Given $\Omega \in \mathcal{C}$, $M \geq 0$, $\Lambda \in L_+^\infty(\Omega)$, $1 < p, q \leq +\infty$, $w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty} \in L_+^\infty(\Omega)$, $\tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty} \geq 0$ such that $M \geq \text{ess sup}_\Omega \Lambda$, define

$$\begin{aligned} \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}} \\ = \left\{ (A, v, V) \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R})) \times L^\infty(\Omega, \mathbb{R}^n) \times L^\infty(\Omega); \right. \\ \left. \begin{aligned} & \|A\|_{W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))} \leq M, \quad A \geq \Lambda \text{Id a.e. in } \Omega, \\ & \|w_{1,p} v\|_p \leq \tau_{1,p}, \quad \|w_{1,\infty} v\|_\infty \leq \tau_{1,\infty}, \quad \|w_{2,q} V\|_q \leq \tau_{2,q}, \quad \|w_{2,\infty} V\|_\infty \leq \tau_{2,\infty} \end{aligned} \right\}, \end{aligned}$$

and

$$\left\{ \begin{aligned} \underline{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) \\ &= \inf_{(A, v, V) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}} \lambda_1(\Omega, A, v, V), \\ \bar{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) \\ &= \sup_{(A, v, V) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}} \lambda_1(\Omega, A, v, V). \end{aligned} \right. \quad (6.1)$$

Observe that $\mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}} \neq \emptyset$, that

$$\underline{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) \geq -\frac{\tau_{2,\infty}}{\text{ess inf}_\Omega w_{2,\infty}}$$

and that

$$\bar{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) < +\infty$$

by [14].

Our first result deals with the optimization problem for (6.1):

Theorem 6.1 *Let $\Omega \in \mathcal{C}$, $M \geq 0$, $\Lambda \in L_+^\infty(\Omega)$, $1 < p, q \leq +\infty$, $w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty} \in L_+^\infty(\Omega)$ and $\tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty} \geq 0$ be given such that $M \geq \text{ess sup}_\Omega \Lambda$. Then,*

(1) *there exists $(\underline{A}, \underline{v}, \underline{V}) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}$ such that, if $\underline{\varphi} = \varphi_{\Omega, \underline{A}, \underline{v}, \underline{V}}$:*

(a) $\underline{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) = \lambda_1(\Omega, \underline{A}, \underline{v}, \underline{V})$,

(b) $\underline{v} \cdot \nabla \underline{\varphi} = -|\underline{v}| \times |\nabla \underline{\varphi}|$ a.e. in Ω ,

(c) $\underline{V}(x) \leq 0$ a.e. in Ω ,

(d) $\|w_{1,p}\underline{v}\|_p = \tau_{1,p}$ or $\|w_{1,\infty}\underline{v}\|_\infty = \tau_{1,\infty}$, $\|w_{2,q}\underline{V}\|_q = \tau_{2,q}$ or $\|w_{2,\infty}\underline{V}\|_\infty = \tau_{2,\infty}$.

Moreover, if $(A, v, V) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}$ is such that

$$\underline{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) = \lambda_1(\Omega, A, v, V)$$

and if $\varphi = \varphi_{\Omega, A, v, V}$, then properties (b), (c) and (d) hold with φ, v, V instead of $\underline{\varphi}, \underline{v}$ and \underline{V} ,

(2) there exists $(\overline{A}, \overline{v}, \overline{V}) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}$ such that, if $\overline{\varphi} = \varphi_{\Omega, \overline{A}, \overline{v}, \overline{V}}$:

(a) $\overline{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) = \lambda_1(\Omega, \overline{A}, \overline{v}, \overline{V})$,

(b) $\overline{v} \cdot \nabla \overline{\varphi} = |\overline{v}| \times |\nabla \overline{\varphi}|$ a.e. in Ω ,

(c) $\overline{V}(x) \geq 0$ a.e. in Ω ,

(d) $\|w_{1,p}\overline{v}\|_p = \tau_{1,p}$ or $\|w_{1,\infty}\overline{v}\|_\infty = \tau_{1,\infty}$, $\|w_{2,q}\overline{V}\|_q = \tau_{2,q}$ or $\|w_{2,\infty}\overline{V}\|_\infty = \tau_{2,\infty}$.

Moreover, if $(A, v, V) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}$ is such that

$$\overline{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) = \lambda_1(\Omega, A, v, V)$$

and if $\varphi = \varphi_{\Omega, A, v, V}$, then properties (b), (c) and (d) hold with φ, v, V instead of $\overline{\varphi}, \overline{v}$ and \overline{V} .

We will use several times in the proof the following comparison result:

Lemma 6.2 *Let $\Omega \in \mathcal{C}$, $\mu \in \mathbb{R}$, $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$ with $A \geq \gamma \text{Id}$ a.e. in Ω for some $\gamma > 0$, $v \in L^\infty(\Omega, \mathbb{R}^n)$ and $V \in L^\infty(\Omega)$. Assume that φ and ψ are functions in $W^{2,r}(\Omega)$ for all $1 \leq r < +\infty$, satisfying $\|\varphi\|_\infty = \|\psi\|_\infty$ and $\varphi = \psi = 0$ on $\partial\Omega$. Assume also that $\varphi \geq 0$ in Ω , $\psi > 0$ in Ω and*

$$\begin{cases} -\text{div}(A\nabla\psi) + v \cdot \nabla\psi + V\psi \geq \mu\psi \text{ a.e. in } \Omega, \\ -\text{div}(A\nabla\varphi) + v \cdot \nabla\varphi + V\varphi \leq \mu\varphi \text{ a.e. in } \Omega. \end{cases}$$

Then $\varphi = \psi$ in Ω .

Proof. The proof uses a classical comparison method (this method was used for instance in [14]). We give it here for the sake of completeness. Since $\psi > 0$ in Ω and $\psi = 0$ on $\partial\Omega$, the Hopf lemma yields $\nu(x) \cdot \nabla\psi(x) < 0$ on $\partial\Omega$, where, for all $x \in \partial\Omega$, $\nu(x)$ denotes the outward normal unit vector at x . Since $\varphi \in C^{1,\beta}(\overline{\Omega})$ for all $0 \leq \beta < 1$, $\varphi \geq 0$ in Ω and $\varphi = 0$ on $\partial\Omega$, it follows that there exists $\gamma > 0$ such that $\gamma\psi > \varphi$ in Ω . Define

$$\gamma^* = \inf \{ \gamma > 0, \gamma\psi > \varphi \text{ in } \Omega \}.$$

One clearly has $\gamma^*\psi \geq \varphi$ in Ω , so that $\gamma^* > 0$. Define $w = \gamma^*\psi - \varphi \geq 0$ and assume that $w > 0$ everywhere in Ω . Since

$$-\text{div}(A\nabla w) + v \cdot \nabla w + Vw - \mu w \geq 0 \text{ a.e. in } \Omega \tag{6.2}$$

and $w = 0$ on $\partial\Omega$, the Hopf maximum principle implies that $\nu \cdot \nabla w < 0$ on $\partial\Omega$. As above, this yields the existence of $\kappa > 0$ such that $w > \kappa\varphi$ in Ω , whence $\gamma^*\psi/(1 + \kappa) > \varphi$ in Ω . This is a contradiction with the minimality of γ^* . Thus, there exists $x_0 \in \Omega$ such that $w(x_0) = 0$ (i.e. $\gamma^*\psi(x_0) = \varphi(x_0)$). Since $w \geq 0$ in Ω , it follows from (6.2) and from the strong maximum principle, that $w = 0$ in Ω , which means that φ and ψ are proportional. Since they are non-negative in Ω and have the same L^∞ norm in Ω , one has $\varphi = \psi$, which ends the proof of Lemma 6.2. \square

For the proof of Theorem 6.1, we will treat the minimization problem only, the maximization problem being clearly analogous. It is plain to see that the result is a consequence of the two following lemmata:

Lemma 6.3 *Let $\Omega \in \mathcal{C}$, $M \geq 0$, $\Lambda \in L_+^\infty(\Omega)$, $1 < p, q \leq +\infty$, $w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty} \in L_+^\infty(\Omega)$ and $\tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty} \geq 0$ be given such that $M \geq \text{ess sup}_\Omega \Lambda$. Assume that $(\underline{A}, \underline{v}, \underline{V}) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}$ is such that*

$$\underline{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) = \lambda_1(\Omega, \underline{A}, \underline{v}, \underline{V}), \quad (6.3)$$

and let $\underline{\varphi} = \varphi_{\Omega, \underline{A}, \underline{v}, \underline{V}}$. Then, properties (b), (c) and (d) in Theorem 6.1 hold.

Lemma 6.4 *Let $\Omega \in \mathcal{C}$, $M \geq 0$, $\Lambda \in L_+^\infty(\Omega)$, $1 < p, q \leq +\infty$, $w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty} \in L_+^\infty(\Omega)$ and $\tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty} \geq 0$ be given such that $M \geq \text{ess sup}_\Omega \Lambda$. Then, there exists $(\underline{A}, \underline{v}, \underline{V}) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}$ such that (6.3) holds.*

Proof of Lemma 6.3. Let $\underline{A}, \underline{v}, \underline{V}$ and $\underline{\varphi}$ as in Lemma 6.3. Remember that $\underline{\varphi} \in C^1(\overline{\Omega})$. Define, for a.e. $x \in \Omega$,

$$w(x) = \begin{cases} -|\underline{v}(x)| \frac{\nabla \underline{\varphi}(x)}{|\nabla \underline{\varphi}(x)|} & \text{if } \nabla \underline{\varphi}(x) \neq 0, \\ 0 & \text{if } \nabla \underline{\varphi}(x) = 0 \end{cases}$$

so that $w \cdot \nabla \underline{\varphi} = -|\underline{v}| \times |\nabla \underline{\varphi}| \leq \underline{v} \cdot \nabla \underline{\varphi}$ a.e. in Ω , and set

$$\mu = \lambda_1(\Omega, \underline{A}, w, \underline{V}) \text{ and } \psi = \varphi_{\Omega, \underline{A}, w, \underline{V}}.$$

Notice that $|w| \leq |\underline{v}|$ a.e. in Ω , whence $(\underline{A}, w, \underline{V}) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}$ and

$$\underline{\lambda} := \underline{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}) \leq \mu.$$

Thus, one has

$$\begin{cases} -\text{div}(\underline{A}\nabla \underline{\varphi}) + w \cdot \nabla \underline{\varphi} + \underline{V}\underline{\varphi} \leq -\text{div}(\underline{A}\nabla \underline{\varphi}) + \underline{v} \cdot \nabla \underline{\varphi} + \underline{V}\underline{\varphi} = \underline{\lambda}\underline{\varphi} \leq \mu\underline{\varphi} & \text{a.e. in } \Omega, \\ -\text{div}(\underline{A}\nabla \psi) + w \cdot \nabla \psi + \underline{V}\psi = \mu\psi & \text{a.e. in } \Omega, \end{cases}$$

and Lemma 6.2 yields $\psi = \underline{\varphi}$ and therefore $\mu = \underline{\lambda}$ and $\underline{v} \cdot \nabla \underline{\varphi} = w \cdot \nabla \underline{\varphi} = -|\underline{v}| \times |\nabla \underline{\varphi}|$ a.e. in Ω .

As far as assertion (c) is concerned, define, for a.e. $x \in \Omega$,

$$W(x) = \begin{cases} \underline{V}(x) & \text{if } \underline{V}(x) \leq 0, \\ 0 & \text{if } \underline{V}(x) > 0. \end{cases}$$

Observe that $|W| \leq |\underline{V}|$ a.e. in Ω , whence $(\underline{A}, \underline{v}, W) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}}$. If $\mu = \lambda_1(\Omega, \underline{A}, \underline{v}, W)$ and $\psi = \varphi_{\Omega, \underline{A}, \underline{v}, W}$, one therefore has $\underline{\lambda} \leq \mu$ and

$$-\operatorname{div}(\underline{A}\nabla\psi) + \underline{v} \cdot \nabla\psi + W\psi = \mu\psi \text{ a.e. in } \Omega,$$

while, since $W\underline{\varphi} \leq \underline{V}\underline{\varphi}$ a.e. in Ω ,

$$-\operatorname{div}(\underline{A}\nabla\underline{\varphi}) + \underline{v} \cdot \nabla\underline{\varphi} + W\underline{\varphi} \leq \mu\underline{\varphi} \text{ a.e. in } \Omega.$$

Lemma 6.2 therefore shows that $\psi = \underline{\varphi}$, and it follows that $W\underline{\varphi} = \underline{V}\underline{\varphi}$ a.e. in Ω , which implies $W = \underline{V}$ a.e. in Ω since $\underline{\varphi} > 0$ in Ω , and this is assertion (c).

Assume now that $\|w_{1,p}\underline{v}\|_p < \tau_{1,p}$ and $\|w_{1,\infty}\underline{v}\|_\infty < \tau_{1,\infty}$, and define, for a.e. $x \in \Omega$,

$$w(x) = \begin{cases} -(|\underline{v}(x)| + \varepsilon) \frac{\nabla\underline{\varphi}(x)}{|\nabla\underline{\varphi}(x)|} & \text{if } \nabla\underline{\varphi}(x) \neq 0, \\ 0 & \text{if } \nabla\underline{\varphi}(x) = 0, \end{cases}$$

where $\varepsilon > 0$ is chosen in such a way that $\|w_{1,p}w\|_p \leq \tau_{1,p}$ and $\|w_{1,\infty}w\|_\infty \leq \tau_{1,\infty}$, so that, if $\mu = \lambda_1(\Omega, \underline{A}, w, \underline{V})$, one has $\underline{\lambda} \leq \mu$. Let $\psi = \varphi_{\Omega, \underline{A}, w, \underline{V}}$. Observe that

$$-\operatorname{div}(\underline{A}\nabla\psi) + w \cdot \nabla\psi + \underline{V}\psi \leq \underline{\lambda}\psi \leq \mu\underline{\varphi} \text{ a.e. in } \Omega,$$

since $w \cdot \nabla\underline{\varphi} = -(|\underline{v}| + \varepsilon) |\nabla\underline{\varphi}| \leq -|\underline{v}| \times |\nabla\underline{\varphi}| \leq \underline{v} \cdot \nabla\underline{\varphi}$ a.e. in Ω , while

$$-\operatorname{div}(\underline{A}\nabla\psi) + w \cdot \nabla\psi + \underline{V}\psi = \mu\psi \text{ a.e. in } \Omega.$$

Another application of Lemma 6.2 yields that $\psi = \underline{\varphi}$ and therefore $w \cdot \nabla\underline{\varphi} = \underline{v} \cdot \nabla\underline{\varphi}$, so that $-\varepsilon |\nabla\underline{\varphi}| = 0$ a.e. in Ω , which is impossible. One argues similarly for \underline{V} , using the fact that $\underline{V} \leq 0$ a.e. in Ω . \square

Proof of Lemma 6.4. Write

$$\underline{\lambda} = \underline{\lambda}(\Omega, M, \Lambda, p, q, w_{1,p}, w_{1,\infty}, w_{2,q}, w_{2,\infty}, \tau_{1,p}, \tau_{1,\infty}, \tau_{2,q}, \tau_{2,\infty}).$$

There exist a sequence $(A_k)_{k \in \mathbb{N}} \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$ with $\|A_k\|_{W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))} \leq M$ and $A_k \geq \Lambda \operatorname{Id}$ a.e. in Ω , a sequence $(v_k)_{k \in \mathbb{N}} \in L^\infty(\Omega, \mathbb{R}^n)$ with $\|w_{1,p}v_k\|_p \leq \tau_{1,p}$ and $\|w_{1,\infty}v_k\|_\infty \leq \tau_{1,\infty}$, and a sequence $(V_k)_{k \in \mathbb{N}} \in L^\infty(\Omega)$ with $\|w_{2,q}V_k\|_q \leq \tau_{2,q}$ and $\|w_{2,\infty}V_k\|_\infty \leq \tau_{2,\infty}$, such that

$$\lambda_k := \lambda_1(\Omega, A_k, v_k, V_k) \rightarrow \underline{\lambda} \text{ as } k \rightarrow +\infty.$$

For each $k \in \mathbb{N}$, call $\varphi_k = \varphi_{\Omega, A_k, v_k, V_k}$, so that

$$-\operatorname{div}(A_k \nabla \varphi_k) + v_k \cdot \nabla \varphi_k + V_k \varphi_k = \lambda_k \varphi_k \text{ in } \Omega \text{ and } \varphi_k = 0 \text{ on } \partial\Omega.$$

By the usual elliptic estimates, the φ_k 's are uniformly bounded in $W^{2,r}(\Omega)$ for all $1 \leq r < +\infty$, and therefore in $C^{1,\alpha}(\overline{\Omega})$ for all $0 \leq \alpha < 1$. Therefore, up to a subsequence, one may assume that, for some nonnegative function $\underline{\varphi} \in W^{2,r}(\Omega)$ for all $1 \leq r < +\infty$, $\varphi_k \rightharpoonup \underline{\varphi}$ weakly in $W^{2,r}(\Omega)$ for all $1 \leq r < +\infty$ and $\varphi_k \rightarrow \underline{\varphi}$ strongly in $C^{1,\alpha}(\overline{\Omega})$ for all $0 \leq \alpha < 1$, as $k \rightarrow +\infty$. Similarly, there exists $\underline{A} \in W^{1,\infty}(\Omega, \overline{\mathcal{S}}_n(\mathbb{R}))$ such that (up to extraction), $A_k \rightarrow \underline{A}$ uniformly in $\overline{\Omega}$ and, for each $1 \leq j \leq n$, $\partial_j A_k \rightarrow \partial_j \underline{A}$ in $\sigma(L^\infty(\Omega), L^1(\Omega))$ componentwise. In particular, $\|\underline{A}\|_{W^{1,\infty}(\Omega, \overline{\mathcal{S}}_n(\mathbb{R}))} \leq M$ and $\underline{A} \geq \Lambda \text{ Id}$ a.e. in Ω . Finally, up to extraction again, there exists $\omega \in L^\infty(\Omega)$ such that $|v_k| \rightarrow \omega \geq 0$ in $\sigma(L^r(\Omega), L^{r'}(\Omega))$ for all $1 < r \leq +\infty$ (where $1/r + 1/r' = 1$) and there exists $\underline{V} \in L^\infty(\Omega)$ such that $V_k \rightarrow \underline{V}$ in $\sigma(L^r(\Omega), L^{r'}(\Omega))$ for all $1 < r \leq +\infty$. Since, for all $k \geq 1$, by Cauchy-Schwarz,

$$-\text{div}(A_k \nabla \varphi_k) - |v_k| |\nabla \varphi_k| + V_k \varphi_k \leq \lambda_k \varphi_k \text{ a.e. in } \Omega,$$

one has

$$-\text{div}(\underline{A} \nabla \underline{\varphi}) - \omega |\nabla \underline{\varphi}| + \underline{V} \underline{\varphi} \leq \underline{\lambda} \underline{\varphi} \text{ a.e. in } \Omega.$$

Define now, for a.e. $x \in \Omega$,

$$\underline{v}(x) = \begin{cases} -\omega(x) \frac{\nabla \underline{\varphi}(x)}{|\nabla \underline{\varphi}(x)|} & \text{if } \nabla \underline{\varphi}(x) \neq 0, \\ 0 & \text{if } \nabla \underline{\varphi}(x) = 0, \end{cases}$$

so that $\underline{v} \cdot \nabla \underline{\varphi} = -\omega |\nabla \underline{\varphi}|$ a.e. in Ω . One therefore has

$$-\text{div}(\underline{A} \nabla \underline{\varphi}) + \underline{v} \cdot \nabla \underline{\varphi} + \underline{V} \underline{\varphi} \leq \underline{\lambda} \underline{\varphi} \text{ a.e. in } \Omega \text{ and } \underline{\varphi} = 0 \text{ on } \partial\Omega. \quad (6.4)$$

Observe that $(\underline{A}, \underline{v}, \underline{V}) \in \mathcal{A}_{\Omega, M, \Lambda, p, q, w_1, p, w_1, \infty, w_2, q, w_2, \infty, \tau_1, p, \tau_1, \infty, \tau_2, q, \tau_2, \infty}$. Define now $\mu = \lambda_1(\Omega, \underline{A}, \underline{v}, \underline{V})$ and $\psi = \varphi_{\Omega, \underline{A}, \underline{v}, \underline{V}}$, so that $\underline{\lambda} \leq \mu$. It follows from (6.4) that

$$\begin{cases} -\text{div}(\underline{A} \nabla \psi) + \underline{v} \cdot \nabla \psi + \underline{V} \psi = \mu \psi & \text{a.e. in } \Omega, \\ -\text{div}(\underline{A} \nabla \underline{\varphi}) + \underline{v} \cdot \nabla \underline{\varphi} + \underline{V} \underline{\varphi} \leq \mu \underline{\varphi} & \text{in } \Omega. \end{cases}$$

Moreover, $\underline{\varphi} \geq 0$ in Ω , $\psi > 0$ in Ω , $\underline{\varphi} = \psi = 0$ on $\partial\Omega$ and $\|\underline{\varphi}\|_\infty = \|\psi\|_\infty = 1$. Lemma 6.2 therefore yields $\underline{\varphi} = \psi$, hence $\underline{\lambda} = \mu$. This ends the proof of Lemma 6.4. \square

Remark 6.5 What happens in Theorem 6.1 if one drops the L^∞ bounds for v or V ? Even if one still assumes that v and V are qualitatively in L^∞ (so that the principal eigenvalue of $-\text{div}(A \nabla) + v \cdot \nabla + V$ is well-defined by [14]), it turns out that the infimum or the supremum considered there may not be achieved. For instance, fix $1 < p < n$, $\Omega \in \mathcal{C}$, $\tau > 0$, $A = \text{Id}$ and $V = 0$ in $\overline{\Omega}$, and define

$$\underline{\lambda}(\Omega, \tau) = \inf_{v \in L^\infty(\Omega, \mathbb{R}^n), \|v\|_p \leq \tau} \lambda_1(\Omega, \text{Id}, v, 0).$$

Since the operator $-\Delta + v \cdot \nabla$ satisfies the maximum principle in Ω , its principal eigenvalue is positive, for each $v \in L^\infty(\Omega, \mathbb{R}^n)$, and therefore $\underline{\lambda}(\Omega, \tau) \geq 0$. We claim that $\underline{\lambda}(\Omega, \tau) = 0$. Indeed, fix $\rho_0 > 0$ such that there exists a ball B_0 with radius ρ_0 included in Ω . Call x_0 its

center. For all $A > 0$ large enough, define $\rho_A \in (0, \rho_0)$ such that $A(\alpha_n \rho_A^n)^{1/p} = \tau$ (recall that α_n is the Lebesgue measure of the Euclidean unit ball in \mathbb{R}^n), let $B_A = B(x_0, \rho_A)$ be the ball with the same center x_0 as B_0 and with radius ρ_A , and set $v = A \times \mathbf{1}_{B_A} e_r(\cdot - x_0)$ in Ω , so that $\|v\|_{L^p(\Omega, \mathbb{R}^n)} = \tau$. One has $\lambda_1(\Omega, \text{Id}, v, 0) \leq \lambda_1(B_A, \text{Id}, A e_r(\cdot - x_0), 0)$ since $B_A \subset \Omega$. But

$$\lambda_1(B_A, \text{Id}, A e_r(\cdot - x_0), 0) = \frac{\mu_A}{\rho_A^2},$$

where $\mu_A = \lambda_1(\tilde{B}, \text{Id}, A \rho_A e_r(\cdot - x_0), 0)$ is the principal eigenvalue of $-\Delta + A \rho_A e_r(\cdot - x_0) \cdot \nabla$ on the ball \tilde{B} with center x_0 and with radius 1, under Dirichlet boundary condition. Notice that $A \rho_A \rightarrow +\infty$ as $A \rightarrow +\infty$ since $1 < p < n$. Furthermore, $\mu_A = \lambda_1(B_1^n, A \rho_A e_r)$ under the notation (2.15), where B_1^n is the Euclidean ball of \mathbb{R}^n with center 0 and radius 1. It then follows immediately from Appendix 7.3 that $\log \mu_A \sim -A \rho_A$ when $A \rightarrow +\infty$ (see also [20] for related results under stronger regularity assumptions). As a consequence, there exists $A_0 > 0$ such that, for all $A > A_0$,

$$\lambda_1(B_A, \text{Id}, A e_r(\cdot - x_0), 0) \leq \frac{e^{-A \rho_A/2}}{\rho_A^2} = \alpha_n^{2/n} \tau^{-2p/n} A^{2p/n} e^{-(\alpha_n^{-1/n} \tau^{p/n} A^{1-p/n})/2},$$

and this expression goes to 0 when $A \rightarrow +\infty$, which proves the claim.

Similarly, one can show that, if $\Omega \in \mathcal{C}$, $\tau > 0$ and $1 < q < n/2$ are fixed,

$$\inf_{V \in L^\infty(\Omega); \|V\|_q \leq \tau} \lambda_1(\Omega, \text{Id}, 0, V) = -\infty. \quad (6.5)$$

Indeed, fix ρ_0 as before, and, for all A large enough, let $\rho_A \in (0, \rho_0)$ such that $A^q \alpha_n \rho_A^n = \tau^q$, and set $V = -A \times \mathbf{1}_{B_A}$ where $B_A = B(x_0, \rho_A)$ is defined as previously, so that $\|V\|_{L^q(\Omega)} = \tau$. One has

$$\begin{aligned} \lambda_1(\Omega, \text{Id}, 0, V) &\leq \lambda_1(B_A, \text{Id}, 0, -A) = \lambda_1(B_A, \text{Id}, 0, 0) - A \\ &= \frac{C}{\rho_A^2} - A = C \alpha_n^{2/n} \tau^{-2q/n} A^{2q/n} - A, \end{aligned} \quad (6.6)$$

where $C = \lambda_1(\tilde{B}, \text{Id}, 0, 0) > 0$. The right-hand side of (6.6) goes to $-\infty$ when $A \rightarrow +\infty$, due to the choice of q . This ends the proof of the claim (6.5).

6.1.2 The case of L^∞ constraints

When solving optimization problems for $\lambda_1(\Omega, A, v, V)$ if A is fixed and v, V vary and satisfy L^∞ bounds, we can precise the conclusions of Theorem 6.1. Fix $\Omega \in \mathcal{C}$ and $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$ such that $A \geq \gamma \text{Id}$ in Ω for some positive real number $\gamma > 0$. Given $\tau_1, \tau_2 \geq 0$ and $w_1 \in L_+^\infty(\Omega)$, define

$$\mathcal{A}_{\Omega, A, w_1, \tau_1, \tau_2} = \{(v, V) \in L^\infty(\Omega, \mathbb{R}^n) \times L^\infty(\Omega); \|w_1 v\|_\infty \leq \tau_1, \|V\|_\infty \leq \tau_2, \},$$

and

$$\begin{cases} \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) &= \inf_{(v, V) \in \mathcal{A}_{\Omega, A, w_1, \tau_1, \tau_2}} \lambda_1(\Omega, A, v, V), \\ \bar{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) &= \sup_{(v, V) \in \mathcal{A}_{\Omega, A, w_1, \tau_1, \tau_2}} \lambda_1(\Omega, A, v, V). \end{cases}$$

The optimization results under these constraints are the following ones:

Theorem 6.6 *Let $\Omega \in \mathcal{C}$, $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$ such that $A \geq \gamma \text{Id}$ in Ω for some $\gamma > 0$, $\tau_1, \tau_2 \geq 0$ and $w_1 \in L_+^\infty(\Omega)$ be given. Then,*

- (1) *there exist a unique vector field $\underline{v} \in L^\infty(\Omega, \mathbb{R}^n)$ with $\|w_1 \underline{v}\|_\infty \leq \tau_1$ and a unique function $\underline{V} \in L^\infty(\Omega)$ with $\|\underline{V}\|_\infty \leq \tau_2$, such that*

$$\underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) = \lambda_1(\Omega, A, \underline{v}, \underline{V}).$$

Moreover, if $\underline{\varphi} = \varphi_{\Omega, A, \underline{v}, \underline{V}}$, one has

- (a) $\underline{v} \cdot \nabla \underline{\varphi} = -\tau_1 w_1^{-1} |\nabla \underline{\varphi}|$ a.e. in Ω ,
- (b) $|\underline{v}(x)| w_1(x) = \tau_1$ a.e. in Ω ,
- (c) $\underline{V}(x) = -\tau_2$ a.e. in Ω ,

Furthermore, $\nabla \underline{\varphi}(x) \neq 0$ a.e. in Ω and $\underline{v}(x) = -\tau_1 w_1(x)^{-1} \nabla \underline{\varphi}(x) / |\nabla \underline{\varphi}(x)|$ a.e. in Ω .

- (2) *there exist a unique vector field $\bar{v} \in L^\infty(\Omega, \mathbb{R}^n)$ with $\|w_1 \bar{v}\|_\infty \leq \tau_1$ and a unique function $\bar{V} \in L^\infty(\Omega)$ with $\|\bar{V}\|_\infty \leq \tau_2$, such that*

$$\bar{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) = \lambda_1(\Omega, A, \bar{v}, \bar{V}).$$

Moreover, if $\bar{\varphi} = \varphi_{\Omega, A, \bar{v}, \bar{V}}$, one has

- (a) $\bar{v} \cdot \nabla \bar{\varphi} = \tau_1 w_1^{-1} |\nabla \bar{\varphi}|$ a.e. in Ω ,
- (b) $|\bar{v}(x)| w_1(x) = \tau_1$ a.e. in Ω ,
- (c) $\bar{V}(x) = \tau_2$ a.e. in Ω ,

Furthermore, $\nabla \bar{\varphi}(x) \neq 0$ a.e. in Ω and $\bar{v}(x) = +\tau_1 w_1(x)^{-1} \nabla \bar{\varphi}(x) / |\nabla \bar{\varphi}(x)|$ a.e. in Ω .

Proof. As in the proof of Theorem 6.1, let us focus on the minimization problem. The existence of \underline{v} and \underline{V} such that $\lambda_1(\Omega, A, \underline{v}, \underline{V}) = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2)$ is obtained in the same way as in Lemma 6.4, except that one has to define, for almost every $x \in \Omega$,

$$\underline{v}(x) = \begin{cases} -\tau_1 w_1(x)^{-1} \frac{\nabla \underline{\varphi}(x)}{|\nabla \underline{\varphi}(x)|} & \text{if } \nabla \underline{\varphi}(x) \neq 0, \\ 0 & \text{if } \nabla \underline{\varphi}(x) = 0, \end{cases}$$

and $\underline{V}(x) = -\tau_2$ for all $x \in \Omega$, and that we do not need to introduce the vector field w . To prove the uniqueness of \underline{V} , proceed as in the proof of assertion (c) in Lemma 6.3. We are now left with the task of proving the uniqueness of \underline{v} and the fact that $w_1(x) |\underline{v}(x)| = \tau_1$ for almost every $x \in \Omega$.

First, arguing as in the proof of Lemma 6.3, one shows that, if v and V are such that $\lambda_1(\Omega, A, v, V) = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2)$ and if $\varphi = \varphi_{\Omega, A, v, V}$, then $v \cdot \nabla \varphi = -\tau_1 w_1^{-1} |\nabla \varphi|$ a.e. in Ω .

To conclude, we need the following lemma:

Lemma 6.7 *Let $\lambda \in \mathbb{R}$ and $\psi \in W^{2,r}(\Omega)$ for all $1 \leq r < +\infty$, be such that $\psi = 0$ on $\partial\Omega$, $\psi > 0$ in Ω , $\|\psi\|_\infty = 1$ and*

$$-\operatorname{div}(A\nabla\psi) - \tau_1 w_1^{-1} |\nabla\psi| - \tau_2 \psi = \lambda\psi \text{ in } \Omega.$$

Let $v \in L^\infty(\Omega, \mathbb{R}^n)$ be such that $\|w_1 v\|_\infty \leq \tau_1$ and $\lambda_1(\Omega, A, v, -\tau_2) = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2)$. Then $\lambda = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2)$ and $\psi = \varphi_{\Omega, A, v, -\tau_2}$.

Proof of Lemma 6.7. Let v be as above and set $\varphi = \varphi_{\Omega, A, v, -\tau_2}$, so that

$$-\operatorname{div}(A\nabla\varphi) - \tau_1 w_1^{-1} |\nabla\varphi| - \tau_2 \varphi = -\operatorname{div}(A\nabla\varphi) + v \cdot \nabla\varphi - \tau_2 \varphi = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) \varphi \text{ in } \Omega$$

by what we have just seen. Define also

$$w(x) = \begin{cases} -\tau_1 w_1(x)^{-1} \frac{\nabla\psi(x)}{|\nabla\psi(x)|} & \text{if } \nabla\psi(x) \neq 0, \\ 0 & \text{if } \nabla\psi(x) = 0. \end{cases}$$

One has $\|w_1 w\|_\infty = \tau_1$ and

$$-\operatorname{div}(A\nabla\psi) + w \cdot \nabla\psi - \tau_2 \psi = -\operatorname{div}(A\nabla\psi) - \tau_1 w_1^{-1} |\nabla\psi| - \tau_2 \psi = \lambda\psi \text{ in } \Omega,$$

so that, since $\psi > 0$ in Ω and $\psi = 0$ on $\partial\Omega$, by the characterization of the principal eigenfunction and the normalization $\|\psi\|_\infty = 1$, one has $\psi = \varphi_{\Omega, A, w, -\tau_2}$ and

$$\lambda = \lambda_1(\Omega, A, w, -\tau_2) \geq \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2).$$

As a consequence,

$$-\operatorname{div}(A\nabla\psi) + v \cdot \nabla\psi - \tau_2 \psi \geq -\operatorname{div}(A\nabla\psi) - \tau_1 w_1^{-1} |\nabla\psi| - \tau_2 \psi = \lambda\psi \geq \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) \psi \text{ in } \Omega,$$

while

$$-\operatorname{div}(A\nabla\varphi) + v \cdot \nabla\varphi - \tau_2 \varphi = \lambda_1(\Omega, A, v, -\tau_2) \varphi = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) \varphi \text{ in } \Omega$$

by assumption. Since $\psi > 0$ in Ω , another application of Lemma 6.2 shows that $\psi = \varphi = \varphi_{\Omega, A, v, -\tau_2}$, and that $\lambda = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2)$. \square

With the help of Lemma 6.7, we conclude the proof of Theorem 6.6. Let $v_1 \in L^\infty(\Omega, \mathbb{R}^n)$ and $v_2 \in L^\infty(\Omega, \mathbb{R}^n)$ with $\|w_1 v_1\|_\infty \leq \tau_1$ and $\|w_1 v_2\|_\infty \leq \tau_1$ be such that

$$\underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) = \lambda_1(\Omega, A, v_1, -\tau_2) = \lambda_1(\Omega, A, v_2, -\tau_2),$$

and set

$$\varphi_1 = \varphi_{\Omega, A, v_1, -\tau_2} \text{ and } \varphi_2 = \varphi_{\Omega, A, v_2, -\tau_2}.$$

Since $v_1 \cdot \nabla\varphi_1 = -\tau_1 w_1^{-1} |\nabla\varphi_1|$ and $v_2 \cdot \nabla\varphi_2 = -\tau_1 w_1^{-1} |\nabla\varphi_2|$ a.e. in Ω , one has

$$\begin{cases} -\operatorname{div}(A\nabla\varphi_1) - \tau_1 w_1^{-1} |\nabla\varphi_1| - \tau_2 \varphi_1 = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) \varphi_1 & \text{in } \Omega, \\ -\operatorname{div}(A\nabla\varphi_2) - \tau_1 w_1^{-1} |\nabla\varphi_2| - \tau_2 \varphi_2 = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) \varphi_2 & \text{in } \Omega, \end{cases}$$

with $\varphi_1, \varphi_2 \in \bigcap_{1 \leq r < +\infty} W^{2,r}(\Omega)$, $\varphi_1, \varphi_2 > 0$ in Ω , $\|\varphi_1\|_\infty = \|\varphi_2\|_\infty = 1$, and $\varphi_1 = \varphi_2 = 0$ on $\partial\Omega$. Lemma 6.7 shows that $\varphi_1 = \varphi_{\Omega, A, v_2, -\tau_2} = \varphi_2 := \varphi$, so that $v_1 \cdot \nabla\varphi = v_2 \cdot \nabla\varphi = -\tau_1 w_1^{-1} |\nabla\varphi|$. It follows that $v_1 = v_2$ and $|v_1| = |v_2| = \tau_1 w_1^{-1}$ a.e. on the set $\{x \in \Omega; \nabla\varphi(x) \neq 0\}$. It remains to be observed that $\nabla\varphi(x) \neq 0$ a.e. in Ω . Indeed, if $E = \{x \in \Omega; \nabla\varphi(x) = 0\}$, one has $\operatorname{div}(A\nabla\varphi) = 0$ a.e. in E , so that $-\tau_2\varphi = \lambda_1(\Omega, A, v_1, -\tau_2)\varphi$ in E , and since $\lambda_1(\Omega, A, v_1, -\tau_2) > -\tau_2$, one has $|E| = 0$. \square

If, in Theorem 6.6, we specialize to the case when Ω is a ball and the diffusion matrix A is equal to $\Lambda \operatorname{Id}$ with Λ radially symmetric, we obtain a more complete description of the unique minimizer and maximizer. More precisely, we have:

Theorem 6.8 *Assume that Ω is a Euclidean ball centered at 0 with radius $R > 0$, let $\Lambda \in L_+^\infty(\Omega) \cap W^{1,\infty}(\Omega)$ be radially symmetric, set $A = \Lambda \operatorname{Id}$ and use the same notations as in Theorem 6.6, under the extra assumption that w_1 is radially symmetric. Then, $\underline{v} = \tau_1 w_1^{-1} e_r$, $\bar{v} = -\tau_1 w_1^{-1} e_r$ a.e. in Ω , and $\underline{\varphi}$ and $\bar{\varphi}$ are radially symmetric and decreasing.*

Proof. Let $\varphi = \underline{\varphi} = \varphi_{\Omega, \Lambda \operatorname{Id}, \underline{v}, -\tau_2}$ where \underline{v} is given in Theorem 6.6. One has

$$-\operatorname{div}(\Lambda \nabla\varphi) - \tau_1 w_1^{-1} |\nabla\varphi| - \tau_2\varphi = \underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2)\varphi \quad (6.7)$$

in Ω . If S is any orthogonal transformation in \mathbb{R}^n and $\psi = \varphi \circ S$, then $\psi \in W^{2,r}(\Omega)$ for all $1 \leq r < +\infty$, satisfies (6.7), vanishes on $\partial\Omega$ and is positive in Ω . Lemma 6.7 therefore yields $\varphi = \psi$, which means that φ is radially symmetric, so that there exists a function $u : [0, R] \rightarrow \mathbb{R}$ such that $\varphi(x) = u(|x|)$ for all $x \in \bar{\Omega}$, and u is $C^{1,\alpha}([0, R])$ for all $0 \leq \alpha < 1$. Let $0 \leq r_1 < r_2 < R$. Remind that $\underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) = \lambda_1(\Omega, A, \underline{v}, -\tau_2) > -\tau_2$. Since

$$-\operatorname{div}(\Lambda \nabla\varphi) + \underline{v} \cdot \nabla\varphi = (\underline{\lambda}(\Omega, A, w_1, \tau_1, \tau_2) + \tau_2)\varphi > 0 \quad (6.8)$$

in Ω , the maximum principle applied to φ in B_{r_2} yields that $\varphi \geq u(r_2)$ in B_{r_2} , which means that $u(r_1) \geq u(r_2)$. Moreover, if $u(r_1) = u(r_2)$, the strong maximum principle implies that φ is constant in B_{r_2} , which is impossible because of (6.8). Therefore, $u(r_1) > u(r_2)$. Finally, if $0 \leq r_1 < r_2 = R$, one has immediately $u(r_1) > u(r_2) = 0$. Thus, u is decreasing in $[0, R]$, and this yields at once $\underline{v} = \tau_1 w_1^{-1} e_r$ from Theorem 6.6.

The arguments for $\bar{\varphi}$ and \bar{v} are entirely analogous and this completes the proof of Theorem 6.8. \square

6.2 Faber-Krahn inequalities

Proof of Theorem 2.7. First, since $\|V\|_\infty \leq \tau_2$, it follows from [14] that

$$\lambda_1(\Omega, A, v, V) \geq \lambda_1(\Omega, A, v, -\tau_2) = -\tau_2 + \lambda_1(\Omega, A, v, 0).$$

But $\lambda_1(\Omega, A, v, 0) > 0$. Theorem 2.3 then yields the existence of a positive constant $\theta = \theta(\Omega, n, \bar{M}_A, \underline{m}_\Lambda, \tau_1) > 0$ depending only on $(\Omega, n, \bar{M}_A, \underline{m}_\Lambda, \tau_1)$, and the existence of two radially symmetric $C^\infty(\bar{\Omega}^*)$ fields $\Lambda^* > 0$, $\omega^* \geq 0$ such that, for $v^* = \omega^* e_r$ in Ω^* ,

$$\begin{cases} \operatorname{ess\,inf}_\Omega \Lambda \leq \min_{\Omega^*} \Lambda^* \leq \max_{\Omega^*} \Lambda^* \leq \operatorname{ess\,sup}_\Omega \Lambda, & \|(\Lambda^*)^{-1}\|_{L^1(\Omega^*)} = \|\Lambda^{-1}\|_{L^1(\Omega)}, \\ \|v^*\|_{L^\infty(\Omega^*, \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}, & \| |v^*|^2 \times (\Lambda^*)^{-1} \|_{L^1(\Omega^*)} = \| |v|^2 \times \Lambda^{-1} \|_{L^1(\Omega)}, \end{cases}$$

and

$$\lambda_1(\Omega, A, v, 0) \geq \lambda_1(\Omega^*, \Lambda^* \text{Id}, v^*, 0) \times (1 + \theta).$$

Observe that $\|v^*\|_\infty \leq \tau_1$. It follows from Theorem 6.8 (with $w_1 = 1$) that

$$-\tau_2 + \lambda_1(\Omega^*, \Lambda^* \text{Id}, v^*, 0) = \lambda_1(\Omega^*, \Lambda^* \text{Id}, v^*, -\tau_2) \geq \lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, -\tau_2)$$

and $\lambda_1(\Omega^*, \Lambda^* \text{Id}, v^*, 0) \geq \lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, 0)$. Therefore,

$$\lambda_1(\Omega, A, v, V) \geq \lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, -\tau_2) + \theta \times \lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, 0)$$

since θ and $\lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, 0)$ are positive.

Let us now estimate $\lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, 0)$ from below. Call $\varphi = \varphi_{\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, 0}$, $\lambda = \lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, 0) > 0$ and

$$U(x) = \tau_1 \int_0^{|x|} \Lambda^*(re)^{-1} dr \text{ for all } x \in \overline{\Omega^*},$$

where e is an arbitrary unit vector. Multiply the equation

$$-\text{div}(\Lambda^* \nabla \varphi) + \tau_1 e_r \cdot \nabla \varphi = \lambda \varphi \text{ in } \Omega^*$$

by $\varphi e^{-U} \in H_0^1(\Omega^*) \cap W^{1,\infty}(\Omega^*)$ and integrate by parts over Ω^* . It follows from the definition of U that

$$\int_{\Omega^*} \Lambda^* |\nabla \varphi|^2 e^{-U} = \lambda \int_{\Omega^*} \varphi^2 e^{-U} \leq \lambda \int_{\Omega^*} \varphi^2.$$

The last inequality holds since $\lambda > 0$, and $U \geq 0$ in Ω^* . But $\Lambda^* \geq \text{ess inf}_\Omega \Lambda \geq \underline{m}_\Lambda > 0$, whence $U \leq \tau_1 \underline{m}_\Lambda^{-1} R$ in $\overline{\Omega^*}$, where $R = \alpha_n^{-1/n} |\Omega|^{1/n} > 0$ is the radius of Ω^* . Finally,

$$\lambda \int_{\Omega^*} \varphi^2 \geq \underline{m}_\Lambda e^{-\tau_1 \underline{m}_\Lambda^{-1} \alpha_n^{-1/n} |\Omega|^{1/n}} \int_{\Omega} |\nabla \varphi|^2,$$

whence

$$\lambda \geq \underline{m}_\Lambda e^{-\tau_1 \underline{m}_\Lambda^{-1} \alpha_n^{-1/n} |\Omega|^{1/n}} \times |\Omega|^{-2/n} \alpha_n^{2/n} (j_{n/2-1,1})^2 =: \kappa > 0$$

from (1.3) and (1.4). The conclusion of Theorem 2.7 follows with the choice

$$\eta = \eta(\Omega, n, \overline{M}_A, \underline{m}_\Lambda, \tau_1) = \theta \times \kappa > 0.$$

□

Proof of Corollary 2.8. Assume first that Ω is not a ball. Write

$$\lambda_1(\Omega, A, v, V) \geq \lambda_1(\Omega, A, v, 0) + \text{ess inf}_\Omega V.$$

Under the notations of Corollary 2.8, then Theorem 2.7 applied to $\lambda_1(\Omega, A, v, 0)$ with $\Lambda = \gamma_A$ clearly gives $\Lambda^* = \gamma_A$ in $\overline{\Omega^*}$, so that

$$\lambda_1(\Omega, A, v, 0) > \lambda_1(\Omega^*, \gamma_A \text{Id}, \|v\|_\infty e_r, 0),$$

whence

$$\lambda_1(\Omega, A, v, V) > \lambda_1(\Omega^*, \gamma_A \text{Id}, \|v\|_\infty e_r, 0) + \text{ess inf}_\Omega V = \lambda_1(\Omega^*, \gamma_A \text{Id}, \|v\|_\infty e_r, \text{ess inf}_\Omega V).$$

Assume now that Ω is a ball. From Theorem 2.1 applied to $\lambda_1(\Omega, A, v, 0)$ with $\Lambda = \gamma_A$, there exists $v^* \in L^\infty(\Omega^*, \mathbb{R}^n)$ such that $\|v^*\|_{L^\infty(\Omega^*, \mathbb{R}^n)} \leq \|v\|_{L^\infty(\Omega, \mathbb{R}^n)}$ and

$$\lambda_1(\Omega^*, \gamma_A \text{Id}, v^*, 0) \leq \lambda_1(\Omega, A, v, 0).$$

But $\lambda_1(\Omega^*, \gamma_A \text{Id}, \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} e_r, 0) \leq \lambda_1(\Omega^*, \gamma_A \text{Id}, v^*, 0)$ from Theorem 6.8 with $w_1 = 1$. Therefore,

$$\lambda_1(\Omega, A, v, V) \geq \lambda_1(\Omega^*, \gamma_A \text{Id}, \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} e_r, \text{ess inf}_\Omega V).$$

The conclusion of Corollary 2.8 follows immediately. \square

Remark 6.9 If, in Theorem 2.7, we specialize to the case when $A = \gamma \text{Id}$ and $\Lambda = \gamma > 0$ is a given constant, then we have immediately

$$\lambda_1(\Omega, \gamma \text{Id}, v, V) > \lambda_1(\Omega^*, \gamma \text{Id}, \|v\|_\infty e_r, -\|V\|_\infty)$$

provided that $\Omega \in \mathcal{C}$ is not a ball. Furthermore, if Ω is a ball, say with center x_0 , the uniqueness statement in Theorem 6.8 shows that $\lambda_1(\Omega, \gamma \text{Id}, v, V) \geq \lambda_1(\Omega, \gamma \text{Id}, \|v\|_\infty e_r(\cdot - x_0), -\|V\|_\infty)$, where the inequality is strict if $v \neq \|v\|_\infty e_r(\cdot - x_0)$ or $V \neq -\|V\|_\infty$. Finally, we obtain that, if $\Omega \in \mathcal{C}$, then

$$\lambda_1(\Omega, \gamma \text{Id}, v, V) \geq \lambda_1(\Omega^*, \gamma \text{Id}, \|v\|_\infty e_r, -\|V\|_\infty) \quad (6.9)$$

and the equality holds if and only if, up to translation, $\Omega = \Omega^*$, $v = \|v\|_\infty e_r$ and $V = -\|V\|_\infty$.

A rough parabolic interpretation of inequality (6.9) can be as follows: consider the evolution equation $u_t = \gamma \Delta u - v \cdot \nabla u - V u$ in Ω , for $t > 0$, with Dirichlet boundary condition on $\partial\Omega$, and with an initial datum at $t = 0$. Roughly speaking, minimizing $\lambda_1(\Omega, \gamma \text{Id}, v, V)$ with given measure $|\Omega|$ and with given L^∞ constraints $\|v\|_\infty \leq \tau_1$ and $\|V\|_\infty \leq \tau_2$ can be interpreted as looking for the slowest exponential time-decay of the solution u . The best way to do that is: 1) to try to minimize the boundary effects, namely to have the domain as round as possible, 2) to have $-V$ as large as possible, that is V as small as possible, and 3) it is not unreasonable to say that the vector field $-v$ should as much as possible point inwards the domain to avoid the drift towards the boundary. Of course diffusion, boundary losses, transport and reaction phenomena take place simultaneously, but these heuristic arguments tend to lead to the optimal triple $(\Omega, -v, -V) = (\Omega^*, -\tau_1 e_r, \tau_2)$ (up to translation).

In Theorem 2.7, it follows from the above proofs that the inequality

$$\lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, \tau_3) \leq \lambda_1(\Omega, A, v, V) - \eta$$

holds if the assumption $\|V\|_\infty = \|V\|_{L^\infty(\Omega)} \leq \tau_2$ is replaced by: $\text{ess inf}_\Omega V \geq \tau_3$. Furthermore, since $\lambda_1(\Omega, A, v, V) \geq \lambda_1(\Omega, A, v, \text{ess inf}_\Omega V) = \lambda_1(\Omega, A, v, 0) + \text{ess inf}_\Omega V$ with a strict inequality if V is not constant (see [14]), it is then immediate to check that formula (6.9) still holds when $-\|V\|_\infty$ is replaced by $\text{ess inf}_\Omega V$, and that the case of equality can be extended similarly. The parabolic interpretation is the same as above if the condition $\|V\|_\infty \leq \tau_2$ is replaced by: $\text{ess inf}_\Omega V \geq \tau_3$.

Remark 6.10 If Ω is a general open subset of \mathbb{R}^n with finite measure, and if $A \in W^{1,\infty}(\Omega, \mathcal{S}_n(\mathbb{R}))$, $v \in L^\infty(\Omega, \mathbb{R}^n)$, $V \in L^\infty(\Omega)$ are such that $A \geq \gamma \text{Id}$ in Ω for some constant $\gamma > 0$, then we claim that

$$\begin{aligned} \lambda_1(\Omega, A, v, V) &\geq \lambda_1(\Omega^*, \gamma \text{Id}, \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} e_r, \text{ess inf}_\Omega V) \\ &\geq \lambda_1(\Omega^*, \gamma \text{Id}, \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} e_r, -\|V\|_{L^\infty(\Omega)}). \end{aligned} \quad (6.10)$$

Indeed, given $\varepsilon > 0$, as in Remark 5.6, there exists a non-empty set $\Omega' = \Omega'_\varepsilon \in \mathcal{C}$ such that $\Omega' \subset\subset \Omega$ and $\lambda_1(\Omega', A, v, V) \leq \lambda_1(\Omega, A, v, V) + \varepsilon$. Then the arguments used in the proof of Corollary 2.8 (with γ instead of γ_A) imply that

$$\lambda_1(\Omega', A, v, V) \geq \lambda_1(\Omega_\varepsilon^*, \gamma \text{Id}, \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} e_r, \text{ess inf}_\Omega V),$$

where Ω_ε^* is the ball centered at the origin with the same measure as Ω' . Therefore,

$$\lambda_1(\Omega, A, v, V) + \varepsilon \geq \lambda_1(\Omega_\varepsilon^*, \gamma \text{Id}, \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} e_r, \text{ess inf}_\Omega V) \geq \lambda_1(\Omega^*, \gamma \text{Id}, \|v\|_{L^\infty(\Omega, \mathbb{R}^n)} e_r, \text{ess inf}_\Omega V)$$

since $\Omega_\varepsilon^* \subset \Omega^*$, and (6.10) follows immediately. Notice that (6.10) holds in particular with $\gamma = \text{ess inf}_\Omega \Lambda[A]$.

7 Appendix

7.1 Proof of the approximation lemma 5.1

Fix $k \in \mathbb{N}$. Call

$$r_{i,k} = \frac{iR}{k+1} \quad \text{for } i = 0, \dots, k+1$$

and

$$r_{i+1/2,k} = \left(\frac{r_{i,k}^n + r_{i+1,k}^n}{2} \right)^{1/n} \in (r_{i,k}, r_{i+1,k}) \quad \text{for } i = 0, \dots, k.$$

Remember that

$$|\Omega_{\rho^{-1}(r_{i+1,k}), \rho^{-1}(r_{i,k})}| = \alpha_n (r_{i+1,k}^n - r_{i,k}^n) = |S_{r_{i,k}, r_{i+1,k}}| \quad \text{for all } i \in \{0, \dots, k\}.$$

Let us first define the function g_k almost everywhere in Ω^* : for $i \in \{0, \dots, k\}$ and $x \in S_{r_{i,k}, r_{i+1,k}}$ such that $|x| \neq r_{i+1/2,k}$, set

$$g_k(x) = G_k(|x|),$$

where

$$G_k(r) = \sup \left\{ a \in \mathbb{R}; \frac{1}{2} \left| \{x \in \Omega_{\rho^{-1}(r_{i+1,k}), \rho^{-1}(r_{i,k})}, g(x) > a\} \right| \geq \alpha_n |r_{i+1/2,k}^n - r^n| \right\}$$

for all $r \in (r_{i,k}, r_{i+1/2,k}) \cup (r_{i+1/2,k}, r_{i+1,k})$.³ It then follows by definition that g_k is radially symmetric, nondecreasing with respect to $|x|$ in $S_{r_{i,k}, r_{i+1/2,k}}$ and nonincreasing with respect to

³In each shell $S_{r_{i,k}, r_{i+1,k}}$, the function g_k is then a kind of Schwarz decreasing rearrangement of the function g in $\Omega_{\rho^{-1}(r_{i+1,k}), \rho^{-1}(r_{i,k})}$, with respect to the inner radius $r_{i+1/2,k}$.

$|x|$ in $S_{r_{i+1/2,k}, r_{i+1,k}}$, and that

$$\begin{cases} |\{x \in \Omega_{\rho^{-1}(r_{i+1,k}), \rho^{-1}(r_{i,k})}, g(x) > t\}| = |\{x \in S_{r_{i,k}, r_{i+1,k}}, g_k(x) > t\}| \\ |\{x \in \Omega_{\rho^{-1}(r_{i+1,k}), \rho^{-1}(r_{i,k})}, g(x) \geq t\}| = |\{x \in S_{r_{i,k}, r_{i+1,k}}, g_k(x) \geq t\}| \end{cases}$$

for all $i \in \{0, \dots, k\}$ and $t \in \mathbb{R}$. As a consequence, the restriction of g_k to $S_{r_{i,k}, r_{i+1,k}}$ is in $L^\infty(S_{r_{i,k}, r_{i+1,k}})$,

$$\operatorname{ess\,inf}_\Omega g \leq \operatorname{ess\,inf}_{\Omega_{\rho^{-1}(r_{i+1,k}), \rho^{-1}(r_{i,k})}} g = \operatorname{ess\,inf}_{S_{r_{i,k}, r_{i+1,k}}} g_k \leq \operatorname{ess\,sup}_{S_{r_{i,k}, r_{i+1,k}}} g_k = \operatorname{ess\,sup}_{\Omega_{\rho^{-1}(r_{i+1,k}), \rho^{-1}(r_{i,k})}} g \leq \operatorname{ess\,sup}_\Omega g$$

and

$$\int_{S_{r_{i,k}, r_{i+1,k}}} g_k = \int_{\Omega_{\rho^{-1}(r_{i+1,k}), \rho^{-1}(r_{i,k})}} g = \int_{S_{r_{i,k}, r_{i+1,k}}} h \quad (7.1)$$

for all $i = 0, \dots, k$, by assumption (5.2). Therefore, $g_k \in L^\infty(\Omega^*)$ with

$$\operatorname{ess\,inf}_\Omega g = \operatorname{ess\,inf}_{\Omega^*} g_k \leq \operatorname{ess\,sup}_{\Omega^*} g_k = \operatorname{ess\,sup}_\Omega g$$

and

$$\begin{cases} |\{x \in \Omega, g(x) > t\}| = |\{x \in \Omega^*, g_k(x) > t\}| \\ |\{x \in \Omega, g(x) \geq t\}| = |\{x \in \Omega^*, g_k(x) \geq t\}| \end{cases}$$

for all $t \in \mathbb{R}$.

Let us now define the sequence of functions $(\underline{g}_k)_{k \in \mathbb{N}}$. Fix $k \in \mathbb{N}$. For each $i \in \{0, \dots, k\}$, the function G_k is by construction nondecreasing in the interval $(r_{i,k}, r_{i+1/2,k})$ and nonincreasing in the interval $(r_{i+1/2,k}, r_{i+1,k})$. Furthermore,

$$\operatorname{ess\,inf}_{(r_{i,k}, r_{i+1/2,k})} G_k = \operatorname{ess\,inf}_{(r_{i+1/2,k}, r_{i+1,k})} G_k = \operatorname{ess\,inf}_{\Omega_{\rho^{-1}(r_{i+1,k}), \rho^{-1}(r_{i,k})}} g \geq \operatorname{ess\,inf}_\Omega g.$$

Therefore, in each of the intervals $(r_{i,k}, r_{i+1/2,k})$ and $(r_{i+1/2,k}, r_{i+1,k})$, the function G_k can be approximated uniformly and from below by piecewise constant functions which are larger than or equal to $\operatorname{ess\,inf}_\Omega g$. As a consequence, there exists a piecewise constant function \tilde{G}_k defined in $[0, R]$ such that

$$\operatorname{ess\,inf}_\Omega g \leq \tilde{G}_k(r) \leq G_k(r) (\leq \operatorname{ess\,sup}_\Omega g) \text{ for all } r \in (0, r_{1/2,k}) \cup \dots \cup (r_{k+1/2,k}, R) \quad (7.2)$$

and

$$\|G_k - \tilde{G}_k\|_{L^\infty(0,R)} \leq \frac{1}{k+1}. \quad (7.3)$$

Let $0 = \rho_{0,k} < \rho_{1,k} < \dots < \rho_{N_k+1,k} = R$ be a subdivision adapted to \tilde{G}_k (with $N_k \in \mathbb{N}$), namely \tilde{G}_k is equal to a constant $m_{j,k} \in [\operatorname{ess\,inf}_\Omega g, \operatorname{ess\,sup}_\Omega g]$ in each interval $(\rho_{j,k}, \rho_{j+1,k})$ for $j = 0, \dots, N_k$. Choose a real number $\underline{\rho}_k$ such that

$$0 < \underline{\rho}_k < \min \left(\min_{0 \leq j \leq N_k} \frac{\rho_{j+1,k} - \rho_{j,k}}{2}, \frac{1}{(N_k + 1)(k + 1)} \right).$$

Let ζ be a fixed $C^\infty(\mathbb{R}, \mathbb{R})$ function such that $0 \leq \zeta \leq 1$ in \mathbb{R} , $\zeta = 0$ in $(-\infty, 1/3]$ and $\zeta = 1$ in $[2/3, +\infty)$. Denote \underline{G}_k the function in $[0, R]$ by:

$$\begin{cases} \underline{G}_k(r) = m_{0,k} & \text{for } r \in [0, \rho_{1,k} - \underline{\rho}_k], \\ \underline{G}_k(r) = m_{j,k} & \text{for } r \in [\rho_{j,k} + \underline{\rho}_k, \rho_{j+1,k} - \underline{\rho}_k] \text{ and } 1 \leq j \leq N_k - 1 \text{ (if } N_k \geq 1), \\ \underline{G}_k(r) = m_{N_k,k} & \text{for } r \in [\rho_{N_k,k} + \underline{\rho}_k, R], \end{cases}$$

$$\begin{aligned} \forall 0 \leq j \leq N_k - 1, \\ \forall r \in [\rho_{j+1,k} - \underline{\rho}_k, \rho_{j+1,k}], \quad \underline{G}_k(r) = \begin{cases} m_{j,k} & \text{if } m_{j,k} \leq m_{j+1,k}, \\ m_{j,k} + (m_{j+1,k} - m_{j,k}) \times \zeta \left(\frac{r - \rho_{j+1,k} + \underline{\rho}_k}{\underline{\rho}_k} \right) & \\ m_{j+1,k} & \text{if } m_{j+1,k} < m_{j,k}, \end{cases} \end{aligned}$$

and

$$\begin{aligned} \forall 1 \leq j \leq N_k, \\ \forall r \in [\rho_{j,k}, \rho_{j,k} + \underline{\rho}_k], \quad \underline{G}_k(r) = \begin{cases} m_{j,k} & \text{if } m_{j-1,k} \geq m_{j,k} \\ m_{j,k} + (m_{j-1,k} - m_{j,k}) \times \zeta \left(\frac{\rho_{j,k} + \underline{\rho}_k - r}{\underline{\rho}_k} \right) & \\ m_{j-1,k} & \text{if } m_{j-1,k} < m_{j,k}. \end{cases} \end{aligned}$$

The function \underline{G}_k is well-defined and C^∞ in $[0, R]$ and

$$\operatorname{ess\,inf}_{(0,R)} \tilde{G}_k = \min_{0 \leq j \leq N_k} m_{j,k} \leq \underline{G}_k(r) \leq \tilde{G}_k(r) \text{ for all } r \in (0, \rho_{1,k}) \cup \dots \cup (\rho_{N_k,k}, R),$$

whence

$$\operatorname{ess\,inf}_{\Omega} g \leq \underline{G}_k \leq \tilde{G}_k \leq G_k \leq \operatorname{ess\,sup}_{\Omega} g \text{ almost everywhere in } [0, R]$$

by (7.2). The function defined by

$$\underline{g}_k(x) = \underline{G}_k(|x|) \text{ for all } x \in \overline{\Omega^*}$$

is radially symmetric and of class $C^\infty(\overline{\Omega^*})$ and it satisfies

$$\operatorname{ess\,inf}_{\Omega} g \leq \underline{g}_k \leq g_k \leq \operatorname{ess\,sup}_{\Omega} g \text{ almost everywhere in } \Omega^*.$$

Fix now $q \in [1, +\infty)$ and let us check that $g_k - \underline{g}_k \rightarrow 0$ in $L^q(\Omega^*)$ as $k \rightarrow +\infty$. One has

$$\begin{aligned} \|g_k - \underline{g}_k\|_{L^q(\Omega^*)} &\leq \|G_k(|\cdot|) - \tilde{G}_k(|\cdot|)\|_{L^q(\Omega^*)} + \|\tilde{G}_k(|\cdot|) - \underline{G}_k(|\cdot|)\|_{L^q(\Omega^*)} \\ &\leq \frac{(\alpha_n R^n)^{1/q}}{k+1} + \|\tilde{G}_k(|\cdot|) - \underline{G}_k(|\cdot|)\|_{L^q(\Omega^*)} \end{aligned}$$

by (7.3). On the other hand, the definition of \underline{G}_k and formula (7.2) imply that

$$\|\tilde{G}_k - \underline{G}_k\|_{L^\infty(0,R)} = \max_{0 \leq j \leq N_k} |m_{j,k} - m_{j+1,k}| \leq 2\|\tilde{G}_k\|_{L^\infty(0,R)} \leq 2\|g\|_{L^\infty(\Omega)}.$$

Using once again the definition of \underline{G}_k , it follows that

$$\begin{aligned} \|\tilde{G}_k(|\cdot|) - \underline{G}_k(|\cdot|)\|_{L^q(\Omega^*)} &= \left[n\alpha_n \sum_{j=0}^{N_k} \left(\int_{\rho_{j,k}}^{\rho_{j,k} + \underline{\rho}_k} (\tilde{G}_k(r) - \underline{G}_k(r))^q r^{n-1} dr \right. \right. \\ &\quad \left. \left. + \int_{\rho_{j+1,k} - \underline{\rho}_k}^{\rho_{j+1,k}} (\tilde{G}_k(r) - \underline{G}_k(r))^q r^{n-1} dr \right) \right]^{1/q} \\ &\leq \left[2n\alpha_n (N_k + 1) \underline{\rho}_k R^{n-1} (2\|g\|_{L^\infty(\Omega)})^q \right]^{1/q} \\ &\leq 2\|g\|_{L^\infty(\Omega)} \times \left(\frac{2n\alpha_n R^{n-1}}{k+1} \right)^{1/q} \end{aligned}$$

from the choice of $\underline{\rho}_k$. Thus,

$$\|g_k - \underline{g}_k\|_{L^q(\Omega^*)} \leq \frac{(\alpha_n R^n)^{1/q}}{k+1} + 2\|g\|_{L^\infty(\Omega)} \times \left(\frac{2n\alpha_n R^{n-1}}{k+1} \right)^{1/q} \rightarrow 0 \text{ as } k \rightarrow +\infty.$$

Finally, let us check that the sequences $(g_k)_{k \in \mathbb{N}}$ and $(\underline{g}_k)_{k \in \mathbb{N}}$ converge to h as $k \rightarrow +\infty$ in $L^p(\Omega^*)$ weak for all $1 < p < +\infty$ and in $L^\infty(\Omega^*)$ weak-*. Let ϕ be in $C(\overline{\Omega^*}, \mathbb{R})$ and fix $\varepsilon > 0$. Since the unit sphere \mathbb{S}^{n-1} is compact and ϕ is uniformly continuous in $\overline{\Omega^*}$, there exists $k_0 \in \mathbb{N}$ and a finite family of measurable pairwise disjoint subsets U_1, \dots, U_q of \mathbb{S}^{n-1} with positive area, such that $\mathbb{S}^{n-1} = U_1 \cup \dots \cup U_q$, and

$$|\phi(x) - \phi(y)| \leq \varepsilon \text{ for all } x, y \in \overline{\Omega^*} \setminus \{0\} \text{ such that } \begin{cases} ||x| - |y|| \leq \frac{1}{k_0 + 1} \\ \exists j \text{ such that } \frac{x}{|x|}, \frac{y}{|y|} \in U_j. \end{cases} \quad (7.4)$$

Fix any k such that $k \geq k_0$. Use the notation $x = r\theta$ with $r = |x|$ and $\theta = x/|x|$ for the points of $\Omega^* \setminus \{0\}$. Denote by $d\sigma$ the surface measure on \mathbb{S}^{n-1} . For all $i \in \{0, \dots, k\}$ and $j \in \{1, \dots, q\}$, call

$$\phi_{i,j} = \frac{\int_{r_{i,k}}^{r_{i+1,k}} \int_{U_j} \phi(r\theta) d\sigma(\theta) dr}{\int_{r_{i,k}}^{r_{i+1,k}} \int_{U_j} d\sigma(\theta) dr}.$$

Since g_k and h are radially symmetric and satisfy (7.1) for all $i = 0, \dots, k$, it follows that

$$\int_{r_{i,k}}^{r_{i+1,k}} \int_{U_j} (g_k(r\theta) - h(r\theta)) \phi_{i,j} d\sigma(\theta) dr = 0$$

for all $i \in \{0, \dots, k\}$ and $j \in \{1, \dots, q\}$. Thus,

$$\int_{\Omega^*} g_k \phi - \int_{\Omega^*} h \phi = \sum_{i=0}^k \sum_{j=1}^q \int_{r_{i,k}}^{r_{i+1,k}} \int_{U_j} (g_k(r\theta) - h(r\theta)) (\phi(r\theta) - \phi_{i,j}) d\sigma(\theta) dr$$

and then

$$\begin{aligned} \left| \int_{\Omega^*} g_k \phi - \int_{\Omega^*} h \phi \right| &\leq \sum_{i=0}^k \sum_{j=1}^q \int_{r_{i,k}}^{r_{i+1,k}} \int_{U_j} (\|g\|_{L^\infty(\Omega)} + \|h\|_{L^\infty(\Omega^*)}) \varepsilon \, d\sigma(\theta) \, dr \\ &= \alpha_n R^n (\|g\|_{L^\infty(\Omega)} + \|h\|_{L^\infty(\Omega^*)}) \varepsilon \end{aligned}$$

for all $k \geq k_0$ (remember that $\|g_k\|_{L^\infty(\Omega^*)} = \|g\|_{L^\infty(\Omega)}$). Since $\varepsilon > 0$ was arbitrary, one concludes that

$$\int_{\Omega^*} g_k \phi \rightarrow \int_{\Omega^*} h \phi \text{ as } k \rightarrow +\infty.$$

Since this is true for every $\phi \in C(\overline{\Omega^*}, \mathbb{R})$, standard density arguments imply then that

$$\int_{\Omega^*} g_k \phi \rightarrow \int_{\Omega^*} h \phi \text{ as } k \rightarrow +\infty, \text{ for all } \phi \in L^{p'}(\Omega^*) \text{ and for all } p' \in [1, +\infty), \quad (7.5)$$

namely $g_k \rightharpoonup h$ as $k \rightarrow +\infty$ in $L^p(\Omega^*)$ weak for all $p \in (1, +\infty)$ and in $L^\infty(\Omega^*)$ weak-*. Lastly, since $g_k - \underline{g}_k \rightarrow 0$ as $k \rightarrow +\infty$ in $L^p(\Omega^*)$ for all $p \in [1, +\infty)$ and since the functions \underline{g}_k are uniformly bounded in $L^\infty(\Omega^*)$, it follows from (7.5) and standard density arguments that

$$\int_{\Omega^*} \underline{g}_k \phi \rightarrow \int_{\Omega^*} h \phi \text{ as } k \rightarrow +\infty, \text{ for all } \phi \in L^{p'}(\Omega^*) \text{ and for all } p' \in [1, +\infty).$$

Thus, $\underline{g}_k \rightharpoonup h$ as $k \rightarrow +\infty$ in $\sigma(L^p(\Omega^*), L^{p'}(\Omega^*))$ for all $1 < p \leq +\infty$.

The construction of the functions \bar{g}_k is similar to that of the functions \underline{g}_k , but they approximate the functions g_k from above. \square

7.2 A remark on distribution functions

Let $\alpha \leq \beta \in \mathbb{R}$ and $m > 0$ be fixed. We extend a definition which we used just before Corollary 2.2: $\mathcal{F}_{\alpha,\beta}(m)$ stands for the set of right-continuous non-increasing functions $\mu : \mathbb{R} \rightarrow [0, m]$ such that

$$\mu(t) = m \text{ for all } t < \alpha \text{ and } \mu(t) = 0 \text{ for all } t \geq \beta.$$

In this appendix, we prove the following fact:

Proposition 7.1 *Let $\alpha \leq \beta \in \mathbb{R}$, $m > 0$, $\mu \in \mathcal{F}_{\alpha,\beta}(m)$ and $\Omega \in \mathcal{C}$ such that $|\Omega| = m$. Then, there exists $V \in L^\infty(\Omega)$ such that $\mu = \mu_V$.*

Proof. This fact is rather classical, but we give here a quick proof for the sake of completeness. Let φ be the solution of

$$\begin{cases} -\Delta \varphi = 1 & \text{in } \Omega, \\ \varphi = 0 & \text{on } \partial\Omega. \end{cases}$$

Observe that the function φ belongs to $W^{2,p}(\Omega)$ for all $1 \leq p < +\infty$, to $C^{1,\gamma}(\overline{\Omega})$ for all $0 \leq \gamma < 1$ and is analytic and positive in Ω . Let $M = \max_{\overline{\Omega}} \varphi$ and, for all $0 \leq a < M$, define (as in Section 3)

$$\Omega_a = \{x \in \Omega; \varphi(x) > a\}.$$

Set also $\Omega_M = \emptyset$. Remember that, for all $0 \leq a \leq M$, $|\partial\Omega_a| = 0$.

Define now, for all $x \in \Omega$,

$$V(x) = \sup \{s \in \mathbb{R}; \mu(s) > |\Omega_{\varphi(x)}|\}.$$

Notice first that this supremum is well-defined for all $x \in \Omega$. Indeed, if $x \in \Omega$, one has $\varphi(x) > 0$, therefore $0 \leq |\Omega_{\varphi(x)}| < |\Omega|$.

We now claim that V is measurable and bounded in Ω and that $\mu_V = \mu$. Indeed, let $t \in \mathbb{R}$. By definition of V , for all $x \in \Omega$,

$$V(x) > t \Leftrightarrow (\exists s > t \text{ such that } \mu(s) > |\Omega_{\varphi(x)}|) \Leftrightarrow |\Omega_{\varphi(x)}| < \mu(t),$$

where the last equivalence follows from the right-continuity of μ and the fact that this function is non-increasing. Define now, for all $0 \leq a \leq M$, $F(a) = |\Omega_a|$. The previous equivalence yields

$$\mu_V(t) = |\{x \in \Omega; F(\varphi(x)) < \mu(t)\}|.$$

Since the function $F : [0, M] \rightarrow [0, |\Omega|]$ is decreasing, one-to-one and onto, one obtains that $\{x \in \Omega; V(x) > t\}$ is measurable for all $t \in \mathbb{R}$, and that

$$\mu_V(t) = |\{x \in \Omega; \varphi(x) > F^{-1}(\mu(t))\}| = |\Omega_{F^{-1}(\mu(t))}| = \mu(t),$$

where the last equality uses the definition of F . Finally, $|\{x \in \Omega; V(x) > \beta\}| = \mu(\beta) = 0$ and, for all $s < \alpha$,

$$|\{x \in \Omega; V(x) \leq s\}| = |\Omega| - \mu(s) = 0,$$

which shows that $V \in L^\infty(\Omega)$. □

7.3 Estimates of $\lambda_1(B_R^n, \tau e_r)$ as $\tau \rightarrow +\infty$

We recall that $\lambda_1(\Omega, v)$ is defined as $\lambda_1(\Omega, \text{Id}, v, 0)$ for $v \in L^\infty(\Omega, \mathbb{R}^n)$. We call B_R^n the open Euclidean ball of \mathbb{R}^n with center 0 and radius $R > 0$, and we set

$$G_n(m, \tau) = \lambda_1(B_{(m/\alpha_n)^{1/n}}^n, \tau e_r)$$

for all $n \in \mathbb{N} \setminus \{0\}$, $m > 0$ and $\tau \geq 0$. Notice that $G_n(m, \tau)$ is always positive.

Our goal here is to discuss the behavior of $G_n(m, \tau)$ for large τ . Indeed, if, in Theorem 2.7, Λ is a constant $\gamma > 0$, then, with the same notations as in Theorem 2.7,

$$\begin{aligned} \lambda_1(\Omega, A, v, V) &\geq \lambda_1(\Omega^*, \gamma \text{Id}, \tau_1 e_r, -\tau_2) = \gamma \lambda_1(\Omega^*, \text{Id}, \tau_1 \gamma^{-1}, 0) - \tau_2 \\ &= \gamma G_n(|\Omega|, \tau_1 \gamma^{-1}) - \tau_2. \end{aligned}$$

The constants γ and τ_2 appear as multiplicative and additive constants in the previous formula. The function $[0, +\infty) \ni \tau \mapsto G_n(m, \tau) > 0$ is obviously continuous, and decreasing (as a consequence of Theorem 6.8). However, the behaviour when $\tau \rightarrow +\infty$ is not immediate. It is the purpose of the following lemma, which was used in Remark 6.5. When Λ is not constant in Theorem 2.7 but still satisfies some given lower and upper bounds, the following lemma provides some bounds of $\lambda_1(\Omega^*, \Lambda^* \text{Id}, \tau_1 e_r, -\tau_2)$ when $\tau_1 \rightarrow +\infty$.

Lemma 7.2 For all $m > 0$, $\tau^{-2}e^{\tau m/2}G_1(m, \tau) \rightarrow 1$ as $\tau \rightarrow +\infty$, and one even has

$$\exists C(m) \geq 0, \exists \tau_0 \geq 0, \forall \tau \geq \tau_0, \quad |\tau^{-2}e^{\tau m/2}G_1(m, \tau) - 1| \leq C(m)\tau e^{-\tau m/2}. \quad (7.6)$$

Moreover, for all $n \geq 2$ and $m > 0$, $G_n(m, \tau) > G_1(2(m/\alpha_n)^{1/n}, \tau)$ for all $\tau \geq 0$, and

$$-\tau^{-1} \log G_n(m, \tau) \rightarrow m^{1/n} \alpha_n^{-1/n} \text{ as } \tau \rightarrow +\infty. \quad (7.7)$$

In [20], with probabilistic arguments, Friedman proved some lower and upper logarithmic estimates, as $\varepsilon \rightarrow 0^+$, for the first eigenvalue of general elliptic operators $-a_{ij}\varepsilon^2\partial_{ij} + b_i\partial_i$ with C^1 drifts $-b = -(b_1, \dots, b_n)$ pointing inwards on the boundary (see also [45]). Apart from the fact that the vector field e_r is not C^1 at the origin, the general result of Friedman would imply the asymptotics (7.7) for $\log G_n(m, \tau) = \log \lambda_1(B_{(m/\alpha_n)^{1/n}}^n, \tau e_r)$. For the sake of completeness, we give here a proof of (7.7) with elementary analytic arguments. Lemma 7.2 also provides the precise equivalent of $G_1(m, \tau)$ for large τ . However, giving an equivalent for $G_n(m, \tau)$ when τ is large and $n \geq 2$ is an open question.

Proof of Lemma 7.2. First, to prove (7.6), fix $m > 0$ and $\tau \geq 0$, set $\Omega = (-R, R)$ with $2R = m$, and denote

$$\lambda = \lambda_1(\Omega, \tau e_r)$$

and $\varphi = \varphi_{\Omega, \text{Id}, \tau e_r, 0}$. Theorem 6.8 ensures that φ is an even function, decreasing in $[0, R]$ and that

$$-\varphi''(r) + \tau\varphi'(r) = \lambda\varphi(r) \text{ for all } 0 \leq r \leq R,$$

with $\varphi(R) = 0$, $\varphi > 0$ in $(-R, R)$ and $\varphi'(0) = 0$ (in particular, the above equality holds in the classical sense in $[0, R]$). For all $s \in [0, \tau R]$, define $\psi(s) = \varphi(s/\tau)$, so that ψ satisfies the equation

$$-\psi''(s) + \psi'(s) = \frac{\lambda}{\tau^2}\psi(s) \text{ for all } 0 \leq s \leq \tau R,$$

with $\psi(\tau R) = 0$ and $\psi'(0) = 0$. Notice that λ depends on τ , but since, for all $\tau \geq 0$, $0 < \lambda \leq \lambda_1((-R, R), 0)$, there exists $\tau_0 > 0$ such that $\tau^2 \geq 4\lambda$ for all $\tau \geq \tau_0$, and we will assume that $\tau \geq \tau_0$ in the sequel. The function ψ can be computed explicitly: there exist $A, B \in \mathbb{R}$ such that, for all $0 \leq s \leq \tau R$,

$$\psi(s) = Ae^{\mu_+ s} + Be^{\mu_- s},$$

where $\mu_{\pm} = (1 \pm \sqrt{1 - 4\lambda/\tau^2})/2$. Using the boundary values of ψ and ψ' , one obtains after straightforward computations:

$$\lambda = \frac{\tau^2}{4} \left(1 + \sqrt{1 - \frac{4\lambda}{\tau^2}} \right)^2 e^{-\sqrt{1 - \frac{4\lambda}{\tau^2}}\tau R}.$$

Since λ remains bounded when $\tau \rightarrow +\infty$, it is then straightforward to check that $\lambda \sim \tau^2 e^{-\tau R}$ when $\tau \rightarrow +\infty$, and that (7.6) follows.

We now turn to the proof of assertion (7.7). Let $n \geq 2$, $m > 0$, $\tau \geq 0$ and $\Omega = B_R^n$ be such that $|\Omega| = m$, so that one has $R = (m/\alpha_n)^{1/n}$ and $G_n(m, \tau) = \lambda_1(\Omega, \tau e_r)$. We first claim that

$$G_n(m, \tau) > G_1(2R, \tau).$$

Indeed, write

$$\lambda = \lambda_1(\Omega, \tau e_r) \quad \text{and} \quad \varphi_n = \varphi_{\Omega, \text{Id}, \tau e_r, 0}.$$

Similarly, $G_1(2R, \tau) = \lambda_1((-R, R), \tau e_r)$, and we denote $\mu = \lambda_1((-R, R), \tau e_r)$ and $\varphi_1 = \varphi_{(-R, R), \text{Id}, \tau e_r, 0}$ (where Id is then understood as the 1×1 identity matrix). As before, define $\psi_n(y) = \varphi_n(y/\tau)$ for all $y \in \tau\overline{\Omega} = \overline{B_{\tau R}^n}$ and $\psi_1(r) = \varphi_1(r/\tau)$ for all $r \in [-\tau R, \tau R]$. Finally, since ψ_n is radially symmetric, let $u_n : [0, \tau R] \rightarrow \mathbb{R}$ such that $\psi_n(y) = u_n(|y|)$ for all $y \in \tau\overline{\Omega} = \overline{B_{\tau R}^n}$. One has

$$\begin{cases} -u_n''(r) - \frac{n-1}{r}u_n'(r) + u_n'(r) = \frac{\lambda}{\tau^2}u_n(r) & \text{in } (0, \tau R], \\ -\psi_1''(r) + \psi_1'(r) = \frac{\mu}{\tau^2}\psi_1(r) & \text{in } [0, \tau R], \end{cases} \quad (7.8)$$

with $u_n'(\tau R) = u_n(\tau R) = 0$, $\psi_1'(0) = \psi_1(\tau R) = 0$.

Assume that $\lambda \leq \mu$. Since $u_n' < 0$ in $(0, \tau R]$ and $u_n \geq 0$, one obtains

$$\begin{cases} -u_n''(r) + u_n'(r) \leq \frac{\mu}{\tau^2}u_n(r) & \text{in } [0, \tau R], \\ -\psi_1''(r) + \psi_1'(r) = \frac{\mu}{\tau^2}\psi_1(r) & \text{in } [0, \tau R]. \end{cases} \quad (7.9)$$

Since $\psi_1'(\tau R) < 0$ by Hopf lemma, while $\psi_1(r) > 0$ in $[0, \tau R)$, $u_n(r) > 0$ in $[0, \tau R)$ and the functions u_n and ψ_1 belong (at least) to $C^1([0, \tau R])$, there exists then $\gamma > 0$ such that $\gamma\psi_1(r) > u_n(r)$ for all $0 \leq r < \tau R$. Define γ^* as the infimum of all the $\gamma > 0$ such that $\gamma\psi_1 > u_n$ in $[0, \tau R]$, observe that $\gamma^* > 0$ and define $z = \gamma^*\psi_1 - u_n$ which is non-negative in $[0, \tau R]$ and satisfies

$$-z''(r) + z'(r) - \frac{\mu}{\tau^2}z(r) \geq 0 \quad (7.10)$$

for all $0 \leq r \leq \tau R$ and $z(\tau R) = 0$.

Assume that there exists $0 < r < \tau R$ such that $z(r) = 0$. The strong maximum principle shows that z is identically zero in $[0, \tau R]$, which means that $\gamma^*\psi_1 = u_n$ in $[0, \tau R]$, and even that $\psi_1 = u_n$ because $\psi_1(0) = u_n(0) = 1$. But this is impossible according to (7.8) and (7.9).

Thus, $z > 0$ in $(0, \tau R)$. Furthermore, $z'(0) = 0$, hence $z(0) > 0$ from Hopf lemma. Another application of Hopf lemma shows that $z'(\tau R) < 0$. Therefore, there exists $\kappa > 0$ such that $z > \kappa u_n$ in $[0, \tau R)$, whence

$$\frac{\gamma^*}{1 + \kappa}\psi_1 > u_n \quad \text{in } [0, \tau R),$$

which is a contradiction with the definition of γ^* .

Finally, we have obtained that $\mu < \lambda$, which means that $G_n(m, \tau) > G_1(2R, \tau)$.

We now look for a reverse inequality. To that purpose, let $\varepsilon \in (0, 1)$ and $R_0 > 0$ such that $\frac{n-1}{R_0} < \varepsilon$. In the following computations, we always assume that $\tau R > R_0$. Define u_n and λ as before. Let

$$\mu' = \lambda_1 \left(\left(- \left(R - \frac{R_0}{\tau} \right), \left(R - \frac{R_0}{\tau} \right) \right), \tau(1 - \varepsilon)e_r \right)$$

and w the normalized corresponding eigenfunction, so that

$$\begin{cases} -w''(r) + \tau(1 - \varepsilon)w'(r) = \mu'w(r) & \text{in } \left[0, R - \frac{R_0}{\tau} \right], \\ w'(0) = 0, w > 0 & \text{in } \left[0, R - \frac{R_0}{\tau} \right), w \left(R - \frac{R_0}{\tau} \right) = 0. \end{cases}$$

For all $R_0 \leq x \leq \tau R$, define $v(x) = w \left(\frac{x - R_0}{\tau} \right)$, which satisfies

$$\begin{cases} -v''(r) + (1 - \varepsilon)v'(r) = \frac{\mu'}{\tau^2}v(r) & \text{in } [R_0, \tau R], \\ v'(R_0) = 0, v > 0 & \text{in } [R_0, \tau R), v(\tau R) = 0. \end{cases}$$

Assume that $\lambda \geq \mu'$. Since $(n-1)/R_0 < \varepsilon$ and $u'_n(r) < 0$ in $(0, \tau R]$, one therefore has

$$\begin{cases} -u''_n(r) + (1 - \varepsilon)u'_n(r) \geq \frac{\mu'}{\tau^2}u_n(r) & \text{in } [R_0, \tau R], \\ -v''(r) + (1 - \varepsilon)v'(r) = \frac{\mu'}{\tau^2}v(r) & \text{in } [R_0, \tau R]. \end{cases}$$

Arguing as before, we see that there exists $\gamma > 0$ such that $\gamma u_n > v$ in $[R_0, \tau R]$. Define γ^* (> 0) as the infimum of all such γ 's and define $z = \gamma^* u_n - v$, which is nonnegative in $[R_0, \tau R]$ and satisfies $-z'' + (1 - \varepsilon)z' - (\mu'/\tau^2)z \geq 0$ in $[R_0, \tau R]$.

Assume that $z(r) = 0$ for some $r \in (R_0, \tau R)$. The strong maximum principle ensures that z is 0 in $[R_0, \tau R]$, which means that $u_n = v$ in $[R_0, \tau R]$, which is impossible because $u'_n(R_0) < 0 = v'(R_0)$.

Therefore, $z > 0$ everywhere in $(R_0, \tau R)$. Furthermore, $z'(R_0) < 0$, thus $z(R_0) > 0$. On the other hand, by Hopf lemma, $z'(\tau R) < 0$. Thus, there exists $\kappa > 0$ such that $z > \kappa v$ in $[R_0, \tau R)$, whence $(\gamma^*/(1 + \kappa))u_n > v$ in $[R_0, \tau R)$. This contradicts the definition of γ^* .

Thus, we have established that $\lambda < \mu'$. Straightforward computations (similar to those of the proof of (7.6)) show that

$$\lambda < \mu' = \frac{\tau^2}{4} \left(1 - \varepsilon + \sqrt{(1 - \varepsilon)^2 - \frac{4\mu'}{\tau^2}} \right)^2 e^{-\sqrt{(1 - \varepsilon)^2 - \frac{4\mu'}{\tau^2}}(\tau R - R_0)},$$

and, since $\lambda > G_1(2R, \tau)$, formula (7.6) and the fact that $m = \alpha_n R^n$ end the proof of (7.7). \square

References

- [1] S. AGMON, On positivity and decay of solutions of second order elliptic equations on Riemannian manifolds. In: *Methods of functional analysis and theory of elliptic equations (Naples, 1982)*, Liguori, Naples, 1983, 19-52.
- [2] A. ALVINO and G. TROMBETTI, Equazioni ellittiche con termini di ordine inferiore e riordinamenti, *Atti Accad. Naz. Lincei Rend. Cl. Sci. Fis. Mat. Natur.* **66** (1979), 1-7.
- [3] A. ALVINO and G. TROMBETTI, Sulle migliori costanti di maggiorazione per una classe di equazioni ellittiche degeneri e non, *Ric. Mat.* **30** (1981), 15-33.
- [4] A. ALVINO, G. TROMBETTI, P.-L. LIONS and S. MATARASSO, Comparison results for solutions of elliptic problems via symmetrization, *Ann. Inst. Henri Poincaré* **16** 2 (1999), 167-188.
- [5] M.S. ASHBAUGH and R.D. BENGURIA, A sharp bound for the ratio of the first two eigenvalues of Dirichlet Laplacians and extensions, *Annals of Math.* **135** (1992), 601-628.
- [6] M.S. ASHBAUGH and R.D. BENGURIA, Eigenvalue ratios for Sturm-Liouville operators, *J. Diff. Eqs.* **103** (1993), 205-219.
- [7] M.S. ASHBAUGH and R.D. BENGURIA, Isoperimetric bounds for higher eigenvalue ratios for the n -dimensional fixed membrane problem, *Proc. Royal Soc. Edinburgh A* **123** (1993), 977-985.
- [8] M.S. ASHBAUGH and R.D. BENGURIA, On Rayleigh's conjecture for the clamped plate and its generalization to three dimensions, *Duke Math. J.* **78** (1995), 1-17.
- [9] M.S. ASHBAUGH, E.M. HARRELL and R. SVIRSKY, On minimal and maximal eigenvalue gaps and their causes, *Pacific J. Math.* **147** (1991), 1-24.
- [10] C. BANDLE, Isoperimetric inequality for some eigenvalues of an inhomogeneous, free membrane, *SIAM J. Appl. Math.* **22** (1972), 142-147.
- [11] C. BANDLE, *Isoperimetric inequalities and applications*, Pitman Monographs and Studies in Math. **7**, Boston, 1980.
- [12] D.C. BARNES, Extremal problems for eigenvalues with applications to buckling, vibration and sloshing, *SIAM J. Math. Anal.* **16** (1985), 341-357.
- [13] H. BERESTYCKI, F. HAMEL and N. NADIRASHVILI, Elliptic eigenvalue problems with large drift and applications to nonlinear propagation phenomena, *Commun. Math. Phys.* **253** (2005), 451-480.
- [14] H. BERESTYCKI, L. NIRENBERG and S.R.S. VARADHAN, The principal eigenvalue and maximum principle for second-order elliptic operators in general domains, *Comm. Pure Appl. Math.* **47** (1994), 47-92.

- [15] M. H. BOSSEL, Membranes élastiquement liées: extensions du théorème de Rayleigh-Faber-Krahn et de l'inégalité de Cheeger, *C. R. Acad. Sci. Paris Sér. I Math.* **302** (1986), 47-50.
- [16] F. BROCK, An isoperimetric inequality for eigenvalues of the Stekloff problem, *Z. Angew. Math. Mech.* **81** (2001), 69-71.
- [17] D. BUCUR and A. HENROT, Minimization of the third eigenvalue of the Dirichlet Laplacian, *Proc. Roy. Soc. London Ser. A* **456** (2000), 985-996.
- [18] S.-Y. CHENG and K. ODEN, Isoperimetric inequalities and the gap between the first and second eigenvalues of an Euclidean domain, *J. Geom. Anal.* **7** (1997), 217-239.
- [19] G. FABER, Beweis, dass unter allen homogenen Membranen von gleicher Fläche und gleicher Spannung die kreisförmige den tiefsten Grundton gibt, *Sitzungsberichte der mathematisch-physikalischen Klasse der Bayerischen Akademie der Wissenschaften zu München Jahrgang, 1923*, 169-172.
- [20] A. FRIEDMAN, The asymptotic behavior of the first real eigenvalue of a second order elliptic operator with a small parameter in the highest derivatives, *Indiana Univ. Math. J.* **22** (1972/1973), 1005-1015.
- [21] F. HAMEL, N. NADIRASHVILI, E. RUSS, An isoperimetric inequality for the principal eigenvalue of the Laplacian with drift, *C. R. Acad. Sci. Paris Ser. I* **340** (2005), 347-352.
- [22] F. HAMEL, N. NADIRASHVILI, E. RUSS, A Faber-Krahn inequality with drift, preprint, <http://fr.arxiv.org/abs/math.AP/0607585>
- [23] A. HENROT, Minimization problems for eigenvalues of the Laplacian, *J. Evol. Eq.* **3** (2003), 443-461.
- [24] A. HENROT, *Extremum Problems for Eigenvalues of Elliptic Operators*, Birkhäuser, 2006.
- [25] C.J. HOLLAND, A minimum principle for the principal eigenvalue for second-order linear elliptic equations with natural boundary conditions, *Comm. Pure Appl. Math.* **31** (1978), 509-519.
- [26] S. KARAA, Extremal eigenvalue gaps for the Schrödinger operator with Dirichlet boundary conditions, *J. Math. Physics* **39** (1998), 2325-2332.
- [27] S. KARAA, Sharp estimates for the eigenvalues of some differential equations, *SIAM J. Math. Anal.* **29** (1998), 1279-1300.
- [28] E. KRAHN, Über eine von Rayleigh formulierte Minimaleigenschaft des Kreises, *Math. Ann.* **94** (1925), 97-100.

- [29] E. KRAHN, Über Minimaleigenschaft der Kugel in drei und mehr Dimensionen, Acta Comm. Univ. Tartu (Dorpat) **A9** (1926), 1-44.
- [30] Ü. LUMISTE and J. PEETRE, *Edgar Krahn, 1894-1961, A Centenary Volume.*, IOS Press, Amsterdam, 1994.
- [31] P. LI and S.T. YAU, On the Schrödinger equation and the eigenvalue problem, Commun. Math. Phys. **88** (1983), 309-318.
- [32] P. MARCELLINI, Bounds for the third membrane eigenvalue, J. Diff. Eqs. **37** (1980), 438-443.
- [33] N.S. NADIRASHVILI, Rayleigh's conjecture on the principal frequency of the clamped plate, Arch. Rational Mech. Anal. **129** (1995), 1-10.
- [34] R.D. NUSSBAUM and Y. PINCHOVER, On variational principles for the generalized principal eigenvalue of second order elliptic operators and some applications, J. Anal. Math. **59** (1992), 161-177.
- [35] R. OSSERMAN, The isoperimetric inequality, Bull. Amer. Math. Soc. **84** (1978), 1182-1238.
- [36] L.E. PAYNE, G. PÓLYA and H.F. WEINBERGER, On the ratio of consecutive eigenvalues, J. Math. Phys. **35** (1956), 289-298.
- [37] G. PÓLYA, On the characteristic frequencies of a symmetric membrane, Math. Z. **63** (1955), 331-337.
- [38] G. PÓLYA, On the eigenvalues of vibrating membranes, Proc. London Math. Soc. (3) **11** (1961), 419-433.
- [39] G. PÓLYA and G. SZEGÖ, *Isoperimetric Inequalities in Mathematical Physics*, Annals of Mathematics Studies **27**, Princeton Univ. Press, Princeton, 1951.
- [40] J.W.S. RAYLEIGH, *The Theory of Sound*, 2nd ed. revised and enlarged (in 2 vols.), Dover Publications, New York, 1945 (republication of the 1894/1896 edition).
- [41] J. SOUČEK and V. SOUČEK, Morse-Sard theorem for real-analytic functions, Comment. Math. Univ. Carolinae **13** (1972), 45-51.
- [42] G. SZEGÖ, Inequalities for certain eigenvalues of a membrane of given area, J. Rational Mech. Anal. **3** (1954), 343-356.
- [43] G. TALENTI, Linear Elliptic P.D.E.'s: Level Sets, Rearrangements and a priori Estimates of Solutions, Bollettino U.M.I. **6** 4-B (1985), 917-949.
- [44] G. TROMBETTI and J. L. VAZQUEZ, A symmetrization result for elliptic equations with lower-order terms, Ann. Fac. Sci. Toulouse **7** (1985), 137-150.

- [45] A.D. VENTCEL, On the asymptotic behavior of the greatest eigenvalue of a second-order elliptic differential operator with a small parameter in the higher derivatives, Soviet Math. **13** (1972), 13-17.
- [46] H.F. WEINBERGER, An isoperimetric inequality for the N -dimensional free membrane problem, J. Rational Mech. Anal. **5** (1956), 633-636.
- [47] S.A. WOLF and J.B. KELLER, Range of the first two eigenvalues of the Laplacian, Proc. Royal Soc. London Ser. A **447** (1994), 397-412.