



**HAL**  
open science

# Full Asymptotics for Steady State Voltage Potentials in a Bidimensional Highly Contrasted Medium

Clair Poignard

► **To cite this version:**

Clair Poignard. Full Asymptotics for Steady State Voltage Potentials in a Bidimensional Highly Contrasted Medium. 2006. hal-00083670v2

**HAL Id: hal-00083670**

**<https://hal.science/hal-00083670v2>**

Preprint submitted on 12 Dec 2006 (v2), last revised 27 Mar 2007 (v3)

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# FULL ASYMPTOTICS FOR STEADY STATE VOLTAGE POTENTIALS IN A BIDIMENSIONAL HIGHLY CONTRASTED MEDIUM

CLAIR POIGNARD

ABSTRACT. We study the behavior of steady state voltage potentials in a bidimensional media composed of material of complex permittivity equal to 1 (respectively  $\alpha$ ) surrounded by a thin membrane of thickness  $h$  and of complex permittivity  $\alpha$  (respectively 1). We provide in both cases a rigorous derivation of the asymptotic expansion of steady state voltage potentials at any order as  $h$  tends to zero, when Neumann boundary condition is imposed on the exterior boundary of the thin layer. Our complex parameter  $\alpha$  is bounded but may be very small compared to 1, hence our results describe the asymptotics of steady state voltage potentials in highly contrasted media. The terms of the potential in the membrane are given explicitly in local coordinates in terms of the boundary data and of the curvature of the domain, while these of the cytoplasmic potential are the solutions to the so-called dielectric formulation with appropriate boundary conditions. The error estimates are given explicitly in terms of  $h$  and  $\alpha$ .

## INTRODUCTION

We study the behavior of the steady state voltage potentials in bidimensional highly contrasted media. This work is the generalization to a domain of class  $\mathcal{C}^\infty$  of arbitrary shape of the asymptotic expansion performed by the author in the case of a circular domain (see Chapter I Section 1.1 of [14]) for the so-called dielectric formulation with Neumann boundary condition and is inspired by Chapter I Section 1.2 of [14]. The motivation of the present work comes from numerical problems raised by the researchers in computational electromagnetics of Ampère<sup>1</sup>, who want to compute the quasi-static electric field in the biological cell, and generally speaking in highly contrasted materials. We refer the reader to the introduction of [14] for a precise description of the biological cell.

In this paper, two kind of media are considered. The first one consists of a conducting inner domain (say that its complex permittivity is equal to 1) surrounded by an insulating<sup>2</sup> thin membrane; we denote by  $\alpha$  the membrane complex permittivity. This model corresponds to the electric model of biological cells [10], [12]. The second material considered consists of an insulating inner domain of permittivity  $\alpha$  surrounded by a conducting thin membrane (say that its complex permittivity is

---

*Date:* ...

1991 *Mathematics Subject Classification.* 34E05, 34E10, 35J05.

*Key words and phrases.* Asymptotics, Voltage Perturbations, Laplace equation.

The author thanks Michelle Schatzman for many fertile discussions and good advices. Remarks of the referees are also of much help in improving the final version.

<sup>1</sup>Ampère is the ex-CEGELY, the center of electrical engineering of Lyon (France).

<sup>2</sup>A material is said insulating as soon as its complex permittivity tends to zero.

equal to 1). The membrane relative thickness is equal to  $h$ , while the characteristic length of the inner domains is equal to 1.

The aim of this paper is to derive full rigorous asymptotic expansion of steady state voltage potentials with respect to the small parameter  $h$ , whatever  $\alpha$  is. Let us write mathematically our problem.

Let  $\Omega_h$  be a smooth bounded bidimensional domain (see Fig. 1), composed of a smooth domain  $\mathcal{O}$  surrounded by a thin membrane  $\mathcal{O}_h$  with a small constant thickness  $h$ :

$$\Omega_h = \mathcal{O} \cup \mathcal{O}_h.$$

Let  $\alpha$  be a non null complex with positive real part;  $\alpha$  is bounded but it may be very small. Without loss of generality, we suppose that  $|\alpha| \leq 1$ . Denote by  $q_h$  and  $\gamma_h$  the following piecewise constant functions

$$\forall x \in \Omega_h, \quad q_h(x) = \begin{cases} 1, & \text{if } x \in \mathcal{O}, \\ \alpha, & \text{if } x \in \mathcal{O}_h, \end{cases}$$

$$\forall x \in \Omega_h, \quad \gamma_h(x) = \begin{cases} \alpha, & \text{if } x \in \mathcal{O}, \\ 1, & \text{if } x \in \mathcal{O}_h. \end{cases}$$

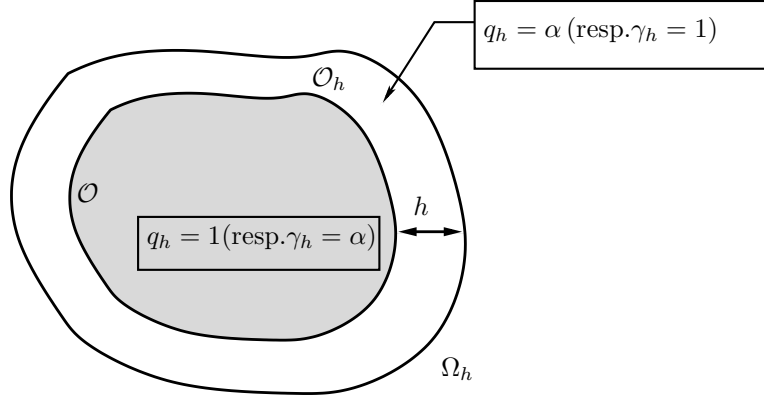


FIGURE 1. Parameters of  $\Omega_h$ .

We would like to understand the behavior for  $h$  tending to zero and uniformly with respect to  $|\alpha| \leq 1$  of  $V$  and  $u$  the respective solutions to Problem (1) and (2) with Neumann boundary condition:

$$(1a) \quad \nabla \cdot (q_h \nabla V) = 0 \text{ in } \Omega_h,$$

$$(1b) \quad \frac{\partial V}{\partial n} = \phi \text{ on } \partial\Omega_h,$$

$$(1c) \quad \int_{\partial\mathcal{O}} V \, d\sigma = 0;$$

and  $u$  satisfies

$$(2a) \quad \nabla \cdot (\gamma_h \nabla u) = 0 \text{ in } \Omega_h,$$

$$(2b) \quad \frac{\partial u}{\partial n} = \phi \text{ on } \partial\Omega_h,$$

$$(2c) \quad \int_{\partial\mathcal{O}} u \, d\sigma = 0.$$

Since we impose a Neumann boundary conditions on  $\partial\Omega_h$  the boundary data  $\phi$  must satisfy the compatibility condition:

$$\int_{\partial\Omega_h} \phi \, d\sigma = 0.$$

The above functions  $V$  and  $u$  are well-defined and belong to  $H^1(\Omega_h)$  as soon as  $\phi$  belongs to  $H^{-1/2}(\partial\Omega_h)$ . For sake of simplicity, we suppose that  $\phi$  is smooth, a weaker assumption will be made precise in our main theorems.

Several authors have worked on similar problems (see for instance Beretta *et al.* [4] and [5]). They compared the exact solution to the so-called background solution defined by replacing the material of the membrane by cytoplasmic material. The difference between these two solutions has then been given through an integral involving the polarization tensor defined for instance in [1], [2], [4], [5], [6], plus some remainder terms. The remainder terms are estimated in terms of the measure of the inhomogeneity. In this paper, we do not use this approach, for several reasons.

The Beretta *et al.* estimate of the remainder terms depends linearly on  $\alpha$  and  $1/\alpha$ : their results are no more valid in a highly contrasted domain (*i.e.* for  $\alpha$  very large or very small). Secondly,  $\alpha$  is complex-valued, hence differential operators involved in our case are not self-adjoint, so that the  $\Gamma$ -convergence techniques of Beretta *et al.* do not apply. Thirdly, the potential in the membrane is not given explicitly in [4], [5] or [6], while we are definitely interested in this potential, in order to obtain the transmembranar potential (see Fear and Stuchly [10]). Finally, the asymptotics of Beretta *et al.* are valid on the boundary of the domain, while we are interested in the potentials in the inner domain.

The heuristics of this work consist in performing a change of coordinates in the membrane  $\mathcal{O}_h$ , so as to parameterize it by local coordinates  $(\eta, \theta)$ , which vary in a domain independently of  $h$ ; in particular, if we denote by  $L$  the length of  $\partial\mathcal{O}$ , the variables  $(\eta, \theta)$  should vary in  $[0, 1] \times \mathbb{R}/L\mathbb{Z}$ . This change of coordinates leads to an expression of the Laplacian in the membrane, which depends on  $h$ . Once the transmission conditions of the new problem are derived, we perform a formal asymptotic expansion of the solution to (1) (respectively (2)) in terms of  $h$ . It remains to validate this expansion. In this paper we work with bidimensional domain and we expect that the same analysis could be performed in higher dimensions.

It has been brought to our attention that a similar approach has been developed in the thesis of Vial [16]. However, our result is more general and more precise. First, our coefficient  $\alpha$  is complex with non-negative real part, and not only strictly positive thus it is slightly more complicated to obtain our estimates. Then our estimates involve appropriate Sobolev norms of the boundary data  $\phi$ , whereas in [16] the constant depends on  $\phi$ . But mainly, our asymptotics are valid for all  $\alpha$ . Actually, Vial estimates cannot be used neither for  $\alpha$  tending to zero nor for  $|\alpha|$  tending to infinity, while ours can. In particular, when the inner domain becomes insulating (*i.e.* if  $\alpha = \beta h^q$ ,  $q \geq 1$ ), the asymptotic expansions of the potentials

begin at the order  $-1$ . In this case, it is a feature of our results that mixed boundary condition or Dirichlet boundary condition have to be imposed to the terms of the asymptotics of the inner potential respectively if  $q = 1$  and  $q \geq 2$ .

This paper is structured as follows. In Section 1, we make precise our geometric conventions. We perform a change of variables in the membrane, and with the help of some differential geometry results, we write Problem (1) and Problem (2) in the language of differential forms. We refer the reader to Flanders [11] or Dubrovin *et al.* [8] (or [7] for the french version) for courses on differential geometry. We derive transmission and boundary conditions in the intrinsic language of differential forms, and we express these relations in local coordinates.

In Section 3 we study Problem (1). In paragraph 3.1 we derive formally all the terms of the asymptotic expansion of the solution to our problem in terms of  $h$ . Paragraph 3.2 is devoted to a proof of the estimate of the error.

In Section 4 Problem (2) is considered. We supposed that  $\alpha$  tends to zero: a boundary layer phenomenon appears. To obtain our error estimates, we link the parameters  $h$  and  $\alpha$ . We introduce a complex parameter  $\beta$  such that

$$\operatorname{Re}(\beta) > 0, \text{ or } (\operatorname{Re}(\beta) = 0, \text{ and } \Im(\beta) \neq 0),$$

and

$$|\beta| = o\left(\frac{1}{h}\right), \quad \text{and} \quad \frac{1}{|\beta|} = o\left(\frac{1}{h}\right).$$

We distinguish two different cases, depending on the convergence of  $|\alpha|$  to zero:  $\alpha = \beta h^q$ , for  $q \in \mathbb{N}^*$  and  $\alpha = o(h^N)$  for all  $N \in \mathbb{N}$ .

For  $q = 1$  we obtain mixed boundary conditions for the asymptotic terms of the inner potential, and as soon as  $q \geq 1$ , appropriate Dirichlet boundary conditions are obtained. We end this section by error estimates.

In Appendix, we give some useful differential geometry formulae.

**Remark 1.** *The use of the formalism of differential forms  $\delta(q_h \mathbf{d})$  could seem futile for the study of the operator  $\nabla \cdot (q_h \nabla)$ . In particular the expression of Laplace operator in local coordinates is well known. However we wanted to present this point of view to show how simple it is to write a Laplacian in curved coordinates once the metric is known.*

*Moreover once this formalism is understood for the functions (or 0-forms), it is easy to study  $\delta(q_h \mathbf{d})$  applied to 1-forms. This leads directly to the study of the operator  $\operatorname{rot}(q_h \operatorname{rot})$ , whose expression in local coordinates is less usual.*

## 1. GEOMETRY

The boundary of the domain  $\mathcal{O}$  is assumed to be smooth. The orientation of the boundary  $\partial\mathcal{O}$  is the trigonometric orientation. To simplify, we suppose that the length of  $\partial\mathcal{O}$  is equal to  $2\pi$ . We denote by  $\mathbb{T}$  the flat torus:

$$\mathbb{T} = \mathbb{R}/2\pi\mathbb{Z}.$$

Since  $\partial\mathcal{O}$  is smooth, we can parameterize it by a function  $\Psi$  of class  $\mathcal{C}^\infty$  from  $\mathbb{T}$  to  $\mathbb{R}^2$  satisfying:

$$\forall \theta \in \mathbb{T}, \quad |\Psi'(\theta)| = 1.$$

Since the boundary  $\partial\Omega_h$  of the cell is parallel to the boundary  $\partial\mathcal{O}$  of the inner domain the following identities hold:

$$\partial\mathcal{O} = \{\Psi(\theta), \theta \in \mathbb{T}\},$$

and

$$\partial\Omega_h = \{\Psi(\theta) + hn(\theta), \theta \in \mathbb{T}\}.$$

Here  $n(\theta)$  is the unitary exterior normal at  $\Psi(\theta)$  to  $\partial\mathcal{O}$ . Therefore the membrane  $\mathcal{O}_h$  is parameterized by:

$$\mathcal{O}_h = \{\Phi(\eta, \theta), (\eta, \theta) \in ]0, 1[ \times \mathbb{T}\},$$

where

$$\Phi(\eta, \theta) = \Psi(\theta) + h\eta n(\theta).$$

Denote by  $\kappa$  the curvature of  $\partial\mathcal{O}$ . Let  $h_0$  belong to  $(0, 1)$  such that:

$$(3) \quad h_0 < \frac{1}{\|\kappa\|_\infty}.$$

Thus for all  $h$  in  $[0, h_0]$ , there exists an open interval  $I$  containing  $(0, 1)$  such that  $\Phi$  is a smooth diffeomorphism from  $I \times \mathbb{R}/2\pi\mathbb{Z}$  to its image, which is a neighborhood of the membrane. The metric in  $\mathcal{O}_h$  is:

$$(4) \quad h^2 d\eta^2 + (1 + h\eta\kappa)^2 d\theta^2.$$

Thus, we use two systems of coordinates, depending on the domains  $\mathcal{O}$  and  $\mathcal{O}_h$ : in the interior domain  $\mathcal{O}$ , we use Euclidean coordinates  $(x, y)$  and in the membrane  $\mathcal{O}_h$ , we use local  $(\eta, \theta)$  coordinates with metric (4).

We translate into the language of differential forms Problem (1) and Problem (2). We refer the reader to Dubrovin, Fomenko and Novikov [8] or Flanders [11] for the definition of the exterior derivative denoted by  $d$ , the exterior product denoted by  $\text{ext}$ , the interior derivative denoted by  $\delta$  and the interior product denoted by  $\text{int}$ . In Appendix we give the formulae describing these operators in the case of a general 2D metric. Our aim, while rewriting our problems (1) and (2) is to take into account nicely the change of coordinates in the thin membrane.

Let  $\mathbb{V}$  be the 0-form on  $\Omega_h$  such that, in the Euclidean coordinates  $(x, y)$ ,  $\mathbb{V}$  is equal to  $V$ , and let  $F$  be the 0-form, which is equal to  $\phi$  on  $\partial\Omega_h$ . We denote by  $\mathbb{N}$  the 1-form corresponding to the inward unit normal on the boundary  $\Omega_h$  (see for instance Gilkey *et al.* [13] p.33):

$$\begin{aligned} \mathbb{N} &= N_x dx + N_y dy, \\ &= N_\eta d\eta. \end{aligned}$$

$\mathbb{N}^*$  is the inward unit normal 1-form. Problem (1) takes now the intrinsic form:

$$(5a) \quad \delta(q_h d\mathbb{V}) = 0, \text{ in } \Omega_h,$$

$$(5b) \quad \text{int}(\mathbb{N}^*) d\mathbb{V} = F, \text{ on } \partial\Omega_h.$$

According to Green's formula (Lemma 1.5.1 of [13]), we obtain the following transmission conditions for  $\mathbb{V}$  along  $\partial\mathcal{O}$ :

$$(5c) \quad \begin{aligned} \text{int}(\mathbb{N}^*)d\mathbb{V}|_{\partial\mathcal{O}} &= \alpha \text{int}(\mathbb{N}^*)d\mathbb{V}|_{\partial\mathcal{O}_h \setminus \partial\Omega_h}, \\ \text{ext}(\mathbb{N}^*)\mathbb{V}|_{\partial\mathcal{O}} &= \text{ext}(\mathbb{N}^*)\mathbb{V}|_{\partial\mathcal{O}_h \setminus \partial\Omega_h}. \end{aligned}$$

Similarly, denoting by  $\mathbb{U}$  the 0-form equal to  $u$  in Euclidean coordinates we rewrite Problem (2) as follows:

$$(6a) \quad \delta(\gamma_h d\mathbb{U}) = 0, \text{ in } \Omega_h,$$

$$(6b) \quad \text{int}(\mathbb{N}^*)d\mathbb{U} = F, \text{ on } \partial\Omega_h;$$

the following transmission conditions hold on  $\partial\mathcal{O}$ :

$$(6c) \quad \begin{aligned} \alpha \text{int}(\mathbb{N}^*)d\mathbb{U}|_{\partial\mathcal{O}} &= \text{int}(\mathbb{N}^*)d\mathbb{U}|_{\partial\mathcal{O}_h \setminus \partial\Omega_h}, \\ \text{ext}(\mathbb{N}^*)\mathbb{U}|_{\partial\mathcal{O}} &= \text{ext}(\mathbb{N}^*)\mathbb{U}|_{\partial\mathcal{O}_h \setminus \partial\Omega_h}. \end{aligned}$$

## 2. STATEMENT OF THE PROBLEM

In this section, we write Problem (5) and Problem (6) in local coordinates, with the help of differential forms. It is convenient to write:

$$\forall \theta \in \mathbb{T}, \quad \Phi_0(\theta) = \Phi(0, \theta), \quad \Phi_1(\theta) = \Phi(1, \theta),$$

and to denote by  $\mathbf{C}$  the cylinder:

$$\mathbf{C} = [0, 1] \times \mathbb{T}.$$

We denote by  $\mathfrak{K}$ ,  $f$  and  $\mathfrak{f}$  the following functions:

$$(7) \quad \forall (x, y) \in \partial\mathcal{O}, \quad \mathfrak{K}(x, y) = \kappa \circ \Phi_0^{-1}(x, y),$$

$$(8) \quad \forall \theta \in \mathbb{T}, \quad f(\theta) = \phi \circ \Phi_1(\theta),$$

$$(9) \quad \forall x \in \partial\mathcal{O}, \quad \mathfrak{f} = f \circ \Phi_0^{-1}(x).$$

Using the expressions of the differential operators  $d$  and  $\delta$ , which are respectively the exterior and the interior derivatives (see Appendix), applied to the metric (4), the Laplacian in the membrane is given in the local coordinates  $(\eta, \theta)$  by:

$$(10) \quad \begin{aligned} \forall (\eta, \theta) \in \mathbf{C}, \\ \Delta|_{\Phi(\eta, \theta)} = \frac{1}{h(1+h\eta\kappa)} \partial_\eta \left( \frac{1+h\eta\kappa}{h} \partial_\eta \right) + \frac{1}{1+h\eta\kappa} \partial_\theta \left( \frac{1}{1+h\eta\kappa} \partial_\theta \right). \end{aligned}$$

Moreover, for a 0-form  $z$  defined in  $\mathcal{O}_h$ , we have:

$$\begin{aligned} \text{int}(\mathbb{N}^*)dz|_{\partial\mathcal{O}} &= \frac{1}{h} \partial_\eta z|_{\eta=0}, \\ \text{int}(\mathbb{N}^*)dz|_{\partial\Omega_h} &= \frac{1}{h} \partial_\eta z|_{\eta=1}. \end{aligned}$$

Denote by

$$\begin{aligned} V^c &= V, \text{ in } \mathcal{O}, \\ V^m &= V \circ \Phi, \text{ in } \mathbf{C}, \end{aligned}$$

and by

$$\begin{aligned} u^c &= u, \text{ in } \mathcal{O}, \\ u^m &= u \circ \Phi, \text{ in } \mathbf{C}. \end{aligned}$$

We infer that Problem (5) may be rewritten as follows:

$$\begin{aligned} (11a) \quad & \Delta V^c = 0, \text{ in } \mathcal{O}, \\ (11b) \quad & \forall (\eta, \theta) \in \mathbf{C}, \quad \frac{1}{h^2} \partial_\eta \left( (1 + h\eta\kappa) \partial_\eta V^m \right) + \partial_\theta \left( \frac{1}{1 + h\eta\kappa} \partial_\theta V^m \right) = 0, \\ (11c) \quad & \partial_n V^c \circ \Phi_0 = \frac{\alpha}{h} \partial_\eta V^m \Big|_{\eta=0}, \\ (11d) \quad & V^c \circ \Phi_0 = V^m \Big|_{\eta=0}, \\ (11e) \quad & \partial_\eta V^m \Big|_{\eta=1} = hf. \\ & \int_{\partial\mathcal{O}} V \, d\sigma = 0. \end{aligned}$$

Similarly the couple  $(u^c, u^m)$  satisfies

$$\begin{aligned} (12a) \quad & \Delta u^c = 0, \text{ in } \mathcal{O}, \\ (12b) \quad & \forall (\eta, \theta) \in \mathbf{C}, \quad \frac{1}{h^2} \partial_\eta \left( (1 + h\eta\kappa) \partial_\eta u^m \right) + \partial_\theta \left( \frac{1}{1 + h\eta\kappa} \partial_\theta u^m \right) = 0, \\ (12c) \quad & \alpha \partial_n u^c \circ \Phi_0 = \frac{1}{h} \partial_\eta u^m \Big|_{\eta=0}, \\ (12d) \quad & u^c \circ \Phi_0 = u^m \Big|_{\eta=0}, \\ (12e) \quad & \partial_\eta u^m \Big|_{\eta=1} = hf, \\ & \int_{\partial\mathcal{O}} u \, d\sigma = 0. \end{aligned}$$

**Remark 2.1.** *In the following, the parameter  $\alpha$  is such that:*

$$\Re(\alpha) > 0 \text{ or } \left\{ \Re(\alpha) = 0 \text{ and } \Im(\alpha) \neq 0 \right\}.$$

*Since  $\alpha$  represents a complex permittivity it may be written (see Balanis and Constantine [3]) as follows:*

$$\alpha = \varepsilon - i\sigma/\omega,$$

*with  $\varepsilon$ ,  $\sigma$ , and  $\omega$  positive. Thus this hypothesis is always satisfied for dielectric materials.*

**Notation 2.2.** *We provide  $\mathbf{C}$  with the metric (4). The  $L^2$  norm of a 0-form  $u$  in  $\mathbf{C}$ , denoted by  $\|u\|_{\Lambda^0 L_m^2(\mathbf{C})}$ , is equal to:*

$$\begin{aligned} \|u\|_{\Lambda^0 L_m^2(\mathbf{C})} &= \left( \int_0^1 \int_0^{2\pi} h(1 + h\eta\kappa) |u(\eta, \theta)|^2 \, d\eta \, d\theta \right)^{1/2}, \\ &= \|u\|_{L^2(\mathcal{O}_h)}, \end{aligned}$$



and the  $L^2$  norm of its exterior derivative  $du$ , denoted by  $\|du\|_{\Lambda^1 L_m^2}$  is equal to

$$\begin{aligned} \|du\|_{\Lambda^1 L_m^2(\mathcal{C})} &= \left( \int_0^1 \int_0^{2\pi} \frac{1+h\eta\kappa}{h} |\partial_\eta u(\eta, \theta)|^2 + \frac{h}{1+h\eta\kappa} |\partial_\theta u(\eta, \theta)|^2 d\eta d\theta \right)^{1/2}, \\ &= \|\text{grad } u\|_{L^2(\mathcal{O}_h)}. \end{aligned}$$

To simplify our notations, for a 0-form  $u$  defined on  $\mathcal{C}$ , we define by  $\|u\|_{H_\mathfrak{q}^1(\mathcal{C})}$  the following quantity

$$\|u\|_{H_\mathfrak{q}^1(\mathcal{C})} = \|u\|_{\Lambda^0 L_m^2(\mathcal{C})} + \|du\|_{\Lambda^1 L_m^2(\mathcal{C})},$$

when the above integrals are well-defined. Observe that for a function  $u \in H^1(\mathcal{O}_h)$ , we have:

$$\|u\|_{H^1(\mathcal{O}_h)} = \|u \circ \Phi\|_{H_\mathfrak{q}^1(\mathcal{C})}.$$

**Remark 2.3** (Poincaré inequality in the thin layer). *Let  $z$  belong to  $H_\mathfrak{q}^1(\mathcal{C})$ , such that*

$$(13) \quad \int_0^{2\pi} z(0, \theta) d\theta = 0.$$

Then, there exists an  $h$ -independent constant  $C_\mathcal{O}$  such that

$$(14) \quad \|z\|_{\Lambda^0 L_\mathfrak{q}^2(\mathcal{C})} \leq C_\mathcal{O} \|dz\|_{\Lambda^1 L_\mathfrak{q}^2(\mathcal{C})}.$$

We prove (14) using Fourier analysis. According to the definition (3) of  $h_0$  there exists two constants  $C_\mathcal{O}$  and  $c_\mathcal{O}$  depending on the domain  $\mathcal{O}$  such that the following inequalities hold:

$$(15a) \quad \|z\|_{\Lambda^0 L_\mathfrak{q}^2(D)}^2 \leq C_\mathcal{O} h \int_0^1 \int_0^{2\pi} |z(\eta, \theta)|^2 d\theta d\eta,$$

$$(15b) \quad \|dz\|_{\Lambda^1 L_\mathfrak{q}^2(D)}^2 \geq c_\mathcal{O} \left( \int_0^1 \int_0^{2\pi} \frac{|\partial_\eta z(\eta, \theta)|^2}{h} + h |\partial_\theta z|^2 d\theta d\eta \right).$$

For  $k \in \mathbb{Z}$ , we denote by  $\widehat{z}_k$  the  $k^{\text{th}}$ -Fourier coefficient (with respect to  $\theta$ ) of  $z$ :

$$\widehat{z}_k = \frac{1}{2\pi} \int_0^{2\pi} z(\theta) e^{-ik\theta} d\theta.$$

Since  $\left(\widehat{\partial_\theta z}\right)_k = ik\widehat{z}_k$ , we infer:

$$\forall k \neq 0, \quad \int_0^1 |\widehat{z}_k(\eta)|^2 d\eta \leq \int_0^1 \left| \left(\widehat{\partial_\theta z}\right)_k(\eta) \right|^2 d\eta.$$

According to gauge condition (13), we have:

$$\widehat{z}_0(0) = 0,$$

thus, using the equality

$$\widehat{z}_0(\eta) = \int_0^\eta \left(\widehat{\partial_\eta z}\right)_0(s) ds,$$

we infer

$$\int_0^1 |\widehat{z}_0(\eta)|^2 d\eta \leq \int_0^1 \left| \left(\widehat{\partial_\eta z}\right)_0(\eta) \right|^2 d\eta.$$

Therefore,

$$\sum_{k \in \mathbb{Z}} \int_0^1 |\widehat{z}_k(\eta, \theta)|^2 d\eta \leq \sum_{k \in \mathbb{Z}} \left\{ \int_0^1 \left| \left( \widehat{\partial_\theta z} \right)_k(\eta) \right|^2 d\eta + \int_0^1 \left| \left( \widehat{\partial_\eta z} \right)_k(\eta) \right|^2 d\eta \right\}.$$

We end the proof of (14) by using Parseval inequality and inequalities (15).

### 3. ASYMPTOTIC EXPANSION OF THE STEADY STATE POTENTIAL FOR AN INSULATING MEMBRANE

We derive asymptotic expansions with respect to  $h$  of the potentials  $(V^c, V^m)$  solution to Problem (11). The membrane is insulating since the modulus of  $\alpha$  is supposed to be smaller than 1. However, our results are still valid if  $|\alpha|$  is bounded by a constant  $C_0$  greater than 1. We emphasize that the following results are valid for  $\alpha$  tending to zero.

**3.1. Formal asymptotic expansion.** We write the following ansatz:

$$(16a) \quad V^c = V_0^c + hV_1^c + h^2V_2^c + \dots,$$

$$(16b) \quad V^m = V_0^m + hV_1^m + h^2V_2^m + \dots.$$

We multiply (11b) by  $h^2(1 + h\eta\kappa)^2$  and we order the powers of  $h$  to obtain:

$$(17) \quad \begin{aligned} & \forall (\eta, \theta) \in [0, 1] \times \mathbb{T}, \\ & \partial_\eta^2 V^m + h\kappa \{3\eta \partial_\eta^2 V^m + \partial_\eta V^m\} + h^2 \{3\eta^2 \kappa^2 \partial_\eta^2 V^m + 2\eta \kappa^2 \partial_\eta V^m + \partial_\theta^2 V^m\} \\ & + h^3 \{\eta^3 \kappa^3 \partial_\eta^2 V^m + \eta^2 \kappa^3 \partial_\eta V^m + \eta \kappa \partial_\theta^2 V^m - \eta \kappa' \partial_\theta V^m\} = 0 \end{aligned}$$

We are now ready to derive formally the terms of the asymptotic expansions of  $V^c$  and  $V^m$  by identifying the terms of the same power in  $h$ .

Recall that for  $(m, n)$  in  $\mathbb{N}^2$ ,  $\delta_{m,n}$  is Kronecker symbol equal to 1 if  $m = n$  and to 0 if  $m \neq n$ . By identifying the powers of  $h$ , we infer that for  $l \in \mathbb{N}$ ,  $V_l^c$  and  $V_l^m$  satisfy the following equations:

$$(18a) \quad \Delta V_l^c = 0, \text{ in } \mathcal{O},$$

for all  $(\eta, \theta) \in \mathcal{C}$ ,

$$(18b) \quad \begin{aligned} \partial_\eta^2 V_l^m = & - \left\{ \kappa \{3\eta \partial_\eta^2 V_{l-1}^m + \partial_\eta V_{l-1}^m\} \right. \\ & + 3\eta^2 \kappa^2 \partial_\eta^2 V_{l-2}^m + 2\eta \kappa^2 \partial_\eta V_{l-2}^m + \partial_\theta^2 V_{l-2}^m \\ & \left. + \eta^3 \kappa^3 \partial_\eta^2 V_{l-3}^m + \eta^2 \kappa^3 \partial_\eta V_{l-3}^m + \eta \kappa \partial_\theta^2 V_{l-3}^m - \eta \kappa' \partial_\theta V_{l-3}^m \right\}, \end{aligned}$$

with transmission conditions

$$(18c) \quad \partial_n V_l^c \circ \Phi_0 = \alpha \partial_\eta V_{l+1}^m \Big|_{\eta=0},$$

$$(18d) \quad V_l^c \circ \Phi_0 = V_l^m \Big|_{\eta=0},$$

with boundary condition

$$(18e) \quad \partial_\eta V_l^m \Big|_{\eta=1} = \delta_{l,1} f,$$

and with gauge condition

$$(18f) \quad \int_{\partial\mathcal{O}} V_l^c d\sigma = 0.$$

In equations (18), we have implicitly imposed

$$(19) \quad \begin{cases} V_l^c = 0, & \text{if } l \leq -1, \\ V_l^m = 0, & \text{if } l \leq -1. \end{cases}$$

The next lemma ensures that for each non null integer  $N$ , the functions  $V_N^c$  and  $V_N^m$  are entirely determined if the boundary condition  $\phi$  is enough regular.

**Notation 3.1.** For  $s \in \mathbb{R}$ , we denote by  $\mathcal{C}^\infty([0, 1]; H^s(\mathbb{T}))$  the space of functions  $u$  defined for  $(\eta, \theta) \in [0, 1] \times \mathbb{T}$ , such that for almost all  $\theta \in \mathbb{T}$ ,  $u(\cdot, \theta)$  belongs to  $\mathcal{C}^\infty([0, 1])$ , and such that for all  $\eta \in [0, 1]$ ,  $u(\eta, \cdot)$  belongs to  $H^s(\mathbb{T})$ .

**Lemma 3.2.** We suppose that  $\partial\mathcal{O}$  is smooth.

For  $N \in \mathbb{N}$  and  $p \geq 0$  we suppose that  $\phi$  belongs to  $H^{N+p-1/2}(\partial\Omega_h)$  and let  $|\alpha| \leq 1$ .

Then the functions  $V_0^m, \dots, V_N^m$  and  $V_0^c, \dots, V_N^c$  are uniquely determined and they belong to the respective functional spaces:

$$\forall k = 0, \dots, N,$$

$$(20a) \quad V_k^m \in \mathcal{C}^\infty([0, 1]; H^{N+p-k+1/2}(\mathbb{T})),$$

$$(20b) \quad V_k^c \in H^{N+p-k+1}(\mathcal{O}).$$

Moreover, there exists a constant  $C_{N,\mathcal{O},p}$  such that:

$$\forall k = 0, \dots, N,$$

$$(21a) \quad \sup_{\eta \in [0,1]} \|V_k^m(\eta, \cdot)\|_{H^{N+p-k+1/2}(\mathbb{T})} \leq C_{N,\mathcal{O},p} \|f\|_{H^{N+p-1/2}(\partial\mathcal{O})},$$

$$(21b) \quad \|V_k^c\|_{H^{N+p-k+1}(\mathcal{O})} \leq |\alpha| C_{N,\mathcal{O},p} \|f\|_{H^{N+p-1/2}(\partial\mathcal{O})}.$$

**Remark 3.3.** To simplify, we suppose that  $|\alpha| \leq 1$ , but the same result may be obtained if there exists  $C_0 > 1$  such that  $|\alpha| \leq C_0$ . In this case, the constant  $C_{N,\mathcal{O},p}$  would also depends on  $C_0$ .

*Proof.* Since  $\partial\mathcal{O}$  is smooth and since  $\phi$  belongs to  $H^{N+p-1/2}(\partial\Omega_h)$ , for  $N \geq 0$  and  $p \geq 0$ , then the functions  $f$  and  $\mathbf{f}$  defined by (8) and by (9) belong respectively to  $H^{N+p-1/2}(\mathbb{T})$  and to  $H^{N+p-1/2}(\partial\mathcal{O})$ . We prove this lemma by recursive process.

- $N = 0$ . Let  $p \geq 0$  and let  $\phi$  belong to  $H^{p-1/2}(\partial\Omega_h)$ .

Thus  $f$  and  $\mathbf{f}$  belong respectively to  $H^{p-1/2}(\mathbb{T})$  and  $H^{p-1/2}(\partial\mathcal{O})$ . Using (18b) and (18e), we infer:

$$(22) \quad \begin{cases} \partial_\eta^2 V_0^m = 0, \\ \partial_\eta V_0^m|_{\eta=1} = 0, \end{cases}$$

hence,  $\partial_\eta V_0^m = 0$ . According to (18b) and to (18e), we straight infer

$$\partial_\eta V_1^m = f.$$

Therefore by (18a) and (18c) the function  $V_0^c$  satisfies the following Laplace problem:

$$(23a) \quad \Delta V_0^c = 0,$$

$$(23b) \quad \partial_n V_0^c|_{\partial\mathcal{O}} = \alpha f,$$

with gauge condition

$$(23c) \quad \int_{\partial\mathcal{O}} V_0^c d\sigma = 0.$$

According to (18d), we infer

$$(24) \quad V_0^m = V_0^c \circ \Phi_0,$$

hence  $V_0^c$  and  $V_0^m$  are entirely determined and they belong to the following spaces:

$$\begin{aligned} V_0^m &\in \mathcal{C}^\infty([0, 1]; H^{p+1/2}(\mathbb{T})), \\ V_0^c &\in H^{p+1}(\mathcal{O}). \end{aligned}$$

Observe also that there exists a constant  $C_{\mathcal{O},p}$  such that

$$\begin{aligned} \sup_{\eta \in [0,1]} \|V_0^m(\eta, \cdot)\|_{H^{p+1/2}(\mathbb{T})} &\leq C_{\mathcal{O},p} \|f\|_{H^{p-1/2}(\partial\mathcal{O})}, \\ \|V_0^c\|_{H^{p+1}(\mathcal{O})} &\leq |\alpha| C_{\mathcal{O},p} \|f\|_{H^{p-1/2}(\partial\mathcal{O})}. \end{aligned}$$

• Induction.

Let  $N \geq 0$ . Suppose that for all  $p \geq 0$ , for all  $\phi \in H^{N+p-1/2}(\partial\Omega_h)$  and for  $M = 0, \dots, N$  the functions  $V_M^c$  and  $V_M^m$  are known. Suppose that they belong respectively to  $H^{N+p-M+1}(\mathcal{O})$  and to  $V_M^m \in \mathcal{C}^\infty([0, 1]; H^{N+p-M+1/2}(\mathbb{T}))$  and that estimates (21) hold.

Let  $\phi$  belong to  $H^{N+p+1/2}(\partial\Omega_h)$ . Therefore, for  $M = 0, \dots, N$ , the functions  $V_M^c$  and  $V_M^m$  are known, they belong respectively to  $H^{N+p-M+2}(\mathcal{O})$  and to  $V_M^m \in \mathcal{C}^\infty([0, 1]; H^{N+p-M+3/2}(\mathbb{T}))$  and the following estimates hold:

$$\begin{aligned} \forall M = 0, \dots, N, \\ \sup_{\eta \in [0,1]} \|V_M^m(\eta, \cdot)\|_{H^{N+p-M+3/2}(\mathbb{T})} &\leq C_{N,\mathcal{O},p} \|f\|_{H^{N+p+1/2}(\partial\mathcal{O})}, \\ \|V_M^c\|_{H^{N+p-M+2}(\mathcal{O})} &\leq |\alpha| C_{N,\mathcal{O},p} \|f\|_{H^{N+p+1/2}(\partial\mathcal{O})}. \end{aligned}$$

We are going to build  $V_{N+1}^c$  and  $V_{N+1}^m$ . From (18b) and (18e), we infer, for all  $(\eta, \theta) \in \mathcal{C}$ ,

$$\begin{aligned} \partial_\eta^2 V_{N+1}^m &= - \left\{ \kappa \{ 3\eta \partial_\eta^2 V_N^m + \partial_\eta V_N^m \} \right. \\ &\quad + 3\eta^2 \kappa^2 \partial_\eta^2 V_{N-1}^m + 2\eta \kappa^2 \partial_\eta V_{N-1}^m + \partial_\theta^2 V_{N-1}^m \\ &\quad \left. + \eta^3 \kappa^3 \partial_\eta^2 V_{N-2}^m + \eta^2 \kappa^3 \partial_\eta V_{N-2}^m + \eta \kappa \partial_\theta^2 V_{N-2}^m - \eta \kappa' \partial_\theta V_{N-2}^m \right\}, \\ \partial_\eta V_{N+1}^m|_{\eta=1} &= 0. \end{aligned}$$

Recall that we use convention (19). Since we have supposed that  $V_M^m$  is known for  $M \leq N$  and belongs to  $\mathcal{C}^\infty([0, 1]; H^{N+1+p-M-1/2}(\mathbb{T}))$ , we infer that:

$$\begin{aligned} & \forall (s, \theta) \in \mathcal{C}, \\ (25) \quad & \partial_\eta V_{N+1}^m(s, \cdot) = \int_s^1 \left\{ \kappa \{ 3\eta \partial_\eta^2 V_N^m + \partial_\eta V_N^m \} \right. \\ & \quad + 3\eta^2 \kappa^2 \partial_\eta^2 V_{N-1}^m + 2\eta \kappa^2 \partial_\eta V_{N-1}^m + \partial_\theta^2 V_{N-1}^m \\ & \quad \left. + \eta^3 \kappa^3 \partial_\eta^2 V_{N-2}^m + \eta^2 \kappa^3 \partial_\eta V_{N-2}^m + \eta \kappa \partial_\theta^2 V_{N-2}^m - \eta \kappa' \partial_\theta V_{N-2}^m \right\} d\eta, \end{aligned}$$

is entirely determined and belongs to  $\mathcal{C}^\infty([0, 1]; H^{p+1/2}(\mathbb{T}))$ . Moreover, since  $\partial_\eta V_{N+1}^m$  is known, we infer exactly by the same way that  $\partial_\eta V_{N+2}^m$  is also determined. Actually, it is equal to

$$\begin{aligned} & \forall (s, \theta) \in \mathcal{C}, \\ & \partial_\eta V_{N+2}^m(s, \cdot) = \int_s^1 \left\{ \kappa \{ 3\eta \partial_\eta^2 V_{N+1}^m + \partial_\eta V_{N+1}^m \} \right. \\ & \quad + 3\eta^2 \kappa^2 \partial_\eta^2 V_N^m + 2\eta \kappa^2 \partial_\eta V_N^m + \partial_\theta^2 V_N^m \\ & \quad \left. + \eta^3 \kappa^3 \partial_\eta^2 V_{N-1}^m + \eta^2 \kappa^3 \partial_\eta V_{N-1}^m + \eta \kappa \partial_\theta^2 V_{N-1}^m - \eta \kappa' \partial_\theta V_{N-1}^m \right\} d\eta, \end{aligned}$$

and it belongs to  $\mathcal{C}^\infty([0, 1]; H^{p+1/2}(\mathbb{T}))$ . According to (18c), the function  $V_{N+1}^c$  is then uniquely determined by

$$(26a) \quad \Delta V_{N+1}^c = 0,$$

$$(26b) \quad \partial_n V_{N+1}^c|_{\partial\mathcal{O}} = \alpha \partial_\eta V_{N+2}^m \circ \Phi_0^{-1},$$

with gauge condition

$$(26c) \quad \int_{\partial\mathcal{O}} V_{N+1}^c d\sigma = 0.$$

Moreover, it belongs to  $H^{p+1}(\mathcal{O})$ . Transmission condition (18d) implies the following expression of  $V_{N+1}^m$ :

$$\forall s \in (0, 1), \quad V_{N+1}^m(s, \cdot) = \int_0^s \partial_\eta V_{N+1}^m(\eta, \cdot) d\eta + V_{N+1}^c \circ \Phi_0,$$

where  $\partial_\eta V_{N+1}^m$  is given by (25) and belongs to  $\mathcal{C}^\infty([0, 1]; H^{p+1/2}(\mathbb{T}))$ . We infer also that there exists  $C_{N+1, \mathcal{O}, p} > 0$  such that

$$\begin{aligned} \sup_{\eta \in [0, 1]} \|V_{N+1}^m(\eta, \cdot)\|_{H^{p+1/2}(\mathbb{T})} & \leq C_{N+1, \mathcal{O}, p} \|f\|_{H^{N+p+1/2}(\partial\mathcal{O})}, \\ \|V_{N+1}^c\|_{H^{p+1}(\mathcal{O})} & \leq |\alpha| C_{N+1, \mathcal{O}, p} \|f\|_{H^{N+p+1/2}(\partial\mathcal{O})}, \end{aligned}$$

hence the lemma.  $\square$

**3.2. Error Estimates.** In the previous section, we have proved that for each non null integer  $N$ , the functions  $V_N^c$  and  $V_N^m$  are well defined by equations (18); more precisely, using convention (19) we get

$$\begin{aligned}\Delta V_0^c &= 0, \\ \partial_n V_0^c|_{\partial\mathcal{O}} &= \alpha f,\end{aligned}$$

with gauge condition

$$\begin{aligned}\int_{\partial\mathcal{O}} V_0^c d\sigma &= 0; \\ \forall(\eta, \theta) \in \mathcal{C}, \quad V_0^m &= V_0^c \circ \Phi_0.\end{aligned}$$

For  $N \geq 0$ , we have

$$\begin{aligned}\forall(s, \theta) \in \mathcal{C}, \\ \partial_\eta V_N^m(s, \cdot) &= \int_s^1 \left\{ \kappa \left\{ 3\eta \partial_\eta^2 V_{N-1}^m + \partial_\eta V_{N-1}^m \right\} \right. \\ &\quad + 3\eta^2 \kappa^2 \partial_\eta^2 V_{N-2}^m + 2\eta \kappa^2 \partial_\eta V_{N-2}^m + \partial_\theta^2 V_{N-2}^m \\ &\quad \left. + \eta^3 \kappa^3 \partial_\eta^2 V_{N-3}^m + \eta^2 \kappa^3 \partial_\eta V_{N-3}^m + \eta \kappa \partial_\theta^2 V_{N-3}^m - \eta \kappa' \partial_\theta V_{N-3}^m \right\} d\eta, \\ \forall(s, \theta) \in \mathcal{C}, \\ \partial_\eta V_{N+1}^m(s, \cdot) &= \int_s^1 \left\{ \kappa \left\{ 3\eta \partial_\eta^2 V_N^m + \partial_\eta V_N^m \right\} \right. \\ &\quad + 3\eta^2 \kappa^2 \partial_\eta^2 V_{N-1}^m + 2\eta \kappa^2 \partial_\eta V_{N-1}^m + \partial_\theta^2 V_{N-1}^m \\ &\quad \left. + \eta^3 \kappa^3 \partial_\eta^2 V_{N-2}^m + \eta^2 \kappa^3 \partial_\eta V_{N-2}^m + \eta \kappa \partial_\theta^2 V_{N-2}^m - \eta \kappa' \partial_\theta V_{N-2}^m \right\} d\eta,\end{aligned}$$

the functions  $V_N^c$  and  $V_N^m$  are then defined by

$$\begin{aligned}\Delta V_N^c &= 0, \\ \partial_n V_N^c|_{\partial\mathcal{O}} &= \alpha \partial_\eta V_{N+1}^m \circ \Phi_0^{-1},\end{aligned}$$

with gauge condition

$$\begin{aligned}\int_{\partial\mathcal{O}} V_N^c d\sigma &= 0, \\ \forall s \in (0, 1), \quad V_{N+1}^m(s, \cdot) &= \int_0^s \partial_\eta V_{N+1}^m(\eta, \cdot) d\eta + V_{N+1}^c \circ \Phi_0.\end{aligned}$$

It remains to prove error estimates.

**Theorem 3.4.** *Let  $h_0$  be defined in (3) and  $h$  belong to  $(0, h_0)$ . The complex parameter  $\alpha$  satisfies*

$$(27) \quad |\alpha| \leq 1,$$

$$(28) \quad \Re(\alpha) > 0 \quad \text{or} \quad \left\{ \Re(\alpha) = 0 \text{ and } \Im(\alpha) \neq 0 \right\}.$$

Let  $N \in \mathbb{N}$  and  $\phi$  belong to  $H^{N+3/2}(\partial\Omega_h)$ . We suppose that  $\partial\mathcal{O}$  is smooth. The function  $f$  is defined by (9). We denote by  $V$  the solution to Problem (1). Let  $R_N^c$

and  $R_N^m$  be the functions defined by:

$$\begin{cases} R_N^c = V - \sum_{k=0}^N V_k^c h^k, \text{ in } \mathcal{O}, \\ R_N^m = V \circ \Phi - \sum_{k=0}^N V_k^m h^k, \text{ in } \mathbf{C}. \end{cases}$$

Then, there exists a constant  $C_{\mathcal{O},N} > 0$  depending only on the domain  $\mathcal{O}$  and on  $N$  such that

$$(29a) \quad \|R_N^c\|_{H^1(\mathcal{O})} \leq C_{\mathcal{O},N} \|f\|_{H^{N+3/2}(\partial\mathcal{O})} |\alpha| h^{N+1/2},$$

$$(29b) \quad \|R_N^m\|_{H_{\mathfrak{g}}^1(\mathbf{C})} \leq C_{\mathcal{O},N} \|f\|_{H^{N+3/2}(\partial\mathcal{O})} h^{N+1/2}.$$

Moreover, if  $\phi$  belongs to  $H^{N+5/2}(\partial\Omega_h)$ , then we have

$$(30a) \quad \|R_N^c\|_{H^1(\mathcal{O})} \leq C_{\mathcal{O},N} \|f\|_{H^{N+5/2}(\partial\mathcal{O})} |\alpha| h^{N+1},$$

$$(30b) \quad \|R_N^m\|_{H_{\mathfrak{g}}^1(\mathbf{C})} \leq C_{\mathcal{O},N} \|f\|_{H^{N+5/2}(\partial\mathcal{O})} h^{N+1/2}.$$

**Remark 3.5.** If  $|\alpha|$  is greater than 1 but bounded by  $C_0$ , the same results hold: the constant  $C_{\mathcal{O},N}$  would also depends on  $C_0$ .

*Proof of Theorem 3.4.* Since  $\phi$  belongs to  $H^{N+3/2}(\partial\Omega_h)$ , according to the previous lemma, the couples of functions  $(R_N^c, R_N^m)$  and  $(R_{N+1}^c, R_{N+1}^m)$  are well defined and belong to  $H^1(\mathcal{O}) \times H_{\mathfrak{g}}^1(\mathbf{C})$ . The Sobolev space  $H_{\mathfrak{g}}^1(\mathbf{C})$  is defined in Notation 2.2.

Denote by  $g_N$  the following function defined on  $\mathbf{C}$ :

$$(31) \quad \begin{aligned} g_N = & \kappa (3\eta \partial_{\eta}^2 V_N^m + \partial_{\eta} V_N^m) + 3\eta^2 \kappa^2 \partial_{\eta}^2 V_{N-1}^m + 2\eta \kappa^2 \partial_{\eta} V_{N-1}^m + \partial_{\theta}^2 V_{N-1}^m \\ & + \eta^3 \kappa^3 \partial_{\eta}^2 V_{N-2}^m + \eta^2 \kappa^3 \partial_{\eta} V_{N-2}^m + \eta \kappa \partial_{\theta}^2 V_{N-2}^m - \eta \kappa' \partial_{\theta} V_{N-2}^m \\ & + h \left( 3\eta^2 \kappa^2 \partial_{\eta}^2 V_N^m + 2\eta \kappa^2 \partial_{\eta} V_N^m + \partial_{\theta}^2 V_N^m \right. \\ & \left. + \eta^3 \kappa^3 \partial_{\eta}^2 V_{N-1}^m + \eta^2 \kappa^3 \partial_{\eta} V_{N-1}^m + \eta \kappa \partial_{\theta}^2 V_{N-1}^m - \eta \kappa' \partial_{\theta} V_{N-1}^m \right) \\ & + h^2 \left( \eta^3 \kappa^3 \partial_{\eta}^2 V_N^m + \eta^2 \kappa^3 \partial_{\eta} V_N^m + \eta \kappa \partial_{\theta}^2 V_N^m - \eta \kappa' \partial_{\theta} V_N^m \right) \end{aligned}$$

According to the previous lemma and since  $\phi$  belongs to  $H^{N+1/2}(\partial\Omega_h)$ , the above function  $g_N$  belongs to  $\mathcal{C}^{\infty}([0, 1]; H^{-1/2}(\mathbb{T}))$  and the function  $\partial_{\eta} V_N^m$  belongs to  $\mathcal{C}^{\infty}([0, 1]; H^{3/2}(\mathbb{T}))$ . Moreover, there exists a constant  $C_{N,\mathcal{O}}$  such that

$$(32) \quad \begin{cases} \sup_{\eta \in [0,1]} \|g_N(\eta, \cdot)\|_{H^{-1/2}(\mathbb{T})} \leq C_{N,\mathcal{O}} \|f\|_{H^{N+1/2}(\mathbb{T})}, \\ \sup_{\eta \in [0,1]} \|\partial_{\eta} V_N^m(\eta, \cdot)\|_{H^{3/2}(\mathbb{T})} \leq C_{N,\mathcal{O}} \|f\|_{H^{N+1/2}(\mathbb{T})}. \end{cases}$$

The functions  $R_N^c$  and  $R_N^m$  satisfy the following problem:

$$\begin{aligned} \Delta R_N^c &= 0, \text{ in } \mathcal{O}, \\ \partial_{\eta} \left( \frac{1+h\eta\kappa}{h} \partial_{\eta} R_N^m \right) + \partial_{\theta} \left( \frac{h}{1+h\eta\kappa} \partial_{\theta} R_N^m \right) &= \frac{-h^N}{(1+h\eta\kappa)} g_N, \end{aligned}$$

with transmission conditions:

$$\begin{aligned}\partial_n R_N^c \circ \Phi_0 &= \frac{\alpha}{h} \left( \partial_\eta R_N^m|_{\eta=0} + h^{N+1} \partial_\eta V_N^m|_{\eta=0} \right), \\ R_N^c \circ \Phi_0 &= R_N^m|_{\eta=0},\end{aligned}$$

with boundary condition

$$\partial_\eta R_N^m|_{\eta=1} = 0,$$

and with gauge condition

$$\int_{\partial\mathcal{O}} R_N^c \, d\sigma = 0.$$

By multiplying the above equality by  $\overline{R_N}$  and by integration by parts, we infer that:

$$\begin{aligned}(33) \quad \|dR_N^c\|_{\Lambda^1 L^2(\mathcal{O})}^2 + \alpha \|dR_N^m\|_{\Lambda^1 L^2_{\mathfrak{g}}(\mathcal{C})}^2 &= -\alpha h^N \int_{\mathcal{C}} g_N(\eta, \theta) \overline{R_N^m}(\eta, \theta) \, d\eta \, d\theta \\ &+ \alpha h^{N+1} \int_{\mathbb{T}} \partial_\eta V_N^m|_{\eta=0} \overline{R_N^m}|_{\eta=0} \, d\theta \\ &- \alpha h^{N+1} \int_{\mathbb{T}} \kappa \partial_\eta V_N^m|_{\eta=1} \overline{R_N^m}|_{\eta=1} \, d\theta.\end{aligned}$$

By hypothesis (28), and using by Cauchy-Schwarz inequality and estimates (32), we infer that there exists a constant  $C_{\mathcal{O},N} > 0$  such that

$$\Re(\alpha) \|dR_N^m\|_{\Lambda^1 L^2_{\mathfrak{g}}(\mathcal{C})}^2 \leq |\alpha| C_{\mathcal{O},N} h^{N-1/2} \|R_N^m\|_{H^1_{\mathfrak{g}}(\mathcal{C})} \|f\|_{H^{N+1/2}(\mathbb{T})} \|R_N^m\|_{H^1_{\mathfrak{g}}(\mathcal{C})},$$

and

$$|\Im(\alpha)| \|dR_N^m\|_{\Lambda^1 L^2_{\mathfrak{g}}(\mathcal{C})}^2 \leq |\alpha| C_{\mathcal{O},N} h^{N-1/2} \|R_N^m\|_{H^1_{\mathfrak{g}}(\mathcal{C})} \|f\|_{H^{N+1/2}(\mathbb{T})} \|R_N^m\|_{H^1_{\mathfrak{g}}(\mathcal{C})},$$

hence

$$\|dR_N^m\|_{\Lambda^1 L^2_{\mathfrak{g}}(\mathcal{C})}^2 \leq C_{\mathcal{O},N} h^{N-1/2} \|f\|_{H^{N+1/2}(\mathbb{T})} \|R_N^m\|_{H^1_{\mathfrak{g}}(\mathcal{C})}.$$

Since  $\int_{\mathbb{T}} R_N^m|_{\eta=0} \, d\theta = 0$ , by Poincaré inequality (14), there exists a strictly positive constant  $C_{\mathcal{O}}$ , which does not depend on  $h$  such that

$$\|R_N^m\|_{\Lambda^0 L^2_{\mathfrak{g}}(\mathcal{C})} \leq C_{\mathcal{O}} \|dR_N^m\|_{\Lambda^1 L^2_{\mathfrak{g}}(\mathcal{C})},$$

hence

$$\|R_N^m\|_{H^1_{\mathfrak{g}}(\mathcal{C})} \leq C_{\mathcal{O},N} \|f\|_{H^{N+1/2}(\mathbb{T})} h^{N-1/2},$$

and therefore we deduce directly from the above estimate and from (33),

$$\|R_N^c\|_{H^1(\mathcal{O})} \leq C_{\mathcal{O},N} \|f\|_{H^{N+1/2}(\mathbb{T})} |\alpha| h^{N-1/2}.$$

The above estimate holds for  $\phi \in H^{N+1/2}(\partial\Omega_h)$ . Since  $\phi$  belongs to  $H^{N+3/2}(\partial\Omega_h)$ , we obtain the same result by replacing  $N$  by  $N+1$ :

$$(34) \quad \begin{cases} \|R_{N+1}^m\|_{H^1_{\mathfrak{g}}(\mathcal{C})} \leq C_{\mathcal{O},N+1} \|f\|_{H^{N+3/2}(\mathbb{T})} h^{N+1/2}, \\ \|R_{N+1}^c\|_{H^1(\mathcal{O})} \leq C_{\mathcal{O},N+1} \|f\|_{H^{N+3/2}(\mathbb{T})} |\alpha| h^{N+1/2}. \end{cases}$$



According to the previous lemma, the functions  $V_{N+1}^c$  and  $V_{N+1}^m$  are well-defined and there exists a constant  $C_{N,\mathcal{O}}$  such that:

$$\begin{aligned}\|V_{N+1}^c\|_{H^1(\mathcal{O})} &\leq |\alpha|C_{N,\mathcal{O}}\|f\|_{H^{N+3/2}(\partial\mathcal{O})}, \\ \|V_{N+1}^m\|_{H_{\frac{1}{6}}^1(\mathcal{C})} &\leq \frac{C_{N,\mathcal{O}}}{\sqrt{h}}\|f\|_{H^{N+3/2}(\partial\mathcal{O})}.\end{aligned}$$

Writing

$$R_N^c = R_{N+1}^c + V_{N+1}^c h^{N+1},$$

and

$$R_N^m = R_{N+1}^m + V_{N+1}^m h^{N+1},$$

we infer that

$$\|R_N^c\|_{H^1(\mathcal{O})} \leq C_{\mathcal{O},N}\|f\|_{H^{N+3/2}(\mathbb{T})}|\alpha|h^{N+1/2},$$

and

$$\|R_N^m\|_{H_{\frac{1}{6}}^1(\mathcal{C})} \leq C_{\mathcal{O},N}\|f\|_{H^{N+3/2}(\mathbb{T})}h^{N+1/2}.$$

If  $\phi$  belongs to  $H^{N+5/2}(\partial\Omega_h)$ , we write

$$R_N^c = R_{N+2}^c + V_{N+1}^c h^{N+1} + V_{N+2}^c h^{N+2},$$

and

$$R_N^m = R_{N+1}^m + V_{N+1}^m h^{N+1} + V_{N+2}^m h^{N+2},$$

to obtain estimates (30), hence the theorem.  $\square$

**Remark 3.6** (The case of an insulating cytoplasm). *Consider Problem (2):*

$$\begin{aligned}\operatorname{div}(\gamma_h \operatorname{grad} u) &= 0 \text{ in } \Omega_h, \\ \frac{\partial u}{\partial n} &= \phi \text{ on } \partial\Omega_h, \\ \int_{\partial\mathcal{O}} u \, d\sigma &= 0.\end{aligned}$$

If the inner domain is perfectly insulating (i.e. if  $\gamma_h$  vanishes in  $\mathcal{O}$ ), the steady state potential in the membrane satisfies:

$$\frac{1}{h^2}\partial_\eta((1+h\eta\kappa)\partial_\eta u^m) + \partial_\theta\left(\frac{1}{1+h\eta\kappa}\partial_\theta u^m\right) = 0, \text{ in } \mathcal{C},$$

with the following boundary conditions:

$$\partial_\eta u^m|_{\eta=0} = 0, \quad \partial_\eta u^m|_{\eta=1} = hf,$$

and with gauge condition

$$\int_0^{2\pi} u^m|_{\eta=0} d\theta = 0.$$

By identifying the terms of the same power of  $h$  we would obtain:  $u_0^m = 0$ , and  $u_1^m$  would satisfy:

$$\begin{aligned}\partial_\eta^2 u_1^m &= 0, \text{ in } \mathcal{C}, \\ \partial_\eta u_1^m|_{\eta=0} &= 0, \quad \partial_\eta u_1^m|_{\eta=1} = f, \\ \int_0^{2\pi} u_1^m|_{\eta=0} d\theta &= 0,\end{aligned}$$

which is a non sense as soon as  $f \neq 0$ . Our ansatz (16) fails. Actually, the asymptotic expansion of  $u^m$  begins at the order  $-1$ : a boundary layer phenomenon appears. This is described in the next Section.

#### 4. ASYMPTOTIC EXPANSION OF THE STEADY STATE POTENTIAL FOR AN INSULATING INNER DOMAIN

Consider now the solution  $u$  to Problem (2). In this section, we suppose

$$(35a) \quad |\alpha| \text{ tends to zero,}$$

$$(35b) \quad \Re(\alpha) > 0 \text{ or } \left\{ \Re(\alpha) = 0 \text{ and } \Im(\alpha) \neq 0 \right\}.$$

Thus the inner domain is insulating. Let  $\beta$  be a complex parameter satisfying:

$$\Re(\beta) > 0, \text{ or } (\Re(\beta) = 0, \text{ and } \Im(\beta) \neq 0).$$

The modulus of  $\beta$  may tend to infinity, or to zero but it must satisfy:

$$|\beta| = o\left(\frac{1}{h}\right), \quad \text{and} \quad \frac{1}{|\beta|} = o\left(\frac{1}{h}\right).$$

We suppose that  $u$  may be written as follows:

$$u = \frac{1}{h}u_{-1} + u_0 + hu_1 + \dots.$$

We denote by  $u^c$  and  $u^m \circ \Phi^{-1}$  the respective restrictions of  $u$  to  $\mathcal{O}$  and to  $\mathcal{O}_h$ .

One of the two following cases holds.

**Hypothesis 4.1** ( $\alpha = \beta h^q$ ). *There exists  $q \geq 1$  such that:*

$$(36) \quad \alpha = \beta h^q.$$

**Hypothesis 4.2** ( $\alpha = o(h^N)$ ,  $\forall N \in \mathbb{N}$ ). *The complex parameter  $\alpha$  satisfies (35) and for all  $N \in \mathbb{N}$ ,*

$$(37) \quad \forall N \in \mathbb{N}, \quad |\alpha| = o(h^N).$$

First we suppose that Hypothesis 4.1 holds: we will discuss on Hypothesis 4.2 later on. We denote by  $(u^{c,q}, u^{m,q})$  the solution to Problem 2 under the Hypothesis 4.1.

According to (12), by ordering and identifying the terms of the same power of  $h$ , for  $k \in \mathbb{N} \cup -1$ , for  $q \in \mathbb{N}^*$ ,  $u_k^{c,q}$  and  $u_k^{m,q}$  satisfy:

$$(38a) \quad \Delta u_l^{c,q} = 0, \text{ in } \mathcal{O},$$

for all  $(\eta, \theta) \in \mathbf{C}$ ,

$$(38b) \quad \partial_\eta^2 u_l^{m,q} = - \left\{ \begin{aligned} &\kappa \{ 3\eta \partial_\eta^2 u_{l-1}^{m,q} + \partial_\eta u_{l-1}^{m,q} \} \\ &+ 3\eta^2 \kappa^2 \partial_\eta^2 u_{l-2}^{m,q} + 2\eta \kappa^2 \partial_\eta u_{l-2}^{m,q} + \partial_\theta^2 u_{l-2}^{m,q} \\ &+ \eta^3 \kappa^3 \partial_\eta^2 u_{l-3}^{m,q} + \eta^2 \kappa^3 \partial_\eta u_{l-3}^{m,q} + \eta \kappa \partial_\theta^2 u_{l-3}^{m,q} - \eta \kappa' \partial_\theta u_{l-3}^{m,q} \end{aligned} \right\},$$

$$(38c) \quad u_l^{c,q} \circ \Phi_0 = u_l^{m,q}|_{\eta=0},$$

$$(38d) \quad \partial_\eta u_l^{m,q}|_{\eta=1} = \delta_{l,1} f,$$

$$(38e) \quad \int_{\partial \mathcal{O}} u_l^{c,q} d\sigma = 0.$$

Transmission condition (12c) coupled with Hypothesis 4.1 implies:

$$(38f) \quad \beta \partial_n u_{l-1-q}^{c,q} \circ \Phi_0 = \partial_\eta u_l^{m,q}|_{\eta=0},$$

In equations (38), we have implicitly imposed

$$(39) \quad \begin{cases} u_l^{c,q} = 0, & \text{if } l \leq -2, \\ u_l^{m,q} = 0, & \text{if } l \leq -2. \end{cases}$$

Let us now derive formal asymptotics of  $u$  when Hypothesis 4.1 holds.

#### 4.1. Formal asymptotics.

- $N = -1$ .

The functions  $u_{-1}^{m,q}$  satisfies

$$\begin{cases} \partial_\eta^2 u_{-1}^{m,q} = 0, & \text{in } \mathbf{C}, \\ \partial_\eta u_{-1}^{m,q}|_{\eta=0} = 0, & \partial_\eta u_{-1}^{m,q}|_{\eta=1} = 0, \end{cases}$$

hence  $u_{-1}^{m,q}$  depends only on the variable  $\theta$ . Observe that we have, for almost all  $\theta \in \mathbb{T}$  the following equality:

$$u_{-1}^{m,q}(\theta) = u_{-1}^{c,q} \circ \Phi_0(\theta).$$

- $N = 0$ .

The function  $u_0^{m,q}$  satisfies:

$$\begin{cases} \partial_\eta^2 u_0^{m,q} = 0, & \text{in } \mathbf{C}, \\ \partial_\eta u_0^{m,q}|_{\eta=0} = 0, & \partial_\eta u_0^{m,q}|_{\eta=1} = 0, \end{cases}$$

hence,  $\partial_\eta u_0^{m,q}$  vanishes identically in  $\mathbf{C}$ .

- $N = 1$ .

The functions  $u_1^{m,q}$  satisfy:

$$\begin{cases} \partial_\eta^2 u_1^{m,q} = -\partial_\theta^2 u_{-1}^{m,q}, & \text{in } \mathbf{C}, \\ \partial_\eta u_1^{m,q}|_{\eta=0} = \beta \partial_n u_{-1-q}^{c,q} \circ \Phi_0, & \partial_\eta u_1^{m,q}|_{\eta=1} = f. \end{cases}$$

Therefore for  $q = 1$  we obtain the following equality:

$$-\partial_\theta^2 u_{-1}^{m,1} + \beta \partial_n u_{-1}^{c,1} \circ \Phi_0 = f,$$

hence the following boundary condition imposed to  $u_{-1}^{c,1}$  on  $\partial\mathcal{O}$ :

$$-\partial_\theta^2 u_{-1}^{c,1} \Big|_{\partial\mathcal{O}} + \beta \partial_n u_{-1}^{c,1} \Big|_{\partial\mathcal{O}} = f.$$

Therefore, the function  $u_{-1}^{c,1}$  is solution to the following problem:

$$(40) \quad \begin{cases} \Delta u_{-1}^{c,1} = 0, & \text{in } \mathcal{O}, \\ -\partial_\theta^2 u_{-1}^{c,1} \Big|_{\partial\mathcal{O}} + \beta \partial_n u_{-1}^{c,1} \Big|_{\partial\mathcal{O}} = f, \\ \int_{\partial\mathcal{O}} u_{-1}^{c,1} d\partial\mathcal{O} = 0. \end{cases},$$

$$(41) \quad \forall(\eta, \theta) \in \mathcal{C}, \quad u_{-1}^{m,1} = u_{-1}^{c,1} \Big|_{\partial\mathcal{O}} \circ \Phi_0.$$

Since  $\Re(\beta) > 0$ , a straight application of Lax-Milgram theorem ensures that  $u_{-1}^{c,1}$  is uniquely determined and belongs to  $H^1(\mathcal{O})$  as soon as the boundary data belongs to  $H^{-3/2}(\partial\mathcal{O})$ .

If  $q \geq 2$ , the function  $u_{-1}^{m,q}$  satisfies:

$$(42) \quad -\partial_\theta^2 u_{-1}^{m,q} = f.$$

Since  $\int_{\mathbb{T}} u_{-1}^{m,q} d\theta = 0$ , equality (42) defines uniquely  $u_{-1}^{m,q}$ . We infer that  $u_{-1}^{c,q}$  is solution to the following problem:

$$(43) \quad \begin{cases} \Delta u_{-1}^{c,q} = 0, & \text{in } \mathcal{O}, \\ -u_{-1}^{c,q} \Big|_{\partial\mathcal{O}} = u_{-1}^{m,q} \circ \Phi_0^{-1}. \end{cases}$$

Hence we have determined  $u_{-1}^{m,q}$  and  $u_{-1}^{c,q}$  for  $q \in \mathbb{N}^*$ . Observe that  $u_{-1}^{c,1}$  is solution to Laplace equation with mixed boundary condition, and for  $q \geq 2$  the potential  $u_{-1}^{c,q}$  is the solution to Laplace equation with Dirichlet boundary condition, while for an insulating membrane, we obtained Neumann conditions for the approximated steady state potentials.

Let us now determined  $u_N^{m,q}$  and  $u_N^{c,q}$  for  $q \in \mathbb{N}^*$  by recurrence.

- Induction.

Suppose that for  $N \geq 0$ , the functions  $u_{N-1}^{m,q}$ ,  $u_{N-1}^{c,q}$ ,  $\partial_\eta u_{N-1}^{m,q}$  and  $\partial_\eta u_{N-1}^{c,q}$  are built.

The function  $u_{N+2}^{m,q}$  satisfies:

$$\begin{cases} \partial_\eta^2 u_{N+2}^{m,q} = -\kappa (3\eta \partial_\eta^2 u_{N+1}^{m,q} + \partial_\eta u_{N+1}^{m,q}) - \partial_\theta^2 u_{N+2}^{m,q} - \eta \kappa \partial_\theta^2 u_{N-1}^{m,q} + \eta \kappa' \partial_\theta u_{N-1}^{m,q}, & \text{in } \mathcal{C}, \\ \partial_\eta u_{N+2}^{m,q} \Big|_{\eta=0} = \beta \partial_n u_{N+1-q}^{c,q} \circ \Phi_0, \quad \partial_\eta u_{N+2}^{m,q} \Big|_{\eta=1} = 0. \end{cases}$$

Denote by  $\phi_N^q$  the following function:

$$\phi_N^q = \int_0^1 \left( \kappa (3\eta \partial_\eta^2 u_{N+1}^{m,q} + \partial_\eta u_{N+1}^{m,q}) + \eta \kappa \partial_\theta^2 u_{N-1}^{m,q} - \eta \kappa' \partial_\theta u_{N-1}^{m,q} \right) d\eta.$$

Since  $\partial_\eta^2 u_{N+1}^{m,q}$  and  $\partial_\eta u_{N+1}^{m,q}$  are supposed to be known, the function  $\phi_N^q$  is entirely determined. Observe that if  $q = 1$ ,  $\partial_\eta u_{N+2}^{m,1} \Big|_{\eta=0}$  is unknown since  $\partial_n u_N^{c,1}$  is not yet determined, while as soon as  $q \geq 2$ ,  $\partial_\eta u_{N+2}^{m,q} \Big|_{\eta=0}$  is known.

Using transmission condition (38c), we infer the following equality satisfied by  $u_N^{m,1}$  in  $\eta = 0$ :

$$-\partial_\theta^2 u_N^{m,1} \Big|_{\eta=0} + \beta \partial_n u_N^{c,1} \circ \Phi_0 = \phi_N^1 - \int_0^1 \eta \partial_\theta^2 \partial_\eta u_N^{m,1} d\eta,$$

hence the boundary condition imposed to  $u_N^{c,1}$  on  $\partial\mathcal{O}$ :

$$\beta \partial_n u_N^{c,1} \Big|_{\partial\mathcal{O}} - \partial_\theta^2 u_N^{c,1} \Big|_{\partial\mathcal{O}} = \left( \phi_N^1 - \int_0^1 \eta \partial_\theta^2 \partial_\eta u_N^{m,1} d\eta \right) \circ \Phi_0^{-1}.$$

Thus the function  $u_N^{c,1}$  is solution to the following problem:

$$(44) \quad \begin{cases} \Delta u_N^{c,1} = 0, & \text{in } \mathcal{O}, \\ -\partial_\theta^2 u_N^{c,1} \Big|_{\partial\mathcal{O}} + \beta \partial_n u_N^{c,1} \Big|_{\partial\mathcal{O}} = \left( \phi_N^1 - \int_0^1 \eta \partial_\theta^2 \partial_\eta u_N^{m,1} d\eta \right) \circ \Phi_0^{-1}, \\ \int_{\partial\mathcal{O}} u_N^{c,1} d\partial\mathcal{O} = 0. \end{cases}$$

In the membrane  $u_N^{m,1}$  is defined by

$$(45) \quad u_N^{m,1} = \int_0^s \partial_\eta u_N^{m,q} d\eta + u_N^{c,q} \circ \Phi_0.$$

If  $q \geq 2$ ,  $u_N^{m,q}|_{\eta=1}$  is entirely determined by the equality:

$$-\partial_\theta^2 u_N^{m,q} \Big|_{\eta=1} = \beta \partial_n u_{N+1-q}^{c,q} \circ \Phi_0 + \phi_N^q - \int_0^1 \eta \partial_\theta^2 \partial_\eta u_N^{m,q} d\eta,$$

hence

$$u_N^{m,q}(s, \theta) = \int_1^s \partial_\eta u_N^{m,q} d\eta + u_N^{m,q} \Big|_{\eta=1}.$$

The potential  $u_N^{c,q}$  satisfies the following boundary value problem:

$$(46) \quad \begin{cases} \Delta u_N^{c,q} = 0, & \text{in } \mathcal{O}, \\ u_N^{c,q} \Big|_{\partial\mathcal{O}} = u_N^{m,q} \circ \Phi_0^{-1}. \end{cases}$$

Observe that for  $q \geq 1$ ,  $\partial_\eta u_{N+2}^{m,q}$  is then entirely determined by:

$$\begin{aligned} \partial_\eta u_{N+2}^{m,q} = \int_1^s \left( -\kappa \left( 3\eta \partial_\eta^2 u_{N+1}^{m,q} + \partial_\eta u_{N+1}^{m,q} \right) \right. \\ \left. - \partial_\theta^2 u_N^{m,q} - \eta \kappa \partial_\theta^2 u_{N-1}^{m,q} + \eta \kappa' \partial_\theta u_{N-1}^{m,q} \right) d\eta. \end{aligned}$$

Therefore, we have proved that for all  $N \geq -1$ , for  $q \in \mathbb{N}^*$ , the functions  $u_N^{c,q}$  and  $u_N^{m,q}$  are uniquely determined.

**4.2. Error estimates.** Let us summarize the expressions of  $(u_k^{c,q}, u_k^{m,q})$ . In the following we use implicitly the convention (39).

- If  $q = 1$ .

‡  $N = -1$ .

$$\begin{cases} \Delta u_{-1}^{c,1} = 0, & \text{in } \mathcal{O}, \\ -\partial_\theta^2 u_{-1}^{c,1} \Big|_{\partial\mathcal{O}} + \beta \partial_n u_{-1}^{c,1} \Big|_{\partial\mathcal{O}} = f, \\ \int_{\partial\mathcal{O}} u_{-1}^{c,1} d\partial\mathcal{O} = 0. \end{cases}$$

$$\forall(\eta, \theta) \in \mathbf{C}, \quad u_{-1}^{m,1} = u_{-1}^{c,1} \Big|_{\partial\mathcal{O}} \circ \Phi_0.$$

Moreover,

$$\partial_\eta u_0^{m,1} = 0, \quad \partial_\eta u_1^{m,1} = -\eta \partial_\theta^2 u_{-1}^{m,1} + f.$$

‡  $N \geq 0$ . Suppose that for  $N \geq 0$ , the functions  $u_{N-1}^{m,1}$ ,  $u_{N-1}^{c,1}$ ,  $\partial_\eta u_N^{m,1}$  and  $\partial_\eta u_{N+1}^{m,1}$  are built. Denote by  $\phi_N^1$  the following function:

$$\phi_N^1 = \int_0^1 \left( \kappa (3\eta \partial_\eta^2 u_{N+1}^{m,q} + \partial_\eta u_{N+1}^{m,q}) + \eta \kappa \partial_\theta^2 u_{N-1}^{m,q} - \eta \kappa' \partial_\theta u_{-1}^{m,q} \right) d\eta.$$

Then  $u_N^{c,1}$  satisfies

$$\begin{cases} \Delta u_N^{c,1} = 0, \text{ in } \mathcal{O}, \\ -\partial_\theta^2 u_N^{c,1} \Big|_{\partial\mathcal{O}} + \beta \partial_n u_N^{c,1} \Big|_{\partial\mathcal{O}} = \left( \phi_N^1 - \int_0^1 \eta \partial_\theta^2 \partial_\eta u_N^{m,1} d\eta \right) \circ \Phi_0^{-1}, \\ \int_{\partial\mathcal{O}} u_N^{c,1} d\partial\mathcal{O} = 0. \end{cases}$$

In the membrane  $u_N^{m,1}$  is defined by

$$u_N^{m,1} = \int_0^s \partial_\eta u_N^{m,q} d\eta + u_N^{c,q} \circ \Phi_0.$$

Observe that  $\partial_\eta u_{N+2}^{m,1}$  is then entirely determined by:

$$\begin{aligned} \partial_\eta u_{N+2}^{m,1} = \int_1^s \left( -\kappa (3\eta \partial_\eta^2 u_{N+1}^{m,1} + \partial_\eta u_{N+1}^{m,1}) \right. \\ \left. - \partial_\theta^2 u_N^{m,1} - \eta \kappa \partial_\theta^2 u_{N-1}^{m,1} + \eta \kappa' \partial_\theta u_{N-1}^{m,1} \right) d\eta. \end{aligned}$$

• If  $q \geq 2$ .

‡  $N = -1$ . The function  $u_1^{m,q}$  is defined by

$$\begin{aligned} \int_{\mathbb{T}} u_{-1}^{m,q} d\theta = 0, \\ -\partial_\theta^2 u_{-1}^{m,q} = f. \end{aligned}$$

$u_{-1}^{c,q}$  is solution to the following problem:

$$\begin{cases} \Delta u_{-1}^{c,q} = 0, \text{ in } \mathcal{O}, \\ -u_{-1}^{c,q} \Big|_{\partial\mathcal{O}} = u_{-1}^{m,q} \circ \Phi_0^{-1}. \end{cases}$$

Moreover,

$$\partial_\eta u_0^{m,q} = 0, \quad \partial_\eta u_1^{m,q} = -\eta \partial_\theta^2 u_{-1}^{m,q} + f.$$

‡  $N \geq 0$ . Suppose that for  $N \geq 0$ , the functions  $u_{N-1}^{m,q}$ ,  $u_{N-1}^{c,q}$ ,  $\partial_\eta u_N^{m,q}$  and  $\partial_\eta u_{N+1}^{m,q}$  are built. Denote by  $\phi_N^q$  the following function:

$$\phi_N^q = \int_0^1 \left( \kappa (3\eta \partial_\eta^2 u_{N+1}^{m,q} + \partial_\eta u_{N+1}^{m,q}) + \eta \kappa \partial_\theta^2 u_{N-1}^{m,q} - \eta \kappa' \partial_\theta u_{-1}^{m,q} \right) d\eta.$$

$u_N^{m,q}|_{\eta=1}$  is entirely determined by the equality:

$$-\partial_\theta^2 u_N^{m,q}|_{\eta=1} = \beta \partial_n u_{N+1-q}^{c,q} \circ \Phi_0 + \phi_N^q - \int_0^1 \eta \partial_\theta^2 \partial_\eta u_N^{m,q} d\eta,$$

hence

$$u_N^{m,q}(s, \theta) = \int_1^s \partial_\eta u_N^{m,q} d\eta + u_N^{m,q}|_{\eta=1}.$$

The potential  $u_N^{c,q}$  satisfies the following boundary value problem:

$$\begin{cases} \Delta u_N^{c,q} = 0, & \text{in } \mathcal{O}, \\ u_N^{c,q}|_{\partial\mathcal{O}} = u_N^{m,q} \circ \Phi_0^{-1}. \end{cases}$$

Observe that for  $q \geq 1$ ,  $\partial_\eta u_{N+2}^{m,q}$  is then entirely determined by:

$$\begin{aligned} \partial_\eta u_{N+2}^{m,q} = \int_1^s & \left( -\kappa (3\eta \partial_\eta^2 u_{N+1}^{m,q} + \partial_\eta u_{N+1}^{m,q}) \right. \\ & \left. - \partial_\theta^2 u_N^{m,q} - \eta \kappa \partial_\theta^2 u_{N-1}^{m,q} + \eta \kappa' \partial_\theta u_{N-1}^{m,q} \right) d\eta. \end{aligned}$$

**Remark 4.3** (Regularity). *We leave the reader verify by induction that the following regularities hold. Let  $q \in \mathbb{N}^*$ ,  $N \geq -1$  and  $p \geq 1$ . Let  $\phi$  belong to  $H^{N+p-3/2}(\partial\Omega_h)$ .*

$$u_{-1}^{c,q} \in H^{1+N+p}(\mathcal{O}),$$

$$u_{-1}^{m,q} \in \mathcal{C}^\infty([0, 1]; H^{1/2+N+p}(\mathbb{T})),$$

$$\forall k = 0, \dots, N,$$

$$(47a) \quad u_k^{c,q} \in H^{1+N+p-2[k/2]}(\mathcal{O}),$$

$$(47b) \quad u_k^{m,q} \in \mathcal{C}^\infty([0, 1]; H^{1/2+N+p-2[(k+1)/2]}(\mathbb{T})).$$

Moreover, there exists a constant  $C_{N,\mathcal{O},p}$  independant on  $h$  and  $\beta$  such that:

$$(48a) \quad \sup_{\eta \in [0,1]} \|u_{-1}^{m,q}(\eta, \cdot)\|_{H^{1/2+N+p}(\mathbb{T})} \leq C_{N,\mathcal{O},p} \|f\|_{H^{N+p-3/2}(\partial\mathcal{O})},$$

$$(48b) \quad \|u_{-1}^c\|_{H^{1+N+p}(\mathcal{O})} \leq C_{N,\mathcal{O},p} \|f\|_{H^{N+p-3/2}(\partial\mathcal{O})},$$

$$\forall k = 0, \dots, N,$$

$$(48c) \quad \sup_{\eta \in [0,1]} \|u_k^{m,q}(\eta, \cdot)\|_{H^{1/2+N+p-2[(k+1)/2]}(\mathbb{T})} \leq C_{N,\mathcal{O},p} \|f\|_{H^{N+p-3/2}(\partial\mathcal{O})},$$

$$(48d) \quad \|u_k^{c,q}\|_{H^{1+N+p-[k/2]}(\mathcal{O})} \leq C_{N,\mathcal{O},p} \|f\|_{H^{N+p-3/2}(\partial\mathcal{O})}.$$

**Theorem 4.4.** *Define  $h_0$  by (3). Let  $h$  belong to  $(0, h_0)$ . Let  $q \in \mathbb{N}^*$  and  $N \in \mathbb{N}$ . The complex parameter  $\alpha$  satisfies (35) with Hypothesis 4.1. Let  $\phi$  belong to  $H^{N+3/2+q}(\partial\Omega_h)$ , and suppose that  $\partial\mathcal{O}$  is smooth. The function  $f$  is defined by (9). We denote by  $u$  the solution to Problem (2). Let  $r_N^{c,q}$  and  $r_N^{m,q}$  be the functions*

defined by:

$$\begin{cases} r_N^{c,q} = u - \sum_{k=-1}^N u_k^{c,q} h^k, & \text{in } \mathcal{O}, \\ r_N^{m,q} = u \circ \Phi - \sum_{k=-1}^N u_k^{m,q} h^k, & \text{in } \mathcal{C}. \end{cases}$$

Then, there exists a constant  $C_{\mathcal{O},N} > 0$  depending only on the domain  $\mathcal{O}$  and on  $N$  such that

$$(49a) \quad \|r_N^{c,q}\|_{H^1(\mathcal{O})} \leq C_{\mathcal{O},N} \|f\|_{H^{N+3/2+q}(\partial\mathcal{O})} \max\left(\sqrt{\frac{h}{|\beta|}}, \sqrt{h}\right) h^{N+1/2},$$

$$(49b) \quad \|r_N^{m,q}\|_{H_{\mathfrak{g}}^1(\mathcal{C})} \leq C_{\mathcal{O},N} \|f\|_{H^{N+3/2}(\partial\mathcal{O})} h^{N+1/2}.$$

If  $\phi$  belongs to  $H^{N+5/2+q}(\partial\Omega_h)$ , we have

$$\|r_N^{c,q}\|_{H^1(\mathcal{O})} \leq C_{\mathcal{O},N} \|f\|_{H^{N+5/2+q}(\mathbb{T})} h^{N+1}.$$

*Proof.* The proof of Theorem 4.4 is similar to the proof of Theorem 3.4. Since  $\phi$  belongs to  $H^{N-1/2}(\partial\Omega_h)$ , according to the previous lemma, the couples of functions  $(r_N^{c,q}, r_N^{m,q})$  and  $(r_{N+1}^{c,q}, r_{N+1}^{m,q})$  are well defined and belong to  $H^1(\mathcal{O}) \times H_{\mathfrak{g}}^1(\mathcal{C})$ .

Denote by  $\tilde{g}_N$  the following function defined on  $\mathcal{C}$ :

$$(50) \quad \begin{aligned} \tilde{g}_N = & \kappa (3\eta \partial_\eta^2 u_N^{m,q} + \partial_\eta u_N^{m,q}) + 3\eta^2 \kappa^2 \partial_\eta^2 u_{N-1}^{m,q} + 2\eta \kappa^2 \partial_\eta u_{N-1}^{m,q} + \partial_\theta^2 u_{N-1}^{m,q} \\ & + \eta^3 \kappa^3 \partial_\eta^2 u_{N-2}^{m,q} + \eta^2 \kappa^3 \partial_\eta u_{N-2}^{m,q} + \eta \kappa \partial_\theta^2 u_{N-2}^{m,q} - \eta \kappa' \partial_\theta u_{N-2}^{m,q} \\ & + h \left( 3\eta^2 \kappa^2 \partial_\eta^2 u_N^{m,q} + 2\eta \kappa^2 \partial_\eta u_N^{m,q} + \partial_\theta^2 u_N^{m,q} \right. \\ & \left. + \eta^3 \kappa^3 \partial_\eta^2 u_{N-1}^{m,q} + \eta^2 \kappa^3 \partial_\eta u_{N-1}^{m,q} + \eta \kappa \partial_\theta^2 u_{N-1}^{m,q} - \eta \kappa' \partial_\theta u_{N-1}^{m,q} \right) \\ & + h^2 \left( \eta^3 \kappa^3 \partial_\eta^2 u_N^{m,q} + \eta^2 \kappa^3 \partial_\eta u_N^{m,q} + \eta \kappa \partial_\theta^2 u_N^{m,q} - \eta \kappa' \partial_\theta u_N^{m,q} \right) \end{aligned}$$

According to the previous lemma and since  $\phi$  belongs to  $H^{N-1/2}(\partial\Omega_h)$ , the above function  $\tilde{g}_N$  belongs to  $\mathcal{C}^\infty([0,1]; H^{-1/2}(\mathbb{T}))$  and the function  $\partial_\eta V_N^{m,q}$  belongs to  $\mathcal{C}^\infty([0,1]; H^{3/2}(\mathbb{T}))$ . Moreover, there exists a constant  $C_{N,\mathcal{O}}$  such that

$$(51) \quad \begin{cases} \sup_{\eta \in [0,1]} \|\tilde{g}_N(\eta, \cdot)\|_{H^{-1/2}(\mathbb{T})} \leq C_{N,\mathcal{O}} \|f\|_{H^{N-1/2}(\mathbb{T})}, \\ \sup_{\eta \in [0,1]} \|\partial_\eta u_N^{m,q}(\eta, \cdot)\|_{H^{3/2}(\mathbb{T})} \leq C_{N,\mathcal{O}} \|f\|_{H^{N-1/2}(\mathbb{T})}. \end{cases}$$

The functions  $r_N^{c,q}$  and  $r_N^{m,q}$  satisfy the following problem:

$$\begin{aligned} \Delta r_N^{c,q} &= 0, & \text{in } \mathcal{O}, \\ \partial_\eta \left( \frac{1+h\eta\kappa}{h} \partial_\eta r_N^{m,q} \right) + \partial_\theta \left( \frac{h}{1+h\eta\kappa} \partial_\theta r_N^{m,q} \right) &= \frac{-h^N}{(1+h\eta\kappa)} \tilde{g}_N, \end{aligned}$$

with transmission conditions:

$$\begin{aligned} \beta h^{1+q} \partial_n r_N^{c,q} \circ \Phi_0 &= \frac{1}{h} \left( \partial_\eta r_N^{m,q} |_{\eta=0} + \beta h^{N+1+q} (\partial_n u_{N-1}^{c,q} \circ \Phi_0 + h \partial_\eta u_N^{m,q} \circ \Phi_0) \right), \\ r_N^{c,q} \circ \Phi_0 &= r_N^{m,q} |_{\eta=0}, \end{aligned}$$



with boundary condition

$$\partial_\eta r_N^{m,q}|_{\eta=1} = 0,$$

and with gauge condition

$$\int_{\partial\mathcal{O}} r_N^{c,q} d\sigma = 0.$$

By multiplying the above equality by  $\overline{r_N}$  and by integration by parts, we infer that:

$$(52) \quad \begin{aligned} \beta h^{1+q} \|dr_N^{c,q}\|_{\Lambda^1 L^2(\mathcal{O})}^2 + \|dr_N^{m,q}\|_{\Lambda^1 L^2_{\mathfrak{g}}(\mathcal{C})}^2 &= -h^N \int_{\mathcal{C}} \tilde{g}_N(\eta, \theta) \overline{r_N^{m,q}}(\eta, \theta) d\eta d\theta \\ &+ \beta h^{N+1+q} \int_{\mathbb{T}} (\partial_\eta u_{N-1}^{c,q} \circ \Phi_0 + h \partial_\eta u_N^{c,q} \circ \Phi_0) \overline{r_N^{m,q}} \Big|_{\eta=0} d\theta. \end{aligned}$$

The end of the proof is similar to Theorem 3.4. Using the positivity of  $\Re(\beta)$  we straight infer estimate (49b) of  $r_N^{m,q}$ . To obtain the estimates of  $r_N^{c,q}$ , we write:

$$r_N^{c,q} = r_{N+q}^{c,q} + \sum_{k=1}^q u_{N+k}^{c,q} h^{N+k}.$$

□

**4.3. The case  $\alpha = o(h^N)$ ,  $\forall N \in \mathbb{N}$ .** Now, we suppose that Hypothesis 4.2 holds. In this case, we prove that  $u^c$  and  $u^m$  may be approximated by  $U^c$  and  $U^m$ , which are solution to:

$$(53a) \quad \Delta U^m = 0, \text{ in } \mathcal{O}_h,$$

$$(53b) \quad \partial_\eta U^m|_{\partial\mathcal{O}} = 0, \quad \partial_\eta U^m|_{\partial\Omega_h} = \phi,$$

$$(53c) \quad \int_{\partial\mathcal{O}} U^m d\sigma = 0.$$

and

$$(54a) \quad \Delta U^c = 0, \text{ in } \mathcal{O},$$

$$(54b) \quad U^c|_{\partial\mathcal{O}} = U^m|_{\partial\mathcal{O}}.$$

Actually, we have the following lemma:

**Lemma 4.5.** *Let  $\phi$  belong to  $H^{-1/2}(\partial\Omega_h)$ . Let  $(u^c, u^m)$  be the solution to Problem (2), and  $U^m$  and  $U^c$  be defined respectively by (53) and (54). Then, we have:*

$$(55) \quad \|u^m - U^m\|_{H^1(\mathcal{O}_h)} \leq C_{\mathcal{O}} |\alpha| |\phi|_{H^{-1/2}(\partial\Omega_h)},$$

$$(56) \quad \|u^c - U^c\|_{H^1(\mathcal{O})} \leq C_{\mathcal{O}} \sqrt{|\alpha|} |\phi|_{H^{-1/2}(\partial\Omega_h)}.$$

*Proof.* Denote by  $w^c$  and  $w^m$  the following functions:

$$w^c = u^c - U^c, \quad w^m = u^m - U^m,$$

and let  $\phi$  belong to  $H^{-1/2}(\partial\Omega_h)$ . We have:

$$\begin{aligned}
(57a) \quad & \Delta w^c = 0, \text{ in } \mathcal{O}, \\
(57b) \quad & \Delta w^m = 0, \text{ in } \mathcal{O}_h, \\
(57c) \quad & \alpha \partial_n w^c|_{\partial\mathcal{O}} = \partial_n w^m|_{\partial\mathcal{O}} - \alpha \partial_n U^c|_{\partial\mathcal{O}}, \\
(57d) \quad & w^c|_{\partial\mathcal{O}} = w^m|_{\partial\mathcal{O}}, \\
(57e) \quad & \partial_\eta w^m|_{\partial\Omega_h} = 0, \\
(57f) \quad & \int_{\partial\mathcal{O}} w^m d\sigma = 0.
\end{aligned}$$

Thus we infer:

$$(58) \quad \alpha \int_{\mathcal{O}} |\nabla w^c|^2 d\text{vol}_{\mathcal{O}} + \int_{\mathcal{O}_h} |\nabla w^m|^2 d\text{vol}_{\mathcal{O}_h} = \alpha \int_{\partial\mathcal{O}} \partial_n U^c|_{\partial\mathcal{O}} \overline{w^m} d\sigma.$$

It is well-known that :

$$\|U^m\|_{H^1(\mathcal{O}_h)} \leq C_{\mathcal{O}} |\phi|_{H^{-1/2}(\partial\Omega_h)},$$

and

$$\|U^c\|_{H^1(\mathcal{O}_h)} \leq C_{\mathcal{O}} |U^m|_{\partial\mathcal{O}}|_{H^{1/2}(\partial\mathcal{O})}.$$

Since  $\alpha$  satisfies (35) we infer,

$$\|w^m\|_{H^1(\mathcal{O}_h)} \leq C_{\mathcal{O}} |\alpha| |\phi|_{H^{-1/2}(\partial\Omega_h)},$$

and thereby

$$\|w^c\|_{H^1(\mathcal{O})} \leq C_{\mathcal{O}} \sqrt{|\alpha|} |\phi|_{H^{-1/2}(\partial\Omega_h)}.$$

□

It remains to derive asymptotics of  $U^m$  and then these of  $U^c$ . They are similar to asymptotics of  $u^{m,q}$  for  $q \geq 2$ : we just have to replace  $\beta$  by zero. We think the reader may easily derive these asymptotics from our previous results.

## CONCLUSION

In this paper, we have studied the steady state potentials in a highly contrasted domain with thin layer when Neumann boundary condition is imposed on the exterior boundary. We derived rigorous asymptotics with respect to the thickness of the potentials in each domain and we gave error estimate in terms of appropriate Sobolev norm of the boundary data, electromagnetic parameters of our domain and a constant depending only on the geometry of the domain. It has to be mentioned that for an insulating inner domain (or equivalently a conducting membrane), the asymptotic expansions start at the order -1 and mixed or Dirichlet boundary conditions has to be imposed on the asymptotic terms of the inner domain.

To illustrate these asymptotics, numerical simulations using FEM with GetDP [9] are forthcoming work with Patrick Dular from Université de Liège. Few results have been shown at the seminar Numelec [15].

## APPENDIX

Let  $\star$  denote the Hodge star operator, which maps 0-forms to 2-forms, 1-forms to 1-forms and 2-forms to 0-forms (see Flanders [11]). We give explicit formulae for the operators  $d$ ,  $\delta$ ,  $\text{ext}$  and  $\text{int}$ . These formulae are straightforward consequences of the definition of the operators  $\star$ ,  $d$  and  $\delta = \star^{-1}d\star$ . We refer the reader to Dubrovin, Fomenko and Novikov [8].

We consider the metric given by the following matrix  $G$

$$(59) \quad G = \begin{pmatrix} g_{11} & g_{12} \\ g_{12} & g_{22} \end{pmatrix}.$$

We denote by  $|G|$  the determinant of  $G$ . The inverse of  $G$  is denoted by  $G^{-1}$

$$G^{-1} = (g^{ij})_{ij},$$

and we suppose that the signature of  $G$  is equal to 1. Thereby, the operator  $\star^2$  is equal to  $\text{Id}$  on the space of 0-forms and 2-forms and it is equal to  $-\text{Id}$  on 1-forms.

#### 4.4. Star operator in $\mathbb{R}^2$ .

4.4.1. *On 0-forms and on 2-forms.* Let  $T$  be a 0-form and let  $S$  be the 2-form  $\nu \, dy^1 dy^2$ . Then  $\star T$  is the 2-form  $\mu \, dy^1 dy^2$  and  $\star S$  is the 0-form  $f$ . The following identities hold:

$$\begin{aligned} \mu &= \sqrt{|G|}T, \\ f &= \frac{1}{\sqrt{|G|}}\nu. \end{aligned}$$

4.4.2. *On 1-forms.* Let  $T$  be the 1-form  $T_1 \, dy^1 + T_2 \, dy^2$ . Then  $\star T$  is the 1-form  $\mu_1 \, dy^1 + \mu_2 \, dy^2$ , and we have the following formulae:

$$\begin{aligned} \mu_1 &= -\sqrt{|G|} (g^{12}T_1 + g^{22}T_2), \\ \mu_2 &= \sqrt{|G|} (g^{11}T_1 + g^{12}T_2). \end{aligned}$$

4.5. **The action of  $d$  acting on 0-forms in  $\mathbb{R}^2$ .** Let  $\mu$  be a 0 form, then  $d\mu$  has the following expression:

$$d\mu = \frac{\partial \mu}{\partial y^1} dy^1 + \frac{\partial \mu}{\partial y^2} dy^2.$$

4.6. **The action of  $\delta$  acting on 1-forms on  $\mathbb{R}^2$ .** Let  $\mu$  be the 1-form  $\mu_1 dy^1 + \mu_2 dy^2$ , and define  $\delta\mu = \alpha$ . The 0-form  $\alpha$  is equal to:

$$\begin{aligned} \alpha &= -\frac{1}{\sqrt{|G|}} \left\{ \frac{\partial}{\partial y^1} \left( \sqrt{|G|} (g^{11}\mu_1 + g^{12}\mu_2) \right) \right. \\ &\quad \left. + \frac{\partial}{\partial y^2} \left( \sqrt{|G|} (g^{12}\mu_1 + g^{22}\mu_2) \right) \right\}. \end{aligned}$$

4.7. **The exterior product of a 1-form with a 0-form.** Let  $N$  be the 1-form  $N_1 dy^1 + N_2 dy^2$  and  $f$  be a 0-form. The exterior product of  $\text{ext}(N)f$  is:

$$\text{ext}(N)f = fN_1 dy^1 + fN_2 dy^2.$$

**4.8. The interior product of a 1-form with a 1-form.** Let  $N$  and  $\mu$  be the 1-forms  $N_1 dy^1 + N_2 dy^2$ , and  $\mu_1 dy^1 + \mu_2 dy^2$ . Then 0-form  $\text{int}(N)\mu$  has the following expression:

$$\text{int}(N)\mu = N_1 (\mu_1 g^{11} + \mu_2 g^{12}) + N_2 (\mu_1 g^{12} + \mu_2 g^{22}).$$

## REFERENCES

- [1] Habib Ammari and Hyeonbae Kang. Properties of the generalized polarization tensors. *Multiscale Model. Simul.*, 1(2):335–348 (electronic), 2003.
- [2] Habib Ammari and Hyeonbae Kang. Reconstruction of conductivity inhomogeneities of small diameter via boundary measurements. In *Inverse problems and spectral theory*, volume 348 of *Contemp. Math.*, pages 23–32. Amer. Math. Soc., Providence, RI, 2004.
- [3] Balanis and Constantine. *Advanced Engineering Electromagnetics*. John Wiley and Sons Ltd, 1989.
- [4] Elena Beretta, Elisa Francini, and Michael S. Vogelius. Asymptotic formulas for steady state voltage potentials in the presence of thin inhomogeneities. A rigorous error analysis. *J. Math. Pures Appl. (9)*, 82(10):1277–1301, 2003.
- [5] Elena Beretta, Arup Mukherjee, and Michael Vogelius. Asymptotic formulas for steady state voltage potentials in the presence of conductivity imperfections of small area. *Z. Angew. Math. Phys.*, 52(4):543–572, 2001.
- [6] Yves Capdeboscq and Michael S. Vogelius. A general representation formula for boundary voltage perturbations caused by internal conductivity inhomogeneities of low volume fraction. *M2AN Math. Model. Numer. Anal.*, 37(1):159–173, 2003.
- [7] B. Doubrovine, S. Novikov, and A. Fomenko. *Géométrie contemporaine. Méthodes et applications. I*. “Mir”, Moscow, 1982. Géométrie des surfaces, des groupes de transformations et des champs. [Geometry of surfaces, groups of transformations and fields], Translated from the Russian by Vladimir Kotliar.
- [8] B. A. Dubrovin, A. T. Fomenko, and S. P. Novikov. *Modern geometry—methods and applications. Part I*, volume 93 of *Graduate Texts in Mathematics*. Springer-Verlag, New York, second edition, 1992. The geometry of surfaces, transformation groups, and fields, Translated from the Russian by Robert G. Burns.
- [9] P. Dular, C. Geuzaine, F. Henrotte, and W. Legros. A general environment for the treatment of discrete problems and its application to the finite element method. 34(5):3395–3398, September 1998.
- [10] E.C. Fear and M.A. Stuchly. Modeling assemblies of biological cells exposed to electric fields. *IEEE Trans.Bio.Eng.*, 45(1):1259–1271 (electronic), 1998.
- [11] Harley Flanders. *Differential forms with applications to the physical sciences*. Academic Press, New York, 1963.
- [12] K.R. Foster and H.P. Schwan. Dielectric properties of tissues and biological materials: a critical review. *CRC in Biomedical Engineering*, 17(1):25–104, 1989.
- [13] Peter B. Gilkey, John V. Leahy, and Jeonghyeong Park. *Spectral geometry, Riemannian submersions, and the Gromov-Lawson conjecture*. Studies in Advanced Mathematics. Chapman & Hall/CRC, Boca Raton, FL, 1999.
- [14] C. Poinard. Méthodes asymptotiques pour le calcul de champs électromagnétiques dans des milieux à couche mince. application aux cellules biologiques. Thesis, November 2006.
- [15] C. Poinard, P. Dular, L. Krähenbühl, L. Nicolas, and M Schatzman. Méthodes asymptotiques pour le calcul de champs. In *5ème Conférence Européenne sur les Méthodes Numériques en Electromagnétisme (NUMELEC 2006)*, 29-30 Nov. 1 Dec. 2006. Lille, FRANCE.
- [16] G. Vial. Analyse multi-échelle et conditions aux limites approchées pour un problème avec couche mince dans un domaine à coin. Thesis, Juin 2003.

INSTITUT CAMILLE JORDAN, UNIVERSITÉ CLAUDE BERNARD LYON1

POIGNARD@MATH.UNIV-LYON1.FR