



HAL
open science

Two-dimensional Keller-Segel model: Optimal critical mass and qualitative properties of the solutions

Adrien Blanchet, Jean Dolbeault, Benoît Perthame

► **To cite this version:**

Adrien Blanchet, Jean Dolbeault, Benoît Perthame. Two-dimensional Keller-Segel model: Optimal critical mass and qualitative properties of the solutions. *Electronic Journal of Differential Equations*, 2006, 44, 32 pp. hal-00021782

HAL Id: hal-00021782

<https://hal.science/hal-00021782v1>

Submitted on 25 Mar 2006

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Two-dimensional Keller-Segel model: Optimal critical mass and qualitative properties of the solutions

Adrien Blanchet,^{*†} Jean Dolbeault,^{*} and Benoît Perthame[‡]

March 25, 2006

Abstract

The Keller-Segel system describes the collective motion of cells which are attracted by a chemical substance and are able to emit it. In its simplest form it is a conservative drift-diffusion equation for the cell density coupled to an elliptic equation for the chemo-attractant concentration. It is known that, in two space dimensions, for small initial mass, there is global existence of solutions and for large initial mass blow-up occurs. In this paper we complete this picture and give a detailed proof of the existence of weak solutions below the critical mass, above which any solution blows-up in finite time in the whole euclidean space. Using hypercontractivity methods, we establish regularity results which allow us to prove an inequality relating the free energy and its time derivative. For a solution with sub-critical mass, this allows us to give for large times an “intermediate asymptotics” description of the vanishing. In self-similar coordinates, we actually prove a convergence result to a limiting self-similar solution which is not a simple reflect of the diffusion.

Résumé

Le système de Keller-Segel décrit le mouvement collectif de cellules qui sont attirées par une substance chimique et qu’elles sont capables d’émettre. Dans sa forme la plus simple, il s’agit d’une équation de dérive-diffusion conservative pour la densité des cellules, couplée à une équation elliptique pour la concentration de chemo-attractant. On sait que, en dimension deux, il y a existence globale de solutions pour une masse initiale petite, alors que pour une masse initiale grande, la solution explose en temps fini dans l’espace euclidien tout entier. Dans cet article, nous précisons cette description et donnons une preuve détaillée de l’existence de solutions faibles avec masse inférieure à la masse critique, au-dessus de laquelle toute solution explose en temps fini. En utilisant des arguments d’hypercontractivité, nous établissons des résultats de régularité qui nous permettent de prouver une inégalité reliant l’énergie libre et sa dérivée en temps. Pour une solution avec masse sous-critique, cela nous permet de donner une description des “asymptotiques intermédiaires” en temps grand, qui décrivent la convergence locale des solutions vers zéro. Dans des coordonnées auto-similaires, nous montrons en fait un résultat de

^{*}CEREMADE (UMR CNRS no. 7534), Université Paris Dauphine, Place de Lattre de Tassigny, 75775 Paris Cédex 16, France. Phone: (33) 1 44 05 46 78, Fax: (33) 1 44 05 45 99.

E-mail: dolbeaul@ceremade.dauphine.fr, Internet: <http://www.ceremade.dauphine.fr/~dolbeaul/>

[†]CERMICS, ENPC, 6–8 avenue Blaise Pascal, Champs-sur-Marne, 77455 Marne-la-Vallée Cédex 2, France. E-mail: blanchet@ceremade.dauphine.fr, Internet: <http://www.ceremade.dauphine.fr/~blanchet/>

[‡]DMA (UMR CNRS no. 8553), Ecole Normale Supérieure, 45 rue d’Ulm, 75005 Paris Cédex 05, France. E-mail: Benoit.Perthame@ens.fr, Internet: <http://www.dma.ens.fr/users/perthame/>

convergence vers une solution auto-similaire limite qui n'est pas simplement donnée par la diffusion.

Key-words. Keller-Segel model, Existence, Weak solutions, Free energy, Entropy method, logarithmic Hardy-Littlewood-Sobolev inequality, Critical Mass, Aubin-Lions compactness method, Hypercontractivity, Large time behavior, Time-dependent rescaling, Self-similar variables, Intermediate asymptotics

MSC (2000). Primary: 35B45, 35B30, 35D05, 35K15, 35B40; Secondary: 35D10, 35K60, 35B32

1 Introduction

The Keller-Segel system for chemotaxis describes the collective motion of cells, usually bacteria or amoebae, that are attracted by a chemical substance and are able to emit it. For a general introduction to chemotaxis, see [42, 40]. Various versions of the Keller-Segel system for chemotaxis are available in the literature, depending on the phenomena and scales one is interested in. We refer the reader to the very nice review papers [29, 30] and references therein. The complete Keller-Segel model is a system of two parabolic equations. In this paper, we consider only the simplified two-dimensional case and assume that the equations take the form

$$\begin{cases} \frac{\partial n}{\partial t} = \Delta n - \chi \nabla \cdot (n \nabla c) & x \in \mathbb{R}^2, t > 0, \\ -\Delta c = n & x \in \mathbb{R}^2, t > 0, \\ n(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2. \end{cases} \quad (1)$$

Here $n(x, t)$ represents the cell density, and $c(x, t)$ is the concentration of chemo-attractant which induces a drift force. A classical parameter of the system is the *sensitivity* $\chi > 0$ of the bacteria to the chemo-attractant which measures the nonlinearity in the system. Here χ is a constant. Such a parameter can be removed by a scaling, to the price of a change of the total mass of the system

$$M := \int_{\mathbb{R}^2} n_0 \, dx .$$

In bounded domains, it is usual to impose no-flux boundary conditions. Here we are not interested in boundary effects and for this reason we are going to consider the system in the full space \mathbb{R}^2 , without boundary conditions. There are related models in gravitation which are defined in \mathbb{R}^3 , see, e.g., [11]. The relevant case for chemotaxis is rather the two-dimensional space, although some three-dimensional versions of the model also make sense. The L^1 -norm is critical in the sense that there exists a critical mass above which all solution blow-up in finite time, see [32], and below which they globally exist (see [21] for the *a priori* estimates and Theorem 1 below for an existence statement). The critical space is $L^{d/2}(\mathbb{R}^d)$ for $d \geq 2$, see [18, 19] and the references therein. In dimension $d = 2$, the Green kernel associated to the Poisson equation is a logarithm, namely $c = -\frac{1}{2\pi} \log |\cdot| * n$. When the Poisson interaction is replaced by a convolution kernel, it is the logarithmic singularity which is critical for the L^1 -norm whatever the dimension is, see [14].

Historically the key papers for this family of models are the original contribution [33] of E. F. Keller and L. A. Segel, and a work by C. S. Patlak, [49]. The rigorous derivation of the Keller-Segel system from an interacting stochastic many-particles system has been done in [57]. Simulations of these can be found in [47]. A very interesting justification of the Keller-Segel model as a diffusion limit of a kinetic model has recently been published, see [16]. Related models with prevention of overcrowding, see [27], volume effects [34, 13, 64], or involving more than one chemo-attractant have also been studied.

As conjectured by S. Childress and J. K. Percus [17] and V. Nanjundiah [46] either the solution of the complete Keller-Segel system globally exists or it blows-up in finite time, a phenomenon called *chemotactic collapse* in the literature. As we shall see, this classification is valid for the simplified Keller-Segel system (1). A large series of results, mostly in the bounded domain case, has been obtained by T. Nagai, T. Senba and T. Suzuki. Many of these results can be found in [55, 58]. Concerning blow-up phenomena, a key contribution has been brought by M. A. Herrero and J. J. L. Velázquez [26, 61]. Also see [41] for numerical computations.

The main tool in this paper is the *free energy*

$$F[n] := \int_{\mathbb{R}^2} n \log n \, dx - \frac{\chi}{2} \int_{\mathbb{R}^2} n c \, dx$$

which provides useful *a priori* estimates. The free energy functional is a well known tool for gravitational models, see [3, 48, 63, 10] and has been introduced for chemotactic models by T. Nagai, T. Senba and K. Yoshida in [44], by P. Biler in [7] and by H. Gajewski and K. Zacharias in [22].

Based on the logarithmic Hardy-Littlewood-Sobolev inequality in its sharp form as established in [15, 4], the free energy is bounded from below if $\chi M \leq 8\pi$, see [21]. Here we use this estimate to prove the global existence of solutions of (1) if $\chi M < 8\pi$. We also prove that for these solutions, the free energy is decaying and use it to study the large time behaviour of the solutions. The limiting case $\chi M = 8\pi$ has recently been studied in the radial case, see [8, 9].

The literature on the Keller-Segel model is huge and it is out of the scope of this paper to give a complete bibliography. Some additional papers will be quoted in the text. Otherwise, we suggest the interested reader to primarily refer to the surveys [51, 29, 30].

Our first main result is the following existence and regularity statement.

Theorem 1 *Assume that $n_0 \in L^1_+(\mathbb{R}^2, (1 + |x|^2) dx)$ and $n_0 \log n_0 \in L^1(\mathbb{R}^2, dx)$. If $M < 8\pi/\chi$, then the Keller-Segel system (1) has a global weak non-negative solution n with initial data n_0 such that*

$$(1 + |x|^2 + |\log n|) n \in L^\infty_{\text{loc}}(\mathbb{R}^+, L^1(\mathbb{R}^2)), \quad \int_0^t \int_{\mathbb{R}^2} n |\nabla \log n - \chi \nabla c|^2 \, dx \, dt < \infty,$$

$$\text{and} \quad \int_{\mathbb{R}^2} |x|^2 n(x, t) \, dx = \int_{\mathbb{R}^2} |x|^2 n_0(x) \, dx + 4M \left(1 - \frac{\chi M}{8\pi}\right) t$$

for any $t > 0$. Moreover $n \in L^\infty_{\text{loc}}((\varepsilon, \infty), L^p(\mathbb{R}^2))$ for any $p \in (1, \infty)$ and any $\varepsilon > 0$, and the following inequality holds for any $t > 0$:

$$F[n(\cdot, t)] + \int_0^t \int_{\mathbb{R}^2} n |\nabla (\log n) - \chi \nabla c|^2 \, dx \, ds \leq F[n_0]. \tag{2}$$

This result was partially announced in [21]. Compared to [21], the main novelty is that we prove the free energy inequality (2) and get the hypercontractive estimate: $n(\cdot, t)$ is bounded in $L^p(\mathbb{R}^2)$ for any $p \in (1, \infty)$ and almost any $t > 0$. The equation holds in the distributions sense. Indeed, writing

$$\Delta n - \chi \nabla \cdot (n \nabla c) = \nabla \cdot [n(\nabla \log n - \chi \nabla c)] ,$$

we can see that the flux is well defined in $L^1(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$ since

$$\iint_{[0, T] \times \mathbb{R}^2} n |\nabla \log n - \chi \nabla c| dx dt \leq \left(\iint_{[0, T] \times \mathbb{R}^2} n dx dt \right)^{1/2} \left(\iint_{[0, T] \times \mathbb{R}^2} n |\nabla \log n - \chi \nabla c|^2 dx dt \right)^{1/2} < \infty .$$

Our second main result deals with large time behavior, intermediate asymptotics and convergence to asymptotically self-similar profiles given in the rescaled variables by the equation

$$u_\infty = M \frac{e^{\chi v_\infty - |x|^2/2}}{\int_{\mathbb{R}^2} e^{\chi v_\infty - |x|^2/2} dx} = -\Delta v_\infty , \quad \text{with} \quad v_\infty = -\frac{1}{2\pi} \log |\cdot| * u_\infty . \quad (3)$$

In the original variables, the self-similar solutions of (1) take the expression:

$$\begin{aligned} n_\infty(x, t) &:= \frac{1}{1+2t} u_\infty(\log(\sqrt{1+2t}), x/\sqrt{1+2t}) , \\ c_\infty(x, t) &:= v_\infty(\log(\sqrt{1+2t}), x/\sqrt{1+2t}) . \end{aligned}$$

This allows us to state our second main result, on intermediate asymptotics.

Theorem 2 *Under the same assumptions as in Theorem 1, there exists a solution of (3) such that*

$$\lim_{t \rightarrow \infty} \|n(\cdot, t) - n_\infty(\cdot, t)\|_{L^1(\mathbb{R}^2)} = 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} \|\nabla c(\cdot, t) - \nabla c_\infty(\cdot, t)\|_{L^2(\mathbb{R}^2)} = 0 .$$

This paper is organized as follows. Section 2 is devoted to the detailed proof of the existence of weak solutions with subcritical mass and without any symmetry assumption. *A priori* estimates have been derived in [21]. The point here is to establish the result with all necessary details: regularized problem, uniform estimates, passage to the limit in the regularization parameter. Compared to [21], we also establish Inequality (2). Proving such an inequality requires a detailed study of the regularity properties of the solutions, which is done in Section 3: By hypercontractivity methods, we prove that the solution is bounded in any L^p space for almost any positive t . Using the free energy we study in Section 4 the asymptotic behavior of the solutions for large times and prove Theorem 2. The main difficulty comes from the fact that the uniqueness of the solutions to (3) for a given $M \in (0, 8\pi/\chi)$ is not known, although many additional properties (radial symmetry, regularity, decay at infinity) of the limiting solution in the self-similar variables are known.

2 Existence for sub-critical masses

We assume that the initial data satisfies the following assumptions:

$$\begin{aligned} n_0 &\in L^1_+(\mathbb{R}^2; (1 + |x|^2) dx) , \\ n_0 \log n_0 &\in L^1(\mathbb{R}^2; dx) . \end{aligned} \tag{4}$$

Because of the divergence form of the right hand side of the equation for n , the total mass is conserved at least for smooth and sufficiently decay solutions

$$M := \int_{\mathbb{R}^2} n_0(x) dx = \int_{\mathbb{R}^2} n(x, t) dx . \tag{5}$$

Our purpose here is first to give a complete existence theory in the subcritical case, *i.e.* in the case

$$M < 8\pi/\chi .$$

This result has been announced in [21], which was dealing only with *a priori* estimates. Here, we give the proofs with all details. More precisely, we prove that under Assumption (4), there are only two cases:

Case 1. Solutions to (1) blow-up in finite time when $M > 8\pi/\chi$,

Case 2. There exists a global in time solution of (1) when $M < 8\pi/\chi$.

The case $M = 8\pi/\chi$ is delicate and for radial solutions, some results have been obtained recently, see [8, 9].

Our existence theory completes the partial picture established in [32]. The solution of the Poisson equation $-\Delta c = n$ is given up to an harmonic function. From the beginning, we have in mind that the concentration of the chemo-attractant is defined by

$$c(x, t) = -\frac{1}{2\pi} \int_{\mathbb{R}^2} \log |x - y| n(y, t) dy . \tag{6}$$

There are other possible solutions, which may result in significantly different qualitative behaviors, as we shall see in Section 4.2. From now on, unless it is explicitly specified, we will only consider concentrations c given by (6). In the following sections, 2.1, 2.2 and 2.3, we closely follow the presentation given in [21].

2.1 Blow-up for super-critical masses

The case $M > 8\pi/\chi$ (Case 1) is easy to understand using moments estimates. The method is classical and has been repeatedly used for various similar problems. See for instance [50] in the similar context of the Euler-Poisson system, [31]. Concerning blow-up, we refer to [18, 19, 31] for recent references on the subject.

Following for instance [53], we can define a notion of weak solution n in $L^\infty_{loc}(\mathbb{R}^+; L^1(\mathbb{R}^2))$ using the symmetry in x, y for the concentration gradient, which has interest in case of blow-up. We shall say that n is a solution to (1) if for all test functions $\psi \in \mathcal{D}(\mathbb{R}^2)$

$$\frac{d}{dt} \int_{\mathbb{R}^2} \psi(x) n(x, t) dx = \int_{\mathbb{R}^2} \Delta \psi(x) n(x, t) dx - \frac{\chi}{4\pi} \int_{\mathbb{R}^2 \times \mathbb{R}^2} [\nabla \psi(x) - \nabla \psi(y)] \cdot \frac{x - y}{|x - y|^2} n(x, t) n(y, t) dx dy .$$

Compared to standard distribution solutions, this is an improved concept that can handle measure solutions because the term $[\nabla\psi(x) - \nabla\psi(y)] \cdot \frac{x-y}{|x-y|^2}$ is continuous.

Lemma 3 *Consider a non-negative L^1 solution n to (1) in the above sense, on an interval $[0, T]$ and assume that n satisfies (5), $\int_{\mathbb{R}^2} |x|^2 n_0(x) dx < \infty$. Then it also satisfies*

$$\frac{d}{dt} \int_{\mathbb{R}^2} |x|^2 n(x, t) dx = 4M \left(1 - \frac{\chi M}{8\pi} \right) .$$

Proof. Consider a smooth function $\varphi_\varepsilon(|x|)$ with compact support that grows nicely to $|x|^2$ as $\varepsilon \rightarrow 0$. Then we use the definition of weak solutions and get

$$\frac{d}{dt} \int_{\mathbb{R}^2} \varphi_\varepsilon n dx = \int_{\mathbb{R}^2} \Delta \varphi_\varepsilon n dx - \frac{\chi}{4\pi} \int_{\mathbb{R}^2} \frac{(\nabla \varphi_\varepsilon(x) - \nabla \varphi_\varepsilon(y)) \cdot (x - y)}{|x - y|^2} n(x, t) n(y, t) dx dy .$$

Because we can always choose $\Delta \varphi_\varepsilon$ bounded and $\nabla \varphi_\varepsilon(x)$ Lipschitz continuous, we deduce that

$$\frac{d}{dt} \int_{\mathbb{R}^2} \varphi_\varepsilon n dx \leq C \int_{\mathbb{R}^2} n_0 dx ,$$

where C is some positive constant. As $\varepsilon \rightarrow 0$ we find that

$$\int_{\mathbb{R}^2} \varphi_\varepsilon n dx \leq c_1 + c_2 t ,$$

where c_1 and c_2 are two positive constants and thus

$$\int_{\mathbb{R}^2} |x|^2 n(x, t) dx < \infty \quad \forall t \in (0, T) .$$

We can pass to the limit using Lebesgue's dominated convergence theorem and thus complete the proof of Lemma 3. \square

As a consequence, we recover the statement of Case 1, namely that for $M > 8\pi/\chi$, there is a finite blow-up time T^* where solutions become singular measures.

Corollary 4 *Consider a non-negative solution n as in Lemma 3 and let $[0, T^*)$ be the maximal interval of existence. Assume that the initial data $n_0 \in L^1(\mathbb{R}^2)$ is such that $I_0 := \int_{\mathbb{R}^2} |x|^2 n_0(x) dx < \infty$. Then either $T^* = \infty$ or $n(\cdot, t)$ converges (up to extraction of sequences) as $t \rightarrow T^*$ to a measure which is not in $L^1(\mathbb{R}^2)$. If $\chi M > 8\pi$, then*

$$T^* \leq \frac{2\pi I_0}{M(\chi M - 8\pi)} .$$

As far as we know, it is an open question to decide whether the solutions of (1) with $\chi M > 8\pi$ and $I_0 = \infty$ also blow-up in finite time. Blow-up statements in bounded domains are available, see [43, 10, 28, 36, 54] and the references therein. When the solution is radially symmetric in x , the second x -moment is not needed and the blow-up profile has been explicated, namely

$$n(x, t) \rightarrow \frac{8\pi}{\chi} \delta + \tilde{n}(x, t) \quad \text{as } t \nearrow T^* ,$$

where \tilde{n} is a $L^1(\mathbb{R}^2 \times \mathbb{R}^+)$ radial function such that $t \mapsto \tilde{n}(\cdot, t)$ is measure valued, see [26, 59]. Except that solutions blow-up for large mass, in the general case very little is known on the blow-up profile (see [53] for concentrations estimates, [41] for numerical computations). Asymptotic expansions at blow-up and continuation of solutions after blow-up have been studied by Velázquez in [61, 60]. The case $\chi M = 8\pi$ has recently been investigated by Biler, Karch, Laurençot and Nadzieja in [8]. In a forthcoming paper, they prove that in the whole space case and $\chi M = 8\pi$, blow-up occurs only for infinite time, [9]. Here we will focus on the subcritical regime and prove that solutions exist and are always asymptotically vanishing for large times.

If the problem is set in dimension $d \geq 3$, the critical norm is $L^p(\mathbb{R}^d)$ with $p = d/2$. In dimension $d = 2$, the value of the mass M is therefore natural to discriminate between super- and sub-critical regimes. However, the limit of the L^p -norm is rather $\int_{\mathbb{R}^2} n \log n \, dx$ than $\int_{\mathbb{R}^2} n \, dx$, which is preserved by the evolution. This explains why it is natural to introduce the entropy, or better, as we shall see below, the *free energy*.

2.2 The usual existence proof for not too large masses

The usual proof of existence is due to W. Jäger and S. Luckhaus in [32]. Here we follow the variant [18, 19] which is based on the following computation. Consider the equation for n and compute $\frac{d}{dt} \int_{\mathbb{R}^2} n \log n \, dx$. Using an integration by parts and the equation for c , we obtain:

$$\begin{aligned} \frac{d}{dt} \int_{\mathbb{R}^2} n \log n \, dx &= -4 \int_{\mathbb{R}^2} |\nabla \sqrt{n}|^2 \, dx + \chi \int_{\mathbb{R}^2} \nabla n \cdot \nabla c \, dx \\ &= -4 \int_{\mathbb{R}^2} |\nabla \sqrt{n}|^2 \, dx + \chi \int_{\mathbb{R}^2} n^2 \, dx . \end{aligned}$$

This shows that two terms compete, namely the diffusion based entropy dissipation term $\int_{\mathbb{R}^2} |\nabla \sqrt{n}|^2 \, dx$ and the hyperbolic production of entropy.

Thus the entropy is nonincreasing if $\chi M \leq 4C_{\text{GNS}}^{-2}$, where $C_{\text{GNS}} = C_{\text{GNS}}^{(4)}$ is the best constant for $p = 4$ in the Gagliardo-Nirenberg-Sobolev inequality:

$$\|u\|_{L^p(\mathbb{R}^2)}^2 \leq C_{\text{GNS}}^{(p)} \|\nabla u\|_{L^2(\mathbb{R}^2)}^{2-4/p} \|u\|_{L^2(\mathbb{R}^2)}^{4/p} \quad \forall u \in H^1(\mathbb{R}^2), \quad \forall p \in [2, \infty). \quad (7)$$

The explicit value of C_{GNS} is not known but can be computed numerically (see [62]) and one finds that the entropy is nonincreasing if $\chi M \leq 4C_{\text{GNS}}^{-2} \approx 1.862... \times (4\pi) < 8\pi$. Such an estimate is therefore not sufficient to cover the whole range of M for global existence in the second case.

In [32] it is also shown that equiintegrability (deduced from the $n \log n$ estimate for instance) is enough to propagate any L^p initial norm. We will come back on this point in Section 2.7 and prove later that due to the regularizing effects, the solution is bounded in time with values in $L^p(\mathbb{R}^2)$ for all positive times.

2.3 A free energy method and a priori estimates

To obtain sharper estimates and prove a global existence result (Case 2), we use the *free energy* which has already been introduced in Section 1:

$$F[n] := \int_{\mathbb{R}^2} n \log n \, dx - \frac{\chi}{2} \int_{\mathbb{R}^2} n c \, dx .$$

See [7, 22, 44] in the case of a bounded domain. The first term in F is the *entropy* and the second one a *potential energy term*. Such a free energy enters in the general notion of entropies, and this is why it is sometimes referred to the method as the “entropy method”, although the notion of *free energy* is physically more appropriate. See [1] for an historical review on these notions. For any solution n of (1), $F[n(\cdot, t)]$ is monotone nonincreasing.

Lemma 5 *Consider a non-negative $C^0(\mathbb{R}^+, L^1(\mathbb{R}^2))$ solution n of (1) such that $n(1 + |x|^2)$, $n \log n$ are bounded in $L^\infty_{\text{loc}}(\mathbb{R}^+, L^1(\mathbb{R}^2))$, $\nabla \sqrt{n} \in L^1_{\text{loc}}(\mathbb{R}^+, L^2(\mathbb{R}^2))$ and $\nabla c \in L^\infty_{\text{loc}}(\mathbb{R}^+ \times \mathbb{R}^2)$. Then*

$$\frac{d}{dt} F[n(\cdot, t)] = - \int_{\mathbb{R}^2} n |\nabla(\log n) - \chi \nabla c|^2 dx. \quad (8)$$

Following a usual denomination in the PDE literature, we will call $\int_{\mathbb{R}^2} n |\nabla(\log n) - \chi \nabla c|^2 dx$ the *free energy production term* or *generalized relative Fisher information*.

Proof. Because the potential energy term $\int_{\mathbb{R}^2} n c dx = \iint_{\mathbb{R}^2 \times \mathbb{R}^2} n(x, t) n(y, t) \log |x - y| dx dy$ is quadratic in n , using Equation (1), the time derivative of $F[n(\cdot, t)]$ is given by

$$\frac{d}{dt} F[n(\cdot, t)] = \int_{\mathbb{R}^2} \left[(1 + \log n - \chi c) \nabla \cdot \left(\frac{\nabla n}{n} - \chi \nabla c \right) \right] dx.$$

An integration by parts completes the proof. \square

From the representation (6) of the solution to the Poisson equation, we deduce that

$$\frac{d}{dt} F[n(\cdot, t)] = \frac{d}{dt} \left[\int_{\mathbb{R}^2} n \log n dx + \frac{\chi}{4\pi} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} n(x, t) n(y, t) \log |x - y| dx dy \right] \leq 0.$$

On the other hand, we recall the logarithmic Hardy-Littlewood-Sobolev inequality.

Lemma 6 [15, 4] *Let f be a non-negative function in $L^1(\mathbb{R}^2)$ such that $f \log f$ and $f \log(1 + |x|^2)$ belong to $L^1(\mathbb{R}^2)$. If $\int_{\mathbb{R}^2} f dx = M$, then*

$$\int_{\mathbb{R}^2} f \log f dx + \frac{2}{M} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) f(y) \log |x - y| dx dy \geq -C(M), \quad (9)$$

with $C(M) := M(1 + \log \pi - \log M)$.

This allows to prove a priori estimates on the two terms involved in the free energy.

Lemma 7 *Consider a non-negative $C^0(\mathbb{R}^+, L^1(\mathbb{R}^2))$ solution n of (1) such that $n(1 + |x|^2)$, $n \log n$ are bounded in $L^\infty_{\text{loc}}(\mathbb{R}^+, L^1(\mathbb{R}^2))$, $\int_{\mathbb{R}^2} \frac{1+|x|}{|x-y|} n(y, t) dy \in L^\infty((0, T) \times \mathbb{R}^2)$, $\nabla \sqrt{n} \in L^1_{\text{loc}}(\mathbb{R}^+, L^2(\mathbb{R}^2))$ and $\nabla c \in L^\infty_{\text{loc}}(\mathbb{R}^+ \times \mathbb{R}^2)$. If $\chi M \leq 8\pi$, then the following estimates hold:*

(i) Entropy:

$$M \log M - M \log[\pi(1 + t)] - K \leq \int_{\mathbb{R}^2} n \log n dx \leq \frac{8\pi F_0 + \chi M C(M)}{8\pi - \chi M}$$

with $K := \max \left\{ \int_{\mathbb{R}^2} |x|^2 n_0(x) dx, \frac{M}{2\pi} (8\pi - \chi M) \right\}$ and $F_0 := F[n_0]$.

(ii) Fisher information: For all $t > 0$, with $C_1 := F_0 + \frac{\chi M}{8\pi} C(M)$ and $C_2 := \frac{\chi M - 8\pi}{8\pi}$,

$$0 \leq \int_0^t ds \int_{\mathbb{R}^2} n(x, s) |\nabla (\log n(x, s)) - \chi \nabla c(x, s)|^2 dx \leq C_1 + C_2 \left[M \log \left(\frac{\pi(1+t)}{M} \right) + K \right]$$

Proof. From (8), with $n(\cdot) = n(\cdot, t)$ for shortness, we get that the quantity

$$(1 - \theta) \int_{\mathbb{R}^2} n \log n dx + \theta \left[\int_{\mathbb{R}^2} n \log n dx + \frac{\chi}{4\pi\theta} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} n(x) n(y) \log |x - y| dx dy \right]$$

is bounded from above by F_0 . We choose

$$\frac{\chi}{4\pi\theta} = \frac{2}{M} \iff \theta = \frac{\chi M}{8\pi}$$

and apply (9):

$$(1 - \theta) \int_{\mathbb{R}^2} n(x, t) \log n(x, t) dx - \theta C(M) \leq F_0.$$

If $\chi M < 8\pi$, then $\theta < 1$ and

$$\int_{\mathbb{R}^2} n(x, t) \log n(x, t) dx \leq \frac{F_0 + \theta C(M)}{1 - \theta}.$$

This estimate proves the upper bound for the entropy. We can also see that $\int_{\mathbb{R}^2} n \log n dx$ is bounded from below. By Lemma 3,

$$\frac{1}{1+t} \int_{\mathbb{R}^2} |x|^2 n(x, t) dx \leq K \quad \forall t > 0.$$

Thus

$$\begin{aligned} \int_{\mathbb{R}^2} n(x, t) \log n(x, t) &\geq \frac{1}{1+t} \int_{\mathbb{R}^2} |x|^2 n(x, t) dx - K + \int_{\mathbb{R}^2} n(x, t) \log n(x, t) dx \\ &= \int_{\mathbb{R}^2} n(x, t) \log \left(\frac{n(x, t)}{e^{-\frac{|x|^2}{1+t}}} \right) dx - K \\ &= \int_{\mathbb{R}^2} n(x, t) \log \left(\frac{n(x, t)}{\mu(x, t)} \right) dx - M \log[\pi(1+t)] - K \\ &= \int_{\mathbb{R}^2} \frac{n(x, t)}{\mu(x, t)} \log \left(\frac{n(x, t)}{\mu(x, t)} \right) \mu(x, t) dx - M \log[\pi(1+t)] - K \end{aligned}$$

with $\mu(x, t) := \frac{1}{\pi(1+t)} e^{-\frac{|x|^2}{1+t}}$. By Jensen's inequality,

$$\int_{\mathbb{R}^2} \frac{n(x, t)}{\mu(x, t)} \log \left(\frac{n(x, t)}{\mu(x, t)} \right) \mu(x, t) dx \geq X \log X \quad \text{where} \quad X = \int_{\mathbb{R}^2} \frac{n(x, t)}{\mu(x, t)} \mu(x, t) dx = M.$$

This gives the lower estimate for the entropy term.

Now, from (8) and (9), we get

$$(1 - \theta) \left[M \log \left(\frac{M}{\pi(1+t)} \right) - K \right] + \theta C(M) + \int_0^t ds \int_{\mathbb{R}^2} n(x, s) |\nabla (\log n(x, s)) - \chi \nabla c(x, s)|^2 dx \leq F_0 .$$

This proves that $\sqrt{n} |\nabla (\log n) - \chi \nabla c|$ is bounded in $L^2_{\text{loc}}(\mathbb{R}^+, L^2(\mathbb{R}^2))$ and gives the estimate on the energy. \square

The *a priori* upper bound on $\int_{\mathbb{R}^2} n \log n dx$ combined with the $|x|^2$ moment bound of Lemma 3 shows that $n \log n$ is bounded in $L^\infty_{\text{loc}}(\mathbb{R}^+, L^1(\mathbb{R}^2))$.

Lemma 8 *For any $u \in L^1_+(\mathbb{R}^2)$, if $\int_{\mathbb{R}^2} |x|^2 u dx$ and $\int_{\mathbb{R}^2} u \log u dx$ are bounded from above, then $u \log u$ is uniformly bounded in $L^\infty(\mathbb{R}^2_{\text{loc}}, L^1(\mathbb{R}^2))$ and*

$$\int_{\mathbb{R}^2} u |\log u| dx \leq \int_{\mathbb{R}^2} u \left(\log u + |x|^2 \right) dx + 2 \log(2\pi) \int_{\mathbb{R}^2} u dx + \frac{2}{e} .$$

Proof. The proof goes as follows. Let $\bar{u} := u \mathbb{1}_{\{u \leq 1\}}$ and $m = \int_{\mathbb{R}^2} \bar{u} dx \leq M$. Then

$$\int_{\mathbb{R}^2} \bar{u} \left(\log \bar{u} + \frac{1}{2} |x|^2 \right) dx = \int_{\mathbb{R}^2} U \log U d\mu - m \log(2\pi)$$

where $U := \bar{u}/\mu$, $d\mu(x) = \mu(x) dx$ and $\mu(x) = (2\pi)^{-1} e^{-|x|^2/2}$. By Jensen's inequality,

$$\int_{\mathbb{R}^2} U \log U d\mu \geq \left(\int_{\mathbb{R}^2} U d\mu \right) \log \left(\int_{\mathbb{R}^2} U d\mu \right) = m \log m ,$$

$$\int_{\mathbb{R}^2} \bar{u} \log \bar{u} dx \geq m \log \left(\frac{m}{2\pi} \right) - \frac{1}{2} \int_{\mathbb{R}^2} |x|^2 \bar{u} dx \geq -\frac{1}{e} - M \log(2\pi) - \frac{1}{2} \int_{\mathbb{R}^2} |x|^2 \bar{u} dx .$$

Using

$$\int_{\mathbb{R}^2} u |\log u| dx = \int_{\mathbb{R}^2} u \log u dx - 2 \int_{\mathbb{R}^2} \bar{u} \log \bar{u} dx ,$$

this completes the proof. \square

2.4 Existence of weak solutions up to critical mass

Using the informations collected in Sections 2.1, 2.2 and 2.3, in the spirit of [18], we can now state, in the subcritical case $M < 8\pi/\chi$, the following existence result of weak solutions, which is essentially the one stated without proof in [21].

Proposition 9 *Under Assumption (4) and $M < 8\pi/\chi$, the Keller-Segel system (1) has a global weak non-negative solution such that, for any $T > 0$,*

$$(1 + |x|^2 + |\log n|) n \in L^\infty(0, T; L^1(\mathbb{R}^2)) \quad \text{and} \quad \iint_{[0, T] \times \mathbb{R}^2} n |\nabla \log n - \chi \nabla c|^2 dx dt < \infty .$$

Proposition 9 strongly relies on the estimates of Lemmata 3 and 7. To establish a complete proof, we need to regularize the problem (Section 2.5) and then prove that the above estimates hold uniformly with respect to the regularization procedure (Section 2.6). This allows to pass to the limit in the regularization parameter (Section 2.7) and proves the existence of a weak solution with a well defined flux. To prove Theorem 1, we need additional regularity properties of the solutions. This is the purpose of Section 3. Hypercontractivity and the free energy inequality (2) will be dealt with in Sections 3.4 and 3.3 respectively.

2.5 A regularized model

The goal of this section is to establish the existence of solutions for a regularized version of the Keller-Segel model, for which the logarithmic singularity of the convolution kernel $\mathcal{K}^0(z) := -\frac{1}{2\pi} \log |z|$ is appropriately truncated.

There are indeed two difficulties when dealing with \mathcal{K}^0 . It is unbounded and has a singularity at $z = 0$. First of all, the unboundedness from above of the kernel is not difficult to handle. For $R > \sqrt{e}$, $R \mapsto R^2 / \log R$ is an increasing function, so that

$$0 \leq \iint_{|x-y|>R} \log |x-y| n(x,t) n(y,t) dx dy \leq \frac{2 \log R}{R^2} M \int_{\mathbb{R}^2} |x|^2 n(x,t) dx .$$

Since $\iint_{1<|x-y|<R} \log |x-y| n(x,t) n(y,t) dx dy \leq M^2 \log R$, we only need to take care of a uniform bound on

$$\iint_{|x-y|<1} \log |x-y| n^\varepsilon(x,t) n^\varepsilon(y,t) dx dy \quad \text{and} \quad \int_{\mathbb{R}^2} n^\varepsilon(x,t) \log n^\varepsilon(x,t) dx .$$

for an approximating family $(n^\varepsilon)_{\varepsilon>0}$.

The other difficulty concerning the convolution kernel \mathcal{K}^0 is the singularity at $z = 0$. This is a much more serious difficulty that we are going to overcome by defining a truncated convolution kernel and deriving uniform estimates in Section 2.6. To do so, we first need to find solutions of the model with a truncated convolution kernel. Let \mathcal{K}^ε be such that

$$\mathcal{K}^\varepsilon(z) := \mathcal{K}^1\left(\frac{z}{\varepsilon}\right)$$

where \mathcal{K}^1 is a radial monotone non-decreasing smooth function satisfying

$$\begin{cases} \mathcal{K}^1(z) = -\frac{1}{2\pi} \log |z| & \text{if } |z| \geq 4, \\ \mathcal{K}^1(z) = 0 & \text{if } |z| \leq 1. \end{cases}$$

Moreover, we can assume without restriction that

$$0 \leq -\nabla \mathcal{K}^1(z) \leq \frac{1}{2\pi |z|}, \quad \mathcal{K}^1(z) \leq -\frac{1}{2\pi} \log |z| \quad \text{and} \quad -\Delta \mathcal{K}^1(z) \geq 0 \quad (10)$$

for any $z \in \mathbb{R}^2$. Since $\mathcal{K}^\varepsilon(z) = \mathcal{K}^1(z/\varepsilon)$, we also have

$$0 \leq -\nabla \mathcal{K}^\varepsilon(z) \leq \frac{1}{2\pi |z|} \quad \forall z \in \mathbb{R}^2. \quad (11)$$

If we replace (1) by the following regularized version

$$\begin{cases} \frac{\partial n^\varepsilon}{\partial t} = \Delta n^\varepsilon - \chi \nabla \cdot (n^\varepsilon \nabla c^\varepsilon) \\ c^\varepsilon = \mathcal{K}^\varepsilon * n^\varepsilon \end{cases} \quad x \in \mathbb{R}^2, t > 0, \quad (12)$$

written in the distribution sense, then we can state the following existence result.

Proposition 10 *For any fixed positive ε , under Assumptions (4), if $n_0 \in L^2(\mathbb{R}^2)$, then for any $T > 0$ there exists $n^\varepsilon \in L^2(0, T; H^1(\mathbb{R}^2)) \cap C(0, T; L^2(\mathbb{R}^2))$ which solves (12) with initial data n_0 .*

To prove Proposition 10, we will first fix a functional framework, then solve a linear problem before using it to make a fixed point argument in order to prove the existence of a solution to the regularized system (12).

2.5.1 Functional framework

We will use the Aubin-Lions *compactness method*, (see [39], Ch. IV, §4 and [2], and [56] for more recent references). A simple statement goes as follows:

Lemma 11 (Aubin Lemma) *Take $T > 0$, $p \in (1, \infty)$ and let $(f_n)_{n \in \mathbb{N}}$ be a bounded sequence of functions in $L^p(0, T; H)$ where H is a Banach space. If $(f_n)_{n \in \mathbb{N}}$ is bounded in $L^p(0, T; V)$, where V is compactly imbedded in H and $\partial f_n / \partial t$ is bounded in $L^p(0, T; V')$ uniformly with respect to $n \in \mathbb{N}$, then $(f_n)_{n \in \mathbb{N}}$ is relatively compact in $L^p(0, T; H)$.*

For our purpose, we fix $T > 0$, $p = 2$ and define $H := L^2(\mathbb{R}^2)$, $V := \{v \in H^1(\mathbb{R}^2) : \sqrt{|x|} v \in L^2(\mathbb{R}^2)\}$, V' its dual, $\mathcal{V} := L^2(0, T; V)$, $\mathcal{H} := L^2(0, T; H)$ and $\mathcal{W}(0, T) := \{v \in L^2(0, T; V) : \partial v / \partial t \in L^2(0, T; V')\}$. In this functional framework, the notion of solution we are looking for is actually more precise than in the distribution sense:

$$0 = \int_0^T \left\{ \langle n_t, \psi \rangle_{V' \times V} + \int_{\mathbb{R}^3} (\nabla n + \chi n \nabla c) \cdot \nabla \psi \, dx \right\} dt \quad \forall \psi \in \mathcal{V}.$$

Notice that V is relatively compact in H , since the bound on $|x| |v|^2$ in $L^1(\mathbb{R}^2)$ allows to consider only compact sets, on which compactness holds by Sobolev's imbeddings: Lemma 11 applies.

2.5.2 Estimates for a linear drift-diffusion equation

We start with the derivation of some *a priori* estimates on the solution of the linear problem

$$\frac{\partial n}{\partial t} = \Delta n - \nabla \cdot (n f) \quad (13)$$

for some function $f \in (L^\infty((0, T) \times \mathbb{R}^2))^2$. We assume in this section that the initial data n_0 belongs to $L^2(\mathbb{R}^2)$. By a fixed-point method, this allows us to prove the

Lemma 12 *Assume that (4) holds and consider $f \in L^\infty((0, T) \times \mathbb{R}^2)$. If $n_0 \in L^2(\mathbb{R}^2)$, for any $T > 0$, there exists $n \in \mathcal{W}(0, T)$ which solves (13) with initial data n_0 .*

Proof. Consider the map $\mathcal{T} : L^\infty(0, T; L^1(\mathbb{R}^2)) \rightarrow L^\infty(0, T; L^1(\mathbb{R}^2))$ defined by

$$\mathcal{T}[n](\cdot, t) := G(\cdot, t) * n_0 + \int_0^t \nabla G(\cdot, t-s) * [n(\cdot, s) f(\cdot, s)] ds \quad \forall (x, t) \in [0, T] \times \mathbb{R}^2,$$

where $*$ denotes the space convolution. Here $G(x, t) := (4\pi t)^{-1} e^{-\frac{|x|^2}{4t}}$ is the Green function associated to the heat equation. Notice that $\|\nabla G(\cdot, s)\|_{L^1(\mathbb{R}^2)} \leq C s^{-1/2}$. We define the sequence $(n_k)_{k \in \mathbb{N}}$ by $n_{k+1} = \mathcal{T}(n_k)$ for $k \geq 1$. For any $t \in [0, T]$, we compute

$$\begin{aligned} \|n_{k+1}(t) - n_k(t)\|_{L^1(\mathbb{R}^2)} &\leq \int_{\mathbb{R}^2} \left| \int_0^t \nabla G(\cdot, t-s) * [(n_k(\cdot, s) - n_{k-1}(\cdot, s)) f(\cdot, s)] ds \right| dx \\ &\leq \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)} \int_0^t \|\nabla G(\cdot, t-s) * (n_k(\cdot, s) - n_{k-1}(\cdot, s))\|_{L^1(\mathbb{R}^2)} ds \\ &\leq \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)} \int_0^t \|\nabla G(\cdot, t-s)\|_{L^1(\mathbb{R}^2)} \|n_k(s) - n_{k-1}(s)\|_{L^1(\mathbb{R}^2)} ds \\ &\leq C \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)} \sqrt{t} \|n_k - n_{k-1}\|_{L^\infty(0, t; L^1(\mathbb{R}^2))}. \end{aligned}$$

For $T > 0$ small enough, $(n_k)_{k \in \mathbb{N}}$ is a Cauchy sequence in $L^\infty(0, T; L^1(\mathbb{R}^2))$, which converges to a fixed point of \mathcal{T} . Iterating the method, we prove the existence of a solution of (13) on an arbitrary time interval $[0, T]$. \square

Next, let us establish some *a priori* estimates. The solution n is bounded in $L^\infty(0, T; L^2(\mathbb{R}^2))$ as a consequence of the following computation:

$$\frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^2} |n(x, t)|^2 dx = - \int_{\mathbb{R}^2} |\nabla n(x, t)|^2 dx + \int_{\mathbb{R}^2} \nabla n(x, t) \cdot n(x, t) f(x, t) dx.$$

The right hand side can be written $\int_{\mathbb{R}^2} a \cdot b f dx$ with $a := \sqrt{1/\lambda} \nabla n$ and $b := \sqrt{\lambda} n$. It is therefore bounded by $(\int_{\mathbb{R}^2} a^2 dx + \frac{1}{4} \int_{\mathbb{R}^2} b^2 dx) \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)}$, which provides the estimate

$$\frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^2} |n|^2 dx \leq \left(-1 + \frac{1}{\lambda} \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)} \right) \int_{\mathbb{R}^2} |\nabla n|^2 dx + \frac{\lambda}{4} \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)} \int_{\mathbb{R}^2} |n|^2 dx.$$

In case $\lambda = \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)}$, we obtain

$$\int_{\mathbb{R}^2} |n|^2 dx \leq \int_{\mathbb{R}^2} |n_0|^2 dx e^{\|f\|_{L^\infty([0, T] \times \mathbb{R}^2)}^2 T/2} \quad \forall t \in (0, T).$$

Hence n is bounded in $L^\infty(0, T; L^2(\mathbb{R}^2)) = \mathcal{H}$. Similarly, in case $\lambda = \frac{3}{2} \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)}$, we obtain

$$\frac{1}{2} \frac{d}{dt} \|n(\cdot, t)\|_{L^2(\mathbb{R}^2)}^2 \leq -\frac{1}{3} \|\nabla n\|_{L^2(\mathbb{R}^2)}^2 + \frac{3}{8} \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)} \|n(\cdot, t)\|_{L^2(\mathbb{R}^2)}^2.$$

This also proves that ∇n is bounded in $L^2((0, T) \times \mathbb{R}^2)$, and n is therefore also bounded in $L^2(0, T; H^1(\mathbb{R}^2))$. Next, we need a moment estimate, which is achieved by

$$\frac{d}{dt} \int_{\mathbb{R}^2} |x|^2 n(x, t) dx \leq 4 \int_{\mathbb{R}^2} n dx + 2 \|f\|_{L^\infty([0, T] \times \mathbb{R}^2)} \left(\int_{\mathbb{R}^2} n dx \right)^{1/2} \left(\int_{\mathbb{R}^2} |x|^2 n(x, t) dx \right)^{1/2}.$$

As a conclusion, this proves that $\int_{\mathbb{R}^2} |x|^2 n(x, t) dx$ is bounded and therefore shows that n is bounded in \mathcal{V} . On the other hand, $\partial n / \partial t$ is bounded in V' as can be checked by an elementary computation. We can therefore apply Aubin's Lemma (Lemma 11) to n :

If $(n_0^k)_{k \in \mathbb{N}}$ is a sequence of initial data with uniform bounds, then the corresponding sequence $(n^k)_{k \in \mathbb{N}}$ of solutions of (13) with f replaced by f_k , for a sequence $(f^k)_{k \in \mathbb{N}}$ uniformly bounded in $(L^\infty([0, T] \times \mathbb{R}^2))^2$, is contained in a relatively compact set in $L^2(0, T; V)$.

We will make use of this property in the next section.

2.5.3 Existence of a solution of the regularized problem

This section is devoted to the proof of Proposition 10, using a fixed point method.

Define the truncation function $h(s) := \min\{1, h_0/s\}$, for some constant $h_0 > 1$ to be fixed later and consider the map $\mathcal{T} : L^2(0, T; H) \rightarrow L^2(0, T; H)$ such that

1. To a function $n \in L^2(0, T; H)$, we associate $\nabla c^\varepsilon := \nabla \mathcal{K}^\varepsilon * n$.
2. With ∇c^ε , we construct the truncated function

$$f := h\left(\|\nabla c^\varepsilon\|_{L^\infty((0, T) \times \mathbb{R}^2)}\right) \nabla c^\varepsilon.$$

3. The function f is bounded in $L^\infty((0, T) \times \mathbb{R}^2)$ by h_0 , so we may apply Lemma 12 and obtain a new function $\tilde{n} =: \mathcal{T}[n]$ which solves (13).

The continuity of \mathcal{T} is straightforward. As noticed in Section 2.5.2, we may apply the Aubin-Lions *compactness method*, which gives enough compactness to apply Schauder's fixed point theorem (Theorem 8.1 p. 199 in [38]) to a ball in $\mathcal{W}(0, T)$. Hence we obtain a solution of

$$\begin{cases} \frac{\partial n^\varepsilon}{\partial t} = \Delta n^\varepsilon - \chi \nabla \cdot (n^\varepsilon f^\varepsilon) \\ f^\varepsilon = h\left(\|\nabla c^\varepsilon\|_{L^\infty((0, T) \times \mathbb{R}^2)}\right) \nabla c^\varepsilon, \quad c^\varepsilon = \mathcal{K}^\varepsilon * n^\varepsilon. \end{cases}$$

Assuming that $h_0 > \|\nabla \mathcal{K}^\varepsilon\|_{L^\infty(\mathbb{R}^2)} \|n_0\|_{L^1(\mathbb{R}^2)}$, we realize that n^ε is a solution of (12). \square

Notice that one can also easily prove a uniqueness result, using an appropriate Gronwall lemma. We refer for instance to [52] for similar results in a ball.

2.6 Uniform a priori estimates

In this section, we prove *a priori* estimates for the regularized problem which are uniform with respect to the regularization parameter ε . These estimates correspond to the formal estimates of Section 2.3.

Lemma 13 *Under Assumption (4), consider a solution n^ε of (12). If $\chi M < 8\pi$ then, uniformly as $\varepsilon \rightarrow 0$, with bounds depending only upon $\int_{\mathbb{R}^2} (1 + |x|^2) n_0 dx$ and $\int_{\mathbb{R}^2} n_0 \log n_0 dx$, we have:*

- (i) *The function $(t, x) \mapsto |x|^2 n^\varepsilon(x, t)$ is bounded in $L^\infty(\mathbb{R}_{\text{loc}}^+; L^1(\mathbb{R}^2))$.*
- (ii) *The functions $t \mapsto \int_{\mathbb{R}^2} n^\varepsilon(x, t) \log n^\varepsilon(x, t) dx$ and $t \mapsto \int_{\mathbb{R}^2} n^\varepsilon(x, t) c^\varepsilon(x, t) dx$ are bounded.*

- (iii) The function $(t, x) \mapsto n^\varepsilon(x, t) \log(n^\varepsilon(x, t))$ is bounded in $L^\infty(\mathbb{R}_{\text{loc}}^+; L^1(\mathbb{R}^2))$.
- (iv) The function $(t, x) \mapsto \nabla \sqrt{n^\varepsilon}(x, t)$ is bounded in $L^2(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$.
- (v) The function $(t, x) \mapsto n^\varepsilon(x, t)$ is bounded in $L^2(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$.
- (vi) The function $(t, x) \mapsto n^\varepsilon(x, t) \Delta c^\varepsilon(x, t)$ is bounded in $L^1(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$.
- (vii) The function $(t, x) \mapsto \sqrt{n^\varepsilon}(x, t) \nabla c^\varepsilon(x, t)$ is bounded in $L^2(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$.

Proof.

(i) The integral $\int_{\mathbb{R}^2} |x|^2 n^\varepsilon(x, t) dx$ can be estimated as in the proof of Lemma 3 because \mathcal{K}^ε is radial and satisfies (11), so

$$\begin{aligned} \frac{d}{dt} \int_{\mathbb{R}^2} |x|^2 n^\varepsilon(x, t) dx &= 4M + 2\chi \int \int_{\mathbb{R}^2 \times \mathbb{R}^2} n^\varepsilon(x, t) n^\varepsilon(y, t) x \cdot \nabla \mathcal{K}^\varepsilon(x - y) dx dy \\ &= 4M + \chi \int \int_{\mathbb{R}^2 \times \mathbb{R}^2} n^\varepsilon(x, t) n^\varepsilon(y, t) (x - y) \nabla \mathcal{K}^\varepsilon(x - y) dx dy \\ &\leq 4M - \frac{\chi}{2\pi} \int \int_{\mathbb{R}^2 \times \mathbb{R}^2} \frac{n^\varepsilon(x, t) n^\varepsilon(y, t)}{|x - y|} dx dy \leq 4M. \end{aligned}$$

(ii) We compute

$$\frac{d}{dt} \left[\int_{\mathbb{R}^2} n^\varepsilon \log n^\varepsilon dx - \frac{\chi}{2} \int_{\mathbb{R}^2} n^\varepsilon c^\varepsilon dx \right] = - \int_{\mathbb{R}^2} n^\varepsilon |\nabla(\log n^\varepsilon) - \chi \nabla c^\varepsilon|^2 dx.$$

Then by (10) and the logarithmic Hardy-Littlewood-Sobolev inequality, see Lemma 6, it follows by Lemma 7 that both terms of the right hand side are uniformly bounded.

(iii) It is a direct consequence of Lemma 8.

(iv) A simple computation shows that

$$\frac{d}{dt} \int_{\mathbb{R}^2} n^\varepsilon \log n^\varepsilon dx \leq -4 \int_{\mathbb{R}^2} |\nabla \sqrt{n^\varepsilon}|^2 dx + \chi \int_{\mathbb{R}^2} n^\varepsilon \cdot (-\Delta c^\varepsilon) dx.$$

Up to the common factor χ , we can write the last term of the right hand side as

$$\int_{\mathbb{R}^2} n^\varepsilon \cdot (-\Delta c^\varepsilon) dx = \int_{\mathbb{R}^2} n^\varepsilon \cdot (-\Delta(\mathcal{K}^\varepsilon * n^\varepsilon)) dx = \text{(I)} + \text{(II)} + \text{(III)}$$

with

$$\text{(I)} := \int_{n^\varepsilon < K} n^\varepsilon \cdot (-\Delta(\mathcal{K}^\varepsilon * n^\varepsilon)), \quad \text{(II)} := \int_{n^\varepsilon \geq K} n^\varepsilon \cdot (-\Delta(\mathcal{K}^\varepsilon * n^\varepsilon)) - \text{(III)} \quad \text{and} \quad \text{(III)} = \int_{n^\varepsilon \geq K} |n^\varepsilon|^2.$$

We define ϕ_1 such that

$$\frac{1}{\varepsilon^2} \phi_1 \left(\frac{\cdot}{\varepsilon} \right) = -\Delta \mathcal{K}^\varepsilon.$$

This gives an easy estimate of (I), namely

$$\text{(I)} \leq \int_{n^\varepsilon < K} K \int_{\mathbb{R}^2} \frac{1}{\varepsilon^2} \phi_1 \left(\frac{x - y}{\varepsilon} \right) n^\varepsilon(y) dy dx = M K.$$

Notice that

$$\frac{1}{\varepsilon^2} \phi_1 \left(\frac{\cdot}{\varepsilon} \right) = -\Delta \mathcal{K}^\varepsilon \rightharpoonup \delta \quad \text{in } \mathcal{D}' , \quad (14)$$

which heuristically explains why (II) should be small. Let us prove that this is indeed the case. By (10), ϕ_1 is non-negative. Using $\|\phi_1\|_{L^1(\mathbb{R}^2)} = 1$, we get

$$\begin{aligned} \text{(II)} &= \int_{n^\varepsilon \geq K} n^\varepsilon(x, t) \int_{\mathbb{R}^2} [n^\varepsilon(x - \varepsilon y, t) - n^\varepsilon(x, t)] \phi_1(y) dy dx \\ &\leq \int_{n^\varepsilon \geq K} n^\varepsilon(x, t) \int_{\mathbb{R}^2} \left[\sqrt{n^\varepsilon(x - \varepsilon y, t)} - \sqrt{n^\varepsilon(x, t)} \right] \sqrt{\phi_1(y)} \\ &\quad \times \left[\sqrt{n^\varepsilon(x - \varepsilon y, t)} - \sqrt{n^\varepsilon(x, t)} + 2\sqrt{n^\varepsilon(x, t)} \right] \sqrt{\phi_1(y)} dy dx . \end{aligned}$$

By the Cauchy-Schwarz inequality and using $(a + 2b)^2 \leq 2a^2 + 8b^2$ we obtain

$$\begin{aligned} \text{(II)} &\leq \int_{n^\varepsilon \geq K} n^\varepsilon(x, t) \left[\|\phi_1\|_{L^\infty(\mathbb{R}^2)} \int_{1/2 \leq y \leq 2} \left| \sqrt{n^\varepsilon(x - \varepsilon y, t)} - \sqrt{n^\varepsilon(x, t)} \right|^2 dy \right]^{1/2} \\ &\quad \cdot \left[\int_{\mathbb{R}^2} \left[2 \left| \sqrt{n^\varepsilon(x - \varepsilon y, t)} - \sqrt{n^\varepsilon(x, t)} \right|^2 + 8 |n^\varepsilon(x, t)| \right] \phi_1(y) dy \right]^{1/2} dx . \end{aligned}$$

Using the Poincaré inequality we get

$$\begin{aligned} \text{(II)} &\leq \int_{n^\varepsilon \geq K} n^\varepsilon(x, t) \|\phi_1\|_{L^\infty(\mathbb{R}^2)}^{1/2} C_P \|\nabla \sqrt{n^\varepsilon}\|_{L^2(\mathbb{R}^2)} \\ &\quad \cdot \left[\sqrt{2} \|\phi_1\|_{L^\infty(\mathbb{R}^2)}^{1/2} C_P \|\nabla \sqrt{n^\varepsilon}\|_{L^2(\mathbb{R}^2)} + 2\sqrt{2} \sqrt{|n^\varepsilon(x, t)|} \|\phi_1\|_{L^1(\mathbb{R}^2)}^{1/2} \right] dx . \end{aligned}$$

Recall the Gagliardo-Nirenberg-Sobolev inequality (7):

$$\int_{n^\varepsilon \geq K} |n^\varepsilon|^2 dx \leq C_{\text{GNS}}^2 \int_{n^\varepsilon \geq K} \left| \nabla \sqrt{n^\varepsilon} \right|^2 dx \int_{n^\varepsilon \geq K} n^\varepsilon dx .$$

The left hand side can therefore be made as small as desired using:

$$\int_{n^\varepsilon \geq K} n^\varepsilon dx \leq \frac{1}{\log K} \int_{n^\varepsilon \geq K} n^\varepsilon \log n^\varepsilon dx \leq \frac{1}{\log K} \int_{\mathbb{R}^2} n^\varepsilon |\log n^\varepsilon| dx =: \eta(K) ,$$

for $K > 1$, large enough. Then

$$\int_{n^\varepsilon \geq K} |n^\varepsilon|^2 dx \leq \eta(K) C_{\text{GNS}}^2 \left\| \nabla \sqrt{n^\varepsilon} \right\|_{L^2(\mathbb{R}^2)}^2 . \quad (15)$$

By the Cauchy-Schwarz inequality

$$\int_{n^\varepsilon \geq K} |n^\varepsilon(x, t)|^{3/2} dx \leq \left(\int_{n^\varepsilon \geq K} |n^\varepsilon| dx \right)^{1/2} \cdot \left(\int_{n^\varepsilon \geq K} |n^\varepsilon|^2 dx \right)^{1/2} \leq \eta(K) C_{\text{GNS}} \|\nabla \sqrt{n^\varepsilon}\|_{L^2(\mathbb{R}^2)} .$$

From this, it follows that

$$\text{(II)} + \text{(III)} \leq B \eta(K) \|\nabla \sqrt{n^\varepsilon}\|_{L^2(\mathbb{R}^2)}^2$$

with

$$B := C_{\text{GNS}}^2 + \sqrt{2} \|\phi_1\|_{L^\infty(\mathbb{R}^2)} C_P^2 + 2\sqrt{2} \|\phi_1\|_{L^\infty(\mathbb{R}^2)}^{1/2} \|\phi_1\|_{L^1(\mathbb{R}^2)}^{1/2} C_P C_{\text{GNS}}.$$

We can choose K large enough such that $\eta(K) < 4/B$. Collecting the estimates, we have shown that

$$\frac{d}{dt} \int_{\mathbb{R}^2} n^\varepsilon \log n^\varepsilon dx \leq M K + (-4 + B \eta(K)) X(t)$$

with

$$X(t) := \|\nabla \sqrt{n^\varepsilon}(t)\|_{L^2(\mathbb{R}^2)}^2,$$

and so

$$(4 - B \eta) \int_0^T X(s) ds \leq M K T + \int_{\mathbb{R}^2} n_0 \log n_0 dx - \int_{\mathbb{R}^2} n^\varepsilon(x, T) \log n^\varepsilon(x, T) dx.$$

(v) follows from the Gagliardo-Nirenberg-Sobolev inequality (7).

(vi) is a straightforward consequence of (iv). Notice that $-\Delta c^\varepsilon$ is non-negative as a convolution of two non-negative functions ϕ_1 and n^ε .

(vii) A computation shows that

$$\frac{d}{dt} \int_{\mathbb{R}^2} \frac{1}{2} n^\varepsilon c^\varepsilon dx = \int_{\mathbb{R}^2} c^\varepsilon (\Delta n^\varepsilon - \chi \nabla \cdot (n^\varepsilon \nabla c^\varepsilon)) dx = \int_{\mathbb{R}^2} n^\varepsilon \Delta c^\varepsilon dx + \chi \int_{\mathbb{R}^2} n^\varepsilon |\nabla c^\varepsilon|^2 dx.$$

This proves that

$$\iint_{[0, T] \times \mathbb{R}^2} n^\varepsilon |\nabla c^\varepsilon|^2 dx dt \leq \frac{1}{2\chi} \left| \int_{\mathbb{R}^2} n^\varepsilon c^\varepsilon dx - \int_{\mathbb{R}^2} n_0 (\mathcal{K}^\varepsilon * n_0) dx \right| + \frac{1}{\chi} \int_0^T \int_{\mathbb{R}^2} n^\varepsilon (-\Delta c^\varepsilon) dx.$$

The last term of the right hand side is controlled by (vi), while the previous one is bounded by (ii). \square

2.7 Passing to the limit

All estimates of Lemma 13 are uniform in the limit $\varepsilon \rightarrow 0$. The fact that n_0 is assumed to be bounded in $L^2(\mathbb{R}^2)$ in Lemma 10 does not play any role. In this section, n_0 is assumed to satisfy Assumption (4) and we consider the solution n^ε of (12) with a non-negative initial data $n_0^\varepsilon = \min\{n_0, \varepsilon^{-1}\}$. We want to pass simultaneously to the limit as $\varepsilon \rightarrow 0$ in $n_0^\varepsilon \rightarrow n_0$ and in $\mathcal{K}^\varepsilon(z) \rightarrow \mathcal{K}^0(z) = -\frac{1}{2\pi} \log |z|$.

Lemma 14 *Assume that n_0 satisfies Assumption (4) and consider the solution n^ε of (12) with a non-negative initial data $n_0^\varepsilon = \min\{n_0, \varepsilon^{-1}\}$. Then up to the extraction of a sequence ε_k of ε converging to 0, n^{ε_k} converges to a function n solution of (1) in the distribution sense. Furthermore the flux $n |\nabla(\log n) - \chi \nabla c|$ is bounded in $L^1([0, T] \times \mathbb{R}^2)$.*

Proof. Assertion (vii) of Lemma 13 allows to give a sense to the equation in the limit $\varepsilon \searrow 0$. The term which is difficult to handle is $n^\varepsilon \nabla c^\varepsilon$. It is first of all bounded in $L^1([0, T] \times \mathbb{R}^2)$ uniformly with respect to ε , as shown by the Cauchy-Schwarz inequality:

$$\left(\iint_{[0, T] \times \mathbb{R}^2} n^\varepsilon |\nabla c^\varepsilon| dx dt \right)^2 \leq \iint_{[0, T] \times \mathbb{R}^2} n^\varepsilon dx dt \iint_{[0, T] \times \mathbb{R}^2} n^\varepsilon |\nabla c^\varepsilon|^2 dx dt = M T \iint_{[0, T] \times \mathbb{R}^2} n^\varepsilon |\nabla c^\varepsilon|^2 dx dt,$$

where the last term is controlled according to (vii) of Lemma 13.

Actually, $n^\varepsilon \nabla c^\varepsilon$ converges to $n \nabla c$ in the sense of distributions. By the Gagliardo-Nirenberg-Sobolev inequality (7), for any $p > 2$, for $t \in \mathbb{R}^+$ a.e.,

$$\int_{\mathbb{R}^2} |n^\varepsilon|^{p/2} dx \leq \left(C_{\text{GNS}}^{(p)} \right)^{p/2} M \left(\int_{\mathbb{R}^2} |\nabla \sqrt{n^\varepsilon}|^2 dx \right)^{\frac{p}{2}-1},$$

which proves that n^ε is bounded in $L^q(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$ for any $p/2 = q \in [1, +\infty)$, and that, up to the extraction of a sequence $(\varepsilon_k)_{k \in \mathbb{N}}$ which converges to 0, n^{ε_k} weakly converges to n in any $L_{\text{loc}}^q(\mathbb{R}^+ \times \mathbb{R}^2)$, $q \geq 1$. Next,

$$\begin{aligned} \nabla c^{\varepsilon_k} - \nabla c &= - \frac{1}{2\pi} \int_{\mathbb{R}^2} \frac{x-y}{|x-y|^2} (n^{\varepsilon_k}(y, t) - n(y, t)) dy \\ &\quad + \int_{|x-y| \leq 2\varepsilon_k} \left(\frac{1}{\varepsilon_k} \nabla \mathcal{K}^1 \left(\frac{x-y}{\varepsilon_k} \right) + \frac{x-y}{2\pi |x-y|^2} \right) n^{\varepsilon_k}(y, t) dy. \end{aligned}$$

Since $\frac{1}{\varepsilon_k} \nabla \mathcal{K}^1 \left(\frac{z}{\varepsilon_k} \right) + \frac{z}{2\pi |z|^2}$ can be bounded by $\frac{1}{2\pi |z|}$, all terms converge to 0 for almost any $(x, t) \in \mathbb{R}^2 \times \mathbb{R}^+$ and the convergence of n^{ε_k} to n is strong in $L_{\text{loc}}^q(\mathbb{R}^+ \times \mathbb{R}^2)$ for any $q \in (2, \infty)$, which is enough to prove that

$$n^{\varepsilon_k} \nabla c^{\varepsilon_k} \rightharpoonup n \nabla c \quad \text{in } \mathcal{D}'(\mathbb{R}^+ \times \mathbb{R}^2).$$

As a consequence, we also get by weak semi-continuity that

$$\begin{aligned} \iint_{[0, T] \times \mathbb{R}^2} n |\nabla c|^2 dx dt &\leq \liminf_{\varepsilon_k \rightarrow 0} \iint_{[0, T] \times \mathbb{R}^2} n^{\varepsilon_k} |\nabla c^{\varepsilon_k}|^2 dx dt, \\ \iint_{[0, T] \times \mathbb{R}^2} n^2 dx dt &\leq \liminf_{\varepsilon_k \rightarrow 0} \iint_{[0, T] \times \mathbb{R}^2} |n^{\varepsilon_k}|^2 dx dt. \end{aligned}$$

Since the functional $n \mapsto \int_{\mathbb{R}^2} |\nabla \sqrt{n}|^2 dx$ is convex, we also get

$$\iint_{[0, T] \times \mathbb{R}^2} |\nabla \sqrt{n}|^2 dx dt \leq \liminf_{\varepsilon_k \rightarrow 0} \iint_{[0, T] \times \mathbb{R}^2} |\nabla \sqrt{n^{\varepsilon_k}}|^2 dx dt.$$

The proof of the convexity goes as follows. Let $n(\tau) = n_0 + \tau \nu$, $\tau > 0$. Then

$$\frac{d^2}{d\tau^2} \int_{\mathbb{R}^2} |\nabla \sqrt{n(\tau)}|^2 dx \Big|_{\tau=0} = \frac{1}{2n_0^3} \int_{\mathbb{R}^2} |\nu \nabla \sqrt{n_0} - n_0 \nabla \sqrt{\nu}|^2 dx \geq 0.$$

See [5, 6] for more details. Now, since

$$\begin{aligned} &\iint_{[0, T] \times \mathbb{R}^2} n^{\varepsilon_k} |\nabla(\log n^{\varepsilon_k}) - \chi \nabla c^{\varepsilon_k}|^2 dx dt \\ &= 4 \iint_{[0, T] \times \mathbb{R}^2} |\nabla \sqrt{n^{\varepsilon_k}}|^2 dx dt - 2\chi \iint_{[0, T] \times \mathbb{R}^2} |n^{\varepsilon_k}|^2 dx dt + \chi^2 \iint_{[0, T] \times \mathbb{R}^2} n^{\varepsilon_k} |\nabla c^{\varepsilon_k}|^2 dx dt \end{aligned}$$

is bounded uniformly with respect to ε_k by (8),

$$\iint_{[0, T] \times \mathbb{R}^2} n |\nabla(\log n) - \chi \nabla c|^2 dx dt$$

is also finite. Notice that this is not enough to prove that (8) holds if n is a solution of (1), even with an inequality instead of the equality. This is however enough to prove that the flux $n |\nabla(\log n) - \chi \nabla c|$ is bounded in $L^1([0, T] \times \mathbb{R}^2)$, simply by using the Cauchy-Schwarz inequality. This concludes the proof of Lemma 14. \square

As a consequence of the approximation procedure and of Lemma 14, we have also proved Proposition 9. To establish Inequality (2) in Theorem 1, we only need to prove that

$$\iint_{[0, T] \times \mathbb{R}^2} n^2 dx dt = \liminf_{\varepsilon_k \rightarrow 0} \iint_{[0, T] \times \mathbb{R}^2} |n^{\varepsilon_k}|^2 dx dt ,$$

but this requires some additional work on the regularity properties of the solutions of (1).

3 Free energy inequality and regularity properties

In this section, we give some additional regularity properties of the solutions when $\chi M < 8\pi$.

3.1 Weak regularity results

The following result is due to Goudon, see [24].

Theorem 15 [24] *Let $n^\varepsilon : (0, T) \times \mathbb{R}^N \rightarrow \mathbb{R}$ be such that for almost all $t \in (0, T)$, $n^\varepsilon(t)$ belongs to a weakly compact set in $L^1(\mathbb{R}^N)$ for almost any $t \in (0, T)$. If $\partial_t n^\varepsilon = \sum_{|\alpha| \leq k} \partial_x^\alpha g_\varepsilon^{(\alpha)}$ where, for any compact set $K \subset \mathbb{R}^n$,*

$$\limsup_{\substack{|E| \rightarrow 0 \\ E \subset \mathbb{R} \text{ is measurable}}} \left(\sup_{\varepsilon > 0} \int \int_{E \times K} |g_\varepsilon^{(\alpha)}| dt dx \right) = 0 ,$$

then $(n^\varepsilon)_{\varepsilon > 0}$ is relatively compact in $C^0([0, T]; L^1_{\text{weak}}(\mathbb{R}^N))$.

This result immediately applies to the solution of (12).

Corollary 16 *Let n^ε be a solution of (12) with initial data $n_0^\varepsilon = \min\{n_0, \varepsilon^{-1}\}$ such that $n_0(1 + |x|^2 + |\log n_0|) \in L^1(\mathbb{R}^2)$. If n is a solution of (1) with initial data n_0 , such that, for a sequence $(\varepsilon_k)_{k \in \mathbb{N}}$ with $\lim_{k \rightarrow \infty} \varepsilon_k = 0$, $n^{\varepsilon_k} \rightharpoonup n$ in $L^1((0, T) \times \mathbb{R}^2)$, then n belongs to $C^0(0, T; L^1_{\text{weak}}(\mathbb{R}^2))$.*

Proof. We are able to apply Theorem 15 to n^ε where $g_\varepsilon^{(1)} := -\chi n^\varepsilon \nabla c^\varepsilon = -\chi \sqrt{n^\varepsilon} \cdot \sqrt{n^\varepsilon} \nabla c^\varepsilon$ and $g_\varepsilon^{(2)} := n^\varepsilon$. Notice indeed that as a consequence of Lemma 18, we have, uniformly with respect to ε ,

$$\limsup_{t_1 \rightarrow t_2} \sup_{\varepsilon} g_\varepsilon^{(1)} \leq \chi \limsup_{t_1 \rightarrow t_2} M(t_2 - t_1) \int_{t_1}^{t_2} \int_{\mathbb{R}^2} n^\varepsilon |\nabla c^\varepsilon|^2 dx ds = 0 ,$$

$$\limsup_{t_1 \rightarrow t_2} \sup_{\varepsilon} g_\varepsilon^{(2)} \leq \limsup_{t_1 \rightarrow t_2} \int_{t_1}^{t_2} \int_{\mathbb{R}^2} n^\varepsilon dx = 0 .$$

\square

3.2 L^p uniform estimates

Here we prove that if the initial data n_0 is bounded in $L^p(\mathbb{R}^2)$, then it is also the case for the solution $n(\cdot, t)$ for any finite positive time t . By uniform, we mean estimates that hold up to $t = 0$.

Proposition 17 *Assume that (4) and $M < 8\pi/\chi$ hold. If n_0 is bounded in $L^p(\mathbb{R}^2)$ for some $p > 1$, then any solution n of (1) is bounded in $L_{\text{loc}}^\infty(\mathbb{R}^+, L^p(\mathbb{R}^2))$.*

Proof. We formally compute

$$\begin{aligned} \frac{1}{2(p-1)} \frac{d}{dt} \int_{\mathbb{R}^2} |n(x, t)|^p dx &= -\frac{2}{p} \int_{\mathbb{R}^2} |\nabla(n^{p/2})|^2 dx + \chi \int_{\mathbb{R}^2} \nabla(n^{p/2}) \cdot n^{p/2} \cdot \nabla c dx \\ &= -\frac{2}{p} \int_{\mathbb{R}^2} |\nabla(n^{p/2})|^2 dx + \chi \int_{\mathbb{R}^2} n^p (-\Delta c) dx \\ &= -\frac{2}{p} \int_{\mathbb{R}^2} |\nabla(n^{p/2})|^2 dx + \chi \int_{\mathbb{R}^2} n^{p+1} dx . \end{aligned}$$

Using the following Gagliardo-Nirenberg-Sobolev inequality:

$$\int_{\mathbb{R}^2} |v|^{2(1+1/p)} dx \leq K_p \int_{\mathbb{R}^2} |\nabla v|^2 dx \int_{\mathbb{R}^2} |v|^{2/p} dx ,$$

or equivalently, with $n = v^{2/p}$,

$$\int_{\mathbb{R}^2} |n|^{p+1} dx \leq K_p \int_{\mathbb{R}^2} |\nabla(n^{p/2})|^2 dx \int_{\mathbb{R}^2} |n| dx ,$$

we get the estimate

$$\frac{1}{2(p-1)} \frac{d}{dt} \int_{\mathbb{R}^2} n^p dx \leq \int_{\mathbb{R}^2} |\nabla(n^{p/2})|^2 dx \left(-\frac{2}{p} + K_p \chi M \right) ,$$

which proves the decay of $\int_{\mathbb{R}^2} n^p dx$ if $M < \frac{2}{pK_p\chi}$. Otherwise, we can rely on the entropy estimate to get a bound: Let $K > 1$ be a constant, to be chosen later.

$$\int_{\mathbb{R}^2} n^p dx = \int_{n \leq K} n^p dx + \int_{n > K} n^p dx .$$

The first term of the right hand side is bounded by $K^{p-1} M$. Concerning the second one, define first

$$M(K) := \int_{n > K} n dx .$$

Using the fact that $|n \log n|$ is bounded in $L^\infty(\mathbb{R}_{\text{loc}}^+; L^1(\mathbb{R}^2))$, we can estimate $M(K)$ by

$$M(K) \leq \frac{1}{\log K} \int_{n > K} n \log n dx \leq \frac{1}{\log K} \int_{\mathbb{R}^2} |n \log n| dx$$

and choose it arbitrarily small on any given time interval $(0, T)$. Following [32], compute now

$$\begin{aligned}
 & \frac{d}{dt} \int_{\mathbb{R}^2} (n - K)_+^p dx + \frac{4}{p} (p - 1) \int_{\mathbb{R}^2} |\nabla((n - K)_+^{p/2})|^2 dx \\
 &= p \int_{\mathbb{R}^2} (n - K)_+^{p-1} [\Delta n - \chi \nabla(n \nabla c)] dx + \frac{4}{p} (p - 1) \int_{\mathbb{R}^2} |\nabla((n - K)_+^{p/2})|^2 dx \\
 &= -p \chi \int_{\mathbb{R}^2} (n - K)_+^{p-1} [\nabla(n - K) \cdot \nabla c + n \Delta c] dx \\
 &= -\chi \int_{\mathbb{R}^2} (n - K)_+^{p-1} [(n - K)_+ (-\Delta c) - p n (-\Delta c)] dx \\
 &= (p - 1) \chi \int_{\mathbb{R}^2} (n - K)_+^{p+1} dx + (2p - 1) \chi K \int_{\mathbb{R}^2} (n - K)_+^p dx + p \chi K^2 \int_{\mathbb{R}^2} (n - K)_+^{p-1} dx
 \end{aligned}$$

The term involving $\int_{\mathbb{R}^2} (n - K)_+^{p-1} dx$ can be estimated as follows:

$$\begin{aligned}
 \int_{\mathbb{R}^2} (n - K)_+^{p-1} dx &= \int_{K < n \leq K+1} (n - K)_+^{p-1} dx + \int_{n > K+1} (n - K)_+^{p-1} dx, \\
 \int_{K < n \leq K+1} (n - K)_+^{p-1} dx &\leq \int_{K < n \leq K+1} 1 dx \leq \frac{1}{K} \int_{K < n \leq K+1} n dx \leq \frac{M}{K}, \\
 \int_{n > K+1} (n - K)_+^{p-1} dx &\leq \int_{n > K+1} (n - K)_+^p dx \leq \int_{\mathbb{R}^2} (n - K)_+^p dx.
 \end{aligned}$$

By choosing K sufficiently large, we obtain

$$-\frac{4}{p} (p - 1) \int_{\mathbb{R}^2} |\nabla((n - K)_+^{p/2})|^2 dx + (p - 1) \chi \int_{\mathbb{R}^2} (n - K)_+^{p+1} dx \leq 0$$

using again the Gagliardo-Nirenberg-Sobolev inequality but with M replaced by $M(K)$, small. Summarizing, for a fixed interval $(0, T)$ with T arbitrarily large, we have found K such that

$$\frac{d}{dt} \int_{\mathbb{R}^2} (n - K)_+^p dx \leq C_1 \int_{\mathbb{R}^2} (n - K)_+^p dx + C_2$$

for some positive constants C_1 and C_2 . A Gronwall estimate shows that $\int_{\mathbb{R}^2} (n - K)_+^p dx$ is finite on $(0, T)$.

To justify this estimate, one has as above to establish it for the regularized problem and then pass to the limit. This is purely technical but not difficult and we leave it to the reader.

To conclude, we still need to check that the bound on $\int_{\mathbb{R}^2} (n - K)_+^p dx$ is enough to control $\int_{n > K} n^p dx$. Using the estimate

$$x^p \leq \left(\frac{\lambda}{\lambda - 1} \right)^{p-1} (x - 1)^p$$

for any $x \geq \lambda > 1$, we get

$$\begin{aligned}
 \int_{n > K} n^p dx &= \int_{K < n \leq \lambda K} n^p dx + \int_{n > \lambda K} n^p dx \\
 &\leq (\lambda K)^{p-1} M + \left(\frac{\lambda}{\lambda - 1} \right)^{p-1} K^p \int_{n > \lambda K} \left(\frac{n}{K} - 1 \right)^p dx \\
 &\leq (\lambda K)^{p-1} M + \left(\frac{\lambda}{\lambda - 1} \right)^{p-1} \int_{\mathbb{R}^2} (n - K)_+^p dx.
 \end{aligned}$$

□

Notice that very similar estimates have been derived, without the knowledge of the optimal bound $\chi M < 8\pi$, by W. Jäger and S. Luckhaus in [32] in \mathbb{R}^d , $d = 2$ (also see [18, 19] if $d \geq 2$), by working directly in an L^p -framework, instead of the free energy framework.

3.3 The free energy inequality in a regular setting

Using the a priori estimates of the previous section for $p = 2 + \varepsilon$, we can prove that the free energy inequality (2) holds.

Lemma 18 *Let n_0 be in a bounded set in $L^1_+(\mathbb{R}^2, (1 + |x|^2)dx) \cap L^{2+\varepsilon}(\mathbb{R}^2, dx)$, for some $\varepsilon > 0$, eventually small. Then n_0 satisfies Assumption (4), the solution n of (1) found in Theorem 1, with initial data n_0 , is in a compact set in $L^2(\mathbb{R}^2_{\text{loc}} \times \mathbb{R}^2)$ and moreover the free energy production estimate (2) holds:*

$$F[n] + \int_0^t \left(\int_{\mathbb{R}^2} n |\nabla (\log n) - \chi \nabla c|^2 dx \right) ds \leq F[n_0].$$

Proof. We split the proof in three steps.

First Step: n is bounded in $L^2(\mathbb{R}^2_{\text{loc}} \times \mathbb{R}^2)$. To apply Theorem 1, we need to prove that $n_0 \log n_0$ is integrable. By Hölder's inequality we have

$$\|u\|_{L^q(\mathbb{R}^2)} \leq \|u\|_{L^p(\mathbb{R}^2)}^\alpha \|u\|_{L^r(\mathbb{R}^2)}^{1-\alpha}$$

with $\alpha = \frac{p}{q} \frac{r-q}{r-p}$, $p \leq q \leq r$. Take the logarithm of both sides:

$$\alpha \log \left(\frac{\|u\|_{L^q(\mathbb{R}^2)}}{\|u\|_{L^p(\mathbb{R}^2)}} \right) + (\alpha - 1) \log \left(\frac{\|u\|_{L^r(\mathbb{R}^2)}}{\|u\|_{L^q(\mathbb{R}^2)}} \right) \leq 0.$$

Since this inequality trivializes to an equality when $q = p$, we may differentiate it with respect to q at $q = p$ and get that for any $u \in L^p(\mathbb{R}^d) \cap L^r(\mathbb{R}^d)$, $1 \leq p < r < +\infty$, we have

$$\int u^p \log \left(\frac{|u|}{\|u\|_{L^p(\mathbb{R}^2)}} \right) dx \leq \frac{r}{r-p} \|u\|_{L^p(\mathbb{R}^2)}^p \log \left(\frac{\|u\|_{L^r(\mathbb{R}^2)}}{\|u\|_{L^p(\mathbb{R}^2)}} \right).$$

With $u = n_0$, $p = 1$ and $r = 2 + \varepsilon$, by applying Lemma 8, we obtain

$$\begin{aligned} & \int_{\mathbb{R}^2} n_0 |\log n_0| dx \\ & \leq \frac{M}{1+\varepsilon} [(2+\varepsilon) \log(\|n_0\|_{L^{2+\varepsilon}(\mathbb{R}^2)}) - \log M + 2 \log(2\pi)] + \int_{\mathbb{R}^2} |x|^2 n_0 dx + \frac{2}{e} < \infty. \end{aligned}$$

Since $n_0 \in L^1 \cap L^{2+\varepsilon}(\mathbb{R}^2)$, by Hölder's inequality, n_0 is initially in any $L^q(\mathbb{R}^2)$ for all $q \in [1, 2 + \varepsilon]$, and as a special case in $L^2(\mathbb{R}^2)$:

$$\|n_0\|_{L^2(\mathbb{R}^2)}^2 \leq \|n_0\|_{L^1(\mathbb{R}^2)}^{\varepsilon/(1+\varepsilon)} \|n_0\|_{L^{2+\varepsilon}(\mathbb{R}^2)}^{1/(1+\varepsilon)}.$$

Hence by Theorem 1, the solution n of (1) is bounded in $L^\infty(\mathbb{R}_{\text{loc}}^+; L^1 \cap L^{2+\varepsilon}(\mathbb{R}^2))$. As a special case n is bounded in $L^2(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$.

Second Step: ∇n is bounded in $L^2(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$. The following computation

$$\frac{d}{dt} \int_{\mathbb{R}^2} n^2 dx = -2 \int_{\mathbb{R}^2} |\nabla n|^2 dx + 2\chi \int_{\mathbb{R}^2} \nabla n \cdot n \nabla c dx$$

shows that $X := \|\nabla n\|_{L^2((0,T) \times \mathbb{R}^2)}$ satisfies the estimate

$$2X^2 - 2\chi \|n \nabla c\|_{L^\infty(0,T;L^2(\mathbb{R}^2))} X \leq \|n\|_{L^\infty(0,T;L^2(\mathbb{R}^2))}^2 + \|n_0\|_{L^2(\mathbb{R}^2)}^2.$$

This implies that X is bounded if $\|n \nabla c\|_{L^\infty(0,T;L^2(\mathbb{R}^2))}$ is bounded. Let us prove that this is indeed the case. The drift force term takes the form

$$\nabla c(x, t) = \frac{1}{2\pi} \int_{\mathbb{R}^2} \frac{x-y}{|x-y|^2} n(y, t) dy.$$

Since $n_0 \in L^{2+\varepsilon}(\mathbb{R}^2)$, by Theorem 1, the solution n of is bounded in $L^\infty(\mathbb{R}_{\text{loc}}^+; L^{2+\varepsilon}(\mathbb{R}^2))$. As a consequence of the Hardy-Littlewood-Sobolev inequality (see below), for any $(p_1, q_1) \in (2, +\infty) \times (1, 2)$ such that $\frac{1}{p_1} = \frac{1}{q_1} - \frac{1}{2}$, there exists a constant $C = C(p_1) > 0$ such that for almost any $t > 0$,

$$\|\nabla c(\cdot, t)\|_{L^{p_1}(\mathbb{R}^2)} \leq C \|n(\cdot, t)\|_{L^{q_1}(\mathbb{R}^2)}.$$

We can indeed evaluate $\|f * |\cdot|^{-\lambda}\|_{L^{p_1}(\mathbb{R}^d)}$ by

$$\|f * |\cdot|^{-\lambda}\|_{L^{p_1}(\mathbb{R}^d)} = \sup_{\substack{g \in L^{q_1}(\mathbb{R}^d) \\ \|g\|_{L^{q_1}(\mathbb{R}^d)} \leq 1}} \int_{\mathbb{R}^d} (f * |\cdot|^{-\lambda}) g dx$$

with $\frac{1}{p_1} + \frac{1}{q_1} = 1$. The right hand side is bounded, up to a multiplicative constant, by $\|f\|_{L^p(\mathbb{R}^d)}$ according to the Hardy-Littlewood-Sobolev inequality, if $\frac{1}{p} + \frac{1}{q_1} + \frac{\lambda}{d} = 2$ and $0 < \lambda < d$. This inequality, see, e.g., [37], indeed states that: *For all $f \in L^p(\mathbb{R}^d)$, $g \in L^q(\mathbb{R}^d)$, $1 < p, q < \infty$ such that $\frac{1}{p} + \frac{1}{q} + \frac{\lambda}{d} = 2$ and $0 < \lambda < d$, there exists a constant $C = C(p, q, \lambda) > 0$ such that*

$$\left| \int_{\mathbb{R}^d \times \mathbb{R}^d} \frac{1}{|x-y|^\lambda} f(x) g(y) dx dy \right| \leq C \|f\|_{L^p(\mathbb{R}^d)} \|g\|_{L^q(\mathbb{R}^d)}.$$

Applied with $\lambda = 1$, $d = 2$, this proves that $\|n \nabla c\|_{L^\infty(0,T;L^2(\mathbb{R}^2))}$ is bounded.

Applying this estimate with $p_1 = 2(1 + 2/\varepsilon)$ and $q_1 = 2 - \varepsilon/(1 + \varepsilon)$, and using Hölder's inequality, we can write

$$\|n(\cdot, t) \nabla c(\cdot, t)\|_{L^2(\mathbb{R}^2)} \leq \|n(\cdot, t)\|_{L^{2+\varepsilon}(\mathbb{R}^2)} \|\nabla c(\cdot, t)\|_{L^{p_1}(\mathbb{R}^2)} \leq C \|n(\cdot, t)\|_{L^{2+\varepsilon}(\mathbb{R}^2)} \|n(\cdot, t)\|_{L^{q_1}(\mathbb{R}^2)}.$$

which is bounded as $q_1 \in [1, 2 + \varepsilon]$. Thus, if n is a solution of (1), $n \nabla c$ is bounded in $L^\infty(\mathbb{R}_{\text{loc}}^+; L^2(\mathbb{R}^2))$.

Third Step: Compactness. As a consequence of Hölder's inequality with $p := (1 + \varepsilon)/\varepsilon$, $q := 1 + \varepsilon$:

$$\int_{\mathbb{R}^2} |x|^{\frac{2\varepsilon}{1+\varepsilon}} n^2 dx = \int_{\mathbb{R}^2} (n |x|^2)^{\frac{\varepsilon}{1+\varepsilon}} \cdot n^{\frac{2+\varepsilon}{1+\varepsilon}} dx \leq \left(\int_{\mathbb{R}^2} n |x|^2 dx \right)^{\frac{\varepsilon}{1+\varepsilon}} \left(\int_{\mathbb{R}^2} n^{2+\varepsilon} dx \right)^{\frac{1}{1+\varepsilon}},$$

the function $(x, t) \mapsto |x|^{\frac{\varepsilon}{1+\varepsilon}} n$ is bounded in $L^\infty(\mathbb{R}_{\text{loc}}^+; L^2(\mathbb{R}^2))$. The imbedding of the set $V := \{u \in H^1 \cap L_+^1(\mathbb{R}^2) : |x|^{\frac{\varepsilon}{1+\varepsilon}} u \in L^1(\mathbb{R}^2)\}$ into $L^2(\mathbb{R}^2) =: H$ is compact and by the Aubin-Lions compactness method (see Lemma 11) as in Section 2.5, it results that n belongs to a compact set of $L^2(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$.

Let $(n_k)_{k \in \mathbb{N}} := (n^{\varepsilon k})_{k \in \mathbb{N}}$ be an approximating sequence defined as in the proof of Theorem 1. Compared to the results of Lemma 14, we have

$$\begin{aligned} \iint_{[0, T] \times \mathbb{R}^2} |\nabla n|^2 dx dt &\leq \liminf_{k \rightarrow \infty} \iint_{[0, T] \times \mathbb{R}^2} |\nabla n_k|^2 dx dt, \\ \iint_{[0, T] \times \mathbb{R}^2} n |\nabla c|^2 dx dt &\leq \liminf_{k \rightarrow \infty} \iint_{[0, T] \times \mathbb{R}^2} n_k |\nabla c_k|^2 dx dt, \\ \iint_{[0, T] \times \mathbb{R}^2} n^2 dx dt &= \liminf_{k \rightarrow \infty} \iint_{[0, T] \times \mathbb{R}^2} |n_k|^2 dx dt, \end{aligned}$$

where the only difference lies in the last equality, a consequence of the above compactness result. This proves the free energy estimate using

$$\iint_{[0, T] \times \mathbb{R}^2} n |\nabla(\log n) - \chi \nabla c|^2 dx dt = 4 \iint_{[0, T] \times \mathbb{R}^2} |\nabla \sqrt{n}|^2 dx dt + \chi^2 \iint_{[0, T] \times \mathbb{R}^2} n |\nabla c|^2 dx dt - 2\chi \iint_{[0, T] \times \mathbb{R}^2} n^2 dx dt.$$

□

3.4 Hypercontractivity

Much more regularity can actually be achieved as follows. All computations are easy to justify for smooth solutions with sufficient decay at infinity. Up to a regularization step, the final estimates certainly hold if the initial data is bounded in $L^\infty(\mathbb{R}^2)$, which is the case for the regularized problem of Section 2.5 with truncated initial data $n_0^\varepsilon = \min\{n_0, \varepsilon^{-1}\}$. However, we will see that the $L^\infty(\mathbb{R}_{\text{loc}}^+; L^p(\mathbb{R}^2))$ -estimates hold for any $p > 1$ independently of ε , so that we may pass to the limit and get the result for any solution of (1) with initial data satisfying only (4) and $\chi M < 8\pi$. To simplify the presentation of the method, we will therefore do the computations only at a formal level, for smooth solutions which behave well at infinity.

Theorem 19 *Consider a solution n of (1) with initial data n_0 satisfying (4) and $\chi M < 8\pi$. Then for any $p \in (1, \infty)$, there exists a continuous function h_p on $(0, \infty)$ such that for almost any $t > 0$, $\|n(\cdot, t)\|_{L^p(\mathbb{R}^2)} \leq h_p(t)$.*

Notice that unless n_0 is bounded in $L^p(\mathbb{R}^2)$, $\lim_{t \rightarrow 0^+} h_p(t) = +\infty$. Such a result is called an *hypercontractivity* result, see [25], since to an initial data which is originally in $L^1(\mathbb{R}^2)$ but not in $L^p(\mathbb{R}^2)$, we associate a solution which at almost any time $t > 0$ is in $L^p(\mathbb{R}^2)$ with p arbitrarily large.

Proof. Fix $t > 0$ and $p \in (1, \infty)$, and consider $q(s) := 1 + (p-1)\frac{s}{t}$, so that $q(0) = 1$ and $q(t) = p$. Exactly as in the proof of Theorem 1, for an arbitrarily small $\eta > 0$ given in advance, we can find $K > 1$ big enough such that $M(K) := \sup_{s \in (0, t)} \int_{n > K} n(\cdot, s) dx$ is smaller than η . It is indeed sufficient to notice that

$$\int_{n > K} n(\cdot, s) dx \leq \frac{1}{\log K} \int_{n > K} n(\cdot, s) \log n(\cdot, s) dx \leq \frac{1}{\log K} \int_{\mathbb{R}^2} |n(\cdot, s) \log n(\cdot, s)| dx.$$

Since $\chi M < 8\pi$, $n |\log n|$ is bounded in $L^\infty(0, t; L^1(\mathbb{R}^2))$. This proves that for K big enough, we may assume

$$\int_{\mathbb{R}^2} (n - K)_+ dx \leq \eta,$$

for an arbitrarily small $\eta > 0$.

Next, we define

$$F(s) := \left[\int_{\mathbb{R}^2} (n - K)_+^{q(s)}(x, s) dx \right]^{1/q(s)}$$

for the function $s \mapsto q(s)$ defined above. A derivation with respect to s gives

$$F' F^{q-1} = \frac{q'}{q^2} \int_{\mathbb{R}^2} (n - K)_+^q \log \left(\frac{(n - K)_+^q}{F^q} \right) + \int_{\mathbb{R}^2} n_t (n - K)_+^{q-1}.$$

If n is a solution to (1), then

$$\int_{\mathbb{R}^2} (n - K)_+^{q-1} n_t dx = -4 \frac{q-1}{q^2} \int_{\mathbb{R}^2} |\nabla((n - K)_+^{q/2})|^2 dx + \chi \frac{q-1}{q} \int_{\mathbb{R}^2} (n - K)_+^{q+1} dx,$$

and we get

$$F' F^{q-1} = \frac{q'}{q^2} \int_{\mathbb{R}^2} (n - K)_+^q \log \left(\frac{(n - K)_+^q}{F^q} \right) - 4 \frac{q-1}{q^2} \int_{\mathbb{R}^2} |\nabla((n - K)_+^{q/2})|^2 + \chi \frac{(q-1)}{q} \int_{\mathbb{R}^2} (n - K)_+^{q+1}.$$

Using the assumption $q' \geq 0$, we can apply the logarithmic Sobolev inequality [25]

$$\int_{\mathbb{R}^2} v^2 \log \left(\frac{v^2}{\int_{\mathbb{R}^2} v^2 dx} \right) dx \leq 2\sigma \int_{\mathbb{R}^2} |\nabla v|^2 dx - (2 + \log(2\pi\sigma)) \int_{\mathbb{R}^2} v^2 dx$$

for any $\sigma > 0$, and the Gagliardo-Nirenberg-Sobolev inequality

$$\int_{\mathbb{R}^2} |v|^{2(1+1/q)} dx \leq \mathcal{K}(q) \|\nabla v\|_{L^2(\mathbb{R}^2)}^2 \int_{\mathbb{R}^2} |v|^{2/q} dx, \quad \forall q \in [2, \infty)$$

to $v := (n - K)_+^{q/2}$, and get

$$F' F^{q-1} \leq \left(\frac{2\sigma q'}{q^2} - 4 \frac{q-1}{q^2} + \chi \frac{q-1}{q} \mathcal{K}(q) \eta \right) \|\nabla v\|_{L^2(\mathbb{R}^2)}^2 - \frac{q'}{q^2} (2 + \log(2\pi\sigma)) F^q.$$

With the specific choice of $\sigma := (q-1)/q'$ and provided η is chosen small enough in order that

$$-2 \frac{q-1}{q^2} + \chi \frac{q-1}{q} \sup_{r \in (1, p)} [\mathcal{K}(r)] \eta \leq 0,$$

this shows that

$$\frac{F'}{F} \leq -\frac{q'}{q^2} (2 + \log(2\pi\sigma)) =: G(t).$$

The function G is integrable on $(0, t)$, which proves that $F(t)$ can be bounded in terms of $F(0)$.

□

3.5 The free energy inequality for weak solutions

As a consequence of Lemma 18 and Theorem 19, we have the following result.

Corollary 20 *Let $(n^k)_{k \in \mathbb{N}}$ be a sequence of solutions of (1) with initial data n_0^k satisfying Assumption (4) with uniform corresponding bounds. For any $t_0 > 0$, $T \in \mathbb{R}^+$ such that $0 < t_0 < T$, $(n^k)_{k \in \mathbb{N}}$ is relatively compact in $L^2((t_0, T) \times \mathbb{R}^2)$, and if n is the limit of $(n^k)_{k \in \mathbb{N}}$, then n is a solution of (1) such that the free energy inequality (2) holds.*

Proof. By Theorem 19, for $t > t_0 > 0$, n^k is bounded in $L^\infty(t_0, t; L^{2+\varepsilon}(\mathbb{R}^2))$, for any $\varepsilon > 0$. We can therefore apply Lemma 18 with initial data $n^k(\cdot, t_0)$ at $t = t_0$:

$$F[n^k(\cdot, t)] + \int_{t_0}^t \left(\int_{\mathbb{R}^2} n^k |\nabla (\log n^k) - \chi \nabla c^k|^2 dx \right) ds \leq F[n^k(\cdot, t_0)].$$

The compactness in $L^2([t_0, t] \times \mathbb{R}^2)$ follows from Lemma 11. Passing to the limit as $k \rightarrow \infty$, we get

$$F[n(\cdot, t)] + \int_{t_0}^t \left(\int_{\mathbb{R}^2} n |\nabla (\log n) - \chi \nabla c|^2 dx \right) ds \leq F[n(\cdot, t_0)].$$

Since, as a function of s , $\int_{\mathbb{R}^2} n(\cdot, s) |\nabla (\log n(\cdot, s)) - \chi \nabla c(\cdot, s)|^2 dx$ is integrable on $(0, t)$, we can pass to the limit $t_0 \rightarrow 0$. By convexity of $n \mapsto n \log n$, it is easy to check that $\lim_{t_0 \rightarrow 0^+} F[n(\cdot, t_0)] \leq F[n_0]$. \square

Apply Corollary 20 with $n_0^k = \min\{n_0, \varepsilon_k^{-1}\}$ as in the regularization procedure of Section 2.5–2.7. This completes the proof of Theorem 1.

4 Intermediate asymptotics and self-similar solutions

In this section, we investigate the behavior of the solutions as time t goes to infinity and prove Theorem 2. The key tool is the free energy written in rescaled variables, F^R , which is defined below. The main difficulty comes from the fact that the uniqueness of the solutions to (3) for a given $M \in (0, 8\pi/\chi)$ is not known. This is not crucial for the proof of Theorem 2 because, in the self-similar variables, the decay of the entropy selects a unique solution to (3). In this section, we will anyway prove several additional properties (radial symmetry, regularity, decay at infinity) of the solution of (3) and comment on related issues.

4.1 Self-similar variables

Assume that $\chi M < 8\pi$, consider a solution of (1) and define the rescaled functions u and v by:

$$n(x, t) = \frac{1}{R^2(t)} u\left(\frac{x}{R(t)}, \tau(t)\right) \quad \text{and} \quad c(x, t) = v\left(\frac{x}{R(t)}, \tau(t)\right) \quad (16)$$

with

$$R(t) = \sqrt{1 + 2t} \quad \text{and} \quad \tau(t) = \log R(t).$$

The rescaled system is

$$\begin{cases} \frac{\partial u}{\partial t} = \Delta u - \nabla \cdot (u(x + \chi \nabla v)) & x \in \mathbb{R}^2, t > 0, \\ v = -\frac{1}{2\pi} \log |\cdot| * u & x \in \mathbb{R}^2, t > 0, \\ u(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2. \end{cases} \quad (17)$$

The free energy now takes the form

$$F^R[u] := \int_{\mathbb{R}^2} u \log u \, dx - \frac{\chi}{2} \int_{\mathbb{R}^2} u v \, dx + \frac{1}{2} \int_{\mathbb{R}^2} |x|^2 u \, dx .$$

If (u, v) is a smooth solution of (17) which decays sufficiently at infinity, then

$$\frac{d}{dt} F^R[u(\cdot, t)] = - \int_{\mathbb{R}^2} u |\nabla \log u - \chi \nabla v + x|^2 \, dx .$$

Because of the hypercontractivity, the above inequality holds as an inequality for the solution of Theorem 1 after rescaling:

$$\frac{d}{dt} F^R[u(\cdot, t)] \leq - \int_{\mathbb{R}^2} u |\nabla \log u - \chi \nabla v + x|^2 \, dx .$$

For a rigorous proof, one has to redo the argument of Section 3.4. Since there is no additional difficulty this is left to the reader.

4.2 The self-similar solution

System (17) has the interesting property that for $\chi M < 8\pi$, it has a stationary solution which minimizes the free energy.

Lemma 21 *The functional F^R is bounded from below on the set*

$$\left\{ u \in L^1_+(\mathbb{R}^2) : |x|^2 u \in L^1(\mathbb{R}^2), \int_{\mathbb{R}^2} u \log u \, dx < \infty \right\}$$

if and only if $\chi \|u\|_{L^1(\mathbb{R}^2)} \leq 8\pi$.

Proof. If $\chi \|u\|_{L^1(\mathbb{R}^2)} \leq 8\pi$, the result is a straightforward consequence of Lemma 6. Notice that by Lemma 13, (iii), $u \log u$ is then bounded in $L^1(\mathbb{R}^2)$.

The functional $F^R[u]$ has an interesting scaling property. For a given u , let $u_\lambda(x) = \lambda^{-2} u(\lambda^{-1}x)$. It is straightforward to check that $\|u_\lambda\|_{L^1(\mathbb{R}^2)} =: M$ does not depend on $\lambda > 0$ and

$$F^R[u_\lambda] = F^R[u] - 2M \left(1 - \frac{\chi M}{8\pi} \right) \log \lambda + \frac{\lambda - 1}{2} \int_{\mathbb{R}^2} |x|^2 u \, dx .$$

As a function of λ , $F^R[u_\lambda]$ is clearly bounded from below if $\chi M < 8\pi$, and not bounded from below if $\chi M > 8\pi$, which completes the proof. \square

The free energy has a minimum which is a radial stationary solution of (17), see [12]. Such a solution is of course a natural candidate for the large time asymptotics of any solution of (17). In [12], there are also indications that (3) should have a unique solution for any given M , and there are strong numerical evidences supporting this fact. However, we are not able to discard the possibility that more than one solution to (3) exists for any given M .

Lemma 22 *Let $\chi M < 8\pi$. If u is a solution of (17), with initial data u_0 satisfying Assumptions (4), corresponding to a solution of (1) as given in Theorem 1, then as $t \rightarrow \infty$, $(s, x) \mapsto u(x, t+s)$ converges in $L^\infty(0, T; L^1(\mathbb{R}^2))$ for any positive T to a solution of (3) which is a stationary solution of (17) and moreover satisfies:*

$$\lim_{t \rightarrow \infty} \int_{\mathbb{R}^2} |x|^2 u(x, t) dx = \int_{\mathbb{R}^2} |x|^2 u_\infty dx = 2M \left(1 - \frac{\chi M}{8\pi}\right). \quad (18)$$

Proof. We use the free energy production term:

$$F^R[u_0] - \liminf_{t \rightarrow \infty} F^R[u(\cdot, t)] = \lim_{t \rightarrow \infty} \int_0^t \left(\int_{\mathbb{R}^2} u |\nabla \log u - \chi \nabla v + x|^2 dx \right) ds.$$

As a consequence,

$$\lim_{t \rightarrow \infty} \int_t^\infty \left(\int_{\mathbb{R}^2} u |\nabla \log u - \chi \nabla v + x|^2 dx \right) ds = 0, \quad (19)$$

which shows that, up to the extraction of subsequences, the limit u_∞ of $u(\cdot, t+\cdot)$, which exists for the same reasons as in the proof of Theorem 1, satisfies

$$\nabla \log u_\infty - \chi \nabla v_\infty + x = 0, \quad v_\infty = -\frac{1}{2\pi} \log |\cdot| * u_\infty,$$

where the first equation holds at least a.e. in the support of u_∞ . This is equivalent to write that (u_∞, v_∞) solves (3). Notice that the limit is unique because of (19) even if the uniqueness of the solutions of (3) is not established. Because of (19), we also know that u_∞ does not depend on the choice of the subsequence.

As in the proof of Lemma 3, consider a smooth function $\varphi_\varepsilon(|x|)$ with compact support that grows nicely to $|x|^2$ as $\varepsilon \rightarrow 0$. If (u, v) is a solution to (17), we compute

$$\begin{aligned} & \frac{d}{dt} \int_{\mathbb{R}^2} \varphi_\varepsilon u dx \\ &= \int_{\mathbb{R}^2} \Delta \varphi_\varepsilon u dx - \frac{\chi}{4\pi} \int_{\mathbb{R}^2} \frac{(\nabla \varphi_\varepsilon(x) - \nabla \varphi_\varepsilon(y)) \cdot (x - y)}{|x - y|^2} u(x, t) u(y, t) dx dy - 2 \int_{\mathbb{R}^2} |x|^2 u dx. \end{aligned}$$

As ε vanishes we may pass to the limit and obtain

$$\frac{d}{dt} \int_{\mathbb{R}^2} |x|^2 u dx = 4M \left(1 - \frac{\chi M}{8\pi}\right) - 2 \int_{\mathbb{R}^2} |x|^2 u dx.$$

This proves that for any $t > 0$,

$$\int_{\mathbb{R}^2} |x|^2 u(x, t) dx = \int_{\mathbb{R}^2} |x|^2 n_0 dx e^{-2t} + 2M \left(1 - \frac{\chi M}{8\pi}\right) (1 - e^{-2t}).$$

Passing to the limit $t \rightarrow \infty$, we get

$$\int_{\mathbb{R}^2} |x|^2 u_\infty dx \leq 2M \left(1 - \frac{\chi M}{8\pi}\right).$$

However, u_∞ is a solution of Equation (17), which satisfies the same assumptions as n_0 . Since it is a stationary solution with finite second moments, we have

$$\int_{\mathbb{R}^2} |x|^2 u_\infty dx = 2M \left(1 - \frac{\chi M}{8\pi}\right).$$

□

Notice that under the constraint $\|u_\infty\|_{L^1(\mathbb{R}^2)} = M$, u_∞ is a critical point of the free energy. If we knew that (3) has at most one solution for a given $M > 0$, u_∞ would automatically be the unique minimizer of the free energy. This result is not known although one can establish that u_∞ is radially symmetric. This is done using the two following results, Lemmata 23 and 24.

Lemma 23 *Let $u \in L^1_+(\mathbb{R}^2, (1 + |x|^2) dx)$ with $M := \int_{\mathbb{R}^2} u dx$, such that $\int_{\mathbb{R}^2} u \log u dx < \infty$, and define*

$$v(x) := -\frac{1}{2\pi} \int_{\mathbb{R}^2} \log |x - y| u(y) dy.$$

Then there exists a positive constant C such that, for any $x \in \mathbb{R}^2$ with $|x| > 1$,

$$\left| v(x) + \frac{M}{2\pi} \log |x| \right| \leq C.$$

Notice that as a straightforward consequence, v is non-positive outside of a ball.

Proof. We estimate

$$\left| v(x) + \frac{M}{2\pi} \log |x| \right| = \left| \frac{1}{2\pi} \int_{\mathbb{R}^2} \log \left(\frac{|x - y|}{|x|} \right) u(y) dy \right| \leq \text{(I)} + \text{(II)} + \text{(III)}$$

by

$$\begin{aligned} \text{(I)} &:= -\frac{1}{2\pi} \int_{\Omega_{\text{I}}} \log \left(\frac{|x - y|}{|x|} \right) u(y) dy \quad \text{with } \Omega_{\text{I}} := \left\{ (x, y) \in \mathbb{R}^2 : \frac{|x - y|}{|x|} \leq \frac{1}{2} \right\} \\ \text{(II)} &:= \left| \frac{1}{2\pi} \int_{\Omega_{\text{II}}} \log \left(\frac{|x - y|}{|x|} \right) u(y) dy \right| \quad \text{with } \Omega_{\text{II}} := \left\{ (x, y) \in \mathbb{R}^2 : \frac{1}{2} < \frac{|x - y|}{|x|} \leq 2 \right\} \\ \text{(III)} &:= \frac{1}{2\pi} \int_{\Omega_{\text{III}}} \log \left(\frac{|x - y|}{|x|} \right) u(y) dy \quad \text{with } \Omega_{\text{III}} := \left\{ (x, y) \in \mathbb{R}^2 : \frac{|x - y|}{|x|} > 2 \right\}. \end{aligned}$$

Using $|x - y|^2 \leq 2(|x|^2 + |y|^2)$ and $\log(1 + t) \leq t$, we get

$$4\pi \text{(III)} = \int_{\Omega_{\text{III}}} \log \left(\frac{|x - y|^2}{|x|^2} \right) u(y) dy \leq \int_{\Omega_{\text{III}}} \log \left(2 + 2 \frac{|y|^2}{|x|^2} \right) u(y) dy \leq M + \frac{2}{|x|^2} \int_{\Omega_{\text{III}}} |y|^2 u(y) dy.$$

On Ω_{II} , $|\log(|x - y|/|x|)|$ is bounded by $\log 2$: $\text{(II)} \leq M \frac{\log 2}{2\pi}$. For the last term, denote $z_x(y) = \frac{|x|}{|x - y|}$:

$$\text{(I)} = \frac{1}{2\pi} \int_{\Omega_{\text{I}}} \log(z_x(y)) u(y) dy.$$

By Jensen's inequality

$$\int_{\Omega_1} u(y) \log \left(\frac{u(y)}{z_x(y)} \right) dy \geq \int_{\Omega_1} u(y) \log \left(\frac{\int_{\Omega_1} u(y) dy}{\int_{\Omega_1} z_x(y) dy} \right) dy ,$$

we get

$$2\pi \text{ (I)} \leq \int_{\Omega_1} u(y) \log(u(y)) dy - \int_{\Omega_1} u(y) \log \left(\frac{\int_{\Omega_1} u(y) dy}{\int_{\Omega_1} z_x(y) dy} \right) dy .$$

The right hand side is bounded since $u \log u$ is bounded in $L^1(\mathbb{R}^2)$ by Lemma 8,

$$\int_{\Omega_1} z_x(y) dy = \int_{\Omega_1} \frac{|x|}{|x-y|} dy = \pi |x|^2 ,$$

and

$$\int_{\Omega_1} u(y) dy \leq \frac{4}{|x|^2} \int_{\Omega_1} |y|^2 u(y) dy .$$

Hence we can control (I) because $\int_{\Omega_1} u(y) dy \log \left(\int_{\Omega_1} z_x(y) dy \right) \leq \frac{4}{|x|^2} \log(\pi |x|^2)$. \square

This is enough to prove that the solution is radially symmetric, see [45].

Lemma 24 [45] *Assume that V is a non-negative non-trivial radial function on \mathbb{R}^2 such that $\lim_{|x| \rightarrow \infty} |x|^\alpha V(x) < \infty$ for some $\alpha \geq 0$. If u is a solution of*

$$\Delta u + V(x) e^u = 0 \quad x \in \mathbb{R}^2$$

such that $u_+ \in L^\infty(\mathbb{R}^2)$, then u is radially symmetric about the origin and $x \cdot \nabla u(x) < 0$ for any $x \in \mathbb{R}^2$.

Notice here that because of the asymptotic logarithmic behavior of v_∞ , the result of Gidas, Ni and Nirenberg, [23], does not directly apply. The boundedness from above is essential, otherwise non-radial solutions can be found, even with no singularity. Consider for instance the perturbation $\Theta(x) = \frac{1}{2} \theta (x_1^2 - x_2^2)$ for any $x = (x_1, x_2)$, for some fixed $\theta \in (0, 1)$, and define the potential $\phi(x) = \frac{1}{2} |x|^2 - \Theta(x)$. By a fixed-point method we can find a solution of

$$w(x) = -\frac{1}{2\pi} \log |\cdot| * M \frac{e^{\chi w - \phi(x)}}{\int_{\mathbb{R}^2} e^{\chi w(y) - \phi(y)} dy}$$

since, as $|x| \rightarrow \infty$, $\phi(x) \sim \frac{1}{2} [(1-\theta)x_1^2 + (1+\theta)x_2^2] \rightarrow +\infty$. This solution is such that $w(x) \sim -\frac{M}{2\pi} \log |x|$ for reasons similar to the ones of Lemma 23. Hence $v(x) := w(x) + \Theta(x)/\chi$ is a non-radial solution of the above equation with $\log V(x) = -\frac{1}{2} |x|^2$, which behaves like $\Theta(x)/\chi$ as $|x| \rightarrow \infty$ with $|x_1| \neq |x_2|$. This gives a non radial solution of Equation (3).

Lemma 25 *If $\chi M > 8\pi$, Equation (17) has no stationary solution (u_∞, v_∞) such that $\|u_\infty\|_{L^1(\mathbb{R}^2)} = M$ and $\int_{\mathbb{R}^2} |x|^2 u_\infty dx < \infty$. If $\chi M < 8\pi$, Equation (17) has at least one radial stationary solution given by (3). This solution is C^∞ and u_∞ is dominated as $|x| \rightarrow \infty$ by $e^{-(1-\varepsilon)|x|^2/2}$ for any $\varepsilon \in (0, 1)$.*

Proof. The existence of a stationary solution if $\chi M < 8\pi$ is easy. It follows from Lemma 22 but can also be achieved by minimizing the free energy, see [12]. If the initial condition is radial or if the minimization is done among radial solutions, then the stationary solution is also radial. Direct approaches (fixed-point methods, ODE shooting methods) can also be used.

If $\chi M > 8\pi$ and if there was a stationary solution with finite second moment, we could write

$$0 = \frac{d}{dt} \int_{\mathbb{R}^2} |x|^2 u_\infty dx = 4M \left(1 - \frac{\chi M}{8\pi} \right) - 2 \int_{\mathbb{R}^2} |x|^2 u_\infty dx .$$

Since the right hand side is negative, this is simply impossible. \square

In the rescaled variables, the solution of (17) converges to a radial stationary solution u_∞ of (3). It is not difficult to check that $\bar{n}(x, t) := \frac{1}{2t} u_\infty \left(\frac{1}{2} \log(2t), x/\sqrt{2t} \right)$ and $\bar{c}(x, t) := v_\infty \left(\frac{1}{2} \log(2t), x/\sqrt{2t} \right)$ gives a self-similar solution of (1), which is supposed to describe the large time asymptotics of (1), and this is what we are going to clarify in the last section.

4.3 Intermediate asymptotics

Lemma 26 *Under the assumptions of Lemma 22,*

$$\lim_{t \rightarrow \infty} F^R[u(\cdot, \cdot + t)] = F^R[u_\infty] .$$

Proof. By (18), we already know that $\lim_{t \rightarrow \infty} \int_{\mathbb{R}^2} |x|^2 u(x, t) dx = \int_{\mathbb{R}^2} |x|^2 u_\infty dx$. Using the estimates of Sections 2.5–2.7 and Lemma 11, we know that $u(\cdot, \cdot + t)$ converges to u_∞ in $L^2((0, 1) \times \mathbb{R}^2)$ and that $\int_{\mathbb{R}^2} u(\cdot, \cdot + t) v(\cdot, \cdot + t) dx$ converges to $\int_{\mathbb{R}^2} u_\infty v_\infty dx$. Concerning the entropy, it is sufficient to prove that $u(\cdot, \cdot + t) \log u(\cdot, \cdot + t)$ weakly converges in $L^1((0, 1) \times \mathbb{R}^2)$ to $u_\infty \log u_\infty$. By Lemma 8, there is a uniform L^1 bound. Concentration is prohibited by the convergence in $L^2((0, 1) \times \mathbb{R}^2)$. Vanishing or dichotomy cannot occur either: Take indeed $R > 0$, large, and compute $\int_{|x|>R} u |\log u| = \text{(I)} + \text{(II)}$, with

$$\begin{aligned} \text{(I)} &= \int_{|x|>R, u \geq 1} u \log u dx \leq \frac{1}{2} \int_{|x|>R, u \geq 1} |u|^2 dx , \\ \text{(II)} &= - \int_{|x|>R, u < 1} u \log u dx \leq \frac{1}{2} \int_{|x|>R, u < 1} |x|^2 u dx - m \log \left(\frac{m}{2\pi} \right) . \end{aligned}$$

In the first case, we have used the inequality $u \log u \leq u^2/2$ for any $u \geq 1$, while the second estimate is based on Jensen's inequality in the spirit of the proof of Lemma 8:

$$m := \int_{|x|>R, u < 1} u dx \leq \frac{1}{R^2} \int_{|x|>R, u < 1} |x|^2 u dx .$$

Because of the convergence of the two quantities $\int_{|x|>R, u < 1} |u|^2 dx$ and $\int_{|x|>R, u < 1} |x|^2 u dx$ to 0 as $R \rightarrow \infty$, we have the uniform estimate

$$\lim_{R \rightarrow \infty} \int_{|x|>R} u |\log u| = 0 ,$$

which completes the proof. \square

The result we have shown above is actually slightly better, since it proves that all terms in the free energy, namely the entropy, the energy corresponding to the potential $\frac{1}{2}|x|^2$ and the self-consistent potential energy, converge to the corresponding values for the limiting stationary solution.

As noted above, u_∞ is a critical point of F^R under the constraint $\|u\|_{L^1(\mathbb{R}^2)} = M$. We can therefore rewrite $F^R[u] - F^R[u_\infty]$ as

$$F^R[u] - F^R[u_\infty] = \int_{\mathbb{R}^2} u \log\left(\frac{u}{u_\infty}\right) dx - \frac{\chi}{2} \int_{\mathbb{R}^2} |\nabla v - \nabla v_\infty|^2 dx,$$

and both terms in the above expression converge to 0 as $t \rightarrow \infty$, if u is a solution of (1). Since for any nonnegative functions $f, g \in L^1(\mathbb{R}^2)$ such that $\int_{\mathbb{R}^2} f dx = \int_{\mathbb{R}^2} g dx = M$,

$$\|f - g\|_{L^1(\mathbb{R}^2)}^2 \leq \frac{1}{4M} \int_{\mathbb{R}^2} f \log\left(\frac{f}{g}\right) dx$$

by the Csiszár-Kullback inequality, [20, 35], this proves the

Corollary 27 *Under the assumptions of Lemma 22,*

$$\lim_{t \rightarrow \infty} \|u(\cdot, \cdot + t) - u_\infty\|_{L^1(\mathbb{R}^2)} = 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} \|\nabla v(\cdot, \cdot + t) - \nabla v_\infty\|_{L^2(\mathbb{R}^2)} = 0.$$

Undoing the change of variables (16), this proves Theorem 2.

Acknowledgments. All authors are partially supported by the EU network HPRN-CT-2002-00282. © 2006 by the authors. This paper may be reproduced, in its entirety, for non-commercial purposes.

References

- [1] A. ARNOLD, J. A. CARRILLO, L. DESVILLETES, J. DOLBEAULT, A. JÜNGEL, C. LEDERMAN, P. A. MARKOWICH, G. TOSCANI, AND C. VILLANI, *Entropies and equilibria of many-particle systems: an essay on recent research*, Monatsh. Math., 142 (2004), pp. 35–43.
- [2] J.-P. AUBIN, *Un théorème de compacité*, C. R. Acad. Sci. Paris, 256 (1963), pp. 5042–5044.
- [3] F. BAVAUD, *Equilibrium properties of the Vlasov functional: the generalized Poisson-Boltzmann-Emden equation*, Rev. Modern Phys., 63 (1991), pp. 129–148.
- [4] W. BECKNER, *Sharp Sobolev inequalities on the sphere and the Moser-Trudinger inequality*, Ann. of Math. (2), 138 (1993), pp. 213–242.
- [5] R. BENGURIA, *PhD Thesis*, PhD thesis, Princeton University, 1979.
- [6] R. BENGURIA, H. BRÉZIS, AND E. H. LIEB, *The Thomas-Fermi-von Weizsäcker theory of atoms and molecules*, Comm. Math. Phys., 79 (1981), pp. 167–180.
- [7] P. BILER, *Local and global solvability of some parabolic systems modelling chemotaxis*, Adv. Math. Sci. Appl., 8 (1998), pp. 715–743.

- [8] P. BILER, G. KARCH, AND P. LAURENÇOT, *The 8π -problem for radially symmetric solutions of a chemotaxis model in a disc*, Top. Meth. Nonlin. Anal. to appear, (2005).
- [9] ———, *The 8π -problem for radially symmetric solutions of a chemotaxis model in the plane*. Preprint, 2006.
- [10] P. BILER AND T. NADZIEJA, *A class of nonlocal parabolic problems occurring in statistical mechanics*, Colloq. Math., 66 (1993), pp. 131–145.
- [11] ———, *Global and exploding solutions in a model of self-gravitating systems*, Rep. Math. Phys., 52 (2003), pp. 205–225.
- [12] E. CAGLIOTI, P.-L. LIONS, C. MARCHIORO, AND M. PULVIRENTI, *A special class of stationary flows for two-dimensional Euler equations: a statistical mechanics description*, Comm. Math. Phys., 143 (1992), pp. 501–525.
- [13] V. CALVEZ AND J. A. CARRILLO, *Volume effects in the Keller-Segel model: energy estimates preventing blow-up*, tech. rep., Preprint ICREA, 2005.
- [14] V. CALVEZ, B. PERTHAME, AND M. SHARIFI TABAR, *Modified Keller-Segel system and critical mass for the log interaction kernel*. In Preparation.
- [15] E. CARLEN AND M. LOSS, *Competing symmetries, the logarithmic HLS inequality and Onofri's inequality on S^n* , Geom. Funct. Anal., 2 (1992), pp. 90–104.
- [16] F. A. C. C. CHALUB, P. A. MARKOWICH, B. PERTHAME, AND C. SCHMEISER, *Kinetic models for chemotaxis and their drift-diffusion limits*, Monatsh. Math., 142 (2004), pp. 123–141.
- [17] S. CHILDRESS AND J. K. PERCUS, *Nonlinear aspects of chemotaxis*, Math. Biosci., 56 (1981), pp. 217–237.
- [18] L. CORRIAS, B. PERTHAME, AND H. ZAAG, *A chemotaxis model motivated by angiogenesis*, C. R. Math. Acad. Sci. Paris, 336 (2003), pp. 141–146.
- [19] ———, *Global solutions of some chemotaxis and angiogenesis systems in high space dimensions*, Milan J. Math., 72 (2004), pp. 1–29.
- [20] I. CSISZÁR, *Information-type measures of difference of probability distributions and indirect observations*, Studia Sci. Math. Hungar., 2 (1967), pp. 299–318.
- [21] J. DOLBEAULT AND B. PERTHAME, *Optimal critical mass in the two-dimensional Keller-Segel model in \mathbb{R}^2* , C. R. Math. Acad. Sci. Paris, 339 (2004), pp. 611–616.
- [22] H. GAJEWSKI AND K. ZACHARIAS, *Global behaviour of a reaction-diffusion system modelling chemotaxis*, Math. Nachr., 195 (1998), pp. 77–114.
- [23] B. GIDAS, W. M. NI, AND L. NIRENBERG, *Symmetry and related properties via the maximum principle*, Comm. Math. Phys., 68 (1979), pp. 209–243.
- [24] T. GOUDON, *Hydrodynamic limit for the Vlasov-Poisson-Fokker-Planck system: analysis of the two-dimensional case*, Math. Models Methods Appl. Sci., 15 (2005), pp. 737–752.

- [25] L. GROSS, *Logarithmic Sobolev inequalities*, Amer. J. Math., 97 (1975), pp. 1061–1083.
- [26] M. A. HERRERO AND J. J. L. VELÁZQUEZ, *Singularity patterns in a chemotaxis model*, Math. Ann., 306 (1996), pp. 583–623.
- [27] T. HILLEN AND K. PAINTER, *Global existence for a parabolic chemotaxis model with prevention of overcrowding*, Adv. in Appl. Math., 26 (2001), pp. 280–301.
- [28] D. HORSTMANN, *The nonsymmetric case of the Keller-Segel model in chemotaxis: some recent results*, NoDEA Nonlinear Differential Equations Appl., 8 (2001), pp. 399–423.
- [29] ———, *From 1970 until present: the Keller-Segel model in chemotaxis and its consequences. I*, Jahresber. Deutsch. Math.-Verein., 105 (2003), pp. 103–165.
- [30] ———, *From 1970 until present: the Keller-Segel model in chemotaxis and its consequences. II*, Jahresber. Deutsch. Math.-Verein., 106 (2004), pp. 51–69.
- [31] D. HORSTMANN AND M. WINKLER, *Boundedness vs. blow-up in a chemotaxis system*, J. Differential Equations, 215 (2005), pp. 52–107.
- [32] W. JÄGER AND S. LUCKHAUS, *On explosions of solutions to a system of partial differential equations modelling chemotaxis*, Trans. Amer. Math. Soc., 329 (1992), pp. 819–824.
- [33] E. F. KELLER AND L. A. SEGEL, *Initiation of slide mold aggregation viewed as an instability*, J. Theor. Biol., 26 (1970).
- [34] R. KOWALCZYK, *Preventing blow-up in a chemotaxis model*, J. Math. Anal. Appl., 305 (2005), pp. 566–588.
- [35] S. KULLBACK, *On the convergence of discrimination information*, IEEE Trans. Information Theory, IT-14 (1968), pp. 765–766.
- [36] P. LAURENÇOT AND D. WRZOSEK, *From the nonlocal to the local discrete diffusive coagulation equations*, Math. Models Methods Appl. Sci., 12 (2002), pp. 1035–1048. Special issue on kinetic theory.
- [37] E. H. LIEB, *Sharp constants in the Hardy-Littlewood-Sobolev and related inequalities*, Ann. of Math. (2), 118 (1983), pp. 349–374.
- [38] G. M. LIEBERMAN, *Second order parabolic differential equations*, World Scientific Publishing Co. Inc., River Edge, NJ, 1996.
- [39] J.-L. LIONS, *Équations différentielles opérationnelles et problèmes aux limites*, Die Grundlehren der mathematischen Wissenschaften, Bd. 111, Springer-Verlag, Berlin, 1961.
- [40] P. K. MAINI, *Applications of mathematical modelling to biological pattern formation*, in Coherent structures in complex systems (Sitges, 2000), vol. 567 of Lecture Notes in Phys., Springer, Berlin, 2001, pp. 205–217.
- [41] A. MARROCCO, *Numerical simulation of chemotactic bacteria aggregation via mixed finite elements*, M2AN Math. Model. Numer. Anal., 37 (2003), pp. 617–630.

- [42] J. D. MURRAY, *Mathematical biology. II*, vol. 18 of Interdisciplinary Applied Mathematics, Springer-Verlag, New York, third ed., 2003. Spatial models and biomedical applications.
- [43] T. NAGAI AND T. SENBA, *Global existence and blow-up of radial solutions to a parabolic-elliptic system of chemotaxis*, Adv. Math. Sci. Appl., 8 (1998), pp. 145–156.
- [44] T. NAGAI, T. SENBA, AND K. YOSHIDA, *Application of the Trudinger-Moser inequality to a parabolic system of chemotaxis*, Funkcial. Ekvac., 40 (1997), pp. 411–433.
- [45] Y. NAITO, *Symmetry results for semilinear elliptic equations in \mathbf{R}^2* , in Proceedings of the Third World Congress of Nonlinear Analysts, Part 6 (Catania, 2000), vol. 47, 2001, pp. 3661–3670.
- [46] V. NANJUNDIAH, *Chemotaxis, signal relaying and aggregation morphology*, Journal of Theoretical Biology, 42 (1973), pp. 63–105.
- [47] H. G. OTHMER AND A. STEVENS, *Aggregation, blowup, and collapse: the ABCs of taxis in reinforced random walks*, SIAM J. Appl. Math., 57 (1997), pp. 1044–1081.
- [48] T. PADMANABHAN, *Statistical mechanics of gravitating systems*, Phys. Rep., 188 (1990), pp. 285–362.
- [49] C. S. PATLAK, *Random walk with persistence and external bias*, Bull. Math. Biophys., 15 (1953), pp. 311–338.
- [50] B. PERTHAME, *Nonexistence of global solutions to Euler-Poisson equations for repulsive forces*, Japan J. Appl. Math., 7 (1990), pp. 363–367.
- [51] ———, *PDE models for chemotactic movements: parabolic, hyperbolic and kinetic*, Appl. Math., 49 (2004), pp. 539–564.
- [52] C. ROSIER, *Problème de Cauchy pour une équation parabolique modélisant la relaxation des systèmes stellaires auto-gravitants*, C. R. Acad. Sci. Paris Sér. I Math., 332 (2001), pp. 903–908.
- [53] T. SENBA AND T. SUZUKI, *Weak solutions to a parabolic-elliptic system of chemotaxis*, J. Funct. Anal., 191 (2002), pp. 17–51.
- [54] ———, *Blowup behavior of solutions to the rescaled Jäger-Luckhaus system*, Adv. Differential Equations, 8 (2003), pp. 787–820.
- [55] ———, *Applied analysis*, Imperial College Press, London, 2004. Mathematical methods in natural science.
- [56] J. SIMON, *Compact sets in the space $L^p(0, T; B)$* , Ann. Mat. Pura Appl. (4), 146 (1987), pp. 65–96.
- [57] A. STEVENS, *The derivation of chemotaxis equations as limit dynamics of moderately interacting stochastic many-particle systems*, SIAM J. Appl. Math., 61 (2000), pp. 183–212 (electronic).

- [58] T. SUZUKI, *Free energy and self-interacting particles*, Progress in Nonlinear Differential Equations and their Applications, 62, Birkhäuser Boston Inc., Boston, MA, 2005.
- [59] J. J. L. VELÁZQUEZ, *Stability of some mechanisms of chemotactic aggregation*, SIAM J. Appl. Math., 62 (2002), pp. 1581–1633 (electronic).
- [60] ———, *Point dynamics in a singular limit of the Keller-Segel model. I. Motion of the concentration regions*, SIAM J. Appl. Math., 64 (2004), pp. 1198–1223 (electronic).
- [61] ———, *Point dynamics in a singular limit of the Keller-Segel model. II. Formation of the concentration regions*, SIAM J. Appl. Math., 64 (2004), pp. 1224–1248 (electronic).
- [62] M. I. WEINSTEIN, *Nonlinear Schrödinger equations and sharp interpolation estimates*, Comm. Math. Phys., 87 (1982/83), pp. 567–576.
- [63] G. WOLANSKY, *Comparison between two models of self-gravitating clusters: conditions for gravitational collapse*, Nonlinear Anal., 24 (1995), pp. 1119–1129.
- [64] D. WRZOSEK, *Long time behaviour of solutions to a chemotaxis model with volume filling effect*, Hyke preprint server, 166 (2004).