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Static and Dynamic Magnetic Properties of [Fe_{13}]

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Introduction

In certain cases the lattice geometry gives rise to the interesting phenomenon of spin frustration, of which one famous example is the Kagome lattice.[1] Central to such lattices is the presence of triangular exchange interaction geometries. The studied systems are generally extended two-dimensional lattices. If the lattice were curved in both dimensions, as it were projected on a sphere, a highly symmetric zero-dimensional object would result, which can be expected to show interesting properties due to that high symmetry and also due to the finite size of the lattice. An example of a highly symmetric, frustrated molecular systems is the polyoxomolybdate cluster $[Fe_{30}Mo_{72}]$ which contains 30 iron(III) ions (s = 5/2) in an icosadodecahedral arrangement.[2] This system suffers from the drawback that the effective isotropic superexchange interaction is rather weak, of the order of J = 1.5 cm^{-1}, due to the long and inefficient superexchange pathways. Recently a new iron(III) system was published,[3] which can be expected to show stronger exchange interaction due to the presence of monoatomic bridges between the paramagnetic ions which can efficiently transmit superexchange interactions, and in addition to show a high degree of spin frustration. This tridecanuclear Fe^{III} compound, $(C_5H_6N^+)_5[Fe_{13}F_{24}(OCH_3)_{12}O_4].CH_3OH.4H_2O$ (≡ $[Fe_{13}]$) crystallizes in the cubic space group F–43m. Due to its high symmetry there are only two independent Fe^{III} ions in the molecule, *viz*. one in the center and one other, which is reproduced twelve times due to the symmetry elements of the space group. Although only four pyridinium cations were found in the x-ray diffractions studies, the presence of a fifth one was assumed, since an overall charge of –4 of the cluster anion would suggest an unrealistic +4 oxidation state of the central iron ion: the presence of five pyridinium cations was as also corroborated by elemental analysis and volumetric titration excluded the presence of Fe(IV) ions. Preliminary magnetic measurements indicated a complicated dependence of the magnetic susceptibility on temperature and suggested a high-spin ground state. In iron clusters, Mössbauer spectroscopy is an excellent local probe for the study of its magnetic properties. Preliminary data supported the +3 oxidation state of all iron ions, but no magnetic structure was observed.[3]

The other point of interest in such exchange-coupled clusters is the fact that many of them show the phenomenon of slow relaxation of the magnetization due to the zero-field splitting of the spin ground state. This zero field splitting lifts the degeneracy of the spin ground state, effectively leading to an energy barrier towards the relaxation of the magnetization. In addition, transverse terms in the spin Hamiltonian can lead to tunnelling of the magnetic moment under the barrier. Such terms are not allowed by symmetry in second order if the symmetry is tetragonal or higher. In cubic symmetry, in fact, no ZFS is expected at all below fourth order. Recently, it has been shown that local lowering of the symmetry due to disorder in the system can lead to the presence of lower order transverse terms. It is clear that disorder can play a major role in highly symmetric systems.

Here we report the study of the static and dynamic magnetic properties of the mentioned cubic symmetry [Fe_{13}] cluster, using a number of experimental techniques, which shows that especially the spin dynamics is mainly governed by disorder.

Experimental Section

$(C_5H_6N^+)_5[Fe_{13}F_{24}(OCH_3)_{12}O_4]\cdot CH_3OH\cdot 4\ H_2O$ ([Fe_{13}]) was synthesized as published recently. [3]. The deuterated derivative for INS experiments was prepared and stocked under inert atmosphere by using fully deuterated methanol.

dc and *ac* magnetic susceptibilities and the magnetization were measured on powder samples, using a Cryogenics S600 SQUID magnetometer and an Oxford VSM with a 12T magnet. The data were corrected for the diamagnetic contribution using Pascal's constants. Low temperature *ac* magnetic susceptibility was measured in a custom-designed susceptometer based on an Oxford ^{3}He Heliox cryostat, a SR830 Lock-In serving also as *ac* current source, an Oxford ITC 503 temperature controller, and a home-made probe built by Prof. M. A. Novak, Universidade Federal do Rio de Janeiro, Brazil. Temperature was measured with a calibrated RuO_2 probe below 4K, magnetic calibration was made with a spherical ferrimagnetic sample previously calibrated with SQUID *dc* measurements.

Inelastic Neutron Scattering (INS) measurements have been performed at the Institute Laue-Langevin (ILL) in Grenoble (France), with the direct-geometry time-of-flight spectrometer IN5. About 1.7 grams of deuterated polycrystalline sample were placed in a hollow aluminum cylinder (0.7 mm wall thickness, 20 mm inner diameter), and inserted into a standard ILL ^{4}He cryostat allowing a base temperature of 2 K. The spectrometer was operated in high and medium energy resolution configurations, with incident wavelengths λ = 11 Å and λ = 6 Å, corresponding to energy resolution at the elastic peak of 12 μeV and 60 μeV, respectively. Solid angle and detector efficiency calibrations have been performed using the spectrum of a vanadium standard.

The low energy transfer region has been explored with the IN16 spectrometer at the ILL. In this instrument high energy resolution is achieved by deflecting the neutrons at the monochromator and at the analyzer with a Bragg angle of 90 degrees. The energy of the incident beam is varied by Doppler motion of the monochromator. By using a spherically shaped Si(111) monochromator and deformed Si(111) single crystals as analyzers, an energy resolution of 0.9 μeV is obtained over an energy transfer range from -15 to 15 μeV. For this experiment, a cylindrical aluminum holder containing 1.2 grams of material was used inside a dilution refrigerator allowing the sample to be cooled at a minimum temperature of 50 mK.

HF-EPR spectra of [Fe_{13}] pellets were recorded on a home-made spectrometer,[4, 5] operating in the single-pass transmission mode. The main magnetic field is supplied by a superconducting

magnet (Cryogenics Consultants) with a maximum field of 12 T. Gunn-diodes and their multipliers (Radiometer Physics) were used as frequency sources, with a basic frequency of 95 GHz or 115 GHz .

For the Mössbauer experiments powdered [Fe_{13}] was mixed with eicosane (Aldrich) and compressed inside a disk-shaped holder. The effective thickness of the iron was 4.1 mg cm^{-2}. The spectra were determined by using a conventional sinusoidal acceleration spectrometer that operated in a multichannel scaling mode. The γ-ray source, which consisted of 25 mCi of ^{57}Co in a rhodium metal matrix, was maintained at room temperature. A natural abundance iron foil was used to calibrate the spectrometer. The measurements were performed in a standard exchange helium cryostat (Oxford, Inc.), using a control program under LabView.

Results

X-Ray diffraction. The structure of [Fe_{13}] as determined previously, is shown in Fig. 1. The cluster itself is highly symmetric, with the central iron ion lying on a T_d symmetry site, and only one independent further iron ion. Closer inspection of the crystal structure revealed that the cluster is surrounded by 12 pyridinium cations, bound by hydrogen bonds between the pyridinium ions and the cluster (Fig. 1). The aromatic hydrogen atoms are relatively weak donors, while the N^+–H groups are expected to be much stronger, of the order of 4 kcal mol^{-1}. [6] Both terminal Fe–F and bridging Fe–F–Fe fluorine atoms can act as acceptors. Since the hydrogen bonds formed with N^+–H donors have been found to be much more directional with angles close to linear, [7] it can be assumed that the N^+–H groups of the pyridinium ions will preferentially bind to the bridging Fe–F–Fe groups. Two factors are expected to lead to local distortions of the lattice, leading to local symmetry lowering of the cluster molecules. First of all, of the five pyridinium cations that are required for charge neutrality only four could be found in the single crystal X-Ray analysis, the fifth being disordered to such an extent that is does not show up in the electron density difference maps. That means that the tetrahedral lattice formed by the hydrogen bonded network of pyridinium ions and $[Fe_{13}]^{5-}$ clusters contains inclusions of the fifth pyridinium ion at non-specific sites, leading to local distortion of the lattice. The same argument holds for the co-crystallized water and methanol molecules which are both good hydrogen donors. Secondly, the strong hydrogen bond formed between the N^+–H moiety and the Fe–F–Fe bridge can induce a local distortion of that bridge, leading to a local lowering of the cluster symmetry, because each N^+–H group can bind to only one cluster

dc magnetic susceptibility. Preliminary susceptibility measurements suggested the presence of strong antiferromagnetic intracluster exchange interactions. In fact, the observed room temperature χT product of 46.5 cm^3 K mol^{-1} (Fig. 2) is much lower than that expected for 13 noninteracting

iron(III) ions ($\chi T = N \times 0.125 \times g^2 \times S(S+1) = 56.9$ cm^3 K mol^{-1}). The rise of χT on lowering the temperature to a maximum value of 54.5 cm^3 K mol^{-1} at 109.5 K suggests ferrimagnetic interactions, which is supported by the observation that the susceptibility does not tend towards zero at low temperatures. A Curie-Weiss plot of the low temperature part (T < 6 K) yields a Curie constant of C = 39.8±0.3 cm^3 K mol^{-1}, close to that expected for S = 17/2 and g = 2.0, assuming only the ground spin multiplet is populated. The Weiss temperature is θ = –0.41±0.04 K, which could indicate small intermolecular antiferromagnetic interactions. The value of the magnetic moment of μ = 20 μ_B at H = 12 T (not shown) also indicates a high spin ground state, while the fact that the magnetization does not saturate even at these fields suggests the presence of low-lying excited spin multiplets.

In spite of the high symmetry, the spin topology is quite complicated and quite frustrated, due to the presence of eight triangular interactions, as shown schematically in Figure 3a. In oxo-bridged iron(III) dimers it has been found that the exchange interaction decreases with increasing Fe–O distance, while the Fe–O–Fe angle dependence is less clear.[8] In polynuclear complexes a definite angle dependence was observed with the exchange interaction becoming more antiferromagnetic with increasing Fe–O–Fe angle.[9] The central iron ion (Fe1) interacts with all other iron ions (Fe2) through μ^4-oxo bridges (O1). These bridges are asymmetric since the Fe1–O1 distance is short (1.888(6) Å), while Fe2–O1 is 2.093(3) Å. The short Fe1–O1 distance, combined with ∠Fe1–O1–Fe2 = 120.77(13)° suggests that there should be a relatively strong antiferromagnetic interaction between Fe1 and Fe2. The outer iron ions interact in two different ways. The first is through a combination of said μ^4-oxo bridge, with ∠Fe2–O1–Fe2 = 96.16(18)°, combined with a methoxide bridge (O2), with ∠Fe2–O2–Fe2 = 104.60(18)°, and a Fe2–O2 distance of 1.968 Å. The second is through a fluoride bridge, with ∠Fe2–F–Fe2 = 146.97(16)°. The exchange path through O1 is expected to give rise to a weakly antiferromagnetic or even ferromagnetic exchange interaction in view of the small ∠Fe2–O1–Fe2 of 96.16(18)°.[9] Both other exchange paths between the different Fe2 atoms are expected to give rise to similar exchange interaction strengths, since the Fe2–O2–Fe2 angle of 104.60(18)° and the Fe2–O2 distance of 1.968 Å favor an antiferromagnetic exchange interaction between the Fe2 atoms, as do the ∠Fe2–F–Fe2 angle of 146.97(16)° and the Fe2–F distance of 1.979(1) Å. For example, for the Fe–F–Fe interaction in the rhombohedral phase of FeFe3, with an Fe–F distance of 1.92 Å and ∠Fe–F–Fe = 152°, a superexchange interaction of J = 17 – 20 cm^{-1} was estimated.[10]

It is immediately clear that a stronger Fe1…Fe2 exchange interaction will lead to a high spin ground state (up to S = 55/2), while a stronger Fe2…Fe2 exchange interaction will give rise to a low spin system. The susceptibility data point to the former case. However, despite the high

symmetry of the system, its size precludes using numeric diagonalization methods to fit the susceptibility curve to an isotropic exchange interaction Hamiltonian. Even when making full use of the symmetry of the system using irreducible tensor operator methods, the dimension of the largest matrix is still some 80 million. The only possibility to obtain insight into the strengths of the exchange interactions is by neglecting one of the two Fe2…Fe2 exchange interactions. While such a procedure is not expected to lead to an accurate description of the spin state energies, it will give information on the relative strengths of the Fe1…Fe2 and Fe2…Fe2 exchange interactions as well as the approximate value of the spin ground state.

After neglecting one of the Fe2…Fe2 interactions the spin topology simplifies enormously (Fig. 3b). In fact, since the system can be divided into two subsystems (one consisting of four triangles and one of the central spin) where the interactions between all the members of one subsystem and all the members of the other subsystem are equal, the problem has an analytical solution. [11] This approach, called the Kambe approach after its inventor,[12] has been used in several instances for clusters of up to 9 spins.[13, 14] The spin Hamiltonian is:

$$\hat{H} = J_2(S_1S_2 + S_2S_3 + S_3S_1 + S_4S_5 + S_5S_6 + S_6S_4 + S_7S_8 + S_8S_9 + S_9S_7 + S_{10}S_{11} + S_{11}S_{12} + S_{12}S_{10}) + J_1(S_{13}\sum_{i=1}^{12} S_i) \quad (1)$$

where S_1 through S_{12} correspond to the symmetry generated crystallographic Fe2 type ions, and S_{13} corresponds to the Fe1 ion. J_1 and J_2 describe the Fe1...Fe2 exchange interaction and the intratriangle Fe2...Fe2 one, respectively. Following the Kambe approach, the system under study is divided into two subsystems, one being the central iron atom (S_{13}) and the other being the outer iron atoms ($S_{1–12}$). All spins in the first subsystem are connected to those in the second subsystem by equivalent coupling constants J_1. The $S_{1–12}$ system is then divided into the systems $S_{1–6}$ and $S_{7–12}$, which are again connected by equivalent (in this case zero) interactions. Finally the two subsystems $S_{1–6}$ and $S_{7–12}$ are divided into the triangles $S_{1–3}$, $S_{4–6}$, $S_{7–9}$ and $S_{10–12}$. For these latter triangular systems the analytical solution of the SH is known. Eventually the exact solution of the SH of the system under study is obtained:

$$E = J_1/2\,[S_{1–13}(S_{1–13} + 1) – S_{1–12}(S_{1–12} + 1)] + J_2/2\,[S_{1–3}(S_{1–3} + 1) + S_{4–6}(S_{4–6} + 1) + S_{7–9}(S_{7–9} + 1) + S_{10–12}(S_{10–12} + 1)] \quad (2)$$

The total spin function is: $|S_1\ S_2\ S_{1–2}\ S_3\ S_{1–3}\ S_4\ S_5\ S_{4–5}\ S_6\ S_{4–6}\ S_7\ S_8\ S_{7–8}\ S_9\ S_{7–9}\ S_{10}\ S_{11}\ S_{10–11}\ S_{12}\ S_{10–12}\ S_{1–6}\ S_{7–12}\ S_{1–12}\ S_{13}\ S_T\rangle$. From this formula, it is clear that the energy of a spin state only depends on the spin states of the four triangles, the way they are coupled and the way they are coupled to the central Fe atom. Figure 4 shows the energies of the ~250 lowest spin states as a function of the ratio between the intratriangle exchange interaction and the central-peripheral exchange interaction, with the consecutive spin ground states as bold lines. In this model the degeneracies of the low lying

excited states is very large, which would partially be removed by the inclusion of the second type of Fe2…Fe2 interaction. The possible spin states of one triangle are S = 15/2, 13/2, 11/2, ..., 1/2. In the spin ground states as a function of J_2/J_1 all for triangles have the same spin, and the resulting intermediate spin S_{1-12} is coupled antiferromagnetically to the central ion spin. Hence, the consecutive spin ground states of the cluster are expected to be S_T = 55/2, 47/2, 43/2,...,1/2.

The magnetic susceptibility can be calculated from the summation over all spin states i of total spin S_T^i and energy E_i:

$$\chi T = \frac{Ng^2\mu_B^2}{3k} \frac{\sum_i S_T^i(S_T^i+1)(2S_T^i+1)\exp(-E_i/kT)}{\sum_i (2S_T^i+1)\exp(-E_i/kT)} \qquad (3)$$

Using a home-written program, χT was fitted to the experimental data (Figure 2). The resulting J_1 and J_2 values are about 43 cm^{-1} and 30 cm^{-1}, respectively, although the J_2/J_1 ratio (0.70) is probably more reliable than the actual values. This J_2/J_1 ratio would then suggest within our model that the spin ground state of the Fe_{13} cluster is either S_T = 23/2 or S_T = 15/2. Those values are an indication that the Fe2···Fe2 and Fe1···Fe2 interactions are of comparable magnitude, which leads to a large spin ground state. Exactly at J_2/J_1 = 0.70, the spin ground state is S_T = 23/2, while the excited states have progressively lower spin values of S_T = 21/2, 19/2, 17/2, and 15/2 with the separation between the states being $\Delta E = 0.05J_1$, corresponding to ΔE = 2.15 cm^{-1} for J_1 = 43 cm^{-1} (inset Fig. 4)

Inelastic neutron scattering (INS). INS is the zero-field technique of choice to obtain information on the total spin energy levels in exchange coupled clusters because of the ΔS=0, ±1 selection rule for INS transitions.[15] In this study we focus especially on the low energy region to verify the presence of low-lying spin excited states, in a series of different measurements. First of all, a very-high resolution experiment on the IN16 backscattering spectrometer at temperatures of T = 55 mK and T = 500 mK showed no inelastic transitions in the accessible energy transfer range of up to 15 μeV. Secondly, the magnetic response up to 1.5 meV has been explored with lower resolution using the IN5 chopper spectrometer. Fig. 5 shows the INS intensity distribution recorded with the sample kept at T = 2 K, and an incident wavelength of λ= 11 Å. Counts in individual detectors at different scattering angles were summed before the time-of-flight to energy conversion. Background, detector efficiency and absorption corrections have been applied following standard procedures. Two peaks are clearly resolved both on the energy-gain and the energy-loss side of the spectrum. They can be fitted to Gaussian lineshapes centered at ± 50(5) μeV and ± 99(6) μeV (0.40 cm^{-1}, and 0.79 cm^{-1}, respectively) and having a Full Width at Half Maximum (FWHM) about three and six times the instrumental resolution (12 μeV), respectively. These wide peaks can arise from several types of transitions. Firstly, they can be due to excitations between ground spin manifolds (formed

by different equivalent topological configurations) and close-lying spin multiplets described by a density of states peaked at 50 and 99 μeV. Secondly, transitions across a small gap created by static geometric distortions were recently shown to be the cause of the observation of a 35 μeV gap in [V_{15}] in INS measurements, and similar phenomena could occur in the present systemm.[16] Finally, attribution to intra-multiplet transitions between different M_S states is a possibility (see HFEPR results below).

Two further excitation bands are observed at higher energy transfer, as shown in Fig. 6, where the results of measurements performed with IN5 and an incident wavelength of $\lambda = 6$ Å are reported. At $T = 2$ K, a broad peak is observed in the neutron energy-loss side at about 1.0 meV (8.0 cm^{-1}). On warming, the intensity of this excitation band decreases, whilst a second band increases in intensity at lower energy. At $T = 10$ K, the INS response can be fitted to the sum of a gaussian centered at 0.7 meV (5.6 cm^{-1}), with FWHM of 0.5 meV (almost an order of magnitude larger than the resolution), and one centered at 1.0 meV, having a FWHM of 0.25 meV (four times the resolution). These results suggest the presence of many spin energy levels closely spaced around a mean of 1 meV and 1.7 meV. These bands undoubtedly belong to $\Delta S = 0, \pm 1$ transitions from the ground state (1.0 meV) and between excited states (0.7 meV).

High-field EPR (HFEPR). Fig. 7 shows the HFEPR spectra recorded on a pressed powder sample at 285 GHz at various temperatures. The intensities have been normalized to make comparison easier. Further HFEPR spectra were recorded at 190 GHz. At high temperatures, a broad resonance line is observed which is centered around $g = 2.01 \pm 0.01$. On lowering the temperature, first the peak-to-peak separation increases and the line shifts to lower field (higher g value). At the lowest temperature of $T = 1.9$ K, the peak-to-peak separation again decreases somewhat, and $g = 2.14 \pm 0.01$, but no well-resolved structure can be observed. The resonance lines can be attributed to the sum of excitations within several total spin multiplets. At high temperatures many of these multiplets are thermally populated, while lowering of the temperature leads to preferential population of the lowest spin multiplets, which can explain the shift in g value. The plateau between the two peaks as well as the temperature dependence of the peak-to-peak separation suggest the presence of zero-field splitting. For strict cubic symmetry, ZFS can only be present from fourth order, where the spin Hamiltonian will include a term like $\Sigma \hat{S}_\alpha^4$ (α=x,y,z), which leads to the presence of three perpendicular unique axes. However, small deviations from perfect cubic symmetry are expected to lead to sizable second order ZFS. For instance, it was shown that small deviations from tetragonal symmetry can lead to significant second order transverse anisotropy.[Sieber, 2005 #1023; Cornia, 2002 #925] If such small distortions are indeed present, then from the HFEPR spectra an estimate of the ZFS D-parameter can be obtained: Assuming that

at the lowest temperature only the lowest spin multiplet is populated, in the strong field limit, the resonance field difference between lowest and highest field transition cannot be more than $(2S–1)D/g\mu_B$ for half-integer spins. The peak-to-peak separation at $T = 1.9$ K and $\nu = 285$ GHz is $\Delta H \approx 1.5$ T. For $S = 15/2$ (the low end of the possible S values for the total spin of the ground state) this leads to a maximum axial ZFS of $|D| \approx 0.1$ cm^{-1}.

ac susceptibility. In agreement with the small ZFS estimate obtained from HFEPR, no χ" (out-of-phase or imaginary component of the ac susceptibility) signal is observed in ac susceptibility measurements down to $T = 2.0$ K, showing that at those temperatures the magnetization dynamics are fast. At temperatures below 1 K, however, a small but reproducible χ" signal is observed, which goes through a maximum as a function of temperature (Fig. 8). The temperature corresponding to the maximum of this χ" signal is frequency dependent ranging from $T_{max} = 560$ mK for $\nu = 15$ Hz to $T_{max} = 800$ mK for $\nu = 17329$ Hz. The real part (χ') also shows a maximum as a function of temperature, which, in contrast, is virtually frequency independent over the frequency range employed ($T_{max} = 700$ mK). The magnitude of the χ' signal does depend on frequency, increasing by almost 20% going from the highest to the lowest frequency. Finally, χ" is an order of magnitude smaller than χ'. At lower temperatures ($T \approx 420$ mK) there seems to be an additional feature in both χ' and χ", which will not be discussed here.

Clearly, there is some slow magnetization dynamics in this sample. Several physical phenomena can give rise to such slow dynamics. Non mi piace molto questo lungo paragrafo. Stai parlando dei vari phenomeni che possono dare una dinamica lenta osservabile in Xac ma dopo un po' si perde il filo e sembra che si parli del campione..(First of all, a long range order phase transition to a ferromagnetic or antiferromagnetic state gives a frequency-independent χ" signal at the phase transition temperature. Secondly, superparamagnetic or single-molecule-magnet behavior is characterized by an out-of-phase component χ", and a decrease in χ' on lowering the temperature, with the maximum in χ" corresponding to the inflection point in χ'. Furthermore both χ' and χ" are strongly temperature dependent. Since there is usually one characteristic relaxation time, the values of χ" and χ' are the same at the temperature corresponding to the maximum in χ".[17] The relaxation time as a function of temperature can be described well by the Arrhenius law. Thirdly, in spin glasses the broad distribution in relaxation times leads to a much smaller χ" signal compared to χ' (typically one to two orders of magnitude).[17] The χ' signal shows a maximum ("cusp") as a function of temperature, which shifts to lower temperature with decreasing frequency. The frequency shift for spin glasses is defined as $\Delta T_f/T_f\Delta(\log\omega)$, where T_f is the temperature at the maximum in χ", and is typically of the order of a few percent per decade in frequency, reaching 6

% for insulating spin-glass-like materials. Around the temperature of the maximum in χ', χ" rises quickly, and the temperature at which the slope in χ" is largest corresponds to the maximum in χ'. This means that also in spin glasses the frequency shifts of χ' and χ" are equal. On further lowering of the temperature, χ" starts to decrease again, but more slowly. The relaxation time as a function of temperature is not well described by the Arrhenius law)

Although in the [Fe_{13}] measurements, the much smaller χ" signal compared to χ' agrees with a spin-glass interpretation, the lack of frequency dependence of χ' does not. In addition the χ" signal as a function of temperature is more symmetric (Lorentzian) than in a typical spin glass, and the frequency-shift is also larger, viz. $\Delta T_f/T_f\Delta(\log\omega) = 0.11$. Fitting the temperatures of the maxima in χ" at the various frequencies (Fig. 9) to the Arrhenius equation, $\tau = \tau_0\exp(-\Delta E/kT)$, yields an energy barrier of ΔE = 13.10±0.12 K, and τ_0 = (7.5±1.4) x 10^{-13} s. The prefactor τ_0 of 7.5 x 10^{-13} s shows that the slow relaxation is not due to strong intermolecular interactions, which would lead to much smaller values for τ_0, like in spin glasses, [17] but nevertheless it is smaller than typical values for single-molecule magnets (τ_0 is 10^{-7} to 10^{-9} s). This combined with the fact that the Arrhenius law gives a good description of the data agrees with superparamagnetic-like behavior being the cause of the slow dynamics, as evidenced by a non-zero χ" signal. Apparently the local disorder in [Fe_{13}] as suggested by the crystal structure, causes deviations from cubic symmetry leading to finite zero-field splitting of the ground state (in agreement with the HFEPR measurements), but the random intermolecular interactions are too weak to cause spin-glass-like behavior. Since the real part of the ac susceptibility, χ', is much larger than the imaginary part, it must be a minority species that shows slow relaxation of the magnetization. The majority species clearly shows no barrier toward relaxation, therefore no slow magnetization dynamics. As mentioned above, the magnitude of the χ' signal does depend on the measurement frequency. The explanation which agrees with both observations is that the majority part of the sample shows fast underbarrier tunneling of the magnetization with a wide distribution in tunneling frequencies, which would account for the lack of temperature dependent out-of-phase ac susceptibility signal for this part. A certain temperature-independent contribution to χ" (as expected for a tunneling phenomenon) can be observed in Fig. 8, since χ" does not tend to zero at the high- and low-temperature ends. The increase of χ' with decreasing frequencies reflects the fact that at lower frequencies a larger fraction of the molecules have tunneling frequencies such that they can tunnel. In terms of a spin Hamiltonian, the necessary transverse terms will then originate in local distortions of the symmetry as found in the X-ray studies. In short, the magnetization dynamics are very much determined by the disorder in the system, leading to a part that relaxes over a barrier, and another part that tunnels under it.

Mössbauer spectroscopy. In order to characterize the low temperature magnetic properties further, Mössbauer spectroscopy, which is a much shorter-timescale technique than ac susceptibility, was employed. Six spectra of 1024 channels in the ±11 mm s^{-1} velocity range were collected between 2.55 and 13 K (Fig. 11). The 2.55 K spectrum displays well-resolved magnetic hyperfine structure that diminishes as the temperature increases and becomes almost completely absent in the 13 K spectrum. This is indicative of a rapid increase in the spin fluctuation rate with temperature. Although the crystal structure suggests the presence of only two independent iron sites, the analysis of the 2.55 K Mössbauer spectrum showed that it is the sum of four contributions in a 1:4:4:4 ratio. This can be rationalized as follows: due to the frustration among the peripheral iron ions many equivalent topological configurations contribute to the ground state band. If the transitions between these configurations is slow on the Mössbauer timescale, in other words, if the spin fluctuations are slower than the Larmor frequency of the nuclear spin around the instantaneous hyperfine field, the peripheral iron nuclei may experience different hyperfine fields, and consequently, the Mössbauer spectrum would contain a contribution from each of the peripheral sites as well as that from the central iron ion. The relaxation effects on the 2.55 K spectrum are fairly small, so that the hyperfine parameters can be obtained with a good precision (Table 1). The values obtained for these parameters were then used throughout the fitting processes of the other spectra. The isomer shifts of 0.40 mm s^{-1} and negligible quadrupole splittings are as expected for iron(III) ions in relatively symmetric surroundings. Interestingly, the magnetic hyperfine fields are very similar for the central and two of the peripheral sites, while a much smaller value was obtained for the third one. In the absence of an external magnetic field, the hyperfine field experienced by a high spin iron(III) nucleus (B_0^{hf}) is mainly determined by the Fermi contact term, i.e. the polarization of the *s* electrons at the nucleus by the electrons in the partially filled *d* shell, which is of the order of 55 T [18, 19]. This field is modified by the exchange field in an exchange coupled system to $B_{hf} = (\langle S \rangle / S) B_0^{hf}$, where <S> is the mean spin value. From these hyperfine field values, an estimate of the spin ground state value can be obtained, because the $S_z = S$ component of the cluster total spin is given by the algebraic sum of the highest values of the mean spin components s_z of the single ions. As the central-peripheral exchange interaction was calculated to be stronger than the peripheral-peripheral one, we can assume the central spin to be opposite to the resultant peripheral one and also to the two largest peripheral spins. The total cluster spin is then estimated by:

$$S = \frac{2.5}{B_0^{hf}} \left[4\left(B^{(1)} + B^{(2)} - B^{(3)} \right) - B^{(0)} \right] \qquad (4)$$

where B_0^{hf} = 50-55 T (see above), and the superscripts (*i*) (*i*=0,1,2,3) label the central site and the peripheral ones in the same order of Table 1, respectively, resulting in a cluster total spin of S = 12-13 in good agreement with the value *S*=23/2 calculated from the dc susceptibility measurements.

For the peripheral ions, the ratio of the areas of the intermediate and inner lines, is found to be 2.7 instead of 2, as would be expected for isotropic *f*-factors [20], which can be explained by different stiffness of the elastic forces in the radial and tangential directions. The same ratio is 2 for central site, as is expected by considering the high symmetry of its surroundings.

The line shape of each sub-spectrum was estimated by using the simple Blume-Tjon (BT) theory.[21] In the simple BT theory, the effect of spin fluctuations on spectra is simulated by stochastic inversions of the hyperfine field, as would happen during inversion of the magnetization in a superparamagnet. Although the actual values for the linewidths (Γ) and BT transition rates (W) are only estimates (Table 2), it proved to be necessary to take the BT transition rates and linewidths as free parameters for each site (Table 2). This indicates that the spin dynamics cannot be simply described in terms of motion of the total spin, and that the correlations between the motions of the different sites have to be taken into account.

In the 5 K spectrum, the central-site sub-spectrum again displays six lines (Fig. 11*b*), while the inner lines of the peripheral-site 1 sites sub-spectrum and all the lines of the other peripheral sub-spectra are collapsed into central enlarged lines. That is, the fluctuation rate of the central spin is slower than that of the peripheral spins, as we can expect in view of the peripheral spin frustration. In the range from 8 to 13 K, the W values increase considerably with temperature: the peripheral-2 sub-spectrum collapses into an enlarged central line; instead in the sub-spectra of the central and peripheral-1 sites, only the inner lines collapse, while the intermediate and external lines are present but in a very broadened form. This reflects the increase of the spin dynamics rate.

Interestingly, using the values for the energy barrier and prefactor obtained from the Arrhenius fit of the ac susceptibility data, and the typical Mössbauer timescale of 10^{-8} s a blocking temperature of 1.70K is found, which is in good agreement with the observation that at 2.55K most of the spin dynamics in the Mössbauer spectrum has been frozen.

Conclusions.

Competing isotropic exchange interactions lead to a large density of close-lying spin states as shown by dc susceptibility and INS measurements. At low temperatures below 1 K, ac susceptibility measurements show that the magnetization dynamics slow down. Part of the sample shows superparamagnetic-like behaviour evidenced by a frequency-dependent out-of-phase susceptibility signal, while for the other part of the sample the dynamics are limited by the tunnelling frequency. By the Mössbauer spectra, the dynamics were shown to involve mainly the peripheral iron ions.

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Tables

Table 1. Hyperfine interaction parameters obtained from the fit of the 2.55 K Mössbauer spectrum. *IS* is the isomer shift, *QS* the quadrupolar splitting, *B* the hyperfine field, and I_2/I_1 the ratio of the intermediate and inner lines intensities. The numbers in parentheses denote errors in the preceding digit.

Site	*IS* (mm s^{-1})	*QS* (10^{-3} mm s^{-1})	*B* (T)	I_2/I_1
central	0.41(4)	-24(3)	46.5(7)	1.8(8)
peripheral 1	0.40(4)	-26(3)	50(1)	2.7(2)[a]
peripheral 2	0.40(2)	-0.8(29)	43.5(3)	2.7(2)[a]
peripheral 3	0.39(4)	37(3)	9.3(5)	2.7(2)[a]

[a]Constrained to the same value for all three peripheral sites

Table 2. Relaxation parameters (linewidth Γ and Blume-Tjon transition probability per unit time *W*) for the four sites obtained from the best fit of the spectra. Since both *W* and Γ contribute to the linewidths of the spectra, we have expressed also *W* in mm s^{-1} (see text for details). For the ^{57}Fe isotope, the equivalence of the units is: 1 mm s^{-1}=11.6 MHz. For the parameter values without the error indication, only the order of magnitude could be obtained from the fit.

T (K)	c		p 1		p 2		p 3	
	Γ(mms^{-1})	*W* (mms^{-1})	Γ(mms^{-1})	*W* (mms^{-1})	Γ(mms^{-1})	*W* (mms^{-1})	Γ(mms^{-1})	*W* (mm s^{-1})
2.55	0.2(9)	0.1(3)	0.2(8)	0.5	0.5(9)	0.3(5)	0.8	0.3(7)
3	0.2(8)	0.3(3)	0.2(6)	0.7(8)	0.4(9)	0.8(4)	0.5	0.5(3)
4	0.2(8)	0.3(6)	0.2(9)	1.0	0.4	0.7(5)	0.5	0.5(5)
5	0.2(9)	0.3(4)	0.2	1.5	0.4	1.6(8)	0.5	0.8(6)
7.9	1.	1.5	0.2(5)	1.0(8)	0.7(2)	21.	1.0	1.6
13	0.8	2.	0.6	2.0	0.8	93.	1.0	16

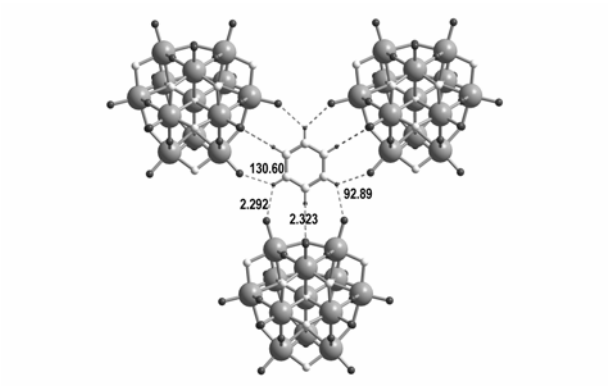


Fig. 1. View along the [111] axis, showing one pyridinium cation and the three surrounding clusters, emphasizing the crystal packing. The big spheres in the clusters are Fe^{III} ions, the small, light spheres are oxygen atoms, the small, dark spheres are fluoride ions. The methanolate methyl groups have been omitted for clarity.

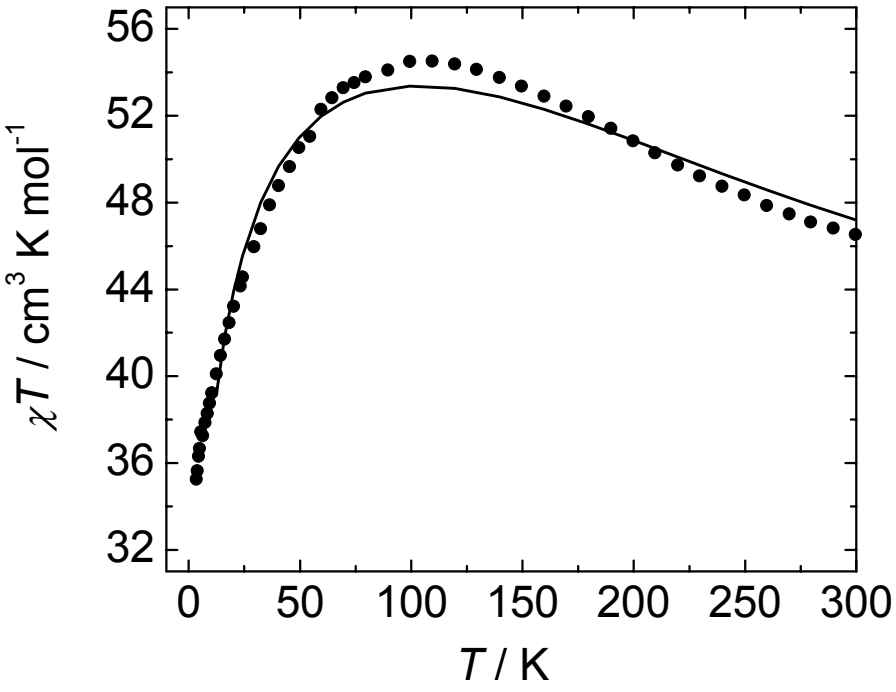


Fig. 2. Observed (●) and calculated (—) χT values. The molar magnetic susceptibility data were obtained at $H = 0.1$ T. The fit was obtained with $J_1 = 43$ cm^{-1} and $J_2 = 30$ cm^{-1}.

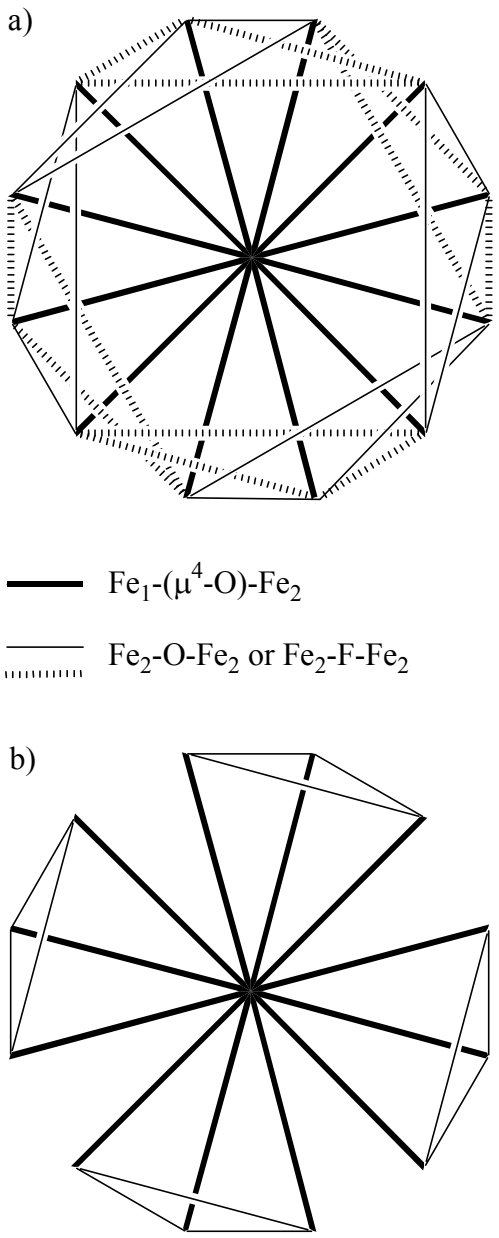


Fig. 3. a) Spin topology of Fe13 b) Spin topology after neglecting one of the intra-triangle interactions.

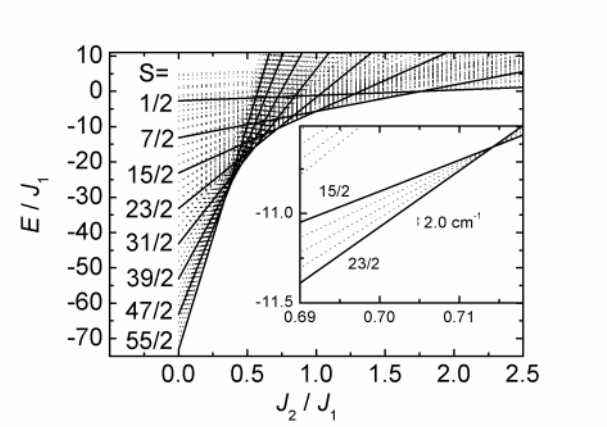


Fig. 4. Plot of the energies of the various spin levels as a function of the ratio of the exchange interaction constants, J_2/J_1. The inset shows the low-lying spin states in the J_2/J_1 = 0.70 region resulting from the fit. The arrow indicates the energy scale for the fitted value of $J_1 = 43$ cm^{-1}.

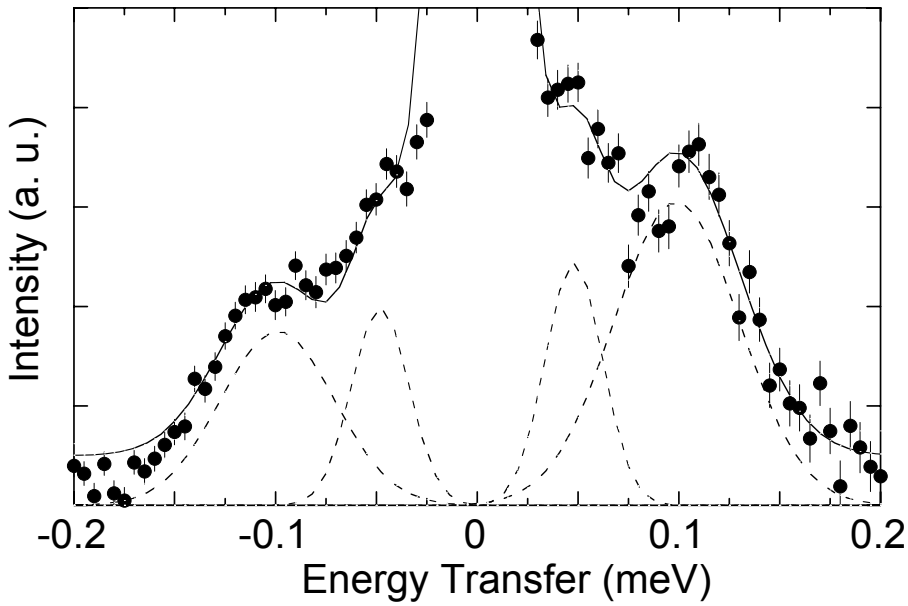


Fig. 5 . INS spectra recorded on IN5 using incident neutrons with λ = 11 Å (energy resolution 12 μeV). Data have been collected at T = 2 K. Solid lines represent the superposition of the elastic peak and four gaussian lineshapes centered at ± 0.05 and ± 0.10 meV. Individual components are shown by broken lines.

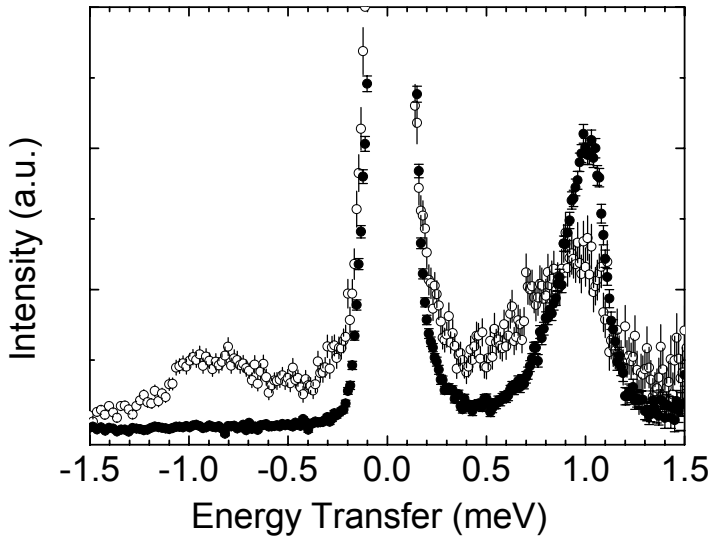


Fig. 6. INS spectra recorded on IN5 using incident neutrons with λ = 6 Å (energy resolution 60 μeV). Data have been collected at T = 2 K (closed circles) and T = 15 K (open circles). The intensity results from the sum over the whole detector bank, covering an angular range from 14.5 to 132.5 degrees.

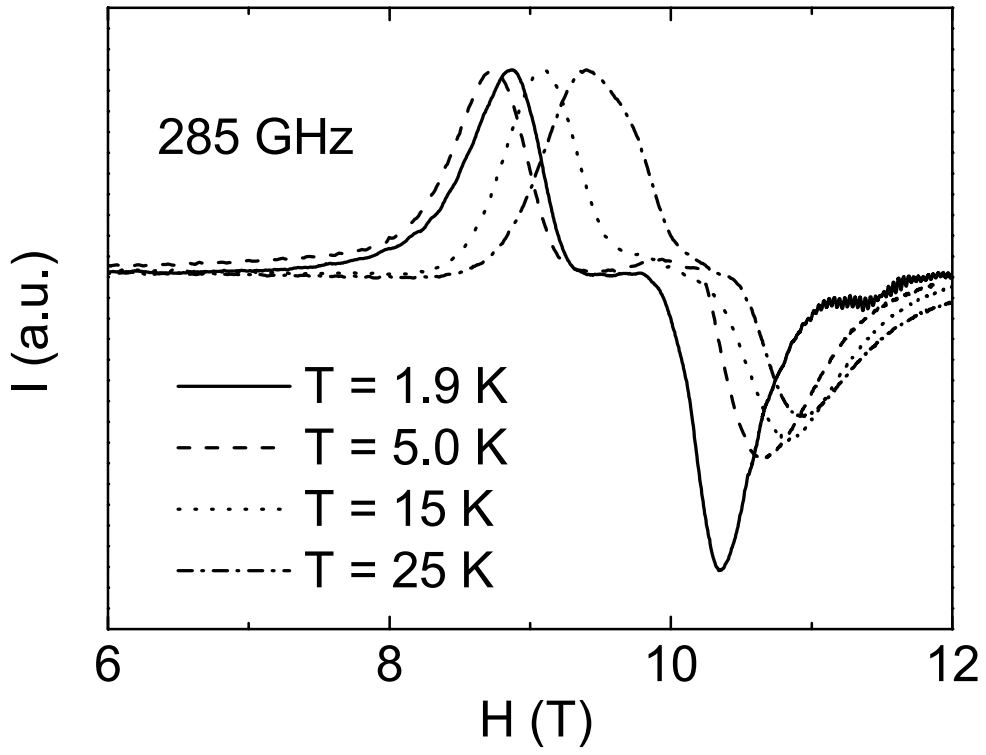


Fig. 7. High-frequency EPR spectra recorded on a pressed powder sample of [Fe_{13}] at ν = 285 GHz at various temperatures as indicated in the Figure.

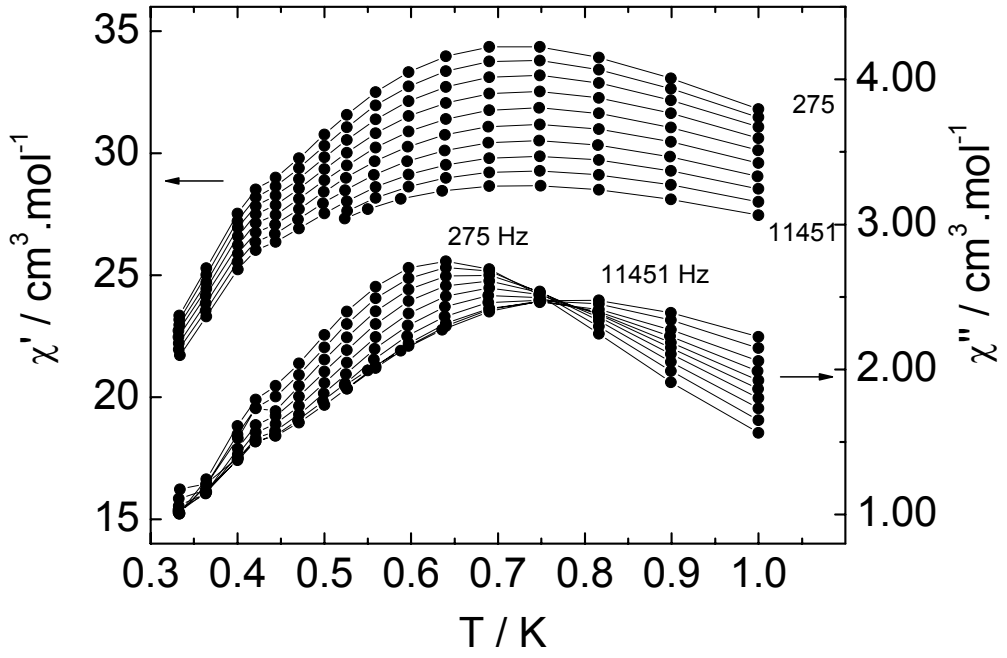


Fig. 8. The real (χ', top) and imaginary (χ", bottom) parts of the ac susceptibility recorded in zero applied field at various driving frequencies ranging between 275 Hz and 11451 Hz as indicated in the Figure.

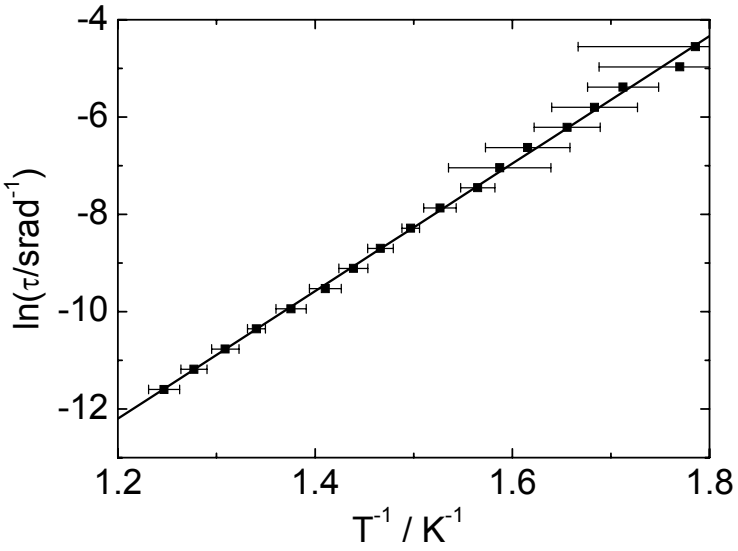


Fig. 9 Measurement time scale as the inverse of the driving frequency as function of the inverse temperature corresponding to the maximum in χ" at that frequency (symbols) and fit (drawn line) according to the Arrhenius law, with $\tau_0 = (7.5 \pm 1.4) \times 10^{-13}$ s, and $\Delta E = 13.10 \pm 0.12$ K.

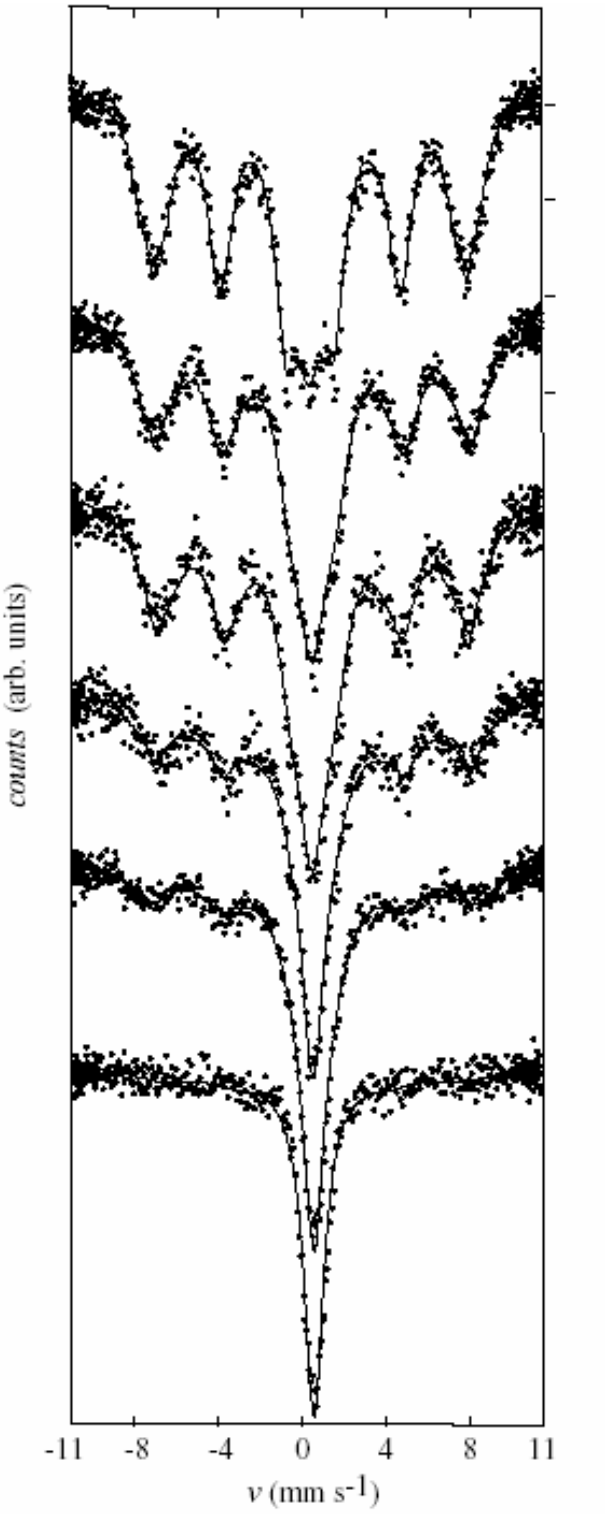


Fig. 10. Mössbauer spectra collected at (starting from the top) T = 2.55, 3, 4, 5, 7.9 and 13 K. Solid lines represent fits obtained as described in the text.

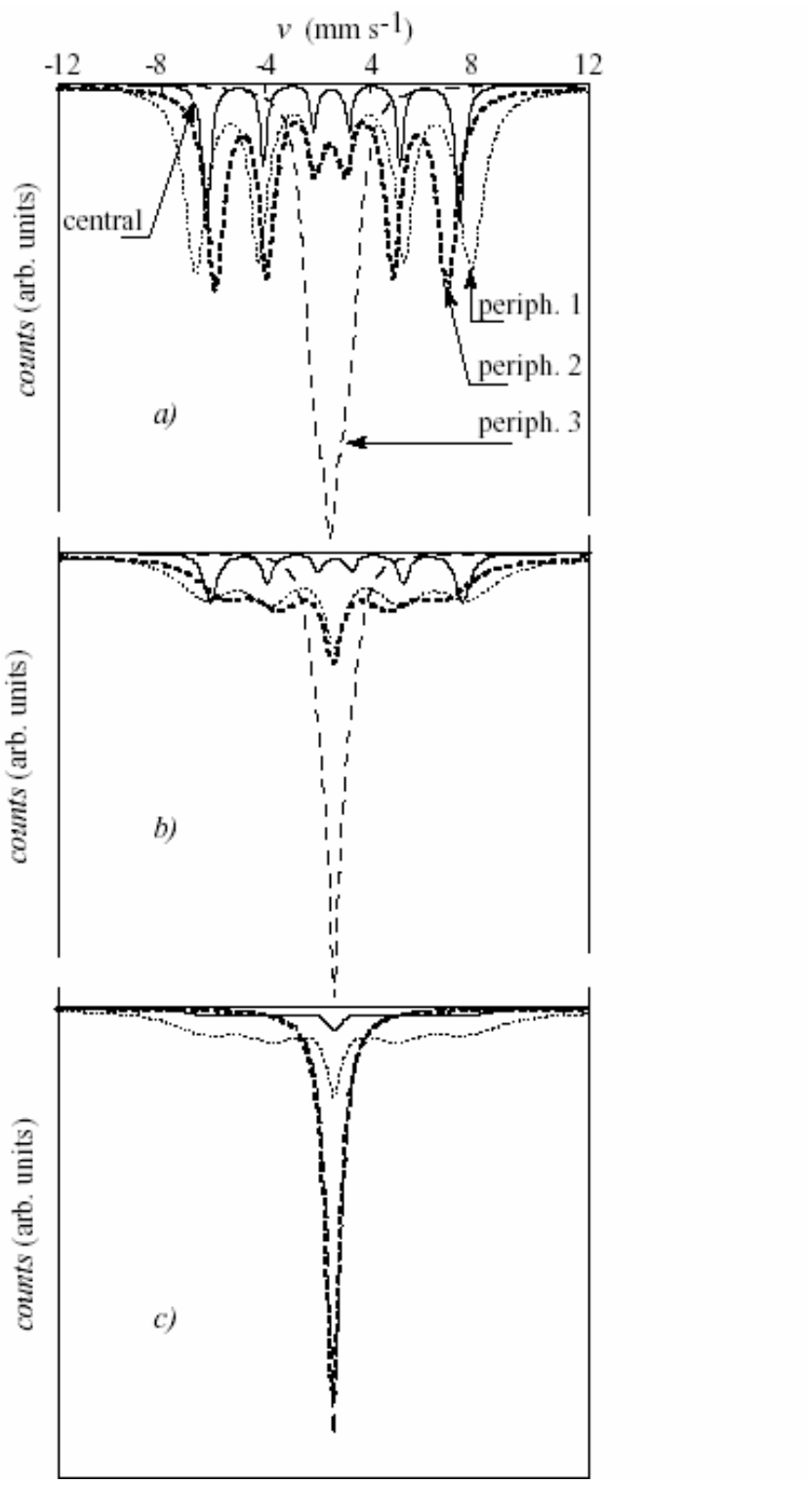


Fig. 11. Spectrum components relative to the four sites, at *a) T*= 2.55 K; *b) T*= 5 K; *c) T*= 13 K.

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