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RHEOLOGICAL STUDY OF AN HYDRATE SLURRY FOR AIR CONDITIONNING APPLICATION

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ABSTRACT

Under atmospheric pressure condition and temperatures between 0°C and +12°C, Tetra-*n*-ButylAmmonium Bromide (TBAB) aqueous solutions crystallise into hydrate slurries. These slurries seems to be well appropriate for cold storage and transportation in the case of air-conditioning applications.

We focus here on the crystallisation and the rheological properties of TBAB hydrates slurries. Once the crystallisation conditions described, we propose an experimental analysis of their flow behaviour. The experimental device is made up of a brushed surface heat exchanger in which the hydrates slurry is generated, and of a measurement loop. The flow behaviour of the slurry is characterised through flow rates and pressure drops measurements. We obtain flow curves of hydrates slurries depending on the volume fraction of hydrate solid particles. In the range of shear rate investigated, hydrate slurries behave as Bingham fluids. We propose a method of determining their apparent viscosities and yield shear stress versus the volume fraction of hydrates.

Keywords : Rheology, Two-phase flow, Heat transportation, Air-conditioning

NOMENCLATURE

A : constant
 D : diameter of the pipe [m]
 D_p : particle diameter [m]
 L : length of the pipe [m]
 M : molecular weight [g.mol⁻¹]
 N : hydration number [-]
 ΔP : pressure drop [Pa]
 Q : volume flow rate [m³.s⁻¹]
 x_{TBAB}, x_H : TBAB, Hydrate, mass fraction [-]
 $\dot{\gamma}, \dot{\gamma}_w$: shear rate, wall shear rate [s⁻¹]
 μ_0, μ_l : Apparent, liquid viscosity [Pa.s]
 ρ_l, ρ_H : liquid, hydrate density [kg.m⁻³]
 τ_w, τ_0 : Wall and yield shear stress [Pa]
 $\overline{\Phi}, \Phi_m$: mean, maximum volume fraction [-]
 ψ : shape factor [-]

INTRODUCTION

In the context of the classical refrigerants phase out, phase change materials (PCM) arouse great interests. Indeed, for very low temperature application such as refrigeration, ice-slurries are spread as secondary refrigerants. Concerning hydrates, they have been used as PCM for cold storage for years. The challenge is now to extend this application to cold transportation in a secondary refrigerant loop. In Japan, JFE marketed the first hydrate slurry for air-conditioning.

The fluid studied is an aqueous solution of Tetra-ButylAmmonium Bromide ((C₄H₉)₄NBr) (TBAB). At low temperature this solution crystallises and forms a suspension of semi-clathrate hydrates. After a presentation of the physical properties of such hydrates, we propose hereafter a rheological study of these TBAB hydrates slurries.

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PHYSICAL PROPERTIES OF TBAB HYDRATES

Thermodynamic properties

In a first time, the existence field of TBAB hydrates was determined experimentally. The experiment consisted in measuring the melting temperature of TBAB hydrates for aqueous solutions of TBAB of different initial concentrations and at atmospheric pressure. Results are presented on figure 1 which represents the melting temperature of hydrates depending on the total concentration of TBAB. This figure highlights the existence of two different hydrate types, called Type A and Type B. These hydrates differ one from the other with their hydration numbers. Tanasawa *et al.* [2] calculated these hydration numbers and gave respectively the values of 26 and 36 for Type A and type B hydrates.

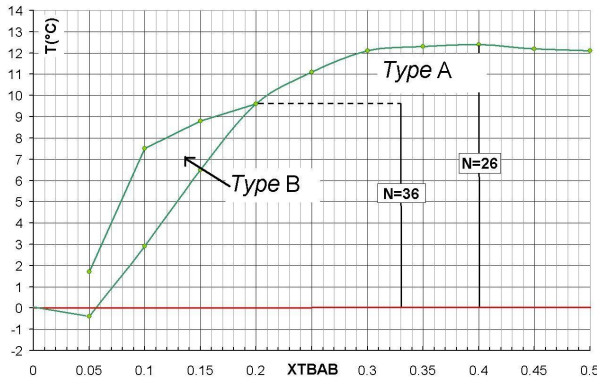


Figure 1: Experimental equilibrium diagram

It is particularly interesting to underline the fact that, at atmospheric pressure, TBAB hydrates are stable for temperatures between 0°C and +12°C. Indeed, in current air-conditioning applications, the cold fluid enters the heat exchanger with a temperature close to +6°C and flow out at about +12°C. This hydrate is therefore a good candidate for air conditioning applications, from a thermodynamic point of view.

From the equilibrium diagram, assuming that the thermodynamics equilibrium is reached, one can deduced from temperature the TBAB mass fraction in the liquid phase. The so-called “moment law”, gives then a mean of calculating the hydrate mass fraction in the suspension (1).

$$x_H = \frac{x_0 - x_l}{x_N - x_l} \quad (1)$$

Figure 2 presents an example of the evolution of the slurry composition versus temperature for a

aqueous solution of 30% mass TBAB. Equations (2) and (3) give the exchange laws between liquid and solid phase for type A and type B TBAB hydrates.

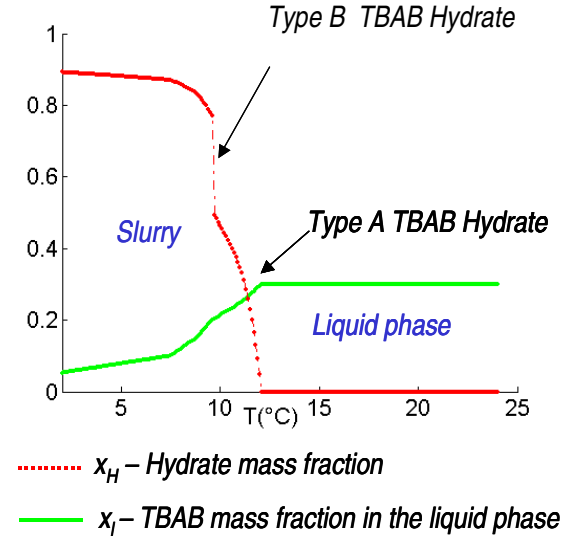


Figure 2: Temperature dependence of the slurry composition ($x_0=0,30$)

TBAB hydrates density

In a second time, the densities of type A and type B hydrates were measured with a pycnometry method consisting in fully crystallising a known mass of TBAB solution. Once hydrates crystallised, hexane was added at the same temperature until the mixture had a precise volume of 10 ml. The mixture is weighted and the hexane volume calculated. The hydrates volume was obtained deducting the hexane volume added.

The densities measured for each hydrate type are presented in Table 1. It also gives the mass fraction of the TBAB in each hydrate crystal depending on its hydration number. Type B hydrate with the highest hydration number is the lightest.

Hydrate type	Hydration number (N)	TBAB mass fraction in the hydrate (x_N)	Density (kg.m^{-3})
A	26	0.41	1082
B	36	0.33	1067

Table 1: Physical properties of type A and type B hydrates.

TBAB hydrates structure

Concerning the crystal structure, type A only is well determined and represented on Figure 2. Davidson [4] first gave a description of this semi-clathrate structure. Then Shimada [3] published figure 3. It is a semi-clathrate hydrate as Nitrogen and Bromine atoms replace some water molecules. The butyl chains are guests of unclosed cavities : 3 tetrakaidecahedrons and one pentakaidecahedron. The dotted lines represent bonds which have disappeared from the classical hydrate structure.

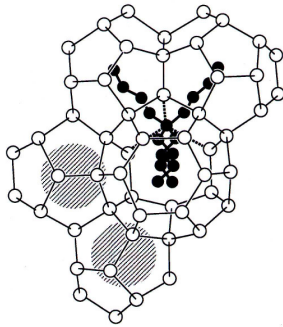


Figure 3: Structure of type A hydrate. [3]

Kinetics effects during TBAB hydrates crystallisation

Decreasing the temperature of TBAB aqueous solutions, hydrates crystals do not appear as soon as temperature reaches their crystallisation temperature : this is the supercooling phenomenon. Figure 3 presents the evolution with time of the temperature of a 20% mass TBAB aqueous solution. The successive apparitions of the two hydrate types are evidenced by two successive temperature increases. The slurry leaves then its metastable state for a stable one. On figure 3 are also represented the different fluid compositions depending on time and temperature (liquid, type A or type B slurries). When type B hydrates appear, type A hydrates definitively disappear.

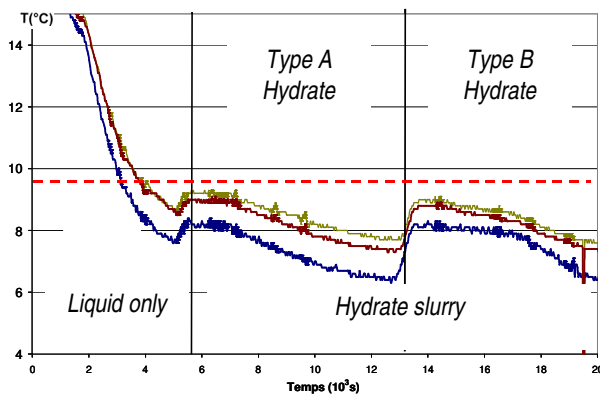


Figure 4: Evidence of the supercooling phenomena

The magnitude of the supercooling defined as the difference between the theoretical temperature of crystallisation and the minimum temperature reached depends on many factors. For instance it is more likely to be great in a small reactor, or if there is no impurity in the fluid.

This supercooling phenomenon is to be avoided or minimized in an industrial context as it represents an additional energy consumption for slurry production. One technique used for example is seeding.

EXPERIMENTAL DEVICE

Set-up description

An “*in-situ*” experimental method is run here in order to avoid some difficulties due to *ex-situ* methods. It is an indirect method based on pressure drop and flow rate measurements. The matter of sliding at the wall still remains but we consider here the global apparent behaviour of the fluid which will be the same as in an industrial installation. The experimental set-up is shown in Figure 5. The hydrate slurry is generated in a brushed surface heat exchanger and drained away with the aid of a gear pump specially developed by GOTECH® for ice-slurry technology. After a 0.5 m long pipe, designed to reach an established hydrodynamic regime, the fluid flows through the “rheometer pipe”. This is a horizontal cylindrical pipe ($L = 1\text{ m}$; $D_{int} = 1.6 \cdot 10^{-2}\text{ m}$) on which the pressure drop between the inlet and outlet is measured with a differential pressure sensor from KELLER® [$0-500\text{ mbar}$] $\pm 0.1\text{ mbar}$. The volumetric flow rate is given by a Coriolis flowmeter from MICRO MOTION® [$0-500\text{ l/h}$] (Accuracy : $\pm 0.5\%$). The loop is also equipped with temperature sensors Pt 100 from PROSENSOR (Accuracy : $\pm 0.1^\circ\text{C}$), and with a transparent section in order to observe the flow regime and the slurry appearance. The rotation frequency of the pump can span the range [$10\text{ Hz}-50\text{ Hz}$] by the use of a frequency controller.

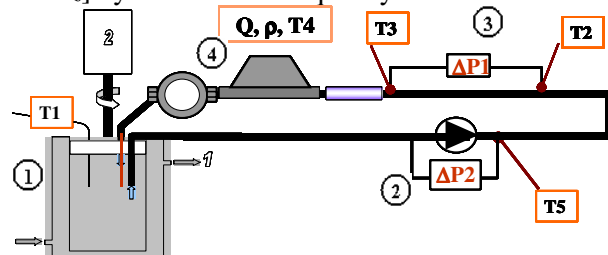


Figure 5: Experimental set-up

Procedure description

Experimental tests were performed on TBAB hydrate slurries to characterise their rheological behaviour. These tests aimed at determining the slurry apparent viscosity for different slurry temperatures, (*i.e.* for different hydrate contents). At a given temperature, the rotation speed of the pump motor is varied between 10 Hz and 50 Hz following regular steps. These steps are 5 Hz high and last about 30 s. Temperature, volumetric flow rate Q and pressure loss ΔP are continuously recorded. The same procedure is repeated for different slurry temperatures, as the hydrate content slowly increases as temperature decreases.

RHEOLOGICAL BEHAVIOUR OF HYDRATE SLURRIES

In this part we propose to study the rheological behaviour of TBAB hydrate slurries produced in the reactor (fig. 5). We first remind some useful tools for the rheological study and then present the results of the experiments.

Theoretical analysis

We consider a laminar and established flow in a horizontal cylindrical pipe. The pressure drop is linked to the wall shear stress through equation (4), derived from a momentum balance.

$$\tau_w = \frac{D\Delta P}{4L} \quad (4)$$

Concerning the shear rate, it can be expressed introducing the volume flow rate thanks to the Rabinovitch relation.

$$\dot{\gamma}_w = \frac{32Q}{\pi D^3} \left(\frac{3n+1}{4n} \right) \text{ with } n = \frac{d \ln(\tau_w)}{d \ln \left(\frac{32Q}{\pi D^3} \right)} \quad (5)$$

In our case, curves representing pressure drop versus flow rate are straight lines which do not intercept the origin. That why, hereafter we describe the rheological behaviour of slurries through an extrapolated shear stress and an apparent viscosity. The behaviour law of such slurries is then similar to a Bingham fluid law :

$$\frac{8\bar{U}}{D} = \frac{1}{\mu_0} \left[\tau_w - \frac{4}{3} \tau_0 \right] \quad (6)$$

$$\begin{cases} \tau - \tau_0 = \mu_0 \dot{\gamma} & \text{if } \tau > \tau_0 \\ \dot{\gamma} = 0 & \text{if } \tau < \tau_0 \end{cases} \quad (7)$$

The actual existence of such a yield shear stress is not obvious, indeed, one can imagine a very rapid decrease of the viscosity with shear rate for very low shear rate. Nevertheless, in the range of mean velocity explored, the slurry behaves just as this yield stress did exist.

Results

Figure 6 presents some examples of flow curves ($\dot{\gamma}_w, \tau_w$) obtained with TBAB hydrate slurries.

These are straight lines which do not intercept the origin. From equation (6), one can say that the intercept point is proportional to the extrapolated or yield shear stress. Concerning the slope it is equal to the apparent viscosity of the slurry. This way, one can associate to each volume fraction of hydrate, an apparent viscosity and a yield shear stress.

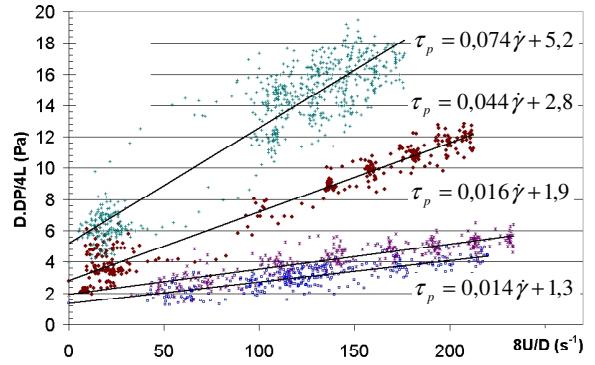


Figure 6: flow curves of hydrate slurries

Apparent viscosity

Figure 7 present the evolution of the relative viscosity of the slurry ($\mu_r = \mu_0/\mu_l$) with hydrate volume fraction. This volume fraction is derived from the mass fraction thanks to equation (8).

$$\bar{\Phi} = \frac{1}{\frac{\rho_H}{\rho_l} \left(\frac{1-x_H}{x_H} \right) + 1} \quad (8)$$

The apparent viscosity of the slurry increases progressively with its load and tends to diverge as the volume fraction of hydrate approaches 60% mass. The apparent viscosity of slurries of type B hydrates seems to diverge for slightly lower mass fraction than type A hydrate.

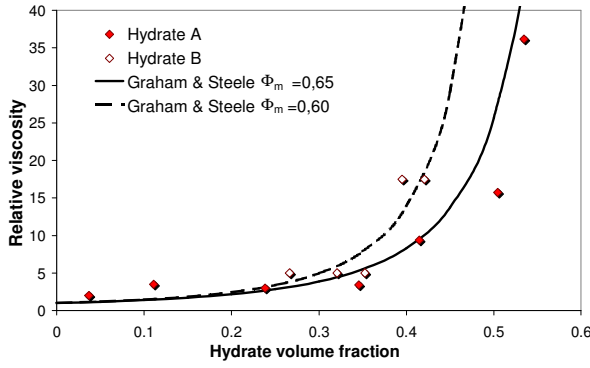


Figure 7: Influence of the hydrate structure on viscosity

The experimental viscosities obtained are then compared with the Graham, Steele and Bird [5] model of viscosity which introduces a maximum packing fraction Φ_m . When particles in the suspension reach this volume fraction Φ_m , one can consider that all the particles are in contact in a more or less porous network. This maximum packing fraction depends on the shape of the crystals and their size distribution.

The expression of suspension viscosity proposed by Graham *et al.* considers an additional volume of the liquid phase V_0 trapped by the particles which can form clusters.

$$V_0 = 1 + \left(\frac{1}{\Phi_m} - 1 \right) \left[1 - \left(\frac{\Phi_m - \bar{\Phi}}{\Phi_m} \right)^2 \right]^{1/2} \quad (9)$$

The corresponding viscosity of the suspension is given by (10).

$$\mu_0 = \mu_s (1 - V_0 \bar{\Phi})^{-2.5} \quad (10)$$

On Figure 7, the results given with this Graham *et al.* model for $\Phi_m = 0.60$ and $\Phi_m = 0.65$ are superposed to the experimental values of apparent viscosity respectively for type A and type B hydrates slurries. It seems that type B hydrates form a suspension less compactable than type A hydrates. Consequently, the structure of the crystals seem to play an important role in the rheological behaviour of their slurries. To give an idea of what represent such maximum packing fractions, one can say that, with monodisperse spheres, $\Phi_m = 0.64$ for a random stack and $\Phi_m = 0.61$ is obtained with an hexagonal arrangement.

EXTRAPOLATED SHEAR STRESS

Figure 8 presents the dependence of the extrapolated shear stress with the volume hydrate fraction for the two hydrate types.

Once again type B hydrate slurries are less transportable as its yield stress are much higher than the ones of type A hydrate slurries.

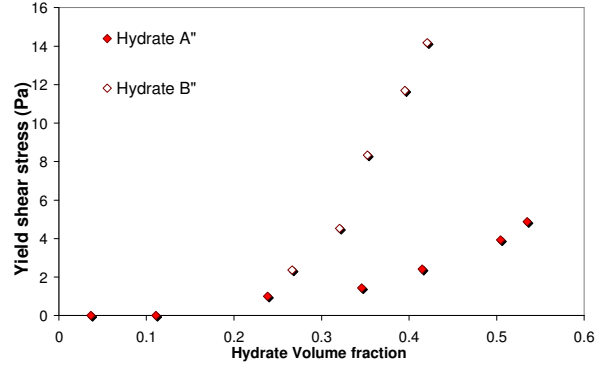


Figure 8: Influence of the hydrate structure on yield shear stress

This yield shear stress is the minimum shear stress which leads to the flow of the fluid. In the case of suspensions it is interpreted as the stress necessary to overcome the interaction forces between particles and permit the flow of the suspension. Therefore, the higher volume fraction of particles, the higher yield stress; and the higher strength of interparticle interactions, the higher yield stress.

With Figure 9 the dependence of the yield shear stress with the volume hydrate fraction is highlighted with a log-log representation. We obtain some straight lines of slopes close to 3. It means that the yield shear stress varies like $\bar{\Phi}^3$. This phenomenon was already observed and evidenced by Thomas [6] for different metal oxides suspensions. The empirical relations he proposed take into account a shape factor ψ (for non-spherical particles), a particle equivalent diameter D_p , and the particle volume fraction:

$$\tau_0 = \frac{A\psi}{D_p^2} \bar{\Phi}^3 \quad (11)$$

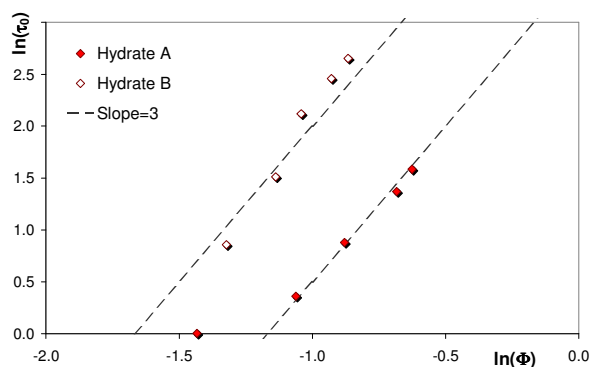


Figure 9 : Dependence of the yield shears stress with hydrate volume fraction.

CONCLUSION

Under atmospheric pressure and temperatures between 0°C and +12°C, aqueous solutions of TBAB can form two types of semi-clathrate TBAB hydrate slurries. Each TBAB hydrate crystal appear after a variable period of metastability and then the thermodynamic equilibrium is regained. In the range of shear rate explored, the rheological behaviour of TBAB hydrate slurries can be described with an apparent viscosity and a extrapolated yield shear stress. This study evidences that TBAB hydrate slurries are transportable until rather high loads of particles (~ 40% volume). Moreover the Graham *et al.* model of viscosity for suspensions seem to be appropriate to describe the evolution of the viscosity of the slurries with the volume fraction of particles if two different maximum packing fractions are introduced for the two hydrate types. Concerning the extrapolated yield shear stress, it varies like $\overline{\Phi}^3$ for the two hydrate types. An additional study of the interparticle interaction

forces should be interesting in order to go further in the interpretation of the yield shear stress.

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