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# QUANTISATION OF LIE-POISSON MANIFOLDS 

SÉBASTIEN RACANIÈRE


#### Abstract

In quantum physics, the operators associated with the position and the momentum of a particle are unbounded operators and $C^{*}$-algebraic quantisation does therefore not deal with such operators. In the present article, I propose a quantisation of the Lie-Poisson structure of the dual of a Lie algebroid which deals with a big enough class of functions to include the above mentioned example.

As an application, I show with an example how the quantisation of the dual of the Lie algebroid associated to a Poisson manifold can lead to a quantisation of the Poisson manifold itself. The example I consider is the torus with constant Poisson structure, in which case I recover its usual $C^{*}$-algebraic quantisation.


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## 1. Introduction

In his PhD thesis, B.Ramazan [6] (see also N. P. Landsman and B. Ramazan [4]) proved a conjecture of Landsman which roughly speaking states that the quantised, that is deformed, algebra of functions on the dual of a Lie algebroid in the direction of its natural Lie-Poisson bracket is the $C^{*}$-algebra of the Lie groupoid integrating

[^0]the Lie algebroid ${ }^{1}$ The type of quantisations that Ramazan considers are deformation quantisations in the sense of M. A. Rieffel [7]. Not all functions are quantised in this way, in fact only functions whose Fourier transform is compactly supported (with respect to a given family of measures) are quantised.

If $M$ is a Riemannian manifold, then its tangent bundle is a Lie algebroid which integrates to the pair groupoid $M \times M$. The induced Lie-Poisson structure on $T^{*} M$ is the usual symplectic structure on a cotangent bundle and Landsman-Ramazan's quantisation can be carried over. Nevertheless, this example shows an important limitation of this procedure: functions that are polynomials in the fibres of $T^{*} M \longrightarrow M$ are not quantised, whereas functions giving the position or the momentum of a particle are of this type. Moreover, it is well-known by physicists that the quantisation of such functions are unbounded operators, whereas LandsmanRamazan's quantisation only gives elements of $C^{*}$-algebras, that is bounded operators on a Hilbert space.

In the present article, I wish to propose a quantisation of the dual of an integrable Lie algebroids $A \longrightarrow M$ which can be used on a wide class of functions. This class contains in particular functions which are polynomial in the fibres of $A^{*} \longrightarrow M$. This will be done in Sections 2 and 3 where Theorem 2.14 is the main result. In Section 4 I show that Theorem 2.14 can be used to recover the physicists's position and momentum operators of a particle moving in $\mathbb{R}^{n}$.

If $M$ is a Poisson manifold, then its cotangent bundle is naturally a Lie algebroid whose dual can be quantised using Theorem 2.14 One might then hope that this quantisation will help finding a quantisation of the original Poisson manifold $M$. This slightly naive idea is shown to work on an example, the torus with constant Poisson structure, in Section 5] and Section 6 There, I recover the usual $C^{*}$-algebraic quantisation of a constant Poisson structure on a torus (see X. Tang and A. Weinstein [8], and A. Weinstein [9]). Part of the strategy of Section 6consists in finding a Poisson map between $T M$ and $M$. Such maps are solutions to a partial differential equation derived in Section 5 In Appendix A, I show how to find a solution to this equation in the case of the sphere in $\mathbb{R}^{3}$.

## 2. Method of quantisation

Let $G \rightrightarrows M$ be a groupoid. Denote by $s$ and $t$ the respective source and target maps of the groupoid $G \rightrightarrows M$. Also, denote by $G^{(2)}=\left\{(x, y) \in G^{2} \mid s(x)=\right.$ $t(y)\}$ the set of pairs of composable arrows in $G$. Let $\tau: A \longrightarrow M$ be Lie algebroid of $G \rightrightarrows M$. I will use the same letter $\tau$ to denote the projection $A^{*} \longrightarrow$ $M$ of the dual of $A$. Choose a Riemannian metric on $A \longrightarrow M$. By duality, this also gives a Riemannian metric on $A^{*}$. I will denote by $X, Y$ or $Z$ elements in $A^{*}$ and by $\xi$ or $\zeta$ elements in $A$.

Definition 2.1. Let $E$ be a s-family of operators on $G$, that is a map $q \longmapsto E_{q}$ from $M$ to the linear forms on $C_{c}^{\infty}\left(s^{-1}(q)\right)$. I will denote $C_{c}^{\infty}(G) \otimes_{s} C_{c}^{\infty}(M)$ the vector space of such operators which in addition satisfy: for all smooth family of

[^1]functions $H$ on $G$ with compact support, that is for all compactly supported smooth function $H$ on $N \times G$ for some manifold $N$, the function
\[

$$
\begin{array}{ccc}
N \times M & \longrightarrow & \mathbb{C} \\
(u, q) & \longmapsto & E_{q}(x \longmapsto H(u, x))
\end{array}
$$
\]

is smooth and compactly supported.
Also I will denote $\mathfrak{O p}(G)$ the vector space of s-family of operators $E$ which in addition satisfy: for all smooth family of functions $H$ on $G$ with compact support, that is for all compactly supported smooth function $H$ on $N \times G$ for some manifold $N$, the function

$$
\begin{array}{ccc}
N \times G & \longrightarrow & \mathbb{C} \\
(u, z) & \longmapsto & E_{t(z)}(x \longmapsto H(u, x z))
\end{array}
$$

is smooth and compactly supported.
Notice that since $M$ is a closed sub-manifold of $G$, the space $\mathfrak{O p}(G)$ is included in $C_{c}^{\infty}(G) \otimes_{s} C_{c}^{\infty}(M)$. On the contrary,
Proposition 2.2. Let $D$ be in $C_{c}^{\infty}(G) \otimes_{s} C_{c}^{\infty}(M)$. For any compactly supported smooth function $H$ on $N \times M$, the map

$$
\begin{array}{ccc}
N \times G & \longrightarrow & \mathbb{C} \\
(u, z) & \longmapsto & D_{t(z)}(H(u, \cdot z))
\end{array}
$$

is smooth. Nevertheless, it might fail to be compactly supported.
Proof. Let $H$ be a compactly supported smooth function on $N \times G$, where $N$ is a manifold. I wish to prove that the map

$$
\begin{array}{ccc}
N \times G & \longrightarrow & \mathbb{C} \\
(u, z) & \longmapsto & \left.D_{t(z)}(H(u, \cdot z))\right)
\end{array}
$$

is smooth. Let $\left(u_{0}, z_{0}\right)$ be a point in $N \times G$. Let $\varphi$ be a compactly supported smooth function on $G$ such that $\varphi \equiv 1$ on a neighbourhood of $z_{0}$. Consider

$$
\begin{array}{cccc}
\widetilde{H}: & \begin{array}{c}
N \times G^{(2)} \\
\\
(u, x, z)
\end{array} & \longmapsto & \mathbb{C} \\
& H(u, x \cdot z) \varphi(z) .
\end{array}
$$

If

$$
K^{\prime}=\left\{(u, x, z) \in N \times G^{(2)} \mid(u, x z) \in \operatorname{supp} H, z \in \operatorname{supp} \varphi\right\},
$$

then the support of $\widetilde{H}$ is a closed subset of $K^{\prime}$ and since one easily checks that $K^{\prime}$ is compact, it follows that $\widetilde{H}$ has compact support. Using the fact that $G^{(2)}$ is a closed sub-manifold of $G \times G$, I extend $\widetilde{H}$ to a function, still denoted $\widetilde{H}$, in $C_{c}^{\infty}(N \times G \times G)$. Let $\widetilde{N}=N \times G$. By interpreting the extended version of $\widetilde{H}$ as a function

$$
\begin{array}{ccc}
\widetilde{N} \times G & \longrightarrow & \mathbb{C} \\
((u, z), x) & \longmapsto & \widetilde{H}(u, x, z),
\end{array}
$$

I can apply $D$ and obtain a function

$$
\begin{array}{ccc}
\tilde{N} \times M & \longrightarrow & \mathbb{C} \\
((u, z), q) & \longmapsto & D_{q}(\widetilde{H}(u, \cdot, z))
\end{array}
$$

in $C_{c}^{\infty}(\widetilde{N} \times M)$. The closed sub-manifold $\{x, z, t(z)\}$ of $\widetilde{N} \times G$ is diffeomorphic to $N \times G$. Therefore $(u, z) \longmapsto D_{t(z)}(H(u, \cdot z) \varphi(z))$ is smooth. This map is equal to $(u, z) \longmapsto D_{t(z)}(H(u, \cdot z))$ in a neighbourhood of $\left(u_{0}, z_{0}\right)$. This can be done for any choice of $\left(u_{0}, z_{0}\right)$, it follows that $(u, z) \longmapsto D_{t(z)}(H(u, \cdot z))$ is smooth.

Nevertheless, the map $(u, z) \longmapsto D_{t(z)}(H(u, \cdot z))$ needs not be compactly supported. Indeed, choose $M$ to be a point, that is $G$ is a genuine group. Let $f$ be a smooth function on $G, \mu=\mathrm{d} g$ be a right invariant measure on $G$ and $D=D_{f}$ be defined as in Proposition 2.3 Then for $h$ a function on $G$, the map $z \longmapsto \int_{G} f(g) h(g z) \mathrm{d} g$ is certainly not compactly supported in general. For example, if $f$ is the constant function equal to 1 then the above map is the constant function equal to $\int_{G} h \mathrm{~d} g$.

Recall that $A \longrightarrow M$ is the Lie algebroid of $G \rightrightarrows M$. Let $T^{s} G \longrightarrow G$ be the vector bundle whose fibre above $x$ in $G$ is $\operatorname{Ker} s_{*, x}$. Denote by $|\Omega|^{1}\left(T^{s} G\right)$ the vector bundle of 1-densities on the fibres of $T^{s} G \longrightarrow G$. Assume that we have a right invariant everywhere positive section $\mu$ of $|\Omega|^{1}\left(T^{s} G\right)$; it defines a right invariant smooth Haar system on $G \rightrightarrows M$. This section is entirely determined by its value along $M$ in $G$, which is a section, denoted by $\mathrm{d} \mu$, of $|\Omega|^{1}(A)$. Equivalently, $\mathrm{d} \mu$ is a smooth family of Lebesgue measures on the fibres of $A \longrightarrow$ $M$. By taking the dual, we obtain a family of Lebesgue measures on $A^{*} \longrightarrow M$, the dual of $A \longrightarrow M$.

Integration provides a way of embedding $C^{\infty}(G)$ in $C_{c}^{\infty}(G) \otimes_{s} C_{c}^{\infty}(M)$.
Proposition 2.3. Let $f$ be in $C^{\infty}(G)$. For each $q$ in $M$, consider the following linear form on $C_{c}^{\infty}\left(s^{-1}(q)\right)$

$$
D_{f, q}: h \longmapsto \int_{s^{-1}(q)} f h \mu .
$$

Then $D_{f}$ is in $C_{c}^{\infty}(G) \otimes_{s} C_{c}^{\infty}(M)$.
If moreover $f$ has compact support, then $D_{f}$ is in $\mathfrak{O p}(G)$.
Proof. Let $N$ be a manifold and $H$ a compactly supported smooth function on $N \times$ $G$. It is clear that the map $(u, q) \longmapsto D_{f, q}(H(u))$ has support in $\left(\operatorname{Id}_{N} \times s\right)\left(\operatorname{supp}_{H}\right)$ which is compact.

Moreover, because $f H$ has compact support, it is a finite sum of functions with support contained in open local charts. Writing things in these local coordinates, it becomes obvious that the map $(u, q) \longmapsto D_{f, q}(H(u))$ is smooth.

In addition, when $f$ has compact support, the map

$$
\begin{aligned}
N \times G^{(2)} & \longrightarrow \mathbb{C} \\
(u, x, z) & \longmapsto f(x) H(u, x z)
\end{aligned}
$$

has compact support. Hence

$$
\begin{aligned}
N \times G & \longrightarrow \mathbb{C} \\
(u, z) & \longmapsto \int_{s^{-1}(t(z))} f(x) H(u, x z)
\end{aligned}
$$

has compact support and $D_{f}$ is in $\mathfrak{O p}(G)$.

The Lie algebroid $\tau: A \longrightarrow M$ is in particular a vector bundle and one can construct a Lie groupoid $A \rightrightarrows M$ with both the source and the target map equal to the projection $\tau: A \longrightarrow M$. In particular, each smooth function on $A$ gives an element of $C_{c}^{\infty}(A) \otimes_{\tau} C_{c}^{\infty}(M)$. Let $f$ be a smooth function on $A^{*}$. Unless the restriction of $f$ to each fibre of $A^{*}$ is $L^{1}$, the Fourier transform of $f$ is not defined. Nevertheless, the Fourier transform of $D_{f}$ is defined for a much larger class of functions.

Definition 2.4. Let $f$ be a smooth function on $A^{*}$. Say that $f$ has polynomial controlled growth if

- for every $q$ in $M$,
- every smooth multi-vector field $v$ on $M$,
- every non-negative integer $k$ and every section $\delta$ of $S^{k} A^{*}$, and
- every trivialisation $\phi:\left.A\right|_{B^{\prime}} \longrightarrow A_{q} \times B^{\prime}$ in a neighbourhood $B^{\prime}$ of $q$, there exists a smaller neighbourhood $B \subset B^{\prime}$ of $q$, a non negative constant $C$ and an integer $m$ such that

$$
\begin{equation*}
(\Upsilon \Delta \cdot f)(Y) \leq C\left(1+\|Y\|^{2}\right)^{m}, \text { for all } Y \text { in }\left.A^{*}\right|_{B} \tag{1}
\end{equation*}
$$

where $\Upsilon$ is the multi-vector field defined on $A^{*}$ using $v$ and the trivialisation $\phi$, and $\Delta$ is the multi-vector field on $A^{*}$ defined using $\delta$ and the vector space structure on the fibres of $A^{*}$.

Denote by $C_{p g}^{\infty}\left(A^{*}\right)$ the set of smooth functions on $A^{*}$ with polynomial controlled growth.

Notice that the above Definition remains unchanged if one replaces $\Upsilon \Delta$ by $\Delta \Upsilon$ in (11). Also, to check if $f$ has polynomial controlled growth, it is enough to check (11) for only one particular choice of trivialisation $\phi$.

The interesting thing about functions with polynomial controlled growth is that one can define the Fourier transform of the operator $D_{f}$ associated to them. This will be a consequence of the following easy Lemma.

Lemma 2.5. Let $(t, \xi)$ be coordinates on $\mathbb{R} \times \mathbb{R}^{n}$ and $K$ a compactly supported smooth function on $\mathbb{R} \times \mathbb{R}^{n}$. If $P$ is any polynomial function on $\mathbb{R}^{n}$ then the map

$$
\begin{aligned}
\mathbb{R} \times \mathbb{R}^{n} & \longrightarrow \mathbb{C} \\
(t, X) & \longmapsto P(X) \int_{\mathbb{R}^{n}} d \xi e^{-i\langle X, \xi\rangle} K(t, \xi)
\end{aligned}
$$

is bounded.
Proof. This is just a simple application of the fact that the Fourier transform takes multiplication by a variable to differentiation with respect to that variable.

If $h$ is a $L^{1}$ function on a vector space $E$ with measure $\mathrm{d}(X)$, set its Fourier transform to be the following function on the dual $E^{*}$ of $E$

$$
\mathfrak{F}(h)(\xi)=\int_{E} \mathrm{~d}(X) e^{-i\langle\xi, X\rangle} h(X) .
$$

Corollary 2.6. If $f$ has polynomial controlled growth, set the Fourier transform of $D_{f}$ to be

$$
\mathfrak{F}\left(D_{f}\right)_{q}(h)=\int_{A_{q}^{*}} d \mu(X) f(X) \mathfrak{F}(h)(X), \forall q \in M, h \in C_{c}^{\infty}\left(\tau^{-1}(q)\right) .
$$

This is a well-defined element of $C_{c}^{\infty}(A) \otimes_{\tau} C_{c}^{\infty}(M)$.
Proof. If $H$ is a compactly supported smooth function on $A$, then using Lemma 2.5 it is easy to prove that $D_{f}(H)$ is a well-defined smooth function on $Q$.

In addition, the support of $D_{f}(H)$ is included in the image of the support of $H$ under the projection $\tau: A \longrightarrow Q$ and is therefore compact.

In Definition 2.11 I define the set of functions acceptable for quantisation as a subset of the set of functions with polynomial controlled growth. One can then apply Proposition 2.13 to see that the set of functions with polynomial controlled growth is big enough for our purpose. Moreover it is a Poisson algebra as the next Lemma shows.

Lemma 2.7. The set of functions with polynomial controlled growth forms a Poisson sub-algebra of $C^{\infty}\left(A^{*}\right)$.

Proof. This is a simple consequence of Lemma 3.6
I now wish to put a structure of algebra on $\mathfrak{O p}(G)$.
Proposition 2.8. Let $D$ and $E$ be two elements of $\mathfrak{O p}(G)$. For $q$ in $M$ and $h$ in $C_{c}^{\infty}\left(s^{-1}(q)\right)$, set

$$
(D \star E)_{q}(h)=E_{q}\left(z \longmapsto D_{t(z)}\left(R_{z}^{*} h\right)\right) .
$$

The operator $D \star E$ lies in $\mathfrak{O p}(G)$.
Proof. Let $N$ be a manifold and $H$ a compactly supported smooth function on $N \times G$. The function $F$ on $N \times G$ defined by $F(u, z)=D_{t(z)}(H(u, \cdot z))$ is smooth and compactly supported because $D$ is in $\mathfrak{O p}(G)$; therefore the function $(u, z) \longmapsto E_{t(z)}(F(u, \cdot z))$ is smooth and compactly supported.

Let $\widetilde{G}$ be the tangent groupoid of $G$ (see M. Hilsum and G. Skandalis [2], A. Weinstein [10] and [11], and A. Connes [1]) with respective source and target maps $\widetilde{s}$ and $\widetilde{t}$. Let $\widetilde{\tau}: \widetilde{A}=\mathbb{R} \times A \longrightarrow \mathbb{R} \times M$ be the Lie algebroid of the tangent groupoid (this is the tangent algebroid, see V. Nistor, A. Weinstein and P. Xu [5]). I will present a method to construct a map from $C_{c}^{\infty}(A) \otimes_{\tau} C_{c}^{\infty}(M)$ to $C_{c}^{\infty}(\widetilde{G}) \otimes_{\tilde{s}} C_{c}^{\infty}(\mathbb{R} \times M)$. Let $\alpha$ be a diffeomorphism from an open neighbourhood $W$ of $M$ in $A$ to an open neighbourhood $V$ of $M$ in $G$ such that

- $\alpha(q)=q$ for $q$ in $M$,
- $s \circ \alpha=\tau$, in particular $\alpha$ sends $A_{q}$ to $s^{-1}(q)$,
- the differential at zero of the restriction of $\alpha$ to $A_{q}$ is the identity map from $A_{q}$ to $A_{q}$.

For example such an $\alpha$ can be obtained from the choice of an exponential map. Let

$$
\widetilde{W}=\{(\hbar, X) \in \mathbb{R} \times A \mid \hbar X \in W\}
$$

be an open subset in $\widetilde{A}$. On it, the map

$$
\widetilde{\alpha}(\hbar, X)=\left\{\begin{array}{l}
(\hbar, \alpha(\hbar X)) \text { for } \hbar \neq 0 \\
(0, X) \text { for } \hbar=0
\end{array}\right.
$$

is a diffeomorphism onto an open neighbourhood $\widetilde{V}$ of $\mathbb{R} \times M$ in $\widetilde{G}$. Choose a smooth function $\psi$ on $A$ with support in $W$ such that $\left.\psi\right|_{\frac{1}{2} W} \equiv 1$. Define $\widetilde{\psi}$ in $C^{\infty}(\mathbb{R} \times A)$ by $\widetilde{\psi}(\hbar, X)=\psi(\hbar X)$.

Proposition 2.9. Let $D$ be in $C_{c}^{\infty}(A) \otimes_{\tau} C_{c}^{\infty}(M)$. For $(\hbar, q)$ in $\mathbb{R} \times M$ and $h$ in $C_{c}^{\infty}\left(\widetilde{s}^{-1}(\hbar, q)\right)$, set

$$
\widetilde{D}_{(\hbar, q)}(h)=D_{q}(X \longmapsto \widetilde{\psi}(\hbar, X) h \circ \widetilde{\alpha}(\hbar, X)) .
$$

The operator $\widetilde{D}$ lies in $C_{c}^{\infty}(\widetilde{G}) \otimes_{\tilde{s}} C_{c}^{\infty}(\mathbb{R} \times M)$.
Proof. Let $N$ be a smooth manifold and $H$ a compactly supported smooth function on $N \times \widetilde{G}$. The function $\widetilde{\psi} \circ \widetilde{\alpha}^{-1}$ defined on $\widetilde{V}$ can be extended, by zero, to a smooth function on the whole of $\widetilde{G}$. The product of this function with $H$ is of course with compact support in $\widetilde{V}$; hence its pull back by $\widetilde{\alpha}$ is compactly supported. It follows that the function

$$
\begin{array}{ccc}
(\mathbb{N} \times \mathbb{R}) \times A & \longrightarrow & \mathbb{C} \\
((u, \hbar), X) & \longmapsto & \psi(\hbar X) H(u, \widetilde{\alpha}(\hbar, X))
\end{array}
$$

is well-defined, smooth and compactly supported. Therefore, I can apply the operator $D$ to it and get a compactly supported smooth function on $N \times \mathbb{R} \times M$. This proves that $\widetilde{D}$ is in $C_{c}^{\infty}(\widetilde{G}) \otimes_{\tilde{s}} C_{c}^{\infty}(\mathbb{R} \times M)$.

Notice that in the above Proof, the function $\psi$ is used to make sense of expressions of the type $\psi(\hbar X) H(u, \widetilde{\alpha}(\hbar, X))$ even when $(\hbar, X)$ is not in $\widetilde{W}$, the domain of definition of $\widetilde{\alpha}$.

Let us see what happens to the product of two operators constructed as in the previous Proposition at $\hbar=0$.
Lemma 2.10. Let $D_{1}$ and $D_{2}$ be in $C_{c}^{\infty}(A) \otimes_{\tau} C_{c}^{\infty}(M)$ such that $\widetilde{D_{1}}$ and $\widetilde{D_{2}}$ are in $\mathfrak{O p}(G)$. Let $H$ be a compactly supported smooth function on $\widetilde{G}$ and $q$ be a point of $M$. If $H_{0}$ denotes the restriction of $H$ to $A \subset \widetilde{G}$ then

$$
\left(\widetilde{D_{1}} \star \widetilde{D_{2}}\right)_{(0, q)}(H)=\left(D_{1} \star D_{2}\right)_{q}\left(H_{0}\right) .
$$

In particular, if $f_{i=1,2}$ are functions on $A^{*}$ such that $\mathfrak{F}\left(D_{f_{1}}\right)$ and $\mathfrak{F}\left(D_{f_{2}}\right)$, respectively $\widetilde{\mathfrak{F}\left(D_{f_{1}}\right)}$ and $\widetilde{\mathfrak{F}\left(D_{f_{2}}\right)}$, are in $\mathfrak{O p}(A)$, respectively $\mathfrak{O p}(\widetilde{G})$, then

$$
\left(\widetilde{\mathfrak{F}\left(D_{f_{1}}\right)} \star \widetilde{\mathfrak{F}\left(D_{f_{2}}\right)}\right)_{(0, q)}(H)=\mathfrak{F}\left(D_{f_{1} f_{2}}\right)_{q}\left(H_{0}\right) .
$$

Proof. The first claim is true because

$$
\begin{aligned}
\left(\widetilde{D_{1}} \star \widetilde{D_{2}}\right)_{(0, q)}(H) & =D_{2, q}\left(Y \longmapsto D_{1, q}(X \longmapsto H(0, X+Y))\right) \\
& =\left(D_{1} \star D_{2}\right)_{q}\left(H_{0}\right) .
\end{aligned}
$$

The second claim is true because

$$
\begin{aligned}
& \left(\widetilde{\mathfrak{F}\left(D_{f_{1}}\right)} \star \widetilde{\mathfrak{F}\left(D_{f_{2}}\right)}\right)_{(0, q)}(H) \\
= & \left(\mathfrak{F}\left(D_{f_{1}}\right) \star \mathfrak{F}\left(D_{f_{2}}\right)\right)_{q}\left(H_{0}\right) \\
= & \int_{A_{q}^{*}} \mathrm{~d} \mu(X) f_{2}(X) \int_{A_{q}} \mathrm{~d} \mu(\xi) e^{-\langle\xi, X\rangle} \int_{A_{q}^{*}} \mathrm{~d} \mu(Y) f_{1}(Y) \int_{A_{q}} \mathrm{~d} \mu(\zeta) e^{-\langle\zeta, Y\rangle} H_{0}(\xi+\zeta) \\
= & \int_{A_{q}^{*}} \mathrm{~d} \mu(X) f_{2}(X) \int_{A_{q}} \mathrm{~d} \mu(\xi) e^{-\langle\xi, X\rangle} \int_{A_{q}^{*}} \mathrm{~d} \mu(Y) f_{1}(Y) \int_{A_{q}} \mathrm{~d} \mu(\zeta) e^{-\langle\zeta-\xi, Y\rangle} H_{0}(\zeta) \\
= & \int_{A_{q}^{*}} \mathrm{~d} \mu(X) f_{2}(X) \mathfrak{F}\left(\mathfrak{F}^{-1}\left(f_{1} \mathfrak{F}\left(H_{0}\right)\right)\right)(X) \\
= & \int_{A_{q}^{*}} \mathrm{~d} \mu(X) f_{2}(X) f_{1}(X) \mathfrak{F}\left(H_{0}\right)(X) \\
= & \mathfrak{F}\left(D_{f_{1} f_{2}}\right)_{q}\left(H_{0}\right) .
\end{aligned}
$$

Of course, not every element $D$ of $C_{c}^{\infty}(A) \otimes_{\tau} C_{c}^{\infty}(M)$ gives an element of $\mathfrak{O p}(\widetilde{G})$ and an important problem is to be able to determine when does $\widetilde{D}$ lie in $\mathfrak{O p}(\widetilde{G})$ ? More precisely, for $f$ in $C^{\infty}\left(A^{*}\right)$, I want to know when does $\widetilde{\mathfrak{F}\left(D_{f}\right)}$ lie in $\mathfrak{O p}(\widetilde{G})$ ?

Definition 2.11 gives an answer to this question.
Let $f$ be a smooth function on $A^{*}$ such that for any $q$ in $M$ and any compactly supported smooth function $h$ on $A_{q}$, the product of the restriction $f_{q}$ of $f$ to $A_{q}^{*}$ by the Fourier transform $\mathfrak{F}(h)$ is again the Fourier transform of a compactly supported smooth function denoted by $m_{f}(q) h$

$$
f_{q} \mathfrak{F}(h)=\mathfrak{F}\left(m_{f}(q) h\right) .
$$

For $N$ a smooth manifold and $\theta: N \longrightarrow M$ a smooth map, denote by $\Theta$ the induced bundle morphism $\theta^{*} A \longrightarrow A$.

Definition 2.11. A smooth function $\widetilde{H}$ on $\theta^{*} A$ is said to be sufficiently compact if
(i) it is in $C_{v c}^{\infty}\left(\theta^{*} A\right)$, the set of vertically compactly supported smooth functions; and
(ii) for any subset $K$ of $A$ which is compact modulo $M$ (that is $K$ is closed and $K \backslash M$ has compact closure $)^{2}$, the set

$$
\left(\operatorname{supp}(\widetilde{H})+\Theta^{-1} K\right) \cap N
$$

is relatively compact. This requirement says that 'small vertical perturbations of the support of $\widetilde{H}$ meet $N$ in a compact set'.
The set of sufficiently compact functions on $\theta^{*} A$ is denoted by $C_{s c}^{\infty}\left(\theta^{*} A\right)$.
A smooth function $f$ on $A^{*}$ is said to be acceptable for quantisation if
(1) $f$ has polynomial controlled growth,
(2) $m_{f}$ preserves $C_{c}^{\infty}(A)$,

[^2](3) $m_{f}$ preserves $C_{s c}^{\infty}\left(\theta^{*} A\right)$ for all manifolds $N$ and smooth functions $\theta$ : $N \longrightarrow M$,
where the action of $m_{f}$ in (2) and (3) is defined fibrewise. The set of smooth functions acceptable for quantisation is denoted by $\mathfrak{Q}\left(A^{*}\right)$.

One reason to state the rather technical above Definition is:
Proposition 2.12. If $f$ is acceptable for quantisation then $\widetilde{\mathfrak{F}\left(D_{f}\right)}$ lies in $\mathfrak{O p}(\widetilde{G})$.
Proof. Let $f$ be acceptable for quantisation. Let $N$ be a manifold and $\widetilde{H}$ a compactly supported smooth function on $N \times \widetilde{G}$. I need to prove that the function

\[

\]

is compactly supported.
Define $\theta: N \times \widetilde{G} \longrightarrow M$ by

$$
\theta(u, \hbar, z)=t(z) .
$$

Let $F$ be the function

$$
\begin{aligned}
\theta^{*} A & \longrightarrow \mathbb{C} \\
(u, \hbar, z, \xi) & \longmapsto \psi(\hbar \xi) \widetilde{H}(u, \widetilde{\alpha}(\hbar, \xi)(\hbar, z)) .
\end{aligned}
$$

I need to prove that $m_{f} F\left(u, \hbar, z, 0_{t(z)}\right)$ is compactly supported in $(u, \hbar, z)$. It will be enough to prove that $F$ is sufficiently compact on $\theta^{*} A$.

Fix $(u, \hbar, z)$ in $N \times \widetilde{G}$. Because $\widetilde{H}$ is compactly supported and because multiplication on the right in a groupoid is a diffeomorphism between two fibres of the source map, the map

$$
\begin{aligned}
N \times \widetilde{s}^{-1}(\hbar, t(z)) & \longrightarrow \mathbb{C} \\
(u, \hbar, x) & \longmapsto \widetilde{H}(u,(\hbar, x)(\hbar, z))
\end{aligned}
$$

is compactly supported. The function $\widetilde{\psi} \circ \widetilde{\alpha}^{-1}$ is defined on an open subset of $\widetilde{s}^{-1}(\hbar, t(z))$ and can be extended by zero to a smooth function on the whole of $\widetilde{s}^{-1}(\hbar, t(z))$. Its product with $\widetilde{H}(u,(\hbar, x)(\hbar, z))$ is compactly supported. This product composed with $\widetilde{\alpha}$ is a compactly supported function on $\theta^{*} A_{(u, \hbar, t(z))}$. This proves that $F$ has vertical compact support.

Let $K$ be a compact modulo $M$ in $A$. I am interested in

$$
\begin{aligned}
& \left(\operatorname{supp} F+\Theta^{-1} K\right) \cap N \times \widetilde{G} \\
\subset & \left\{\left(u, \hbar, z, 0_{t(z)}\right) \mid \exists \xi \in A_{t(z)}, \hbar \xi \in \operatorname{supp} \psi,(u, \widetilde{\alpha}(\hbar, \xi)(\hbar, z)) \in \operatorname{supp} \widetilde{H},-\xi \in K\right\} .
\end{aligned}
$$

Let $\left(u_{j}, \hbar_{j}, z_{j}, 0\right)$ be a sequence in the set on the right hand side of the above inclusion. For each $j$, choose an element $\xi_{j}$ of $-K$ such that $\left(u_{j}, \widetilde{\alpha}\left(\hbar_{j}, \xi_{j}\right)\left(\hbar_{j}, z_{j}\right)\right)$ is in the support of $\widetilde{H}$. This sequence satisfies
(1) $-\xi_{j}$ is a sequence in $K$,
(2) $\hbar_{j} \xi_{j}$ is a sequence in the support of $\psi$,
(3) $\left(u_{j}, \widetilde{\alpha}\left(\hbar_{j}, \xi_{j}\right)\left(\hbar_{j}, z_{j}\right)\right)$ is a sequence in the support of $\widetilde{H}$.

Because the support of $\widetilde{H}$ is compact, we can find a subsequence such that $u_{j_{k}}, \hbar_{j_{k}}$ and $\widetilde{\alpha}\left(\hbar_{j_{k}}, \xi_{j_{k}}\right)\left(\hbar_{j_{k}}, z_{j_{k}}\right)$ converge. Also, since $K$ is compact modulo $M$, we can extract a subsequence such that either $\xi_{j_{k}}$ converges or $\xi_{j_{k}}$ lies in $M$. In the former case, $\hbar_{j_{k}} \xi_{j_{k}}$ converges in supp $\psi$ and $\widetilde{\alpha}\left(\hbar_{j_{k}}, \xi_{j_{k}}\right)$ admits a limit, therefore $\left(\hbar_{j_{k}}, z_{j_{k}}\right)$ converges. In the latter case, $\widetilde{\alpha}\left(\hbar_{j_{k}}, \xi_{j_{k}}\right)\left(\hbar_{j_{k}}, z_{j_{k}}\right)=\left(\hbar_{j_{k}}, z_{j_{k}}\right)$ converges as well. It follows that $\left(\operatorname{supp} F+\Theta^{-1} K\right) \cap N \times \widetilde{G}$ is compact.

## Proposition 2.13. If $f$ is either

- the Fourier transform of a compactly supported smooth function on $A$,
- polynomial in the fibres, that is $f$ is a smooth section of $\bigoplus_{k} S^{k} A$,
- a compactly supported character, i.e. it is of the type $X_{q} \longmapsto e^{i\langle\ell(q), X\rangle}$, where $\ell$ is a compactly supported smooth section of $A$,
then $f$ is in $\mathfrak{Q}\left(A^{*}\right)$.
This Proposition shows that $\mathfrak{Q}\left(A^{*}\right)$ contains indeed many interesting functions.
Proof. Let $f$ be the Fourier transform of a compactly supported smooth function $g$. For each $q$ in $M$, choose a local chart together with a trivialisation of $A$ and $A^{*}$ over it. Writing things in these local chart and local trivialisation, to prove that $f$ has polynomial controlled growth is a simple matter of differentiating under the integral sign in the definition of the Fourier transform.

If $f$ is either polynomial in the fibres or a compactly supported character, it is even more immediate to prove that $f$ has polynomial controlled growth.

Fix a compactly supported smooth function $H$ on $A$, a manifold $N$, a smooth $\operatorname{map} \theta: N \longrightarrow M$ and a sufficiently compact smooth function $\widetilde{H} \in C_{s c}^{\infty}\left(\theta^{*} A\right)$.

Firstly, assume that $f$ is the Fourier transform of a compactly supported smooth function $g$ on $A$. The support of $m_{f} H$ is included in the sum of supp $g$ and the support of $H$, hence it is compactly supported. In the same way, the support of $m_{f} \widetilde{H}$ is included in the sum of the support of $\widetilde{H}$ and $\Theta^{-1} \operatorname{supp} g$. It easily follows that $m_{f} \widetilde{H}$ is again sufficiently compact.

Secondly, assume that $f$ is a smooth section of $\bigoplus_{k} S^{k} A$. For such a function, the operator $m_{f}$ is given by a differential operator $\partial_{f}$ and $\partial_{f} H$ has support included in the support of $H$, therefore $\partial_{f} H$ is in $C_{c}^{\infty}\left(A^{*}\right)$. In the same way, $m_{f}=\partial_{f}$ preserves $C_{s c}^{\infty}\left(\theta^{*} A\right)$. Hence $f$ is acceptable for quantisation.

Finally, assume that $f$ is of the type $f\left(X_{q}\right)=e^{i\langle\ell(q), X\rangle}$. For $\xi$ in $A_{q}, m_{f} H(\xi)=$ $H(\xi+\ell(q))$. Hence $m_{f}$ translates the support of $H$ by $\ell$ on each fibre of $A \longrightarrow M$. Therefore $m_{f} H$ has also compact support. For the same reason, $m_{f} \widetilde{H}$ is also still vertically compactly supported. The support of $m_{f} \widetilde{H}$ is equal to $\operatorname{supp} \widetilde{H}+$ $\Theta^{-1} \operatorname{Im}(-\ell)$. Since $l$ is compactly supported, its image is compact modulo $M$ and $m_{f} \widetilde{H}$ is again sufficiently compact.

Theorem 2.14. Let $G \rightrightarrows M$ be a groupoid with Lie algebroid $\tau: A \longrightarrow M$. Define a quantisation map

$$
\begin{aligned}
\mathcal{Q}: C_{p g}^{\infty}\left(A^{*}\right) & \longrightarrow C_{c}^{\infty}(\widetilde{G}) \otimes_{\tilde{s}} C_{c}^{\infty}(\mathbb{R} \times M) \\
f & \longmapsto \widetilde{\mathfrak{F}\left(D_{f}\right)} .
\end{aligned}
$$

This defines a quantisation of the Poisson manifold $A^{*}$ in the sense that $\mathcal{Q}$ sends the set of functions acceptable for quantisation $\mathfrak{Q}\left(A^{*}\right)$ into $\mathfrak{O p}(\widetilde{G})$; moreover, if $f$ and $g$ are two functions acceptable for quantisation then

$$
\mathcal{Q}(f)^{\star} \mathcal{Q}(g)_{(0, q)}=\mathcal{Q}(f g)_{(0, q)},
$$

and the operator $D=\frac{1}{i \hbar}[\mathcal{Q}(f), \mathcal{Q}(g)]$ is in $C_{c}^{\infty}(\widetilde{G}) \otimes_{\tilde{s}} C_{c}^{\infty}(\mathbb{R} \times M)$ with

$$
D_{(0, q)}=\mathcal{Q}(\{f, g\})_{(0, q)},
$$

for every $q$ in $M$.
Notice that along a non zero $\hbar, \mathcal{Q}(f)$ restricts to an operator

$$
\mathcal{Q}(f)_{\hbar}: C_{c}^{\infty}(G) \longrightarrow C_{c}^{\infty}(G),
$$

while for $\hbar=0$ it restricts to an operator $C_{c}^{\infty}(A) \longrightarrow C_{c}^{\infty}(A)$ which is the Fourier transform of the operator multiplication by $f$ on $C_{c}^{\infty}\left(A^{*}\right)$.

Theorem 2.14 is a consequence of Lemma 2.10 and Corollary 3.8 The proof of this Corollary will take up the whole of next Section.

In Section 4 I show how by applying Theorem [2.14] one recovers the quantisation of the position and momentum operators used by physicists. In Section 6 I will show how to use it to recover the usual strict deformation quantisation of the torus with constant Poisson structure.

## 3. Computation in local coordinates

Let $m=\operatorname{dim} M, U$ be an open subset of $\mathbb{R}^{m}$ and $\varphi$ a diffeomorphism between $U$ and an open subset of $M$

$$
\varphi: U \longrightarrow \varphi(U) \subset M
$$

Let

$$
\begin{aligned}
U \times \mathbb{R}^{n} & \left.\longrightarrow A\right|_{U} \\
(u, \xi) & \longmapsto \gamma(u, \xi)
\end{aligned}
$$

be a trivialisation above $\varphi(U)$ of the vector bundle $A \longrightarrow M$, read in the local chart $(U, \varphi)$. I identify $\mathbb{R}^{n}$ with its dual using the usual euclidean structure of $\mathbb{R}^{n}$. Therefore $\gamma$ also defines a trivialisation $\delta$ of the restriction of $A^{*} \longrightarrow M$ to $\varphi(U)$. This trivialisation is characterised by

$$
\langle\delta(u, X), \gamma(u, \xi)\rangle=\langle X, \xi\rangle,
$$

where $\langle$,$\rangle denotes both the pairing between A$ and $A^{*}$, and the euclidean product on $\mathbb{R}^{n}$.

Choose an open neighbourhood $V_{U}$ of $U \times\{0\}$ in $\mathbb{R}^{m} \times \mathbb{R}^{n}$, such that $\alpha$ is defined on $\gamma\left(V_{U}\right)$. I can define a local chart for $G$

$$
\begin{aligned}
\theta: V_{U} & \longrightarrow G \\
(u, v) & \longmapsto \alpha \circ \gamma(u, v) .
\end{aligned}
$$

Let $V_{U}^{\prime}$ be an open neighbourhood of $U \times\{0\}$ in $V_{U}$ verifying
(1) for each $(u, v)$ in $V_{U}^{\prime}$, there exists $u$ in $U$ such that $t \circ \theta(u, v)=\varphi(u)$,
(2) for each $\left(u_{1}, v_{1}\right)$ and $\left(u_{2}, v_{2}\right)$ in $V_{U}^{\prime}$ with $t \circ \theta\left(u_{2}, v_{2}\right)=\varphi\left(u_{1}\right)$, the product $\theta\left(u_{1}, v_{1}\right) \theta\left(u_{2}, v_{2}\right)$ is in $\theta\left(V_{U}\right)$.
Let $\sigma: V_{U}^{\prime} \longrightarrow U$ be given by

$$
\sigma(u, v)=\varphi^{-1} \circ t \circ \theta(u, v),
$$

(notice that $\varphi^{-1} \circ s \circ \theta(u)=u$ ). Let

$$
V_{U}^{\prime} \operatorname{pr}_{U} \times{ }_{\sigma} V_{U}^{\prime}=\left\{\left(u, v_{1}, v_{2}\right) \mid\left(u, v_{2}\right) \in V_{U}^{\prime},\left(\sigma\left(u, v_{2}\right), v_{1}\right) \in V_{U}^{\prime}\right\}
$$

and define

$$
\begin{aligned}
p: V_{U}^{\prime} \mathrm{pr}_{U} \times_{\sigma} V_{U}^{\prime} & \longrightarrow V_{U} \\
\left(u, v_{1}, v_{2}\right) & \longmapsto \theta^{-1}\left(\theta\left(\sigma\left(u, v_{2}\right), v_{1}\right) \theta\left(u, v_{2}\right)\right) .
\end{aligned}
$$

Ramazan [6 Proposition 2.2.5] proved

$$
\begin{aligned}
p\left(u, v_{1}, v_{2}\right) & =\left(u, v_{1}+v_{2}+B\left(u, v_{1}, v_{2}\right)+\mathrm{O}_{3}\left(u, v_{1}, v_{2}\right)\right), \\
\theta(u, v)^{-1} & =\theta\left(\sigma(u, v),-v+B(u, v, v)+\mathrm{O}_{3}(u, v)\right),
\end{aligned}
$$

where $B\left(u, v_{1}, v_{2}\right)$ is bilinear in $\left(v_{1}, v_{2}\right)$ and $\mathrm{O}_{3}\left(u, v_{1}, v_{2}\right)$, respectively $\mathrm{O}_{3}(u, v)$, is of degree of homogeneity at least 3 in $v_{1}$ and $v_{2}$, respectively $v$.

For $\xi$ in $\mathbb{R}^{n}$

$$
\frac{\partial \theta}{\partial v}(u, v) \xi=\mathrm{d}_{\gamma(u, v)} \alpha \circ \frac{\partial \gamma}{\partial v}(u, v) \xi=\mathrm{d}_{\gamma(u, v)} \alpha \circ \gamma(u, \xi),
$$

because $\gamma$ is linear in $v$. In particular

$$
\frac{\partial \theta}{\partial v}(u, 0) \xi=\gamma(u, \xi),
$$

because $\mathrm{d} \alpha$ is the identity along $M$. Moreover $\varphi(u)=\theta(u, 0)$.
The map

$$
\begin{aligned}
\widetilde{\gamma}: \mathbb{R} \times U \times \mathbb{R}^{n} & \left.\longrightarrow \widetilde{A}\right|_{\mathbb{R} \times U} \\
(\hbar, u, \xi) & \longmapsto(\hbar, \gamma(u, \xi))
\end{aligned}
$$

gives a local trivialisation of the Lie algebroid $\widetilde{A} \longrightarrow \mathbb{R} \times M$ over $\mathbb{R} \times U$. Let

$$
\widetilde{V_{U}}=\left\{(\hbar, u, v) \mid(u, \hbar v) \in V_{U}\right\}
$$

and

$$
\widetilde{V_{U}^{\prime}}=\left\{(\hbar, u, v) \mid(u, \hbar v) \in V_{U}^{\prime}\right\}
$$

I obtain local coordinates on $\widetilde{G}$ by taking $\widetilde{\theta}=\widetilde{\alpha} \circ \widetilde{\gamma}$

$$
\begin{aligned}
\widetilde{\theta}: \widetilde{V_{U}} & \longrightarrow \widetilde{G} \\
(\hbar, u, v) & \longmapsto \widetilde{\alpha}(\hbar, \gamma(u, v))
\end{aligned}
$$

Let $q$ be in $M$ and $\xi$ be in $A_{q}$. Assume that $\xi$ is in the domain of $\alpha$. The map

$$
\begin{aligned}
\mathcal{T}_{\xi}: A_{q} & \longrightarrow A_{t o \alpha(\xi)} \\
\zeta & \left.\longmapsto \frac{\mathrm{d}}{\mathrm{~d} r}\right|_{r=0} \alpha(\xi+r \zeta) \alpha(\xi)^{-1}
\end{aligned}
$$

defines an isomorphism between $A_{q}$ and $A_{t o \alpha(\xi)}$.
Lemma 3.1. Let $u$ be in $U$. Let $\hbar \in \mathbb{R}$ and $\zeta$, $\xi$ in $\mathbb{R}^{n}$

$$
\left.\frac{d}{d \hbar}\right|_{\hbar=0} \widetilde{\alpha}\left(\hbar, \mathcal{T}_{\hbar \gamma(u, \xi)} \circ \gamma(u, \zeta)\right) \widetilde{\alpha}(\hbar, \gamma(u, \xi))=d_{(0, u, \zeta+\xi)} \widetilde{\theta}(1,0,0)
$$

## Proof. I first compute

$$
\begin{aligned}
& \widetilde{\alpha}\left(\hbar,\left.\frac{\mathrm{d}}{\mathrm{~d} r}\right|_{r=0} \alpha \circ \gamma(u, \hbar \xi+r \zeta) \alpha \circ \gamma(u, \hbar \xi)^{-1}\right) \widetilde{\alpha}(\hbar, \gamma(u, \xi)) \\
= & \widetilde{\alpha}\left(\hbar,\left.\frac{\mathrm{d}}{\mathrm{~d} r}\right|_{r=0} \theta(u, \hbar \xi+r \zeta) \theta(u, \hbar \xi)^{-1}\right) \widetilde{\theta}(\hbar, u, \xi) \\
= & \widetilde{\alpha}\left(\hbar,\left.\frac{\mathrm{d}}{\mathrm{~d} r}\right|_{r=0} \theta(u, \hbar \xi+r \zeta) \theta\left(\sigma(u, \hbar \xi),-\hbar \xi+\mathrm{O}\left(\hbar^{2}\right)\right)\right) \widetilde{\theta}(\hbar, u, \xi) \\
= & \widetilde{\alpha}\left(\hbar,\left.\frac{\mathrm{d}}{\mathrm{~d} r}\right|_{r=0} \theta\left(\sigma(u, \hbar \xi), r \zeta+\mathrm{O}\left(\hbar^{2}\right)+B\left(\sigma(u, \hbar \xi), \hbar \xi+r \zeta,-\hbar \xi+\mathrm{O}\left(\hbar^{2}\right)\right)\right.\right. \\
& \left.\left.+\mathrm{O}_{3}\left(\sigma(u, \hbar \xi), \hbar \xi+r \zeta,-\hbar \xi+\mathrm{O}\left(\hbar^{2}\right)\right)\right)\right) \widetilde{\theta}(\hbar, u, \xi) \\
= & \widetilde{\theta} \circ \widetilde{\gamma}^{-1}\left(\hbar, \frac{\partial \theta}{\partial v}(\sigma(u, \hbar \xi), 0)\left(\zeta+B\left(\sigma(u, \hbar \xi), \zeta,-\hbar \xi+\mathrm{O}\left(\hbar^{2}\right)\right)+\mathrm{O}\left(\hbar^{2}\right)\right)\right) \widetilde{\theta}(\hbar, u, \xi)^{3} \\
= & \widetilde{\theta}\left(\hbar, \sigma(u, \hbar \xi), \zeta-\hbar B(\sigma(u, \hbar \xi), \zeta, \xi)+\mathrm{O}\left(\hbar^{2}\right)\right) \widetilde{\theta}(\hbar, u, \xi) \\
= & \left(\hbar, \theta\left(\sigma(u, \hbar \xi), \hbar \zeta-\hbar^{2} B(\sigma(u, \hbar \xi), \zeta, \xi)+\mathrm{O}\left(\hbar^{3}\right)\right) \theta(u, \hbar \xi)\right)^{4} \\
= & \left(\hbar, \theta\left(u, \hbar \zeta-\hbar^{2} B(\sigma(u, \hbar \xi), \zeta, \xi)+\mathrm{O}\left(\hbar^{3}\right)+\hbar \xi+\right.\right. \\
& \left.\left.+B\left(u, \hbar \zeta-\hbar^{2} B(\sigma(u, \hbar \xi), \zeta, \xi)+\mathrm{O}\left(\hbar^{3}\right), \hbar \xi\right)+\mathrm{O}\left(\hbar^{3}\right)\right)\right) \\
= & \left(\hbar, \theta\left(u, \hbar \zeta+\hbar \xi-\hbar^{2} B(\sigma(u, \hbar \xi), \zeta, \xi)+B(u, \hbar \zeta, \hbar \xi)+\mathrm{O}\left(\hbar^{3}\right)\right)\right) \\
= & \left.\widetilde{\theta}\left(\hbar, u, \zeta+\xi-\hbar B(\sigma(u, \hbar \xi), \zeta, \xi)+\hbar B(u, \zeta, \xi)+\mathrm{O}\left(\hbar^{2}\right)\right)\right)
\end{aligned}
$$

The Lemma follows by differentiation with respect to $\hbar$ at 0 .

[^3]Lemma 3.2. Let $f$ and $g$ be in $\mathfrak{Q}\left(A^{*}\right)$. Let $q$ be a point in $M$ and $H$ a compactly supported smooth function on $\widetilde{G}$. Let

$$
N_{f g}(\hbar)=\widetilde{\mathfrak{F}\left(D_{f}\right)} \star \widetilde{\mathfrak{F}\left(D_{g}\right)_{\hbar, q}}(H),
$$

then

$$
\frac{d N_{f g}}{d \hbar}(0)=\mathfrak{F}\left(D_{f^{\prime} g}\right)_{q}\left(H_{0}\right)+\mathfrak{F}\left(D_{f g}\right)_{q}\left(\left.\frac{\partial H}{\partial \hbar}\right|_{\hbar=0}\right),
$$

where

$$
f^{\prime}(Y)=\left.\frac{d}{d \hbar}\right|_{\hbar=0} f\left(Y \circ \mathcal{T}_{\hbar \xi}^{-1}\right),
$$

and $H_{0}=\left.H\right|_{\hbar=0}$.
The term $\left.\frac{\partial H}{\partial \hbar}\right|_{\hbar=0}$ is defined by first pulling back $H$ to a neighbourhood of $\{0\} \times$ $A$ in $\mathbb{R} \times A$ via $\widetilde{\alpha}$, then differentiating with respect to $\hbar$ and finally pushing forward the result via $\widetilde{\alpha}$ again. This definition is actually independent of the choice of $\alpha$.

Proof. We have

$$
\begin{aligned}
N_{f g}(\hbar)= & \int_{A_{q}^{*}} \mathrm{~d} \mu(X) g(X) \int_{A_{q}} \mathrm{~d} \mu(\xi) e^{-i\langle X, \xi\rangle} \int_{A_{t \circ \alpha(\hbar \xi)}^{*}} \mathrm{~d} \mu(Y) f(Y) \\
& \int_{A_{t \circ \alpha(\hbar \xi)}} \mathrm{d} \mu(\zeta) e^{-i\langle Y, \zeta\rangle} \psi(\hbar \xi) \psi(\hbar \zeta) H(\widetilde{\alpha}(\hbar, \zeta) \widetilde{\alpha}(\hbar, \xi))
\end{aligned}
$$

The following change of variables,

- replace $Y$ by $Y \circ \mathcal{T}_{\text {hॄ }}^{-1}$ with $Y \in A_{q}^{*}$,
- replace $\zeta$ by $\mathcal{T}_{\hbar \xi}(\zeta)$ with $\zeta \in A_{q}$,
gives

$$
\begin{aligned}
N_{f g}(\hbar)= & \int_{A_{q}^{*}} \mathrm{~d} \mu(X) g(X) \int_{A_{q}} \mathrm{~d} \mu(\xi) e^{-i\langle X, \xi\rangle} \int_{A_{q}^{*}} \mathrm{~d} \mu(Y) f\left(Y \circ \mathcal{T}_{\hbar \xi}^{-1}\right) \\
& \int_{A_{q}} \mathrm{~d} \mu(\zeta) e^{-i\langle Y, \zeta\rangle} \psi(\hbar \xi) \psi\left(\hbar \mathcal{T}_{\hbar \xi}(\zeta)\right) H\left(\widetilde{\alpha}\left(\hbar, \mathcal{T}_{\hbar \xi}(\zeta)\right) \widetilde{\alpha}(\hbar, \xi)\right) .
\end{aligned}
$$

The changes of variables require to introduce the terms $\operatorname{det} \mathcal{T}_{\hbar \xi}$ and $\operatorname{det} \mathcal{T}_{\hbar \xi}^{-1}$ in the above integral; but these two terms cancel each other since their product is 1 .

Leaving out the justification for it for later, I differentiate the above expression under the integral signs. Since $\psi$ is constant and equal to 1 in a neighbourhood of $M$, it follows that

- $\psi\left(0_{q}\right)=1$ and
- $\left.\frac{\mathrm{d}}{\mathrm{d} \hbar}\right|_{\hbar=0} \psi(\hbar \xi)=\left.\frac{\mathrm{d}}{\mathrm{d} \hbar}\right|_{\hbar=0} \psi\left(\hbar \mathcal{T}_{\hbar \xi}(\zeta)\right)=0$.

Because of Lemma 3.1 and by definition of $\frac{\partial H}{\partial \hbar}$

$$
\left.\frac{\mathrm{d}}{\mathrm{~d} \hbar}\right|_{\hbar=0} H\left(\widetilde{\alpha}\left(\hbar, \mathcal{T}_{\hbar \xi}(\zeta)\right) \widetilde{\alpha}(\hbar, \xi)\right)=\frac{\partial H}{\partial \hbar}(0, \zeta+\xi) .
$$

The Lemma follows since

$$
\begin{aligned}
& \int_{A_{q}^{*}} \mathrm{~d} \mu(X) g(X) \int_{A_{q}} \mathrm{~d} \mu(\xi) e^{-i\langle X, \xi\rangle} \int_{A_{q}^{*}} \mathrm{~d} \mu(Y) f(Y) \int_{A_{q}} \mathrm{~d} \mu(\zeta) e^{-i\langle Y, \zeta\rangle} \frac{\partial H}{\partial \hbar}(0, \zeta+\xi) \\
= & \mathfrak{F}\left(D_{f}\right) \star \mathfrak{F}\left(D_{g}\right)_{q}\left(\left.\frac{\partial H}{\partial \hbar}\right|_{\hbar=0}\right) \\
= & \mathfrak{F}\left(D_{f g}\right)_{q}\left(\left.\frac{\partial H}{\partial \hbar}\right|_{\hbar=0}\right),
\end{aligned}
$$

where the last line is true by Lemma 2.10
There now remains to justify differentiation below the integral signs in

$$
\begin{aligned}
& \int_{A_{q}^{*}} \mathrm{~d} \mu(X) g(X) \int_{A_{q}} \mathrm{~d} \mu(\xi) e^{-i\langle X, \xi\rangle} \int_{A_{q}^{*}} \mathrm{~d} \mu(Y) f\left(Y \circ \mathcal{T}_{\hbar \xi}^{-1}\right) \\
& \int_{A_{q}} \mathrm{~d} \mu(\zeta) e^{-i\langle Y, \zeta\rangle} \psi(\hbar \xi) \psi\left(\hbar \mathcal{T}_{\hbar \xi}(\zeta)\right) H\left(\widetilde{\alpha}\left(\hbar, \mathcal{T}_{\hbar \xi}(\zeta)\right) \widetilde{\alpha}(\hbar, \xi)\right) .
\end{aligned}
$$

Let $\theta$ be the map

$$
\begin{aligned}
\widetilde{V_{U}} & \longrightarrow M \\
(\hbar, \xi) & \longmapsto t \circ \alpha(\hbar \xi) .
\end{aligned}
$$

Define a function $\widetilde{H}$ on $\theta^{*} A$ by

$$
\widetilde{H}(\hbar, \xi, \zeta)=\psi(\hbar \xi) \psi(\hbar \zeta) H(\widetilde{\alpha}(\hbar, \zeta) \widetilde{\alpha}(\hbar, \xi))
$$

Let

$$
\begin{aligned}
S_{1}(\hbar, \xi, Y) & =\int_{A_{q}} \mathrm{~d} \mu(\zeta) e^{-i\langle Y, \zeta\rangle} \widetilde{H}(\hbar, \xi, \zeta) \\
S_{2}(\hbar, \xi) & =\int_{A_{q}^{*}} \mathrm{~d} \mu(Y) f\left(Y \circ \mathcal{T}_{\hbar \xi}^{-1}\right) S_{1}(\hbar, \xi, Y), \\
S_{3}(\hbar, X) & =\int_{A_{q}} \mathrm{~d} \mu(\xi) e^{-i\langle X, \xi\rangle} S_{2}(\hbar, \xi)
\end{aligned}
$$

where $S_{1}$ is defined on $\theta^{*} A^{*}, S_{2}$ is defined on $\widetilde{V_{U}}$ and extended by zero to $\mathbb{R} \times A$ and $S_{3}$ is defined on $\mathbb{R} \times A^{*}$.

I claim that $\widetilde{H}$ is in $C_{s c}^{\infty}\left(\theta^{*} A\right)$. The proof of this claim is similar to that for $F$ in the proof of Proposition 2.12 and will not be reproduced here. It follows that

$$
S_{2}(\hbar, \xi)=m_{f} \widetilde{H}(\hbar, \xi, 0)
$$

is compactly supported in $(\hbar, \xi)$.
For $\xi$ fixed, the function $\widetilde{H}$ is compactly supported in $(\hbar, \zeta)$, hence derivation below the integral sign in $S_{1}$ is possible.

Since $f$ has polynomial controlled growth, for $\xi$ fixed, there exists a positive constant $C$, an $\epsilon>0$ and an integer $m$ such that

$$
\left|\frac{\partial}{\partial \hbar}\left(f\left(Y \circ \mathcal{T}_{\hbar \xi}^{-1}\right) S_{1}(\hbar, \xi, Y)\right)\right| \leq C\left(1+\|Y\|^{2}\right)^{m}\left(\left|S_{1}(\hbar, \xi, Y)\right|+\left|\frac{\partial}{\partial \hbar} S_{1}(\hbar, \xi, Y)\right|\right)
$$

By Lemma 2.5 both terms on the right hand side of the above inequality are bounded by a smooth $L^{1}$ function independent of $\hbar$. Differentiation below the integral sign in $S_{2}$ is therefore possible.

Since $S_{2}$ is compactly supported, differentiation below the integral sign in $S_{3}$ is possible.

To finish, since $g$ has polynomial controlled growth and by Lemma 2.5 differentiation below the integral sign in $N_{f g}$ is possible.
Lemma 3.3. Let $\xi$ and $\zeta$ be in $\mathbb{R}^{n}$, then

$$
\left.\frac{d}{d \hbar}\right|_{\hbar=0} \mathcal{T}_{\gamma(u, \hbar \xi)} \circ \gamma(u, \zeta)=\frac{\partial \gamma}{\partial u}(u, \zeta) \circ \frac{\partial \sigma}{\partial v}(u, 0) \xi-\gamma(u, B(u, \zeta, \xi)) .
$$

Proof. First, I compute

$$
\begin{aligned}
& \mathcal{I}_{\gamma(u, \hbar \xi)} \circ \gamma(u, \zeta) \\
= & \left.\frac{\mathrm{d}}{\mathrm{~d} r}\right|_{r=0} \theta(u, \hbar \xi+r \zeta) \theta\left(\sigma(u, \hbar \xi),-\hbar \xi+\mathrm{O}\left(\hbar^{2}\right)\right) \\
= & \left.\frac{\mathrm{d}}{\mathrm{~d} r}\right|_{r=0} \theta\left(\sigma(u, \hbar \xi), r \zeta+B\left(\sigma(u, \hbar \xi), \hbar \xi+r \zeta,-\hbar \xi+\mathrm{O}\left(\hbar^{2}\right)\right)+\right. \\
& \left.+\mathrm{O}_{3}\left(\sigma(u, \hbar \xi), \hbar \xi+r \zeta,-\hbar \xi+\mathrm{O}\left(\hbar^{2}\right)\right)\right) \\
= & \frac{\mathrm{d} \theta}{\mathrm{~d} v}(\sigma(u, \hbar \xi), 0)\left(\zeta-\hbar B(\sigma(u, \hbar \xi), \zeta, \xi)+\mathrm{O}\left(\hbar^{2}\right)\right) \\
= & \gamma\left(\sigma(u, \hbar \xi), \zeta-\hbar B(\sigma(u, \hbar \xi), \zeta, \xi)+\mathrm{O}\left(\hbar^{2}\right)\right) .
\end{aligned}
$$

The result follows by differentiation and because $\gamma$ is linear in the second variable.

The map $\gamma: U \times\left.\mathbb{R}^{n} \longrightarrow A\right|_{U}$ is a local trivialisation of $A$. The induced local trivialisation of $A^{*}$

$$
\delta: U \times\left.\mathbb{R}^{n} \quad \longrightarrow \quad A^{*}\right|_{U}
$$

is characterised by

$$
\langle\delta(u, X), \gamma(u, \zeta)\rangle=\langle X, \zeta\rangle,
$$

where $\langle$,$\rangle denotes the euclidean product on \mathbb{R}^{n}$.
Lemma 3.4. Let $e_{1}, \ldots, e_{n}$ be a basis of $\mathbb{R}^{n}$, then for $(u, \xi)$ in $U \times \mathbb{R}^{n}$ and $Y$ in $A_{\varphi(u)}^{*}$

$$
\left.\frac{d}{d \hbar}\right|_{\hbar=0} \delta^{-1}\left(Y \circ \mathcal{T}_{\hbar \gamma(u \xi)}^{-1}\right)=\left(\frac{\partial \sigma}{\partial v}(u, 0) \xi, \sum_{k} Y\left(\gamma\left(u, B\left(u, e_{k}, \xi\right)\right)\right) e_{k}\right) .
$$

Proof. In the proof of Lemma 3.3, I showed that

$$
\mathcal{T}_{\gamma(u, \hbar \xi)} \circ \gamma(u, \zeta)=\gamma\left(\sigma(u, \hbar \xi), \zeta-\hbar B(\sigma(u, \hbar \xi), \zeta, \xi)+\mathbf{O}\left(\hbar^{2}\right)\right),
$$

therefore

$$
\mathcal{T}_{\gamma(u, \hbar \xi)}^{-1} \circ \gamma(\sigma(u, \hbar \xi), \zeta)=\gamma\left(u, \zeta+\hbar B(u, \zeta, \xi)+\mathrm{O}\left(\hbar^{2}\right)\right) .^{5}
$$

The Lemma is proved by using this formula when differentiating

$$
\delta^{-1}\left(Y \circ \mathcal{T}_{\hbar \gamma(u, \xi)}^{-1}\right)=\left(\sigma(u, \hbar \xi), \sum_{k} Y\left(\mathcal{T}_{\hbar \gamma(u \xi)}^{-1} \circ \gamma\left(\sigma(u, \hbar \xi), e_{k}\right)\right) e_{k}\right) .
$$

Write

$$
B\left(u, e_{k}, e_{h}\right)=\sum_{j} B_{k, h}^{j} e_{j} .
$$

In particular, the $B_{k h}^{j}$ 's depend on $u$.
Corollary 3.5. For $Y=\sum_{j} Y_{j} e_{j}$ and $\xi=\sum_{h} \xi_{h} e_{h}$ in $\mathbb{R}^{n}$

$$
\left.\frac{d}{d \hbar}\right|_{\hbar=0} \delta^{-1}\left(\delta(u, Y) \circ \mathcal{T}_{\hbar \gamma(u, \xi)}^{-1}\right)=\left(\frac{\partial \sigma}{\partial v}(u, 0) \xi, \sum_{k, h, j} Y_{j} B_{k, h}^{j} \xi_{h} e_{k}\right) .
$$

Let $f$ be a smooth function on $A^{*}$. Define $F=f \circ \delta$, then

$$
\left.\frac{d}{d \hbar}\right|_{\hbar=0} f\left(\delta(u, Y) \circ \mathcal{T}_{\hbar \gamma(u, \xi)}^{-1}\right)=\sum_{k, h} \xi_{k} \frac{\partial F}{\partial u_{h}}(u, Y) \frac{\partial \sigma_{h}}{\partial v_{k}}(u, 0)+\sum_{k, h, j} Y_{j} \xi_{h} \frac{\partial F}{\partial Y_{k}}(u, Y) B_{k, h}^{j} .
$$

Proof. The first formula is just Lemma 3.4 written in local coordinates. The second one is a straightforward computation.

Let us look at the Poisson bracket on $A^{*}$ in local coordinates.
Lemma 3.6. Let $f$ and $g$ be smooth functions on $A^{*}$. Define $F=f \circ \delta$ and $G=g \circ \delta$, smooth functions on $U \times \mathbb{R}^{n}$. Set

$$
\{F, G\}=\{f, g\} \circ \delta .
$$

Let $(u, Z)$ be in $U \times \mathbb{R}^{n}$ and denote $\frac{\partial F}{\partial u_{j}}$ for $\frac{\partial F}{\partial u_{j}}(u, Z)$. I will use similar notations for $\frac{\partial G}{\partial u_{j}}, \frac{\partial F}{\partial Z_{k}}$ and $\frac{\partial G}{\partial Z_{k}}$. Then
$\{F, G\}(u, Z)=\sum_{k, h, j} \frac{\partial F}{\partial Z_{k}} \frac{\partial G}{\partial Z_{h}}\left(B_{h k}^{j}-B_{k h}^{j}\right) Z_{j}+\sum_{k, h}\left(\frac{\partial F}{\partial Z_{k}} \frac{\partial G}{\partial u_{h}}-\frac{\partial F}{\partial u_{h}} \frac{\partial G}{\partial Z_{k}}\right) \frac{\partial \sigma_{h}}{\partial v_{k}}(u, 0)$.
Proof. This is essentially Equation (1.2.6) and Proposition 2.2.6 in Ramazan [6] where it is proved that ${ }^{6}$

$$
\left[\gamma\left(u, e_{k}\right), \gamma\left(u, e_{h}\right)\right]=\sum_{j}\left(B_{k h}^{j}-B_{h k}^{j}\right) \gamma\left(u, e_{j}\right)
$$

and, if $u_{j}^{*}$ is the $j$-th coordinate map on $U$,

$$
\rho\left(e_{k}\right) \cdot u_{h}^{*}(u)=\frac{\partial \sigma_{h}}{\partial v_{k}}(u, 0)
$$

[^4]Proposition 3.7. Let $f$ and $g$ be in $\mathfrak{Q}\left(A^{*}\right)$. Let $H$ be a compactly supported smooth function on $\widetilde{G}$ and $q$ a point in $M$, then

$$
\left.\frac{d}{d \hbar}\right|_{\hbar=0}\left(\widetilde{\mathfrak{F}\left(D_{f}\right)} \star \widetilde{\mathfrak{F}\left(D_{g}\right)}-\widetilde{\mathfrak{F}\left(D_{g}\right)} \star \widetilde{\mathfrak{F}\left(D_{f}\right)}\right)_{(\hbar, q)}(H)=i \widetilde{\mathfrak{F}\left({\left.\widetilde{D_{\{f, g\}}}\right)_{(0, q)}}(H) . . . . . . . .\right.}
$$

Proof. The left hand side of the above equation is equal to

$$
\left.\frac{\mathrm{d}}{\mathrm{~d} \hbar}\right|_{\hbar=0}\left(N_{f g}-N_{g f}\right)
$$

The terms

$$
\mathfrak{F}\left(D_{f g}\right)_{q}\left(\left.\frac{\partial H}{\partial \hbar}\right|_{\hbar=0}\right)
$$

appear in both $\left.\frac{\mathrm{d}}{\mathrm{d} \hbar}\right|_{\hbar=0} N_{f g}$ and $\left.\frac{\mathrm{d}}{\mathrm{d} \hbar}\right|_{\hbar=0} N_{g f}$ in Lemma 3.2, they will therefore cancel each other when taking the difference. The other term in $\left.\frac{\mathrm{d}}{\mathrm{d} \hbar}\right|_{\hbar=0} N_{f g}$, when using Corollary 3.5, becomes a sum of terms. These terms can be dealt with by recalling that the Fourier transform takes the operator 'multiplication by a variable' to the operator 'derivation with respect to this variable'. For example (with some slight abuse of notations)

$$
\begin{aligned}
& \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(X) G(u, X) \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\xi) e^{-i\langle X, \xi\rangle} \xi_{k} \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(Y) \frac{\partial F}{\partial u_{h}}(u, Y) \frac{\partial \sigma_{h}}{\partial v_{k}}(u, 0) \\
& \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\zeta) e^{-i\langle Y, \zeta\rangle} H_{0}(\xi+\zeta) \\
= & \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(X) G(u, X) i \frac{\partial}{\partial X_{k}} \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\xi) e^{-i\langle X, \xi\rangle} \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(Y) \frac{\partial F}{\partial u_{h}}(u, Y) \frac{\partial \sigma_{h}}{\partial v_{k}}(u, 0) \\
& \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\zeta) e^{-i\langle Y, \zeta\rangle} H_{0}(\xi+\zeta) \\
= & -i \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(X) \frac{\partial G}{\partial X_{k}}(u, X) \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\xi) e^{-i\langle X, \xi\rangle} \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(Y) \frac{\partial F}{\partial u_{h}}(u, Y) \frac{\partial \sigma_{h}}{\partial v_{k}}(u, 0) \\
& \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\zeta) e^{-i\langle Y, \zeta\rangle} H_{0}(\xi+\zeta) \\
= & -i \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(Z) \frac{\partial G}{\partial Z_{k}}(u, Z) \frac{\partial F}{\partial u_{h}}(u, Z) \frac{\partial \sigma_{h}}{\partial v_{k}}(u, 0) \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\xi) e^{-i\langle Z, \xi\rangle} H_{0}(\xi),
\end{aligned}
$$

where the last line is true for the same reason that $\mathfrak{F}\left(D_{f}\right) \star \mathfrak{F}\left(D_{g}\right)=\mathfrak{F}\left(D_{f g}\right)$ (see Lemma 2.10). A similar computation leads to

$$
\begin{aligned}
& \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(X) G(u, X) \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\xi) e^{-i\langle X, \xi\rangle} \xi_{h} \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(Y) Y_{j} B_{k h}^{j} \frac{\partial F}{\partial Y_{k}}(u, Y) \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\zeta) e^{-i\langle Y, \zeta\rangle} H_{0}(\xi+\zeta) \\
= & -i \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(Z) \frac{\partial G}{\partial Z_{h}}(u, Z) \frac{\partial F}{\partial Z_{k}}(u, Z) Z_{j} B_{k h}^{j} \int_{\mathbb{R}^{n}} \mathrm{~d} \mu(\xi) e^{-i\langle Z, \xi\rangle} H_{0}(\xi) .
\end{aligned}
$$

These computations together with Lemma 3.6 prove Proposition 3.7

As a corollary, I obtain the following:
Corollary 3.8. Let $f$ and $g$ be two smooth functions on $A^{*}$ acceptable for quantisation. Then

$$
\frac{1}{i \hbar}\left[\widetilde{\mathfrak{F}\left(D_{f}\right)}, \widetilde{\mathfrak{F}\left(D_{g}\right)}\right]
$$

is a well-defined element of $C_{c}^{\infty}(\widetilde{G}) \otimes_{\tilde{s}} C_{c}^{\infty}(\mathbb{R} \times M)$, which along $\hbar=0$ is equal to

$$
\widetilde{\mathfrak{F}\left(D_{\{f, g\}}\right)_{0}} .
$$

## 4. Quantisation of $\mathbb{R}^{2 n}$

In this short section, I will discuss the case of the quantisation of observables on the phase space of a particle moving in $\mathbb{R}^{n}$.

Let $M=\mathbb{R}^{n}$ with its euclidean structure and $A$ be the tangent bundle of $M$, that is $A=\mathbb{R}^{2 n}$. With these conditions, the Lie groupoid $G$ integrating $A$ is the pair groupoid $M \times M=\mathbb{R}^{n} \times \mathbb{R}^{n}$ with source map, tangent map and product

$$
s(p, q)=q, t(p, q)=p,(r, p) \cdot(p, q)=(r, q) .
$$

The euclidean product gives a natural family of measures on the fibres of $A$. I can take $\alpha$ to be defined on the whole of $A$ by

$$
\alpha(q, \xi)=(q+\xi, q) .
$$

Also, I can choose $\psi$ to be equal to the constant function 1 . The space of morphisms of the tangent groupoid is diffeomorphic to $\mathbb{R} \times \mathbb{R}^{n} \times \mathbb{R}^{n}$. Let $H$ be a compactly supported smooth function on $\mathbb{R} \times \mathbb{R}^{n} \times \mathbb{R}^{n}$; it is a a function of $(\hbar, p, q)$. Let $f$ be a function acceptable for quantisation. For $\hbar \neq 0$ and $x=(p, q)$ in $G$

$$
\begin{aligned}
\mathcal{Q}_{\hbar}(f)(H)(x) & =\int_{A_{p}^{*}} \mathrm{~d} X f(p, X) \int_{A_{p}} e^{-i\langle\xi, X\rangle} H(\widetilde{\alpha}(\hbar, p, \xi) \cdot(\hbar, p, q)) \\
& =\int_{\mathbb{R}^{n}} \mathrm{~d} X f(p, X) \int_{\mathbb{R}^{n}} e^{-i\langle\xi, X\rangle} H((\hbar, p+\hbar \xi, p) \cdot(\hbar, p, q)) \\
& =\int_{\mathbb{R}^{n}} \mathrm{~d} X f(p, X) \int_{\mathbb{R}^{n}} e^{-i\langle\xi, X\rangle} H(\hbar, p+\hbar \xi, q) .
\end{aligned}
$$

If $f(p, X)$ is equal to $X_{k}$, the $k$-th coordinate of $X$, its quantisation at a given value of $\hbar$ is

$$
\mathcal{Q}_{\hbar}\left(X_{k}\right)=-i \hbar \frac{\partial}{\partial p_{k}} ;
$$

whereas if $f(p, X)=p_{k}$ then

$$
\mathcal{Q}_{\hbar}\left(p_{k}\right)=p_{k},
$$

the operator multiplication by $p_{k}$.

## 5. SOME GENERAL RESULTS ABOUT PoISSON MANIFOLDS

Given a Poisson manifold $P$, its cotangent bundle is naturally a Lie algebroid $A$. If this Lie algebroid is integrable to a Lie groupoid then Theorem 2.14 gives a quantisation of the Lie-Poisson manifold $A^{*}$. Since the Poisson structure of $A^{*}$ is completely determined by the one of $P$, and vice-versa, one might hope to be able to say something about a quantisation of $P$. One way of doing so might consist in looking for a surjective Poisson map $\pi$ between $A^{*}$ and $P$ and then quantised a function on $P$ by taking the quantisation of the pulled back function on $A^{*}$. In $C^{*}$-algebraic quantisation, such an idea is bound to fail because if $f$ is a function on $P$, then its pull-back $\pi^{*} f$ has little chance of being quantisable. Nevertheless, I will show with an example that this idea can be made to work when using the quantisation defined in Theorem 2.14

The aim of this Section is to derive the partial differential equation that a map $\pi: T P \longrightarrow P$ has to satisfy to be Poisson. This PDE is given in Corollary 5.2

Assume $P$ is a Poisson manifold with Poisson bivector field $\eta$. Denote by $p$ : $\mathrm{T} P \rightarrow P$ the natural projection. The Poisson structure of $P$ induces a Lie algebroid structure on the cotangent space of $P$ with anchor map $\eta: T^{*} P \longrightarrow T P^{7}$. Its dual, the tangent space of $P$, inherits a Poisson structure in the following manner. Let $\alpha$ and $\beta$ be closed 1 -forms on $P$. They naturally define smooth functions, denoted $\widetilde{\alpha}$ and $\widetilde{\beta}$, on $T P$ by, for $v$ in $T_{x} P$,

$$
\widetilde{\alpha}(v)=\alpha(v) \text { and } \widetilde{\beta}(v)=\beta(v) .
$$

Put

$$
\begin{equation*}
\{\widetilde{\alpha}, \widetilde{\beta}\}_{\mathbf{T} P}(v)=[\alpha, \beta](v) \tag{2}
\end{equation*}
$$

where [, ] is the bracket on the Lie algebroid $\mathrm{T}^{*} P$. Let $f$ and $g$ be smooth functions on $P$. Put

$$
\begin{equation*}
\left\{p^{*} f, p^{*} g\right\}_{\mathbf{T} P}=0 . \tag{3}
\end{equation*}
$$

Finally, for $v$ in $T_{x} P$, put

$$
\begin{equation*}
\left\{\widetilde{\alpha}, p^{*} f\right\}_{\mathbf{T}_{P}}(v)=\eta(\alpha(x)) \cdot f . \tag{4}
\end{equation*}
$$

The formulae (2), (3) and (4) completely determine the Poisson structure on $\mathrm{T} P$.
Suppose there is a torsion free connection on $T P$

$$
\begin{aligned}
\nabla: \Gamma(T P) \otimes \Gamma(T P) & \longrightarrow \Gamma(T P) \\
(X, Y) & \longmapsto \nabla_{X} Y .
\end{aligned}
$$

For example, the Levi-Civita connection of a metric would do. In particular, for every $v$ in $\mathrm{T}_{x} P$, there is a splitting of $\mathrm{T}_{v}(\mathrm{~T} P)$ as a direct sum of a horizontal space $\mathbf{H}_{v}(\mathrm{~T} P)$ and a vertical space $\mathbf{V}_{v}(\mathrm{~T} P)$. Denote by $\mathbf{H}$ and $\mathbf{V}$ the projection on respectively $\mathbf{H}(\mathrm{T} P)$ and $\mathbf{V}(\mathrm{T} P)$. Both these spaces are isomorphic to $\mathrm{T}_{x} P$ and

[^5]the projection $\mathbf{H}$ is equal to $p_{*}$ while $\mathbf{V}(\mathrm{T} P)$ is the kernel of $p_{*}$. The isomorphism between $\mathbf{V}_{v}(\mathbf{T} P)$ and $\mathrm{T}_{x} P$ is given by
\[

$$
\begin{aligned}
\mathrm{T}_{x} P & \longrightarrow \mathbf{V}_{v}(\mathrm{~T} P) \\
u & \left.\longmapsto \frac{\mathrm{~d}}{\mathrm{~d} t}\right|_{t=0} v+t u .
\end{aligned}
$$
\]

Let $\mu$ be in $\Lambda^{2} \mathrm{~T}_{x} P$. The connection also defines a splitting of the tangent space of $\bigwedge^{2} \mathrm{~T} P$ at $\mu$ into the direct sum of a horizontal space isomorphic to $\mathrm{T}_{x} P$ and a vertical space isomorphic to $\bigwedge^{2} \mathrm{~T}_{x} P$. Because $p \circ \eta$ is the identity of $P$, the horizontal component of $T \eta: \mathrm{T} P \rightarrow \mathrm{~T} \bigwedge^{2} \mathrm{~T} P$ is the identity. Denote its vertical component by

$$
\mathfrak{D} \eta: \mathrm{T} P \longrightarrow \bigwedge^{2} \mathrm{~T} P
$$

If $f$ is a smooth function on $\mathrm{T} P$, its differential at $v$ in $T_{x} P$ has a horizontal and a vertical component. Denote by $\partial_{2} f: T_{x} P \longrightarrow \mathbb{C}$ its horizontal component and by $\partial_{1} f: T_{x} P \longrightarrow \mathbb{C}$ its vertical one. In the same fashion, for $\pi: \mathrm{T} P \rightarrow P$, denote by $\partial_{1} \pi$ and $\partial_{2} p$ respectively the vertical and horizontal components of $\pi_{*}: \mathrm{T}(\mathrm{T} P) \rightarrow$ $\mathrm{T} P$.

Lemma 5.1. Let $f$ and $g$ be smooth functions on $T P$. For $x$ in $P$ and $v$ in $T_{x} P$
$\{f, g\}_{T P}(v)=\left\langle\mathfrak{D} \eta(v), \partial_{1} f(v) \wedge \partial_{1} g(v)\right\rangle+\left\langle\eta(x), \partial_{1} f(v) \wedge \partial_{2} g(v)-\partial_{1} g(v) \wedge \partial_{2} f(v)\right\rangle$.
Proof. Use the right hand side of the above equation to define a bracket

$$
\{,\}: \mathbb{C}^{\infty}(P) \times \mathbb{C}^{\infty}(P) \longrightarrow \mathbb{C}^{\infty}(P) .
$$

This bracket satisfies the Leibniz identity because the operators $\partial_{1}$ and $\partial_{2}$ do. It is also clearly anti-symmetric. To prove that it is equal to $\{,\}_{\mathrm{T} P}$, it suffices to prove that it satisfies Equations (2), (3) and (4).

Equation (3) is satisfied because $\partial_{1}$ vanishes on pull-backs to $T P$ of functions on $P$.

If $g$ is a function on $P$, then $\partial_{2} p^{*} g=\mathrm{d} g$. If $f$ is equal to $\widetilde{\alpha}$ for some 1 -form $\alpha$ on $P$, then $\partial_{1} \widetilde{\alpha}(v)=\alpha_{x}$ for all $v$ in $T_{x} P$. Hence the bracket $\{$,$\} satisfies$ Equation (4).

The connection on $\mathrm{T} P \rightarrow P$ also defines a connection on its dual bundle and on all bundles one can construct from $\mathrm{T} P$ and $\mathrm{T}^{*} P$ through direct sums, tensor products...

Let $\alpha$ and $\beta$ be closed 1-forms on $P$. Let $v$ be in $\mathrm{T}_{x} P$ and let $\sigma$ be a path in $P$ such that $\sigma(0)=x$ and $\dot{\sigma}(0)=v$; for example, take $\sigma(t)=\operatorname{Exp}(t v)$, where Exp is the exponential map of the connection. Equation (2) gives

$$
\begin{equation*}
\{\widetilde{\alpha}, \widetilde{\beta}\}_{\mathbf{T} P}(v)=\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \eta_{\sigma(t)}(\alpha \wedge \beta) . \tag{5}
\end{equation*}
$$

Let $E \rightarrow P$ be the bundle $\mathrm{T}^{*} P \oplus \mathrm{~T}^{*} P \oplus \bigwedge^{2} \mathrm{~T} P \rightarrow P$. There is a natural map

$$
m: E \longrightarrow \mathbb{R}
$$

$$
(\alpha, \beta, \mu) \longmapsto \mu(\alpha \wedge \beta)
$$

I will compute the differential of this map. Let $w$ be in $\mathrm{T}_{\left(\alpha_{x}, \beta_{x}, \mu_{x}\right)} E$ with horizontal component $\mathbf{H}(w)=v$. Its vertical component $\mathbf{V}(w)=\left(\epsilon_{1}, \epsilon_{2}, \theta\right)$ is in $\mathrm{T}_{x}^{*} P \oplus$ $\mathrm{T}_{x}^{*} P \oplus \wedge^{2} \mathrm{~T}_{x} P$. Let $\phi_{\sigma(t)}: E_{x} \rightarrow E_{\sigma(t)}$ be the parallel transport along the path $\sigma$. Define a path in $E$ by

$$
\gamma(t)=\phi_{\sigma(t)}\left(\alpha_{x}+t \epsilon_{1}, \beta_{x}+t \epsilon_{2}, \mu_{x}+t \theta\right) .
$$

The path $\gamma$ satisfies

$$
\gamma(0)=\left(\alpha_{x}, \beta_{x}, \mu_{x}\right) \text { and } \dot{\gamma}(0)=w .
$$

Notice that, because the connection on $E$ is defined using a single connection on $\mathrm{T} P$, we have

$$
m \circ \phi_{\sigma(t)}=m .
$$

This means

$$
\begin{aligned}
m_{*}(w) & =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} m \circ \phi_{\sigma(t)}\left(\alpha_{x}+t \epsilon_{1}, \beta_{x}+t \epsilon_{2}, \mu_{x}+t \theta\right) \\
& =m\left(\epsilon_{1}, \beta_{x}, \mu_{x}\right)+m\left(\alpha_{x}, \epsilon_{2}, \mu_{x}\right)+m\left(\alpha_{x}, \beta_{x}, \theta\right) .
\end{aligned}
$$

This last computation together with Equation (5) gives

$$
\begin{equation*}
\{\widetilde{\alpha}, \widetilde{\beta}\}_{T P}(v)=m\left(\partial_{1, x} \alpha(v), \beta_{x}, \eta_{x}\right)+m\left(\alpha_{x}, \partial_{1, x} \beta(v), \eta_{x}\right)+m\left(\alpha_{x}, \beta_{x}, \mathfrak{D} \eta(v)\right) \tag{6}
\end{equation*}
$$

Firstly, in this equality, one can replace $\alpha_{x}$ and $\beta_{x}$ by respectively $\partial_{1, v} \widetilde{\alpha}$ and $\partial_{1, v} \widetilde{\beta}$. Secondly, consider a vector field $X$ on $P$. Let

$$
\begin{aligned}
\iota: P & \longrightarrow T^{*} P \oplus T P \\
x & \longmapsto\left(\alpha_{x}, X_{x}\right),
\end{aligned}
$$

and

$$
\begin{aligned}
k: T^{*} P \oplus T P & \longrightarrow \mathbb{C} \\
(\delta, Z) & \longmapsto \delta(Z) .
\end{aligned}
$$

I choose $X$ such that $\partial_{1} X=0$. Then, differentiation of the equality

$$
k \circ \iota=\widetilde{\alpha} \circ X
$$

leads to

$$
\begin{equation*}
\partial_{2, X_{x}} \widetilde{\alpha}(v)=\partial_{1, x} \alpha(v)\left(X_{x}\right) . \tag{7}
\end{equation*}
$$

Since $\alpha$ and $\beta$ are closed 1-forms, and since $\nabla$ is torsion free

$$
\begin{equation*}
\partial_{1, x} \alpha(v)\left(X_{x}\right)=\partial_{1, x} \alpha\left(X_{x}\right)(v) . \tag{8}
\end{equation*}
$$

Equations (6), (7) and (8) put together prove that the bracket $\{$,$\} satisfies Equa-$ tion (2).

I deduce the following corollaries.

Corollary 5.2. Let $\pi$ be a map $T P \rightarrow P$. Let $f$ and $g$ be functions on $P$. Their pull-backs by $\pi$ satisfy
$\left\{\pi^{*} f, \pi^{*} g\right\}_{T_{P}}(v)=\left\langle\left(\partial_{1} \pi(v) \otimes \partial_{1} \pi(v)\right)\left(\mathfrak{D}_{x} \eta(v)\right)+2\left(\partial_{1} \pi(v) \otimes \partial_{2} \pi(v)\right)(\eta(x)), d_{\pi(v)} f \wedge d_{\pi(v)} g\right\rangle$.
In particular, $\pi$ is Poisson if and only if

$$
\frac{1}{2}\left(\partial_{1} \pi(v) \odot \partial_{1} \pi(v)\right)\left(\mathfrak{D}_{x} \eta(v)\right)+\left(\partial_{1} \pi(v) \odot \partial_{2} \pi(v)\right)(\eta(x))=\eta(\pi(v)),
$$

for all $v$ in $T_{x} P$, where $\odot$ means the symmetric product.
Corollary 5.3. Assume the Poisson bivector field is parallel relative to the connection $\nabla$. Then, a map $\pi: T P \rightarrow P$ is Poisson if and only if

$$
\left(\partial_{1} \pi(v) \odot \partial_{2} \pi(v)\right)(\eta(x))=\eta(\pi(v))
$$

for all $v$ in $T_{x} P$.
Proof. Indeed, since the Poisson bivector field is parallel

$$
\mathfrak{D} \eta(v)=0 .
$$

Assume the connection on $P$ is the Levi-Civita connection of a metric on $P$ and assume $P$ is complete. Its exponential map is denoted

$$
\text { Exp : TP } \rightarrow P .
$$

When restricted to a fibre $\mathrm{T}_{x} P$, I will denote it $\operatorname{Exp}_{x}$. I quote here the following Lemma for future reference.

Lemma 5.4. Let v be a tangent vector to $P$ at a point $x$. Let $w$ be a tangent vector to TP at $v$. Its horizontal and vertical components are respectively $\mathbf{H} w$ and $\mathbf{V} w$. Consider the geodesic $\sigma(t)=\operatorname{Exp}(t \mathbf{H} w)$ and the 1-parameter family of geodesics

$$
\gamma_{s}(t)=\operatorname{Exp}_{\sigma(t)}\left(s \phi_{\sigma(t)}(v+t \mathbf{V} w)\right) .
$$

The differential of Exp at $v$ is given by

$$
T_{v} \operatorname{Exp}(w)=\left.\frac{d}{d t}\right|_{t=0} \gamma_{1}(t) .
$$

It follows that $T_{v} \operatorname{Exp}(w)$ is also the value at $t=1$ of the Jacobi field J along the geodesic $\sigma$ with initial value $J(0)=\mathbf{H} w$ and $J^{\prime}(0)=\mathbf{V} w$. In particular, Exp :TP $\rightarrow P$ is a submersion.

Proof. This is a simple exercise in Riemannian geometry.

## 6. The torus with constant Poisson structure

Let $P$ be the $n$-dimensional torus $\mathbb{R}^{n} / 2 \pi \mathbb{Z}^{n}$ with its metric inherited from the euclidean metric $\langle$,$\rangle on \mathbb{R}^{n}$. Consider a constant Poisson structure on $P$ given by a skew-symmetric $n \times n$ matrix $\eta$. Identify $\mathrm{T} P$ with $\mathbb{R}^{n} \times P$ in the obvious way. From Corollary 5.2. I deduce that the map

$$
\begin{aligned}
\pi: \mathbb{R}^{n} \times P & \longrightarrow P \\
(u, p) & \longmapsto \operatorname{Exp}_{p}\left(\frac{1}{2} u\right)=p+\frac{1}{2} u
\end{aligned}
$$

is Poisson. It is also a surjective submersion. Hence, I can hope that a quantisation of T $P$ will lead to a quantisation of $P$.

The dual $A=T^{*} P$ of $\mathrm{T} P$, identified with T $P$ using the euclidean metric, is a Lie algebroid. It can be integrated to a source simply-connected Lie groupoid. The space of morphisms of this groupoid is $\mathbb{R}^{n} \times P$. Notice that $\mathbb{R}^{n}$ is the direct orthogonal sum of $\operatorname{Ker}(\eta)$ and $\operatorname{Im}(\eta)$. Let $\mathrm{pr}_{1}$ be the orthogonal projection on $\operatorname{Ker}(\eta)$ and $\mathrm{pr}_{2}$ the orthogonal projection on $\operatorname{Im}(\eta)$. The source map of the groupoid is

$$
\begin{aligned}
s: \mathbb{R}^{n} \times P & \longrightarrow P \\
(u, p) & \longmapsto p
\end{aligned}
$$

whereas the target map is

$$
\begin{aligned}
t: \mathbb{R}^{n} \times P & \longrightarrow P \\
(u, p) & \longmapsto p+\mathrm{pr}_{2}(u)
\end{aligned}
$$

Given two elements $(u, p)$ and $(v, q)$ in the groupoid, their multiplication $(u, p)$. $(v, q)$ is well-defined if the target of $(u, p)$ is equal to the source of $(v, q)$, that is if $q=p+\mathrm{pr}_{2}(u)$; in this case

$$
(u, p) \cdot(v, q)=(u+v, p)
$$

Assume that $\eta$ is invertible, that is $P$ is symplectic. In this case, $\eta$ defines an isomorphism of Lie algebroids between $T^{*} P$ and $T P$

$$
\begin{aligned}
T^{*} P & \longrightarrow T P \\
(\xi, p) & \longmapsto(\eta(\xi), p)
\end{aligned}
$$

Choose the natural connection on the trivial vector bundle $\mathrm{T} P \simeq \mathbb{R}^{n} \times P \longrightarrow P$. The exponential map Exp for $T P$ is

$$
\begin{aligned}
\mathrm{T} P & \longrightarrow \mathbb{R}^{n} \times P \\
(X, p) & \longmapsto(X, p) .
\end{aligned}
$$

Whereas the $\operatorname{Exp}$ map for $T^{*} P$, that is $\alpha$, is

$$
\begin{aligned}
\mathrm{T}^{*} P & \longrightarrow P \\
(\xi, p) & \longmapsto(\eta(\xi), p)
\end{aligned}
$$

The tangent groupoid is given by $\widetilde{G}=\mathbb{R} \times \mathbb{R}^{n} \times P$ with

$$
s(\hbar, u, q)=(\hbar, q), \text { and } t(\hbar, u, q)=(\hbar, q+\hbar u)
$$

The product is given by

$$
(\hbar, v, q+\hbar u) \cdot(\hbar, u, q)=(\hbar, v+u, q)
$$

With this representation of $\widetilde{G}$, the exponential of the tangent groupoid is

$$
\begin{aligned}
\mathbb{R} \times T P & \longrightarrow \mathbb{R} \times \mathbb{R}^{n} \times P \\
(\hbar, X, q) & \longmapsto(\hbar, X, q)
\end{aligned}
$$

therefore $\widetilde{\alpha}$ is

$$
\begin{aligned}
\mathbb{R} \times T^{*} P & \longrightarrow \mathbb{R} \times \mathbb{R}^{n} \times P \\
(\hbar, \xi, q) & \longmapsto(\hbar, \eta(\xi), q)
\end{aligned}
$$

Let $f$ be a function on $A^{*}=T P$. Assume it is acceptable for quantisation. Let $H$ be a compactly supported smooth function on $\widetilde{G}$. For $(\hbar, q)$ in $\mathbb{R} \times P$,

$$
\begin{equation*}
\mathcal{Q}(f)_{\hbar, q}(H)=\int_{\mathbb{R}^{n}} \mathrm{~d}(X) f(X, q) \int_{\mathbb{R}^{n}} \mathrm{~d}(\xi) e^{-i\langle\xi, X\rangle} H(\hbar, \eta(\xi), q) \tag{9}
\end{equation*}
$$

Let $r$ be a vector in $\mathbb{Z}^{n}$ and define a function

$$
\begin{aligned}
g_{r}: P & \longrightarrow \mathbb{C} \\
q & \longmapsto e^{i\langle r, q\rangle} .
\end{aligned}
$$

The number $e^{i\langle r, q\rangle}$ is well defined because $\langle r, q\rangle$ is well defined modulo $2 \pi$. Set $f_{r}=\pi^{*} g_{r}$, that is

$$
\begin{aligned}
f_{r}: T P & \longrightarrow \mathbb{C} \\
(X, q) & \longmapsto e^{i\langle r, q\rangle} e^{\frac{1}{2} i\langle r, X\rangle} .
\end{aligned}
$$

Proposition 6.1. Let $r$ and $r^{\prime}$ be vectors in $\mathbb{Z}^{n}$, then

$$
\mathcal{Q}\left(f_{r}\right) \star \mathcal{Q}\left(f_{r^{\prime}}\right)=\mathcal{Q}\left(e^{i \frac{\hbar}{2}\langle r, \eta(s)\rangle} f_{r+r^{\prime}}\right)
$$

Proof. With $f_{r}$ instead of $f$, 10) becomes

$$
\begin{aligned}
\mathcal{Q}\left(f_{r}\right)_{\hbar, q}(H) & =\int_{\mathbb{R}^{n}} \mathrm{~d}(X) e^{i\langle r, q\rangle} e^{\frac{1}{2} i\langle r, X\rangle} \int_{\mathbb{R}^{n}} \mathrm{~d}(\xi) e^{-i\langle\xi, X\rangle} H(\hbar, \eta(\xi), q) \\
& =e^{i\langle r, q\rangle} H\left(\hbar, \eta\left(\frac{1}{2} r\right), q\right)
\end{aligned}
$$

Hence, the product $\mathcal{Q}\left(f_{r}\right) \star \mathcal{Q}\left(f_{r^{\prime}}\right)$ is

$$
\begin{aligned}
\mathcal{Q}\left(f_{r}\right) \star \mathcal{Q}\left(f_{r^{\prime}}\right)_{\hbar, q}(H)= & \int_{\mathbb{R}^{n}} \mathrm{~d}(X) f_{r^{\prime}}(X, q) \int_{\mathbb{R}^{n}} \mathrm{~d}(\xi) e^{i\langle\xi, X\rangle} \int_{\mathbb{R}^{n}} \mathrm{~d}(Y) f_{r}(Y, q+\hbar \eta(\xi)) \\
= & \int_{\mathbb{R}^{n}} \mathrm{~d}(\zeta) e^{-i\langle\zeta, Y\rangle} H(\hbar, \eta(\zeta+\xi), q) \\
& \mathrm{d}(X) e^{i\left\langle r^{\prime}, q\right\rangle} e^{\frac{1}{2} i\left\langle r^{\prime}, X\right\rangle} \int_{\mathbb{R}^{n}} \mathrm{~d}(\xi) e^{i\langle\zeta, X\rangle} \int_{\mathbb{R}^{n}} \mathrm{~d}(Y) e^{i\langle r, q+\hbar \eta(\xi)\rangle} e^{\frac{1}{2} i\langle r, Y\rangle} \\
& \int_{\mathbb{R}^{n}} \mathrm{~d}(\zeta) e^{-i\langle\zeta, Y\rangle} H(\hbar, \eta(\zeta+\xi), q) \\
= & e^{i\left\langle r^{\prime}+r, q\right\rangle} \int_{\mathbb{R}^{n}} \mathrm{~d}(Y) e^{\left.i\left\langle r, \hbar \eta\left(\frac{1}{2} r^{\prime}\right)\right\rangle\right\rangle} e^{i\langle r, Y\rangle} \int_{\mathbb{R}^{n}} \mathrm{~d}(\zeta) e^{-i\langle\zeta, Y\rangle} H\left(\hbar, \eta\left(\zeta+\frac{1}{2} r^{\prime}\right), q\right) \\
= & e^{i\left\langle r^{\prime}+r, q\right\rangle} e^{i \frac{\hbar}{2}\left\langle r, \eta\left(r^{\prime}\right)\right\rangle} H\left(\hbar, \frac{1}{2} \eta\left(r+r^{\prime}\right), q\right) \\
= & e^{i \frac{\hbar}{2}\left\langle r, \eta\left(r^{\prime}\right)\right\rangle} \mathcal{Q}\left(f_{\left.r+r^{\prime}\right)}\right), q(H) .
\end{aligned}
$$

Let $\mathcal{P}$ be the algebra of functions on $P$ generated by $\left\{g_{r}, r \in \mathbb{Z}^{n}\right\}$. It is a dense sub-algebra of the $C^{*}$-algebra of continuous functions on $P$. Proposition 6.1 shows that the product on this algebra can be deformed in

$$
g_{r} \star_{\hbar} g_{r^{\prime}}=e^{i \frac{\hbar}{2}\langle r, \eta(s)\rangle} g_{r+r^{\prime}},
$$

for each $\hbar$. With this new product, $\mathcal{P}$ becomes a $*$-algebra which can be completed into a $C^{*}$-algebra $\mathcal{P}_{\hbar}$. The natural family of injections $\mathcal{P} \longrightarrow \mathcal{P}_{\hbar}$ gives the usual quantisation of the torus with constant Poisson structure as defined in Tang and Weinstein [8]. It is a strict deformation quantisation in the sense of Rieffel.

## Appendix: The 2 -sphere in $\mathbb{R}^{3}$

Let $P=S^{2}$ be the 2 -sphere $\left\{(x, y, z) \in \mathbb{R}^{3} / x^{2}+y^{2}+z^{2}=1\right\}$. In this Section, I show how to construct a Poisson map between $T S^{2}$ and $S^{2}$.

Consider the metric on $S^{2}$ given by the restriction of the euclidean metric $\mathrm{d} x^{2}+$ $\mathrm{d} y^{2}+\mathrm{d} z^{2}$ on $\mathbb{R}^{3}$. For $p=(x, y, z)$ in $\mathbb{R}^{3}$, define the endomorphism

$$
J_{p}(u)=u \wedge p, \text { for all } u \in \mathbb{R}^{3} .
$$

The restriction of $J$ to each tangent space of $S^{2}$ defines a complex structure on the sphere. Also

$$
\omega=g(J \cdot, \cdot)
$$

is a symplectic form on $S^{2}$. It is the restriction to $S^{2}$ of the 2 -form $z \mathrm{~d} y \wedge \mathrm{~d} z-$ $y \mathrm{~d} x \wedge \mathrm{~d} z+z \mathrm{~d} x \wedge \mathrm{~d} y$ defined on $\mathbb{R}^{3}$. The geodesics are great circles on the sphere so that the exponential is given by

$$
\operatorname{Exp}_{p}(u)=p \cos (\|u\|)+\frac{\sin (\|u\|)}{\|u\|} u .
$$

The differential $\partial_{1} E x p$ is a map from $T_{p} S^{2}$ to $T_{E x p(u)} S^{2}$ given by

$$
\begin{aligned}
\partial_{1} \operatorname{Exp}_{(u, p)}(h) & =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \operatorname{Exp}_{p}(u+t h) \\
& =\frac{\sin (\|u\|)}{\|u\|}(-\langle u, h\rangle p+h)+\frac{\|u\| \cos (\|u\|)-\sin (\|u\|)}{\|u\|^{2}}\left\langle\frac{u}{\|u\|}, h\right\rangle u
\end{aligned}
$$

In addition,

$$
\partial_{2} \operatorname{Exp}_{(u, p)}(\epsilon)=\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \operatorname{Exp}_{\sigma(t)} \phi_{\sigma(t)}(u),
$$

where $\sigma(t)=\operatorname{Exp}_{p}(t \epsilon)$ and $\phi_{\sigma(t)}$ is the parallel transport along $\sigma(t)$. Without loss of generality, I can assume that $\epsilon$ is a unit vector. In this situation,

$$
\sigma(t)=\cos (t) p+\sin (t) \epsilon .
$$

Also, $\{p, \epsilon, p \wedge \epsilon\}$ forms an orthonormal basis of $\mathbb{R}$ and $\phi_{\sigma(t)}$ is a morphism in $\mathrm{SO}(3)$. It is given by

$$
\begin{aligned}
\phi_{\sigma(t)}(p) & =\sigma(t) \\
\phi_{\sigma(t)}(\epsilon) & =\frac{\mathrm{d}}{\mathrm{~d} t} \sigma(t) \\
& =-\sin (t) p+\cos (t) \epsilon \\
\phi_{\sigma(t)}(p \wedge \epsilon) & =p \wedge \epsilon .
\end{aligned}
$$

It follows that

$$
\phi_{\sigma(t)}(u)=\langle u, \epsilon\rangle(-\sin (t) p+\cos (t) \epsilon)+\langle u, p \wedge \epsilon\rangle p \wedge \epsilon .
$$

The parallel transport preserves the norm, hence $u$ and $\phi_{\sigma(t)}(u)$ have the same norm and

$$
\begin{aligned}
\operatorname{Exp}_{\sigma(t)} \phi_{\sigma(t)}(u)= & \cos (\|u\|) \sigma(t)+\frac{\sin (\|u\|)}{\|u\|} \phi_{\sigma(t)}(u) \\
= & \cos (\|u\|)(\cos (t) p+\sin (t) \epsilon)+ \\
& +\frac{\sin (\|u\|)}{\|u\|}(\langle u, \epsilon\rangle(-\sin (t) p+\cos (t) \epsilon)+\langle u, p \wedge \epsilon\rangle p \wedge \epsilon) .
\end{aligned}
$$

I can now compute

$$
\begin{aligned}
\partial_{2} \operatorname{Exp}_{(u, p)}(\epsilon) & =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \operatorname{Exp}_{\sigma(t)} \phi_{\sigma(t)}(u) \\
& =\cos (\|u\|) \epsilon-\frac{\sin (\|u\|)}{\|u\|}\langle u, \epsilon\rangle p .
\end{aligned}
$$

This last formula is of course still valid when $\epsilon$ is not a unit vector.
Consider a map

$$
\begin{aligned}
& \pi: \mathrm{T} S^{2} \longrightarrow S^{2} \\
& u \longmapsto \operatorname{Exp}(\lambda\|u\|) \\
&=\cos (\lambda\|u\|) p+f(\lambda\|u\|) \lambda u
\end{aligned}
$$

where $\lambda$ is a function of $\|u\|^{2}$ defined for $\|u\|^{2}$ in some neighbourhood of 0 in $\mathbb{R}$.
For $\pi$ to be a Poisson map, I need $\lambda(0)=\frac{1}{2}$.
I will now compute the differential of such a map. Firstly,

$$
\begin{aligned}
\partial_{1,(u, p)} \pi(h)= & \left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \operatorname{Exp}_{p}\left(\lambda\left(\|u+t h\|^{2}\right)(u+t h)\right) \\
= & \left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \cos \left(\lambda\left(\|u+t h\|^{2}\right)\|u+t h\|\right) p+ \\
& +f\left(\lambda\left(\|u+t h\|^{2}\right)\|u+t h\|\right) \lambda\left(\|u+t h\|^{2}\right)(u+t h) \\
= & -\sin (\lambda\|u\|)\left(2 \lambda^{\prime}\langle u, h\rangle\|u\|+\lambda\left\langle h, \frac{u}{\|u\|}\right\rangle\right) p+f(\lambda\|u\|) \lambda h+ \\
& +2 f(\lambda\|u\|)\langle u, h\rangle \lambda^{\prime} u+\left(2\langle u, h\rangle \lambda^{\prime}\|u\|+\lambda\left\langle h, \frac{u}{\|u\|}\right\rangle\right) f^{\prime}(\lambda\|u\|) \lambda u \\
= & -\lambda f(\lambda\|u\|)\left(2 \lambda^{\prime}\langle u, h\rangle\|u\|^{2}+\lambda\langle u, h\rangle\right) \lambda f(\lambda\|u\|) p+f(\lambda\|u\|) \lambda h+ \\
& +\left(2 f(\lambda\|u\|) \lambda^{\prime}+2 \lambda^{\prime} \lambda f^{\prime}(\lambda\|u\|)+\frac{\lambda^{2}}{\|u\|} f^{\prime}(\lambda\|u\|)\right)\langle u, h\rangle u \\
= & -\lambda f(\lambda\|u\|)\left(2 \lambda^{\prime}\langle u, h\rangle\|u\|^{2}+\lambda\langle u, h\rangle\right) \lambda f(\lambda\|u\|) p+f(\lambda\|u\|) \lambda h+ \\
& +\left(2 \lambda^{\prime} \cos (\lambda\|u\|)+\frac{\lambda^{2}}{\|u\|} f^{\prime}(\lambda\|u\|)\right)\langle u, h\rangle u,
\end{aligned}
$$

where for the last equality I have used the relation

$$
t f^{\prime}(t)=\cos (t)-f(t)
$$

Secondly, assuming without loss of generality that $\|\epsilon\|=1$,

$$
\begin{aligned}
\partial_{2,(u, p)} \pi(\epsilon) & =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \pi\left(\phi_{\sigma(t)}(u)\right) \\
& =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \operatorname{Exp}_{\sigma(t)}\left(\lambda\left(\left\|\phi_{\sigma(t)}(u)\right\|^{2}\right) \phi_{\sigma(t)}(u)\right) \\
& =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \operatorname{Exp}_{\sigma(t)}\left(\lambda\left(\|(u)\|^{2}\right) \phi_{\sigma(t)}(u)\right) \\
& =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \cos \left(\lambda\left\|\phi_{\sigma(t)}(u)\right\|\right) \sigma(t)+f\left(\lambda\left\|\phi_{\sigma(t)}(u)\right\|\right) \lambda \phi_{\sigma(t)}(u) \\
& =\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \cos (\lambda\|u\|) \sigma(t)+f(\lambda\|u\|) \lambda \phi_{\sigma(t)}(\langle u, \epsilon\rangle \epsilon+\langle u, p \times \epsilon\rangle p \times \epsilon) \\
& =\cos (\lambda\|u\|) \sigma^{\prime}(0)+\left.\lambda f(\lambda\|u\|) \frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0}\left(\langle u, \epsilon\rangle \sigma^{\prime}(t)+\langle u, p \times \epsilon\rangle p \times \epsilon\right) \\
& =\cos (\lambda\|u\|) \epsilon+\langle u, \epsilon\rangle \lambda f(\lambda\|u\|) \sigma^{\prime \prime}(0) \\
& =\cos (\lambda\|u\|) \epsilon-\langle u, \epsilon\rangle \lambda f(\lambda\|u\|) p .
\end{aligned}
$$

I wish to compute $\|u\|^{2} \partial_{1} \pi \odot \partial_{2} \pi(\eta(p))$. I know

$$
\eta(p)=\frac{1}{\|u\|^{2}} u \wedge(p \times u), \text { whenever } u \neq 0
$$

So, I need to compute

$$
\begin{aligned}
\partial_{1,(u, p)} \pi(u)= & -\lambda f(\lambda\|u\|)\left(2 \lambda^{\prime}\|u\|^{2}+\lambda\right)\|u\|^{2} p+\lambda f(\lambda\|u\|) u+ \\
& +\left(2 \lambda^{\prime} \cos (\lambda\|u\|)+\frac{\lambda}{\|u\|^{2}} f^{\prime}(\lambda\|u\|)\right)\|u\|^{2} u \\
= & -\lambda\|u\|^{2} f(\lambda\|u\|)\left(2 \lambda^{\prime}\|u\|^{2}+\lambda\right) p+(\lambda f(\lambda\|u\|)+ \\
& \left.+2 \lambda^{\prime} \cos (\lambda\|u\|)\|u\|^{2}+\lambda^{2}\|u\| f^{\prime}(\lambda\|u\|)\right) u \\
= & \left(2 \lambda^{\prime}\|u\|^{2}+\lambda\right)\left(\cos (\lambda\|u\|) u-\|u\|^{2} \lambda f(\lambda\|u\|) p\right)
\end{aligned}
$$

and

$$
\partial_{1,(u, p)} \pi(p \times u)=\lambda f(\lambda\|u\|) p \times u
$$

and

$$
\partial_{2,(u, p)} \pi(u)=\cos (\lambda\|u\|) u-\|u\|^{2} \lambda f(\lambda\|u\|) p
$$

and finally

$$
\partial_{2,(u, p)} \pi(p \times u)=\cos (\lambda\|u\|) p \times u
$$

Notice that

$$
\partial_{1,(u, p)} \pi(u)=\left(2 \lambda^{\prime}\|u\|^{2}+\lambda\right) \partial_{2,(u, p)} \pi(u)
$$

I can now compute

$$
\begin{aligned}
\|u\|^{2} \partial_{1} \pi \odot \partial_{2} \pi(\eta(p))= & \partial_{1} \pi(u) \wedge \partial_{2} \pi(p \times u)+\partial_{2} \pi(u) \wedge \partial_{2} \pi(p \times u) \\
= & \partial_{2}(u) \wedge\left(\left(2 \lambda^{\prime}\|u\|^{2}+\lambda\right) \partial_{2,(u, p)} \pi(u)+\partial_{1} \pi(p \times u)\right) \\
= & \left(\cos (\lambda\|u\|) u-\|u\|^{2} \lambda f(\lambda\|u\|) p\right) \wedge\left(\left(2 \lambda^{\prime}\|u\|^{2}+\lambda\right) \cos (\lambda\|u\|)+\right. \\
& +\lambda f(\lambda\|u\|)) p \times u \\
= & \left(\cos (\lambda\|u\|) u-\|u\|^{2} \sin (\lambda\|u\|) p\right) \wedge\left(\left(2 \lambda^{\prime}\|u\|^{2}+\lambda\right) \cos (\lambda\|u\|)+\right. \\
& +\lambda f(\lambda\|u\|)) p \times u
\end{aligned}
$$

On the other hand

$$
\begin{aligned}
\|u\|^{2} \eta(\pi(u, p)) & =\|u\|^{2} \phi_{\sigma(\lambda\|u\|)}(\eta(p)) \quad, \text { with } \sigma(t)=\operatorname{Exp}\left(t \frac{u}{\|u\|}\right)=\cos (t) p+\sin (t) \frac{u}{\|u\|} \\
& =\phi_{\sigma(\lambda\|u\|)}(u) \wedge \phi_{\sigma(\lambda\|u\|)}(p \times u) \\
& =\sigma^{\prime}(\lambda\|u\|) \wedge(p \times u) \\
& =(\cos (\lambda\|u\|) u-\|u\| \sin (\lambda\|u\|) p) \wedge(p \times u)
\end{aligned}
$$

It follows that $\pi$ is a Poisson map if and only if $\lambda$ satisfies the following differential equation

$$
\left(2 \lambda^{\prime}\left(t^{2}\right) t^{2}+\lambda\left(t^{2}\right)\right) \cos \left(\lambda\left(t^{2}\right) t\right)+\lambda\left(t^{2}\right) f\left(\lambda\left(t^{2}\right) t\right)=1
$$

Put $\mu(t)=\lambda\left(t^{2}\right)$ so that $\mu^{\prime}(t)=2 t \lambda^{\prime}\left(t^{2}\right)$. The function $\mu$ satisfies the differential equation

$$
\left(\mu^{\prime}(t) t+\mu(t)\right) \cos (\mu(t) t)+\frac{\sin (\mu(t) t)}{t}=1 .
$$

Put $\alpha(t)=\sin (\mu(t) t)$ so that $\alpha^{\prime}(t)=\left(\mu^{\prime}(t) t+\mu(t)\right) \cos (\mu(t) t)$. The function $\alpha$ satisfies the differential equation

$$
t \alpha^{\prime}(t)+\alpha(t)=t
$$

A general solution of this equation is

$$
\alpha(t)=\frac{a}{t}+\frac{t}{2} \text {, with } a \in \mathbb{R} \text {. }
$$

Hence

$$
\mu(t)=\frac{1}{t} \arcsin \left(\frac{a}{t}+\frac{t}{2}\right) .
$$

Since I want $\mu(0)=\lambda(0)=\frac{1}{2}$, I need $a=0$ and

$$
\mu(t)=\frac{1}{t} \arcsin \left(\frac{t}{2}\right) .
$$

I deduce
Proposition 6.2. The map

$$
\begin{aligned}
\pi: T S^{2} & \longrightarrow S^{2} \\
(u, p) & \longmapsto \operatorname{Exp}_{p}\left(\frac{1}{\|u\|} \arcsin \left(\frac{\|u\|}{2}\right) u\right)
\end{aligned}
$$

## is Poisson

The way it is written in the previous Proposition, the map $\pi$ is only defined on a neighbourhood of $S^{2}$ in $T S^{2}$. Nevertheless, it can easily be extended to a continuous function on the whole of $T S^{2}$, so that it is Poisson wherever it is smooth. The fact that this map is not smooth at all points means that technics used in Section 6 will not carry over here. I nevertheless believe that a modification of these technics will make things work.

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[^0]:    Date: November 3, 2004.

[^1]:    ${ }^{1}$ That is if the Lie algebroid is integrable.

[^2]:    ${ }^{2}$ Any compact set is compact modulo $M$; but there might be other compacts modulo $M: M \subset A$ itself is compact modulo $M$ even if it is not compact

[^3]:    ${ }^{3}$ At $r=0$, we have $\theta(u, \hbar \xi+r \zeta) \theta(u, \hbar \xi)^{-1}=t \circ \theta(u, \xi)$, thus the differential of $\theta(u, \hbar \xi+$ $r \zeta) \theta(u, \hbar \xi)^{-1}$ at $r=0$ is of the type $\frac{\partial \theta}{\partial v}(\sigma(u, \hbar \xi), 0) \phi$ for a certain vector $\phi$.
    ${ }^{4}$ This is true only for $\hbar \neq 0$, nevertheless the final result of the computation is trivially true for $\hbar=0$.

[^4]:    ${ }^{5}$ Here I used $\sigma(u, \hbar \xi)=u+\mathrm{O}(\hbar)$ to simplify the formula.
    ${ }^{6}$ The signs here and in Ramazan [6] do not agree. This is due to different choices in the definition of the map $\alpha$. Essentially, I have $s \circ \alpha$ constant on the fibres of $A \longrightarrow Q$, whereas he has $t \circ \alpha$ constant on the same fibres.

[^5]:    ${ }^{7}$ Here, $\eta$ is understood as an anti-symmetric map between $T^{*} P$ and $T P$. I will use different sorts of interpretations of $\eta$, the precise interpretation depending on the context.

