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## DISCRETE SYMMETRIES AND GENERAL RELATIVITY, THE DARK SIDE OF GRAVITY

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Received ..

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The parity, time reversal and space/time exchange invariant actions, equations and their conjugate metric solutions are obtained in the context of a general relativistic model modified in order to take into account discrete symmetries. The equations are not covariant however the predictions of the model, in particular its Schwarzschild metric solution in vacuum, only start to differ from those of General Relativity at the Post-Post-Newtonian order. No coordinate singularity (black hole) arises in the privileged coordinate system where the energy of gravity is found to vanish. Vacuum energies cancel as gravitational sources and a flat universe accelerated expansion phase is obtained without resorting to inflation nor a cosmological constant. The Pioneer anomalous blue-shift is a straightforward outcome. The context is promising to help us elucidate several outstanding enigmas such as flat galactic rotation curves or the universe voids. A wave solution is obtained leading to the same binary pulsar decay rate as observed and predicted in GR.

*Keywords:* Negative energy; time reversal, tachyon.

### 1. Introduction

We might summarize the point of view developed in Ref. 1 as follows: any physicist would agree that symmetries are fundamental in contemporary physics. They both allow to constrain the form assumed by the actions and the properties of the basic objects such as fields from which the actions are built and the phenomenology is derived. If we miss some fundamental symmetry of mother nature or misunderstand the way it works, the physical world description we will derive might be either totally wrong or in the best case incomplete. As is well known, a discrete operator such as the time reversal one in Quantum Field Theory may be either unitary or anti-unitary. The anti-unitary choice for it is the conventional one, and the commonly accepted derived picture undoubtedly constitutes an important part of our culture. However, our starting point in Ref. 1 was to take serious the other mathematical possibility for the time reversal operator, e.g. that it be unitary and see how far we could follow this theoretical option. Then a very different conception of time reversal arises together with negative energy fields. However, at some point we found that there is no way to reach a coherent description in flat space-time. But fortunately,

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we know that the complete picture should include gravity and this, as we shall show, naturally allows us to achieve our primary goal: understand time reversal and by the way the negative energy representations of the Lorentz group. The reader is referred to Ref. 1 for our investigation of negative energies and time reversal restricted to QFT in flat space-time but an interesting analysis can also be found in Ref. 2.

## 2. Negative Energies in Quantum Field Theory

Let us gather the main information we learned from our investigation in Ref. 1 of negative energies in Relativistic QFT indicating the the correct theoretical framework for handling them should be a modified GR.

- Theoretical Motivations

In second quantification, all relativistic field equations admit genuine negative energy field solutions creating and annihilating negative energy quanta. Unitary time reversal links these fields to the positive energy ones. The unitary choice, usual for all other symmetries in physics, also allows us to avoid the well known paradoxes associated with time reversal. Positive and negative energy fields vacuum divergences we encounter after second quantization are unsurprisingly found to be exactly opposite. The negative energy fields action must be maximised. However, there is no way to reach a coherent theory involving negative energies in flat-spacetime. Indeed, if positive and negative energy scalar fields are time reversal conjugate, also must be their Hamiltonian densities and actions. This is only possible in the context of GR thanks to the metric transformation under discrete symmetries.

- Phenomenological Motivations

In a mirror negative energy world whose fields remain non coupled to our world positive energy fields, stability is insured and the behavior of matter and radiation is as usual. Hence, it is just a matter of convention to define each one as a positive or negative energy world. Otherwise, If they interact gravitationally, promising phenomenology is expected. Indeed, many outstanding enigmas indicate that repelling gravity might play an important role in physics: flat galactic rotation curves, the Pioneer effect, the flatness of the universe, acceleration and its voids, etc... But negative energy states never manifested themselves up to now, suggesting that a barrier is at work preventing the two worlds to interact except through gravity.

- A Modified GR to Circumvent the Main Issues

A trivial cancellation between vacuum divergences is not acceptable since the Casimir effect shows evidence for vacuum fluctuations. But the positive and negative energy worlds could be maximally gravitationally coupled in such a way as to produce at least exact cancellations of vacuum energies gravitational effects. Also, a generic catastrophic instability issue arises whenever quantum positive and negative energy fields are allowed to interact. If we restrict the stability issue to the modified gravity, this disastrous

scenario is avoided. Finally, allowing both positive and negative energy virtual photons to propagate the electromagnetic interaction, simply makes it disappear. The local gravitational interaction is treated very differently in our modified GR so that this unpleasant feature is also avoided.

- Outlooks

A left-handed kinetic and interaction Lagrangian can satisfactorily describe all known physics (except mass terms which anyway remain problematic in modern physics). This strongly supports the idea that the right handed chiral fields might be living in another world (where the 3-volume reversal under parity presumably would make these fields acquire a negative energy density) and may provide as shown in Ref. 1 an interesting explanation for maximal parity violation observed in the weak interaction.

If the connection between the two worlds is fully reestablished above a given energy threshold, then loop divergences naturally would get cancelled thanks to the positive and negative energy virtual propagators compensation. Such reconnection might take place through a new transformation process allowing particles to jump from one metric to the conjugate one presumably at places where the conjugate metrics meet each other.

### 3. conjugate Worlds Gravitational Coupling

From Ref. 3 we learn that if a discrete symmetry (Parity or/and Time reversal) transforms the general coordinates, this will not affect a scalar action however if the inertial coordinates  $\xi^\alpha$  are also transformed in a non-trivial way:

$$\xi^\alpha \rightarrow \tilde{\xi}^\alpha, \quad (1)$$

non-trivial in the sense that in general  $\tilde{\xi}^\alpha \neq \xi_\alpha$ , our metric terms will be affected and our action is not expected to be invariant under P or T. Having two conjugate inertial coordinate systems, we can build, following the usual procedure, two discrete symmetry reversal conjugate metric tensors:

$$g_{\mu\nu} = \eta_{\alpha\beta} \frac{\partial \xi^\alpha}{\partial x^\mu} \frac{\partial \xi^\beta}{\partial x^\nu}, \quad \tilde{g}_{\mu\nu} = \eta_{\alpha\beta} \frac{\partial \tilde{\xi}^\alpha}{\partial x^\mu} \frac{\partial \tilde{\xi}^\beta}{\partial x^\nu}. \quad (2)$$

This means that there exists a coordinate system where the discrete symmetry applies in the trivial way (for instance  $x^0 \rightarrow -x^0$  in case of time reversal) transforming the conjugate metrics into one another. We postulate that this coordinate system is such that  $\tilde{g}_{\mu\nu}$  identifies with  $g^{\mu\nu}$ , a non covariant relation making this frame a privileged one. Then, we should distinguish between those fields following the geodesics of  $\tilde{g}_{\mu\nu}$  and the others following the geodesics of  $g_{\mu\nu}$ . We first build the total action: sum of usual  $I_{RG}$  and the new conjugate one  $\tilde{I}_{RG}$ . The conjugate actions are separately general coordinate scalars and adding the two pieces is necessary to obtain a discrete symmetry reversal invariant total action. In this system, varying our action, applying the extremum action principle and making use of the relation

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$\delta g^{\rho\kappa}(x) = -g^{\rho\mu}(x)g^{\nu\kappa}(x)\delta g_{\mu\nu}(x)$  would lead us to a modified Einstein equation (with  $c = 1$ ) :

$$-8\pi G \left( \sqrt{g(x)}T_{\rho\sigma} - \sqrt{g^{-1}(x)}\tilde{T}^{\rho\sigma} \right) = \sqrt{g(x)} \left( R_{\rho\sigma} - \frac{1}{2}g_{\rho\sigma}R \right) - \sqrt{g^{-1}(x)} \left( R^{\rho\sigma} - \frac{1}{2}g^{\rho\sigma}R \right) \Big|_{g^{\rho\sigma} \rightarrow g_{\rho\sigma}, g_{\rho\sigma} \rightarrow g^{\rho\sigma}}. \quad (3)$$

This equation, only valid as it stands in our privileged working coordinate system, is not general covariant and not intended to be so. Indeed, it follows from both an extremum principle of our general covariant and discrete symmetry invariant action and a noncovariant relation between one metric and its discrete symmetry reversal conjugate. The straightforward interpretation of the Left hand side is that fields living in the reversed metric world are seen from our world as negative energy density fields but living in the conjugate metric prevents them from interacting with our world fields except through gravitation. However coupling in this traditional way metric fields with matter and radiation fields assumes implicitly that we treat and understand in the same way both directions of the bidirectionnal talk between matter and gravity. We shall show why this is not always possible and in general (there will be a noticeable exception) will give up the matter action, keeping only the purely gravitational action  $I_G + \tilde{I}_G$ . The derived equation is neither unic nor valid in such a general form as above. Indeed, symmetry requirements will a priori determine two possible simplified forms for the conjugate metrics in the privileged coordinate system.

#### 4. Isotropy and Space/Time Exchange

The isotropy condition for both conjugate metrics determines their general form:

$$d\tau^2 = B(r, t)dt^2 + A(r, t)d\sigma^2, \quad d\tilde{\tau}^2 = \frac{1}{B(r, t)}dt^2 + \frac{1}{A(r, t)}d\sigma^2. \quad (4)$$

We now want to investigate the space/time exchange symmetry in order to understand the tachyonic Lorentz group representation. Indeed, the natural symmetry linking tachyons to bradyons is the transformation reversing the metric signature, hence, following Ref. 11 12 13 transforming space-like coordinates into time-like ones (we call it x/t symmetry). This can be achieved by Wick rotating them and leads us to introduce the flipped signature conjugate metric:  $\hat{g}_{\mu\nu}(r, t) = -g_{\mu\nu}(r, t)$ . More generally, it is natural to allow our metrics to be complex, the phases corresponding to rotations in the complex plane continuously transforming into one another these conjugate metrics. The trivial B=A=C Euclidian form manifestly involves time and space coordinates in a symmetrical way, whatever the function C(r,t):

$$d\tau^2 = C(r, t) (dt^2 + d\sigma^2), \quad (5)$$

$$d\tilde{\tau}^2 = \frac{1}{C(r, t)} (dt^2 + d\sigma^2). \quad (6)$$

But in general we really need the inverse metric to restore the space/time symmetry. Indeed, the most general expression being

$$d\tau^2 = C(r, t) \left( \frac{1}{A(r, t)} dt^2 + A(r, t) d\sigma^2 \right), \quad (7)$$

$$d\tilde{\tau}^2 = \frac{1}{C(r, t)} \left( A(r, t) dt^2 + \frac{1}{A(r, t)} d\sigma^2 \right). \quad (8)$$

only the presence of  $C(r, t)$  breaks down the space / time symmetry so that in general, as we already noticed, such term is not allowed and the conjugate metrics verify  $A(r, t) = B^{-1}(r, t)$ . However, in the special case where A is a pure phase with angle:  $\theta(r, t) = \pi/2$  it is possible to keep the  $C(r, t)$  term and respect the space/time symmetry thanks to the opposite conjugate metric  $\hat{g}_{\mu\nu}(r, t) = -g_{\mu\nu}(r, t)$ . Then we have four conjugate metrics:

- $g_{\mu\nu}$
- $\tilde{g}_{\mu\nu} = g_{\mu\nu}^{-1}$
- $\hat{g}_{\mu\nu} = -g_{\mu\nu}$
- $\tilde{\hat{g}}_{\mu\nu} = \hat{g}_{\mu\nu}$

Explicitely:

$$\begin{aligned} d\tau^2 &= C(r, t) [dt^2 - d\sigma^2] \\ d\tilde{\tau}^2 &= \frac{1}{C(r, t)} [dt^2 - d\sigma^2] \\ d\hat{\tau}^2 &= C(r, t) [d\sigma^2 - dt^2] \\ d\tilde{\hat{\tau}}^2 &= \frac{1}{C(r, t)} [d\sigma^2 - dt^2] \end{aligned}$$

Thus we could show that appart from the trivial B=A form we only have two kinds of metrics: the "euclidian" ones satisfying B=1/A and the Minkowskian ones with B=-A. Of course, in the above formula, thanks to an additional global  $\theta = \pi/2$  phase, C(r,t) thus the Minkowskian metric can be real. Therefore, in case the metric elements are pure phases, the inverse metric is nothing but the complex conjugate metric and a first  $\theta = \pi/2$  phase rotation can transform the Euclidian into imaginary Minkowskian ones which under the additional  $B = A = C(r, t) = 1/i$  phase rotation reduce to our two real space/time conjugate stationary worlds.

$$\begin{aligned} d\tau^2 &= d\tilde{\tau}^2 = [dt^2 - d\sigma^2] \\ d\hat{\tau}^2 &= d\tilde{\hat{\tau}}^2 = [d\sigma^2 - dt^2] \end{aligned}$$

Then, if we allow modulus terms, two couples of time reversal conjugate non-stationary worlds will emerge as we shall show. We will thus have conjugate tachyonic/bradyonic worlds as well as positive/negative energy worlds.

## 5. The Complex Metrics and their Variations

In the following sections we will require that either the metric elements are a priori linked by  $A(r, t) = -B(r, t)$  or  $A(r, t) = B^{-1}(r, t)$  so that a single degree of freedom

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eventually remains and we shall investigate the form of its solution in each case. The metric elements being related can no longer be varied independently. Therefore, given any tensor  $T_{\mu\nu}$  and for real metric elements making use of  $\frac{\delta g_{i'i'}}{\delta g_{ii}} = \frac{g_{i'i'}}{g_{ii}}$  and  $\frac{\delta g_{tt}}{\delta g_{ii}} = \frac{g_{tt}}{g_{ii}}$  in case B=-A, a typical action variation will be proportional to a scalar trace:

$$\delta g_{rr} T^{rr} + \delta g_{tt} T^{tt} + \delta g_{\theta\theta} T^{\theta\theta} + \delta g_{\phi\phi} T^{\phi\phi} = \frac{\delta g_{rr}}{g_{rr}} T$$

while in case B=1/A,  $\frac{\delta g_{i'i'}}{\delta g_{ii}} = \frac{g_{i'i'}}{g_{ii}}$  and  $\frac{\delta g_{tt}}{\delta g_{ii}} = -\frac{g_{tt}}{g_{ii}}$  and the additional minus sign gives us a modified trace:

$$T_{B=1/A} = g_{rr} T^{rr} + g_{\theta\theta} T^{\theta\theta} + g_{\phi\phi} T^{\phi\phi} - g_{tt} T^{tt}$$

We also need the relation between the relative variations of the inverse conjugate metrics needed to obtain the gravitational equation in term of the components of a single metric:  $\frac{\delta \tilde{g}_{xx}}{g_{xx}} = -\frac{\delta g_{xx}}{g_{xx}}$

But in case the metric and their variations can be complex, we might tentatively write two kinds of metric solutions.

- The metric corresponding to the previous case where the space time intervals are real moduli. Its conjugate is its inverse. The variation must be real.
- The metric corresponding to the case where the space time intervals are pure phases. Its conjugate is its complex conjugated. The variation is imaginary. Then, the conjugate metric variations are opposite in such a way that  $\delta \tilde{g}_{xx} = -\delta g_{xx}$  and  $\delta g_{i'i'} = \delta g_{ii} = \pm \delta g_{tt}$  for phase metrics for B=-A and B=1/A respectively.

Eventually, we will find ourselves with two equations of gravity: one for the modulus and one for the phase in both cases. The moduli describe homogeneous and isotropic backgrounds for both conjugate metrics which necessarily implies their spatial flatness and r independance. The phases describe the perturbations over this background. The conjugate metric solutions should transform into one another under  $r \rightarrow -r$  and  $t \rightarrow -t$ .

## 6. B=1/A, The Phase: The Schwarzschild Solution

With B=1/A pure phase metric elements, defining g to be the modulus of the metric determinant, its squareroot simplifies and we obtain the following equation in vacuum:

$$\frac{1}{A} (g_{rr} T^{rr} + g_{\theta\theta} T^{\theta\theta} + g_{\phi\phi} T^{\phi\phi}) - \frac{1}{B} g_{tt} T^{tt} - inv = 0$$

where inv denotes the same expression with the metric elements being everywhere replaced by their inverse. This yields:

$$4 \frac{A''}{A} - 4 \left( \frac{A'}{A} \right)^2 + 8 \frac{A'}{Ar} = 0$$

We easily get the unic static phase solution

$$1/B = A = e^{\frac{2iMG}{r}} \quad (9)$$

This euclidian metric is not appropriate to get the geodesics followed by test masses. Rather should we first perform a Wick rotation of the spatial coordinates  $r \Rightarrow ir$  to obtain the real Minkowskian metric with components B, -A satisfying:

$$\Rightarrow A = e^{\frac{2MG}{r}} \approx 1 + 2\frac{MG}{r} + 2\frac{M^2G^2}{r^2}, \quad (10)$$

$$\Rightarrow B = \frac{1}{A} = e^{-\frac{2MG}{r}} \approx 1 - 2\frac{MG}{r} + 2\frac{M^2G^2}{r^2} - \frac{4}{3}\frac{M^3G^3}{r^3}. \quad (11)$$

different from the GR one though in good agreement up to Post-Newtonian order. No black hole type singularity arises in our isotropic system. The phenomenology is simple: objects living in the same metric attract each other. Objects living in different metrics repel each other, as if the object living in the conjugate metric contributes as a negative energy source from the point of view of our metric. The conjugate metrics can be transformed into one another through  $r \rightarrow -r$  or  $M \rightarrow -M$ . We could show in Ref. 1 that a left-handed Lagrangian could satisfactorily describe all known physics (at least kinetic and interaction terms) and provide an interesting explanation for maximal parity violation. This strongly supports the idea that the right handed chiral fields are living in another metric and acquire thanks to the 3-volume reversal a negative energy density. Eventually, because the solution is unpropagated in the privileged coordinate system (was Newton right ?) the usual stability issues reviewed in Refs. 5–8 are avoided. The instability is usually seen in the phenomenology of a positive energy mass interacting with a negative energy mass through an interaction propagated by positive energy virtual interaction particles. The negative energy object is being attracted by the positive energy object, the latter being repulsed by the former. They then accelerate together for ever this resulting in an obviously instable picture. But here the gravitational interaction of two masses living in different metrics exhibits no such instability since they just repel each other. Yet, from the point of view of one of the two metrics, this is really the interaction between a positive energy mass and a negative energy mass.

Postulating the existence of many comoving coordinate systems, each system being defined relative to its source with the gravitational interaction being instantaneous only in the source frame is expected to potentially generate severe causality violations if we still transform all space-time intervals according Lorentz. But we believe that the interaction is instantaneous in any frame. Indeed, the physics of the instantaneous interaction takes place in the Euclidian metric where a zero multidimensional interval implies that all individual one-dimensional intervals are separately zero. Usually it's tentative in Minkowskian space-time to interpret non zero 4d space-time intervals (or masses) as contributions from extradimensional space or time intervals neglected in a first stage, but this is only possible thanks



to the mixed signature of the metric. In an Euclidian metric with any number of extradimensions, all sub-intervals are zero, this statement being also invariant under the Wick rotated version of Lorentz transformations. Thus the non-locality of our interaction is probably frame independent, a consequence of the null value of any interval in any frame from the Euclidian world perspective. Of course, after the Euclidian metric is Wick rotated into the final Minkowskian metric we have non zero intervals. But here takes place another physics sector: the physics describing how objects move under the influence of gravity. Eventually, if this is the correct approach, a given interval should be estimated depending on which kind of physics relates the interval events.

Whatever is the correct way to understand the privileged coordinate system, our metric solution can still (as in GR) be exported by a Lorentz transformation into any frame moving at a constant velocity relative to the privileged frame so that this transformation alone should not introduce any additional violation of the equivalence principle. Anyway, the Post-Newtonian formalism considering extended source distributions and propagated metric solutions does not apply here.

Another issue is the impossibility to span 3d-space with isotropic elementary volumes so that having determined the energy momentum tensor of a source distribution inside a given (necessarily non isotropic) 3d-space cell and computed its trace, we must recover isotropy by postulating that this source cell trace contribution is concentrated in a point. We thus have two modes of spacetime: the continuous Minkowskian one where we have to define an energy momentum tensor for our sources and understand how fields move under the influence of gravity, and a discret one where instantaneous gravity in Euclidian space takes place allowing to derive the Schwarzschild potential out of each individual cell. Eventually, whatever the number of privileged coordinate systems, we can divide any source distribution into cells having a mass point source (isotropic and static) behavior where the above Schwarzschild treatment always applies in vacuum.

We can compute gravitational energy following the standard method using the pseudo-energy momentum tensor. It vanishes for each individual metric solution thanks to the relation  $B=1/A$  in the privileged coordinate system. However the pseudo energy-momentum tensor does not happen to be meaningful out of the comoving privileged coordinate system. Indeed, in our framework this object is not even a Lorentz tensor. If we cannot compose several metrics into a single one then gravitational energy is only a valid concept for each individual source which probably defines a stand alone sector of physics.

The weak equivalence principle is obviously not menaced since once the metric field solution is established, matter and radiation will have to follow its geodesics as in GR. But, because of the non covariance of our equation, a violation of the strong equivalence principle arises at the PPN level for a point-like source. Gravitational energy, since it vanishes, of course falls in the same way as any other form of energy in a gravitational field and this case does not menace the strong equivalence principle more than in the point-like source case.

## 7. B=-A=h: The Phase: Gravitational Waves

The metric a priori takes the form:

$$\begin{aligned} d\tau^2 &= h(r, t) [d\sigma^2 - dt^2] \\ d\tilde{\tau}^2 &= h^*(r, t) [d\sigma^2 - dt^2] \end{aligned}$$

And the left hand side of our modified Einstein equation yields:

$$\frac{1}{A} (g_{rr}T^{rr} + g_{\theta\theta}T^{\theta\theta} + g_{\phi\phi}T^{\phi\phi} - g_{tt}T^{tt}) - inv = 6 \left( h'' - \ddot{h} + 2\frac{h'}{r} - h'^2 + \dot{h}^2 \right)$$

Then, we easily obtain the evolution equation for h in vacuum:

$$6 \left( \square^2 h + \dot{h}^2 - h'^2 \right) = 0$$

Only plane wave pure phase conjugate solutions transforming as required into one another through time and space reversal are acceptable. In general a superposition of such plane wave solutions is not a phase and for instance the spherical wave cannot be accepted as a solution for our equation unless the elementary wave solutions are one-dimensional topological defects with arbitrary directions. Each isotropic cluster of such superposition of one dimensional plane wave solutions produced with a common frequency by a quantum of generic impulse source term  $n\pi G\delta T = \frac{\delta(r)\delta(t)}{A}$  can be approximated by a continuous spherical wave on large scales so that the computation of the lost energy will proceed as in GR.

The elementary one dimensional  $h(r,t)$  outgoing plane wave solution of  $\square^2 h(r, t) = n\pi\frac{G}{6}\delta(r)\delta(t)$  makes  $\dot{h}^2 - h'^2$  vanish so that the energy carried by the wave is not a source for gravity. The superposition of this outgoing wave and ingoing wave with the same frequency is a standing wave. Because it is not a phase, it can contribute as a source term in the moduli equation. On large scales the averaged  $\dot{h}^2 - h'^2$  should vanish while on small space-time scales relative to the wavelength, the implied Zitterbewegung of  $\dot{h}^2 - h'^2$  is the perturbation needed to start a non stationary evolution, i.e. for the birth of a couple of time reversal conjugate universes which scale factors are the evolving moduli as we shall show in the next section.

We shall now predict the same energy lost through gravitational waves radiation of the binary pulsar as in GR in good agreement with the observed decay of the orbit period. We follow Weinberg's computation of the power emitted per unit solid angle and adopt the same notations. For any extended non relativistic source the solution is the retarded potential:

$$h(x, t) = n\pi \frac{-G}{4\pi} \frac{1}{6} \int d^3x' \frac{\delta T_{00}(x', t - |x - x'|)}{|x - x'|}$$

The radial momentum component of our gravitational wave energy momentum tensor reads:

$$T_{r0} = \frac{1}{n\pi G} \sum_{\sigma=0,3} \frac{\partial R}{\partial \left( \frac{\partial A_\sigma}{\partial r} \right)} \frac{\partial A_\sigma}{\partial t} = \frac{1}{n\pi G} \frac{\partial R}{\partial \left( \frac{\partial h}{\partial r} \right)} \frac{\partial h}{\partial t}$$

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We have used here an energy-momentum pseudo-tensor different from the usual one used in GR. It exploits the Lorentz invariance of both our action and the relation between conjugate metrics in the privilege coordinate system. Indeed, both the B=-A metric and its inverse are Lorentz invariant starting from the privileged coordinate system.

Replacing by the expression of our wave solution,

$$\forall \sigma, A_\sigma = h = \sum_{\omega, k} h(\omega, k) e^{i(\omega t - kr)} + h^*(\omega, k) e^{-i(\omega t - kr)}$$

$$\langle T_{r0}(\omega, k) \rangle = \frac{6}{n\pi G} (-2) \langle h \dot{h} \rangle_{\omega, k} = \frac{6}{n\pi G} 4\omega^2 |h(\omega, k)|^2$$

We find that the power emitted per unit solid angle in the direction  $k$  is:

$$\frac{dP}{d\Omega}(\omega, k) = r^2 \langle T_{r0}(\omega, k) \rangle = r^2 \frac{24}{n\pi G} \omega^2 |h(\omega, k)|^2$$

$$\begin{aligned} \frac{dP}{d\Omega}(\omega, k) &= 24 \frac{\omega^2}{n\pi G} \left( n\pi \frac{G}{4\pi} \frac{1}{6} \right)^2 \delta T_{00}^2(\omega, k) \\ &= \frac{n}{24} \frac{\omega^2 G}{\pi} \delta T_{00}^2(\omega, k) \\ &= \frac{2}{\pi} G \omega^2 \delta T_{00}^2(\omega, k) \\ &= \frac{G \omega^6}{2\pi} k_i k_j k_l k_m D_{ij}(\omega) D_{lm}(\omega) \end{aligned}$$

Where  $n$  has been given the numerical value 48 taking account of the fact that in the Newtonian limit our equation  $6\Box^2 h = 48\pi G \delta T$  must give again  $\nabla^2 g_{00} = -8\pi G T_{00}$  (the same straightforward reasoning allows to determine  $n$  in all the other gravitational equations of the theory!). Then following Weinberg, we may write in terms of the moment of inertia Fourier transforms in the observer coordinate system:

$$\begin{aligned} P &= \frac{2}{15} G \omega^6 \left[ (D_{ii}(\omega))^2 + 2D_{ij}(\omega) D_{ij}(\omega) \right] \\ P &= \frac{2}{15} G \omega^6 12 D_{11}^2(\omega) \end{aligned}$$

For a rotating body with angular speed  $\Omega$ , equatorial ellipticity  $e$ , moment of inertia  $I$  in the rotating coordinates,  $\omega = 2\Omega$ ,  $D_{11}(\omega) = \frac{eI}{4}$  and the radiated power reads:

$$P = 2 \frac{12}{15} G 64 \Omega^6 e^2 \frac{I^2}{16} = \frac{32}{5} G \Omega^6 e^2 I^2$$

as in General Relativity. The main difference is that our gravitational wave is found to propagate pure monopole modes linked by  $-g_{00} = g_{11} = g_{22} = g_{33} = A$ . But these cannot be excited independently whatever the source configuration since we have the single degree of freedom  $A(r, t)$ . Quantifying the  $h$  field must generate a new gravitational propagated interaction in addition to the Schwarzschild non propagated solution we obtained in the previous section. We cannot add its potential to the exponential Schwarzschild one since this would severely conflict with observations. Thus the two solutions do not cohabit. Rather the Schwarzschild solution is valid up to a critical distance (probably related to the value of the quantum of gravitational energy carried by the waves) where the other one takes over.

### 8. B=-A: The Modulus: Cosmology

The gravity for moduli, i.e. global gravity (this is cosmology) only depends on time. Indeed, there also exists a global privileged coordinate system where a couple of time reversal conjugate purely time dependent solutions can be derived from a new couple of conjugate actions. The existence of a time reversal conjugate universe was also suggested a long time ago in Ref. 8. Notice that even if we did not impose the B=-A condition, the only possible privileged coordinate system with both metrics spatially homogeneous and isotropic would be a flat Cartesian one. When B=-A, our equation for gravity now reads:

$$\sqrt{A^4}R - \frac{1}{\sqrt{A^4}}\tilde{R} = 0$$

with  $\tilde{R} = R_{A \rightarrow 1/A}$ . For a cosmological homogeneous source we will get purely time dependent background solutions and can keep only the time derivatives leading in the cold-cold regime to:

$$3A \left( -\frac{\ddot{A}}{A} + \frac{1}{2} \left( \frac{\dot{A}}{A} \right)^2 \right) - \frac{3}{A} \left( \frac{\ddot{A}}{A} - \frac{3}{2} \left( \frac{\dot{A}}{A} \right)^2 \right) = 0$$

The solutions such that the right hand side always vanishes describe worlds where only light can live (null trace). Its trivial A=1 stationary solution thus describes a self conjugate light world. A mass perturbation  $\delta(t)$  (as a superposition of ingoing and outgoing GW) is needed for the birth of times to take place and see a couple of universes start evolving. The background matter worlds have simple evolution laws in the particular ranges  $A \ll 1$ ,  $A \approx 1$ ,  $A \gg 1$ . Indeed, the scale factor evolution is then driven (here nondimensional time unit is used) by the following differential equations in the three particular domains:

$$a \ll 1 \Rightarrow \ddot{a} \propto \frac{3}{2} \frac{\dot{a}^2}{a} \Rightarrow a \propto 1/t^2 \text{ where } t < 0, \quad (12)$$

$$a \approx 1 \Rightarrow \ddot{a} \propto \frac{\dot{a}^2}{a} \Rightarrow a \propto e^t, \quad (13)$$

$$a \gg 1 \Rightarrow \ddot{a} \propto \frac{1}{2} \frac{\dot{a}^2}{a} \Rightarrow a \propto t^2 \text{ where } t > 0. \quad (14)$$

We can check that  $t \rightarrow -t$  implies  $1/t^2 \rightarrow t^2$  but also  $e^t \rightarrow e^{-t}$  thus  $A \rightarrow 1/A$ ,  $B \rightarrow 1/B$  when t reverses as required. Let us stress that the couple of cosmological metric solutions does not imply any local gravitational interaction between objects but only a global one between the two conjugate universes as in Ref. 9. A striking and very uncommon feature is that the evolution of the scale factor is mostly driven by the gravitational energy exchange between the coupled universes almost independently of the universes matter and radiation content. In particular, the observed flatness can no longer be translated into the usual estimation  $\Omega_m = 1$  from the WMAP data. The  $t^2$  evolution is one of the very few possibilities. Thus, we are most probably

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living in a constantly accelerating universe. Our and the conjugate universe crossed each other and two reversed time parameters appeared at their birthtime. At last, not only our universe is accelerated without any need for a cosmological constant or dark energy component but it is flat without inflation and gets rid of the big-bang singularity. The vanishing of cosmological constant terms provided

$$d^4\xi(x)\Lambda = d^4\tilde{\xi}(x)\tilde{\Lambda}. \quad (15)$$

appears to be only a local issue (for our Schwarzschild solution).

## 9. $B=1/A$ : The Modulus: Pioneer Effect and Pseudo-Horizon

### 9.1. *The Pioneer Effect*

When  $B=1/A$ , our equation for the background takes a very simple form:

$$\frac{3}{2}\left(B - \frac{1}{B}\right)\left(\frac{\dot{B}}{B}\right)^2 = 0$$

implying  $1=B=1/A$ . However, the unit element here remains to be defined. We require the unit element to be such that an object at rest ( $d\sigma = 0$ ) will not feel any gravitational discontinuity when passing from the  $B=-1/A$  regime to the  $B=-A$  regime. Thus the  $B$  element must be the same for both  $B=-A$  and  $B=1/A$  metrics. At a particular distance from a massive body, the cosmological metric flips locally from a  $A=-B$  solution to a  $A=-1/B$  solution. This results in the photons being redshifted compared to wavelengths emitted by atomic references in the  $A=-B$  expanded regime (it is more accurate to say that the reference periods contract while the photon keeps unaffected) while they will be blue shifted with exactly the same magnitude as compared to the same references in the  $A=-1/B$  regime (it is more accurate to say that the reference periods still contract while the photon periods contract twice more) . Then, perhaps we should not be surprised to receive the photons from an object at several Astronomical Units as are the Pioneer aircrafts slightly blue-shifted, an effect which according to Ref. 10 has been measured to a very good precision with the expected magnitude. The space-space metric component discontinuity when passing from the  $B=-A$  to the  $B=-1/A$  regime might transmit and deviate part of the photons and reflect the others as a genuine gravitational mirror. Alternatively, we could have required that the space-space element is the continuous one from one regime to the other and also got (this is easy to check) the correct redshifts and blueshifts provided  $A=-B$  are in contraction. But in this case, we get discontinuities of the metric for both matter and light. A discontinuity of the gravitational field is expected to produce caustics where matter is unexpectedly concentrated along rings.

### 9.2. *The Black Hole Horizon*

Our exponential solution tells us that there is no more BH singularity in our theory. However a test mass approaching the Schwarzschild radius of a massive body at rest

with respect to the global coordinate system is propagating in a modulus and phase originating superposition:

$$d\tau^2 = \frac{1}{A(r)R(t)^2} dt^2 - A(r)R(t)^2 d\sigma^2. \quad (16)$$

But for a strong enough gravitational field, e.g probably not much below the Schwarzschild radius thanks to the exponential regime, the  $A(r)$  term is expected to "cross" the  $1/A(t)$  scale factor term so that the space/time symmetry (hence a more stable configuration) of the euclidian metric could be restored. Reaching this radius the test mass would be transferred in the Euclidian world or possibly via an extra rotation in the space/time conjugate light world. As well as a genuine black hole horizon this mechanism would account for the absence of thermonuclear bursts from BH candidates.

## 10. Test mass motion

The model developed up to now only modified the way matter sources affect geometry. Remains to be clarified how various fields should move under the influence of various metrics. In the  $B=-A$  regime, we can adopt the usual action minimally coupling radiation and matter fields to the non dynamical metric as obtained from an additive superposition of our gravitational equations solutions. In the  $B=1/A$  regime, postulating a multiplicative superposition of local metrics solutions before applying the Wick rotation is only possible if there exists a single privileged coordinate system. Indeed, in general the multiplicative superposition of metrics will only generate a metric provided we have a common privileged system. If this is not the case, perhaps should we give up the hope to combine the metrics and minimally couple the resulting metric with matter and radiation in the usual way.

## 11. Phenomenological Outlooks

We now show that the new context may hopefully provide a very powerful alternative to dark matter models and a seductive common explanation to various outstanding enigmas after careful study of the postulated mechanisms.

### 11.1. *An Alternative to Dark Matter*

It is very tempting to interpret the universe voids as being filled with invisible matter living in the conjugate metric and repelling our matter at the frontier of what eventually appears to us as empty bubbles. On the other hand a massive structure in our metric repels the matter living in the conjugate metric creating there an attracting void for our structure which might help to explain the flat galaxy rotation curves and other gravitational lensing effects.

### 11.2. *Large Scale Structure Formation and the Early Universe*

Of course, the large scale structure formation and evolution as well as nucleosynthesis would need to be completely revisited in this new context where the game rules are drastically modified. We now have to consider:

- The interactions between conjugate density fluctuations favoring the gravitational collapse.
- The absence of gravitational horizon up to a given distance which would account for the large scale homogeneity of the CMB.
- A slower expanding rate in the past thanks to the accelerated regime.
- A twice older universe ( $2/H_0=28$  billion years) to be compared with the oldest galaxies ages ( $z=5$ )  $\approx 17$  billion years in quite a good agreement with the oldest stars ages while providing more time for galaxy formation.

Our hope is to be as successful as the standard model in cosmology while avoiding ad-hoc hypothesis and adjustable parameters such as a cosmological constant, cold dark matter, inflation...

## 12. Conclusion

We could settle down here the foundations for a modified theory of gravitation. This theory is essentially general relativity enriched to take into account the fundamental discrete symmetries involved in the structure of the Lorentz group. Eventually, we find that this allows to solve many long lasting theoretical issues such as negative energies and stability, QFT vacuum divergences and the cosmological constant but also leads to very remarkable phenomenological predictions: Locally, the disagreement with GR only arises at the PPN level and black holes disappear. Globally (this is cosmology), a constantly accelerating necessarily flat universe in good agreement with the present data is a natural outcome of the model. At last, we could also show that the space/time exchange symmetry allows to derive a gravitational wave solution leading to the observed decay of the binary pulsar orbital period and to clarify the status of tachyonic representations.

## References

1. F. Henry-Couannier, gr-qc/0404110.
2. S. Chen, hep-th/0203230.
3. S. Weinberg, *Gravitation and Cosmology* (John Wiley and sons, New York, 1972).
4. A. D. Linde, *Rept. Prog. Phys.* **47**, 925 (1984).
5. R. R. Caldwell, *Phys. Rev. Lett.* **91**, 071301 (2003).
6. P. H. Frampton, *Mod. Phys. Lett.* **A19**, 801 (2004).
7. S. M. Carrol, M. Hoffman and M. Trodden *Phys. Rev.* **D68**, 023509 (2003).
8. A. D. Sakharov, *JETP Lett.* **5**, 24 (1967).
9. A. D. Linde, *Phys. Lett.* **B200**, 272 (1988).
10. J. D. Anderson, *Phys. Rev.* **D65**, 082004 (2002).
11. M. Tegmark, gr-qc/9702052

12. G. Feinberg, *Phys. Rev.* **D17**, No 6.
13. E. Recami, arXiv:physics/0101108.