



SELECTIVE REFLECTION SPECTROSCOPY WITH A HIGHLY PARALLEL WINDOW: PHASE TUNABLE HOMODYNE DETECTION OF THE RADIATED ATOMIC FIELD

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Aram V. Papoyan, G. G. Grigoryan, S. V. Shmavonyan, David Sarkisyan, Jocelyne Guéna, et al.. SELECTIVE REFLECTION SPECTROSCOPY WITH A HIGHLY PARALLEL WINDOW: PHASE TUNABLE HOMODYNE DETECTION OF THE RADIATED ATOMIC FIELD. 8th EPS conference on atomic and molecular physics, 2004, Rennes, France. pp.4-84. hal-00002633v2

HAL Id: hal-00002633

<https://hal.science/hal-00002633v2>

Submitted on 27 Aug 2004

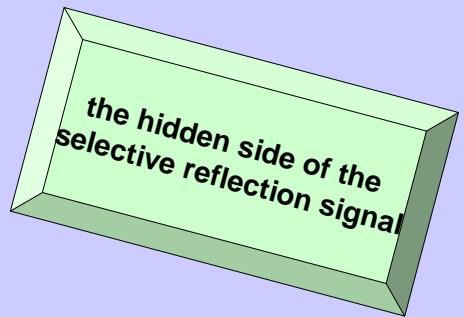
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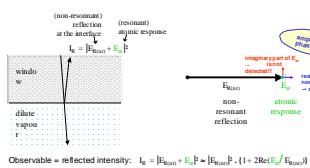
SELECTIVE REFLECTION SPECTROSCOPY WITH A HIGHLY PARALLEL WINDOW: PHASE TUNABLE HOMODYNE DETECTION OF THE RADIATED ATOMIC FIELD

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(to be published in Eur. Phys. J. D)



"ordinary" selective reflection

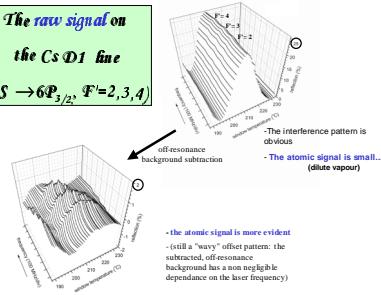


Observable = reflected intensity $I_R = |E_{R(n)}|^2 + |E_{A(\omega)}|^2 \cdot [1 + 2Re(E_A/E_{R(n)})]$

How to detect the imaginary part?? Some proposals have been made:

- Brewster incidence ($E_{R(n)}=0$)? (Gardiner et al. Soc. Opt. Eng. 1989, 416)
- the sub-doppler feature of SR spectroscopy is lost;
- multi dielectric coating? (here, work by Vartanian and Pogor. Opt Commun 110 (1994), 315)
- the coating may be damaged by the atomic vapour
- metallic coating? (Gardiner et al. Phys Rev E 64(2001), 026701)
- considerable attenuation of the atomic signal, due to the required metal thickness

The raw signal on the Cs D1 line ($6S \rightarrow 6P_{3/2}$, $F=2,3,4$)



The model

Continuity equations at the two boundaries between the three media:

- air, $n_s=1$

- (sapphire) window, $n_w=1.76$

- vapour, $n_v=1$

Maxwell equations for the propagation of the backward atomic field in the vapour (without using the slowly varying envelope approximation)

$$\frac{\partial^2 E(z)}{\partial z^2} + 2ik \frac{\partial E(z)}{\partial z} = -\langle k^2 |E\rangle V(z)$$

assuming cell length L as absorption length (no backward beam coming from $z=L$)

$$E_{R(\text{backward})} = \left[\frac{p(z) t(z) \exp(2ikz)}{1 - p(z) t(z) \exp(2ikz)} \right] E_0$$

$$+ \frac{p(z) r(z) \exp(2ikz)}{1 - p(z) r(z) \exp(2ikz)} E_0 = E_{R(\text{vapour})}$$

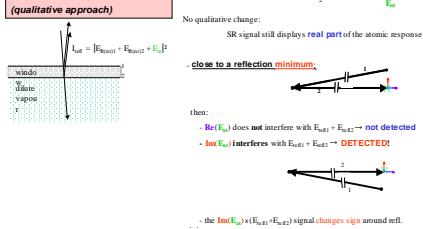
(where the $p(z)$ and $r(z)$ are the amplitude transmission and reflection coefficients and the backward atomic field is generated by the vapour atomic polarisation):

$$E_0 = 2ik \int_0^{L-z} P(z') \exp(2ikz') dz'$$

Defining the atomic response, $I_R = \frac{1}{4} E_R^2 E^* \exp(2ikz)$ and assuming the absence of saturation and non-linearity, we get ($T_c = 190^\circ\text{C}$):

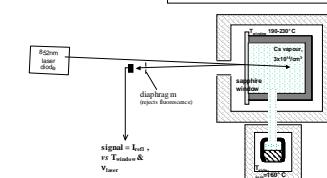
$$I_R = \frac{1}{h} \nabla^2 \nabla^2 \left[\frac{1}{2} \exp(-2kMz) \right]$$

selective reflection with a parallel window (qualitative approach)

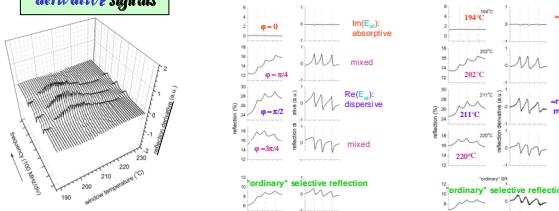


The experiment

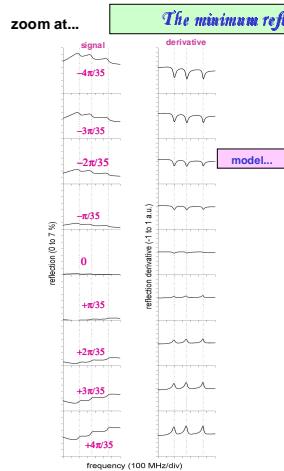
How to change the interference condition in the window?
very easily, by changing the window temperature
For 0.5 mm sapphire window and $\lambda=852\text{nm}$:
 $\Delta T = 30^\circ\text{C} \leftrightarrow 2\pi$ change of the interference
(see Jaksic et al. Appl Phys B 71 (2000), 561 for the use of the "temperature tuning" of the windows for reflection loss free vapour cells)



The raw and derivative signals



The minimum reflection regime



Conclusion

The model and experiment agree very well (no fitted parameter!) on the size and the temperature dependence of the spectra.

By using a "temperature tunable" window, one can detect at will

- the real (dispersive) part
- or the imaginary (absorptive) part of the atomic response.

SN is better near the reflection minimum.

Changing from one regime to the other is obtained very easily, just by changing the window temperature by a few degree C.

Possible application: temperature-tunable locking of a laser frequency on the zero of the derivative signal

