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Epitaxy and growth of titanium buffer layers on $Al_2O_3(0001)$

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Abstract

The structure and growth of thin films of titanium on $\alpha-Al_2O_3$ at room temperature were investigated through *in situ* RHEED observations. Two different structures coexists at low coverage. One corresponds to the $Ti(0001) \parallel Al_2O_3(0001)$, $Ti[1\bar{1}00] \parallel Al_2O_3[2\bar{1}\bar{1}0]$ and $Ti[10\bar{1}0] \parallel Al_2O_3[1\bar{1}00]$ epitaxy of the α phase of titanium reported before for thick films prepared at high temperature. The other structure can be explained by co-existence of α and ω Ti in thin films. It was shown, with the use of tight-binding total energetic calculations that, the ω phase could actually be stabilized by the substrate. In addition, it was demonstrated that the presence of this extra structure has a dramatic effect on the epitaxial growth of the Ag overlayers on the system. This can be the origin of the non-trivial buffer effect of titanium previously observed.

1 Introduction

Finding surface treatments which can force layer-by-layer growth is among the key issues in heteroepitaxy (1). Indeed, many attempts have been made to promote the wetting of materials of different and non interacting nature such as the metal on oxide interface (2; 3). The challenge is not only to grow durable interfaces but also to control the morphology of the metal film as it affects the physical properties of the system. The subject is of increasing importance as many industrial applications rely on the mastering of interfaces such as

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copper interconnects on oxides for microelectronics and silver on oxides for optical purposes. As a trend, adhesion energy decreases from early transition metals to noble metals (4; 5), due to the increased filling of the 3d band (6; 7). These differences have major consequences regarding the wetting properties of metals on oxides and lead to a large variety of growth situations. The thermodynamic condition for two-dimensional growth of a metal film on a substrate requires the adhesion energy β to be equal to or higher than twice the surface energy of the metal σ_M . The way β compares to σ_M is of central importance for the growth of the film. While the condition $\beta < \sigma_M$ implies $\theta > 90^\circ$ and a growth in a three-dimensional fashion up to very high coverage, $\beta > \sigma_M$ implies $\theta < 90^\circ$ and leads to a quickly percolating film. Clearly, noble metals and late transition metals on wide band gap oxides belong to the first category (8) and early transition metals to the second (9). A method often used to improve the wetting of noble metals on oxide surfaces is pre-deposition of early transition metals (10). This so-called buffer effect involves both electronic and geometric contributions (10).

In the present work, attention is focused on the effect of titanium buffers for noble metal growth on alumina, a subject receiving attention because of its application in microelectronic devices (3; 10), and references therein. Suzuki et al. (11) found by ion scattering that there is no interdiffusion at the interface between the titanium overlayer and the Al-terminated α - $\text{Al}_2\text{O}_3(0001)$ substrate. By performing surface diffraction experiments on a 20 monolayer thick titanium film annealed to 1100 K, they showed that the epitaxial relationship was $\text{Ti}(0001) \parallel \text{Al}_2\text{O}_3(0001)$, $\text{Ti}[1\bar{1}00] \parallel \text{Al}_2\text{O}_3[2\bar{1}\bar{1}0]$ and $\text{Ti}[10\bar{1}0] \parallel \text{Al}_2\text{O}_3[1\bar{1}00]$. This means that the triangular lattice of Ti(111) is rotated by 30° with respect to the triangular lattice formed by the outer Al atoms of $\text{Al}_2\text{O}_3(0001)$, see figure 1. Dehm et al. (3) came to similar conclusions through combined transmission electron microscopy (TEM) and reflection high energy electron diffraction (RHEED) observations of a 10-100 nm thick titanium film on $\text{Al}_2\text{O}_3(0001)$. These authors did not investigate in details the crystal structure of the interface Ti/ Al_2O_3 . Recently, a first principle calculation of Ti on $\text{Al}_2\text{O}_3(0001)$ predicted a tendency to intermixing for aluminium and titanium at the interface such that the most stable site for the titanium adatoms is above or in replacement of the surface aluminium atoms (12). These results have still not been experimentally verified.

Concerning the buffer effect it was reported in Ref. (3) that copper grew in an epitaxial manner on thick titanium films (10-100 nm) on $\text{Al}_2\text{O}_3(0001)$. The buffer effect of titanium was indirectly demonstrated on silver for much thinner films through optical reflection *in situ* measurements of the silver plasmon (13). It was found that the use of titanium buffers had a significant influence on the morphology of silver films, although the effect of the titanium film was shown to depend highly on its thickness. Very thin titanium layers (< 0.3 nm) hardly produce any change in the three dimensional growth of silver with respect to

the bare Al_2O_3 surface. In contrast, a titanium layer of 0.6 nm leads to a better wetting of silver between silver and the substrate. However, for higher titanium coverage, the quality of the wetting decreases, although it remains better than on the bare surface. Similar results were obtained by surface diffraction experiments (14). The very intriguing complexity of this scenario can hardly be explained by the surface energy of titanium alone. In the present work the structure of the titanium buffer is investigated as a function of its thickness. The corresponding behavior of a silver overlayer is further discussed.

2 Experimental

Single crystalline $\text{Al}_2\text{O}_3(0001)$ substrates from Mateck GmbH with a $\leq 0.5^\circ$ miscut were prepared using acetone plus ultrasonic bath followed by a cleaning in a 10 percent dilution of a standard buffered basic soap with $\text{PH}=10.6$. After washing in deionized water and drying in isopropanol vapor, an *ex situ* annealing at 1320 K was performed under an atmospheric pressure of oxygen to obtain a good crystallinity. Afterwards, samples were inserted in a MECA2000 molecular beam epitaxy apparatus, annealed under a partial pressure of $5 \cdot 10^{-5}$ torr O_2 to get rid of carbon contamination, first at 820 K over a couple of hours and then at 1070 K for one hour. Samples were cooled under oxygen prior to transfer to the evaporation chamber where the base pressure was $1 \cdot 10^{-10}$ torr. In a similar vacuum chamber equipped with a photoemission analyser, it was checked that this treatment leads to clean $\text{Al}_2\text{O}_3(0001)$ surfaces (13).

Titanium and silver films were evaporated on the samples using a Telemark electron gun. The flux was between 0.08-0.1 nm per min and the samples were held at 400 K. As titanium is a very reactive material special care was taken to keep the chamber pressure within the 10^{-10} torr range during evaporation. The thickness was measured with a quartz microbalance calibrated using *ex situ* electron microprobe measurements to obtain the number of atoms deposited. To eliminate any distortions of the diffraction pattern due to stray fields from the electron guns, special care was taken to align the optical center of the RHEED with the electron guns in operation condition. The evolution of the growth was followed by a Staib RHEED instrument and the diffraction patterns were on a CCD camera. During the deposition any physical movement of the samples used for the lattice measurements were avoided to keep the substrate reference for the diffraction patterns. When analysing short lattice spacings, for which the diffraction patterns would be outside the screen, the scattered beams were deflected on the screen by the means of an electrostatic device. It was verified several times that this method does not introduce any errors on the positions of the diffraction peaks on the bare sapphire substrate.

3 Results

The following section first presents the results and the discussion related to the growth of titanium on Al_2O_3 . The influence of titanium on the growth of silver is then briefly considered. Directions along which RHEED analysis is performed are always given with respect to the crystallographic axes in figure 1. Distances in real and reciprocal space will be labelled l and g respectively.

3.1 Growth of the titanium film

Figure 2 shows the typical evolution of the RHEED pattern of a titanium layer collected along the $\text{Al}_2\text{O}_3[\bar{1}100]$ zone axis which means that the observed surface streaks correspond to the $\text{Al}_2\text{O}_3[2\bar{1}\bar{1}0]$ axis. The film starts growing with a reciprocal lattice parameter close to the substrate lattice, g_{s2} . Up to 1 nm, the broad and not very intense diffraction rods indicate a low degree of order. Upon increasing the titanium coverage, the diffraction pattern globally gains intensity. Between approximately 1.8- 6 nm, two separated lattice parameters $g_{\omega 2}$ and $g_{\alpha 1}$ coexist. At lower coverage, the external rod $g_{\omega 2}$ is the most intense (figure 2, 0.3-3.6 nm.). For increasing coverage the internal rod $g_{\alpha 1}$ increases in intensity to finally dominate the pattern (figure 2, 5.6 nm). Note that the thickness where the switch in intensity between $g_{\omega 2}$ and $g_{\alpha 1}$ occurs depends on the sample. For some crystals it is impossible to get a correct diffraction pattern below 2 nm of Ti - for others the switching between the lattice spacings $g_{\omega 2}$ and $g_{\alpha 1}$ is already observed at a lower coverage. But, in all cases, in-plane rotations of 360° show that the symmetry of the structure is 6 over the whole thickness range.

To collect the peak corresponding to $\text{Ti}[10\bar{1}0]$, the RHEED pattern of the growing titanium was also recorded along the $\text{Al}_2\text{O}_3[2\bar{1}\bar{1}0]$ azimuth. This was delicate because the working distance in the RHEED setup is mainly aimed at analysing large lattice spacings. Small lattice spacings can only be observed by deflection of the diffraction pattern on the CCD camera. Therefore, the diffraction patterns were collected by halves. To get a complete image the two halves were recombined using the (00) streaks as a reference (figure 3). The corresponding lattice parameters were obtained as the mean value found from the half images. The diffraction patterns in the $\text{Al}_2\text{O}_3[2\bar{1}\bar{1}0]$ zone axis were in general of low intensity. Interestingly, the patterns obtained around 2 nm of titanium are close to those of the clean substrate (figure 3). This pattern disappears as the streaks corresponding to the $\text{Ti}[10\bar{1}0]$ appear. A full rotation of the sample reveals a six fold symmetry of the pattern.

The profiles of the diffraction features were extracted to analyse the pattern

in the $\text{Al}_2\text{O}_3[1\bar{1}00]$ zone axis. In order to smooth the noise of the data, three photos recorded for each film thickness were summed and intensities were cumulated along the diffraction peaks. To avoid zones where the diffraction streaks can be curved only the part of the diffraction pattern close to the shadow edge of the sample was used. Intensity profiles are showed in figure 4 (upper right corner). The analysis of this data is complicated by the presence of two overlapping and broad peaks. Nevertheless, since in a wide coverage range, the streaks are split into two lines separated by a valley, it has been chosen to decompose the profiles into two components over the whole titanium coverage range. Diffractions peaks are Lorentzian functions but deviations from this are expected when working with a CCD camera. Therefore several types of functions were tested to obtain the best possible model over as large a range of coverage as possible. It proved to be a Voigt function with a 80% ratio of the Lorentzian to Gaussian. All the diffraction peaks could be independently fitted by means of two Voigt functions of equal width (figure 4). Due to the high background level and the presence of Kikuchi lines, profiles associated with low titanium coverage were difficult to decompose. In particular for coverage below 2 nm, the existence of the $g_{\alpha 1}$ line could be questioned. But at higher coverage of titanium the fits are unambiguous. The fact that the in-plane atomic distances are constant over a wide range of coverage (figure 5) while the residue of the decomposition is very small (figure 4) proves that the decomposition of the diffraction line in two peaks is physically sound.

Figure 5 a) shows the evolution of the area of the fitted peaks. In the 1-3 nm range, the $g_{\alpha 1}$ and $g_{\omega 2}$ peak areas are about constant with $A_{g_{\omega 2}} > A_{g_{\alpha 1}}$, while, above 3 nm, $g_{\alpha 1}$ increases and $g_{\omega 2}$ progressively vanishes. Due to the existence of multiple diffraction effects in RHEED these intensities are only a qualitative indications of the presence of the structures corresponding to the observed peaks in the film. The relative positions of the fitted peaks $g_{\alpha 1}$ and $g_{\omega 2}$ converted into the real space lattice constants l_{α} and l_{ω} are presented in figure 5 b). Titanium first grows with a lattice constant close to that of the substrate l_s and then undergoes a contraction of $\approx 4\%$ to a lattice constant l_{ω} and finally it expands to the value l_{α} close to the interatomic distance of bulk $\text{Ti}[1\bar{1}00]$ which is 7.3% larger than l_s , see ref. (19). The intermediate contraction of the film is very surprising as the titanium lattice is already contracted on the substrate, there should be no gain in energy by a further contraction.

An analysis along the $\text{Al}_2\text{O}_3[2\bar{1}\bar{1}0]$ zone axis also reveals a change in lattice spacing of the titanium film even if the diffraction patterns are more blurred than along the $\text{Al}_2\text{O}_3[1\bar{1}00]$ axis (figure 3). Around 2 nm, $g_{\omega 1}$ streaks are observed in a position which corresponds to a compression by 4% with respect to the substrate. These streaks disappear while the $g_{\alpha 2}$ streaks become visible.

3.2 Titanium as a buffer layer for silver growth

To elucidate the role of the thickness of the titanium buffer effect, 2 nm thick silver films were grown with titanium buffers of different thickness. Figure 6 shows typical RHEED patterns of the silver films with the beam along the $\text{Al}_2\text{O}_3[1\bar{1}00]$ axis. On the bare $\text{Al}_2\text{O}_3(0001)$ substrate, between 300-600 K, silver grows without any epitaxial relation and develops the ring structure typical for polycrystalline films constituted of three dimensional domains (figure 6 a). For a titanium layer of 0.5-1.5 nm, silver deposition leads to streaks indicating a better spread of the film (figure 6 b.). A full rotation in the plane shows that the film has a sixfold symmetry corresponding to an epitaxy of a Ag(111) plane. However, at $\pm 15^\circ$ around the principal axes of the substrate, the diffraction pattern does not change considerably and only decreases in intensity meaning that the epitaxy is loose. On thicker titanium films (2-4 nm), when the two streaks $g_{\omega 2}$ and $g_{\alpha 1}$ are present, the silver layer exhibits a peculiar structure (figure 6 c). The diffraction pattern is composed of dots which suggest a bad wetting of the silver on substrate. The pattern seen in figure 6 c) presents spots which are characteristic of a long range order, but do not correspond to any low index direction of silver. Finally, when silver is deposited on a titanium film thicker than 4 nm and when the $g_{\alpha 1}$ streak dominates, the diffraction streaks are unperturbed by the silver deposition, which means that the silver wets the substrate and grows epitaxially (figure 6d). Therefore, the behavior of the silver film unambiguously shows that an important change occurs in the titanium buffer between 4 and 5 nm.

4 Discussion

4.1 Growth of titanium

The data connected to the growth of the titanium film will first be discussed. At high titanium coverage, the epitaxial relationship is dominated by $\text{Ti}(0001) \parallel \text{Al}_2\text{O}_3(0001)$, $\text{Ti}[1\bar{1}00] \parallel \text{Al}_2\text{O}_3[2\bar{1}\bar{1}0]$ and $\text{Ti}[10\bar{1}0] \parallel \text{Al}_2\text{O}_3[1\bar{1}00]$ in agreement with the epitaxy found by other groups for thicker films at much higher temperature (3; 11). But at lower coverage, extra structures appear in the diffraction diagram in both zone axes in the surface plane (figures 2 and 3). As the $g_{\omega 1}$ and the $g_{\omega 2}$ streaks show up in the same coverage range and behave similarly, it is reasonable to attach them to the same structure. The extra streaks can either be due to multiple scattering between the principal streaks and an overlying domain structure or the existence of another epitaxial relation between the substrate and the titanium film. Both structures exhibit the same 6 fold symmetry as the substrate and they are the only visi-

ble structures during a full rotation of the sample. The dramatic change in the growth mode of the silver overlayers is the indication of a significant difference between the underlying structures of the titanium. This and the absence of extra streaks around the zero order peak and the asymmetry of the $g_{\omega 1}$ and the $g_{\omega 2}$ streaks make the existence of a domain structure unlikely.

If we assume that the observed pattern is the diffraction from a crystal plane then the ratio $g_{\omega 2}/g_{\omega 1} = \sqrt{3}$ suggests a hexagonal structure with a lattice parameter $a \approx 0.455$ nm. The orientation of the hexagon is the same as the hexagon formed by the outer Al atoms on the $\text{Al}_2\text{O}_3(0001)$ surface. This plane does not correspond to any hexagonal plane of the α -Ti but it compares with that of the (001) plane of the high pressure hexagonal ω phase of titanium (21). Therefore, the double peak observed in figure 2 can be explained by the coexistence between two epitaxial hexagonal configurations, one is α phase and the other is considered to be the ω phase of titanium .

The ω -Ti phase is a distorted bcc and one of the first steps in the high pressure hcp to bcc phase transition exhibited in the *IV* metals (20). This phase is frequently encountered in studies of the mechanical properties of titanium. The different high pressure phases of titanium were recently extensively studied both experimentally in a diamond anvil cell (21; 22) and theoretically (23; 24). Under hydrostatic pressure the ω phase coexists with the α phase from around 2 GPa to 9 GPa where the ω phase dominates. The phase is stable up to 116 GPa. Once the ω phase is obtained it can be kept at the ambient pressure although it is metastable. The ω -Ti(001) plane has been schematized (figure 7) using the lattice parameters given in Ref. (23) and the structure description in Ref. (25). The ω phase can be seen as the successive stacking of a hexagonal lattice and a graphite like sheet, the hexagon being the origin of the periodicity of the lattice. The parameter of this hexagon $a=0.4598$ nm (25) is indeed close to what was observed herein.

A few substrate-induced phase transitions have been observed in metal on metal epitaxial systems. One of most studied examples is the fcc - bcc martensitic transition of iron on metallic substrates like Cu and $\text{Cu}_3\text{Au}(001)$. A very elaborated example is the STM and RHEED study of the Bains transition of iron on $\text{Cu}_3\text{Au}(001)$, (26; 27; 28). As in the present case, the authors obtained a critical thickness for the fcc phase between 1 and 10 ML depending on the substrate preparation and the films exhibited a high amount of in-plane disorder.

To evaluate the energetic balance between the α and ω phases under various compressive and extensive conditions a serie of total energy calculations was undertaken. The very small energy difference between the α and ω phase makes such calculations difficult. As a matter of fact, this phase transition is not properly described by first-principles methods by which the equilibrium

ω -Ti phase is found to be slightly lower in energy (about 7 meV) than α -Ti (23). Therefore, a total energy calculations was performed using the set of tight-binding parameters recently developed by Papaconstantopoulos and co-workers (23; 30) to reproduce as accurately as possible the α to ω phase transition. The calculated equilibrium lattice constants of the α -Ti and ω -Ti phases are $a_\alpha = 0.294$ nm (*i.e.*, $l_\alpha = \sqrt{3}a_\alpha = 0.51$ nm) and $a_\omega = 0.459$ nm respectively, showing a lattice mismatch with respect to the substrate of +7.4% and -3.4%. The total energy per atom of α -Ti (*resp.* ω -Ti) has been calculated for hexagonal lattice parameters ranging from perfect matching with the substrate, up (*resp.* down) to the equilibrium lattice constants (figure 8). Each point of the curve is obtained by fixing the hexagonal lattice parameter a and minimizing the total energy with respect to the parameter c . It clearly appears that, at the lattice constant giving the perfect matching with the substrate, ω -Ti is favored with respect to α -Ti, whereas it is the reverse at the equilibrium lattice constant. The small energy difference between the two states makes the coexistence between the two phases over a range of deposited thicknesses very plausible. The presence of defects on the substrate surface can also modify the energetic balance between the two phases. The RHEED data do not give any indication about the arrangement of the two phases which are assumed to coexists within the titanium layer nor does it give insight about the burried part of for the titanium. However, data show clearly that above a certain thickness the surface of the titanium film is exclusively made of the α - Ti with the Ti(0001) \parallel Al₂O₃(0001), Ti[1 $\bar{1}$ 00] \parallel Al₂O₃[2 $\bar{1}$ 10] and Ti[10 $\bar{1}$ 0] \parallel Al₂O₃[1 $\bar{1}$ 00] epitaxy.

4.2 Influence of the cristallography of titanium on the wetting of silver

Concerning the growth of silver on titanium the strength of the metallic bond should clearly favor the good spreading of silver on titanium. Furthermore the lattice spacing the of two most compact hexagonal planes of titanium and silver are the Ti(0001) with $a = 0.295$ nm and the Ag(111) with $a = 0.2889$ nm (19). Therefore, it can be anticipated that the Ag/Ti epitaxy is favored over the Ag/Al₂O₃ epitaxy because it corresponds to a lower mismatch (-2% versus 5.34%). Thus, both the surface energetic and the lattice constants seems to be in favor of a better wetting of silver on titanium on Al₂O₃ at least when bulk values are considered.

The data presented in figure 6 present a more complex picture and suggest that a part of the buffer effect of titanium on silver is linked to the structure and the stress of the interface. On the bare alumina substrate silver grows as three dimensional polycrystals. Silver does not wet the alumina which transfers little order to the metal overlayer. However, the presence of very thin titanium films improves the wetting of alumina by silver (≈ 1 nm, see Ref. (13), the silver film

showing a texture around Ag(111) orientation meaning that the titanium film transfers a part of the substrate order. Taking into account the very diffuse diffraction patterns for thin Ti layers in this region it is possible that the silver film only replicates the underlying Ti structure. At higher coverage, for titanium layers where $g_{\omega 2}$ dominates and the $g_{\alpha 1}$ just starts to become visible, the wetting of the substrate by silver becomes poor, figure 6c. Finally, when the titanium film is thicker and has a nearly relaxed lattice parameter, silver grows epitaxially in a laminar way. The Ag(111)||Ti(0001), Ag[110]||Ti[2 $\bar{1}$ 10] epitaxy corresponds to expectation and it is similar to that found for the Cu/Ti/Al₂O₃ system (3). Such a structure is consistent with the observation of a pure α Ti phase by electron diffraction. The bad wetting of silver on the intermediate titanium coverage is an evidence that a titanium structure drastically different from α Ti is present. Indeed, a like explanation is the presence of the ω Ti at the surface of the titanium film. The change in the interfacial energy is due the large mismatch (9.9%) between ω Ti(0001) and Ag(111) and prevents the spreading of silver.

The present data show that the silver layers exhibit a remarkable amount of different textures. Laminar growth is only obtained when the lattice constant of the titanium film is very close to the silver bulk value. Therefore, the buffer effect of titanium for noble metals is not only the result of a change in surface and interface energies. It is also very dependent on the crystallography of the buffer and therefore on the elastic contributions to the interfacial energy.

5 Conclusion

The room temperature growth of thin titanium films on $\alpha - \text{Al}_2\text{O}_3(0001)$ has been shown be more complex than hereto reported. For films above 4 nm the α -Ti phase is obtained in an epitaxial relation already observed and described in literature for the high temperature growth of thick titanium films on sapphire. For thin films an additional structure is present. This structure is assigned to the (0001) plane of the ω phase where the hexagon is aligned with the one of an Al-terminated $\alpha - \text{Al}_2\text{O}_3(0001)$. The two phases coexist at the beginning of the growth of the titanium film in the range of coverage (≈ 1 -4 nm). The substrate-induced existence of an ω phase is plausible when taking simple lattice considerations into account, but the details behind the growth and stabilisation of this phase still need to be explained.

The structure of the titanium film has a dramatic influence on the so-called buffer effect, which has been tested herein by depositing silver on a $\alpha - \text{Al}_2\text{O}_3(0001)$ surface pre-covered by titanium at various coverage. Indeed, the titanium film structure and strain was proved to strongly affect, not only the structure, but also the wetting of the silver film. The silver film, which clearly

grows in a 3D manner when the film structure contains the structure we assign to the ω -Ti phase, spreads out on the α -Ti phase. Therefore, the growth mode and spreading of noble metals deposited on titanium-covered alumina arises not only from thermodynamics but from a subtle interplay between the structure and the surface energy of the titanium film.

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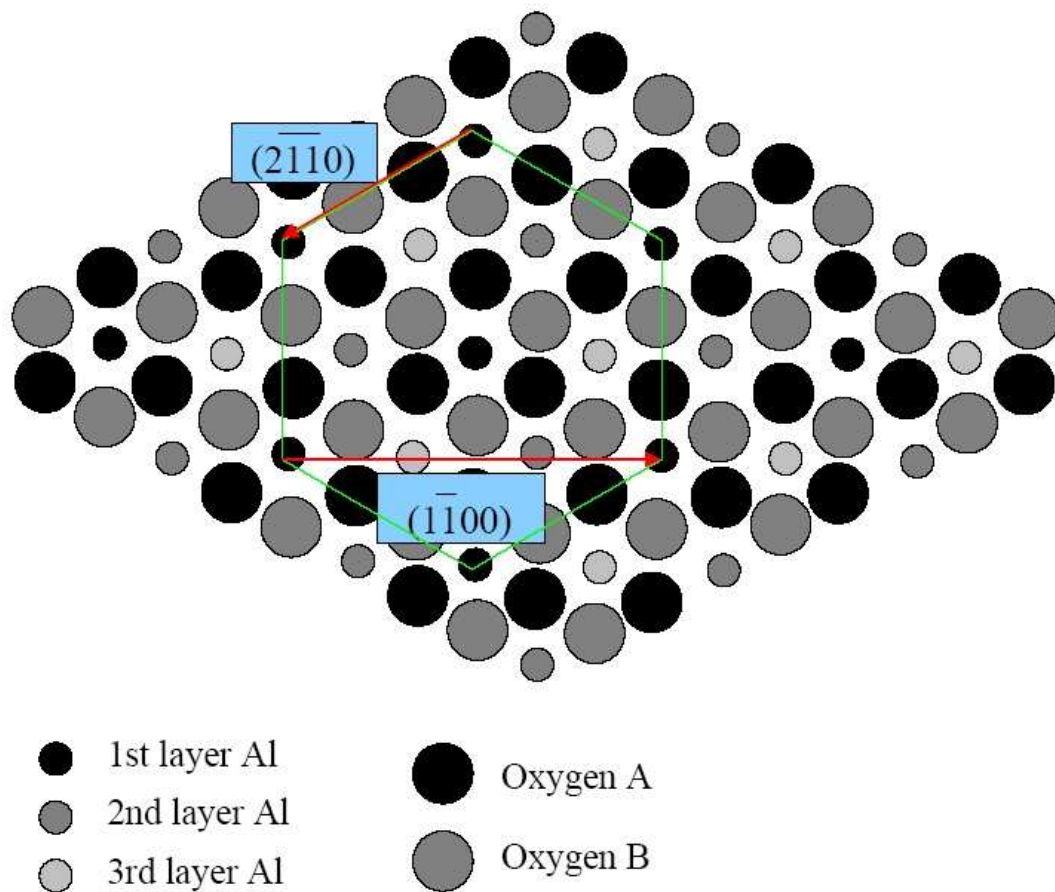


Fig. 1. Schematic Top view of an Al terminated α $\text{Al}_2\text{O}_3(0001)$ surface. The hexagon indicates the position of the outer Al- layer.

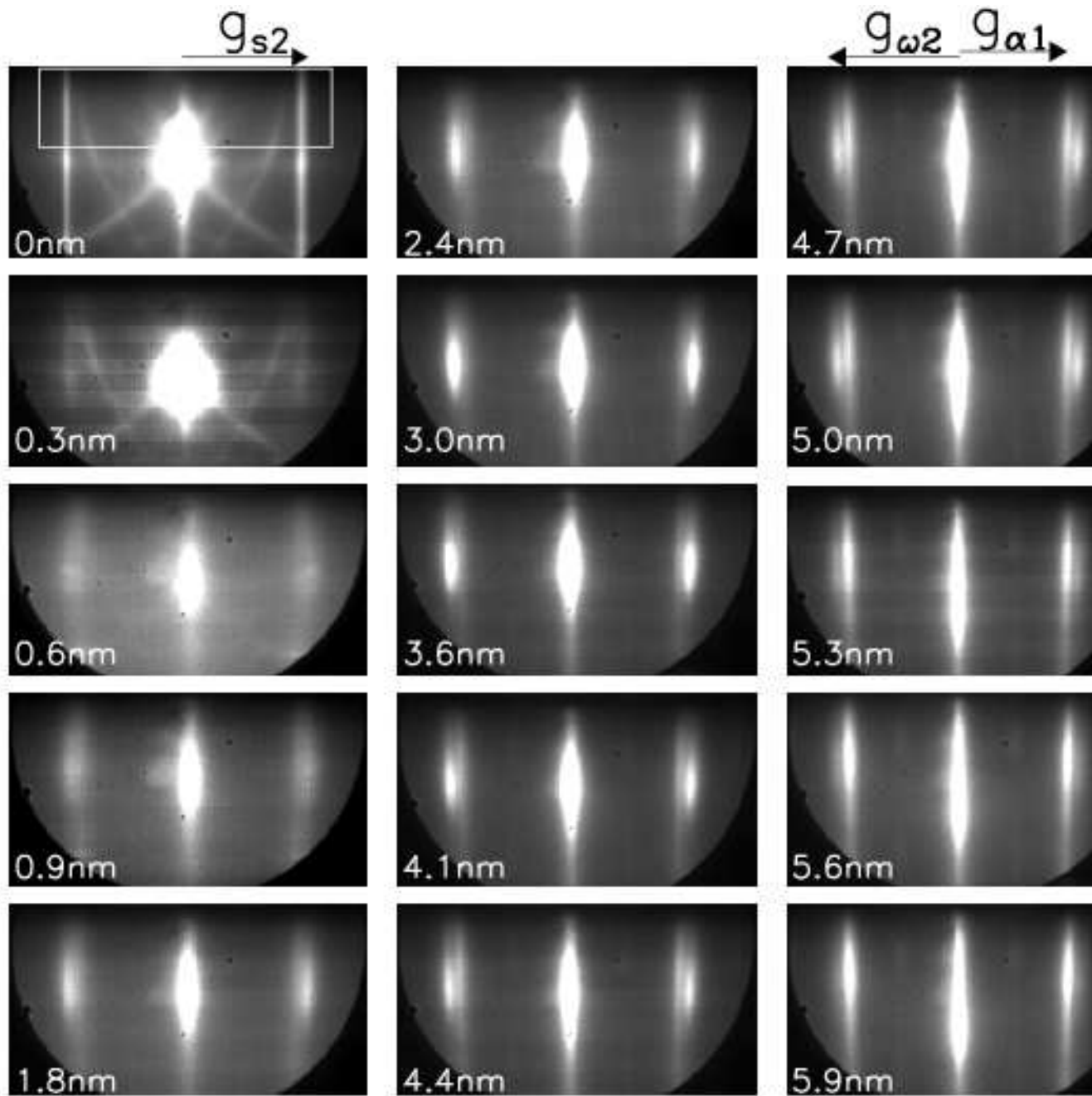


Fig. 2. RHEED images showing the growth of the Ti film with the beam along the $Al_2O_3[1\bar{1}00]$ axis.

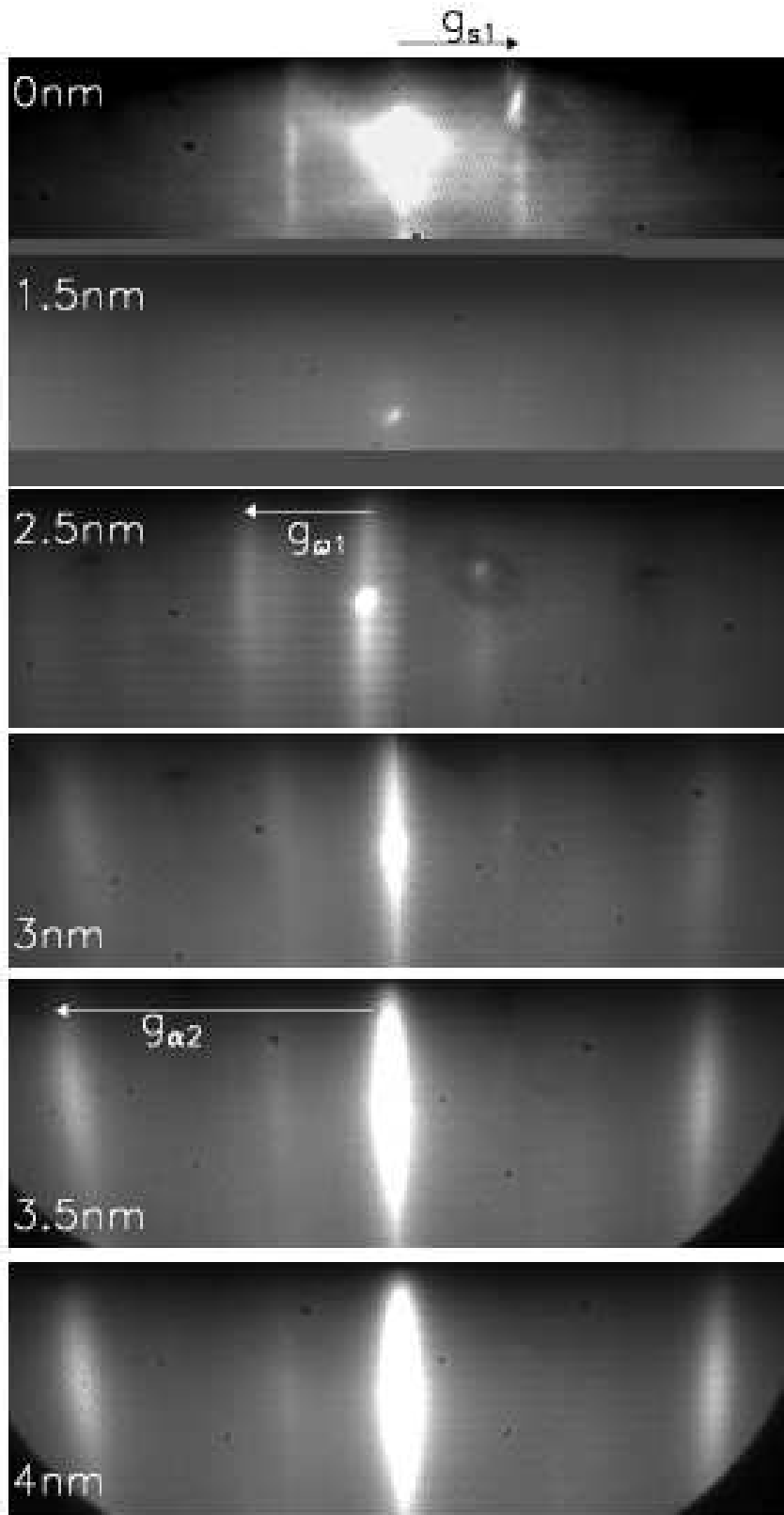


Fig. 3. RHEED images showing the growth of the Ti film with the beam along $Al_2O_3[2\bar{1}\bar{1}0]$. The images presented were obtained as the recombination of two deflected images in order to observe the diffraction peak due to the $Ti[10\bar{1}0]$.

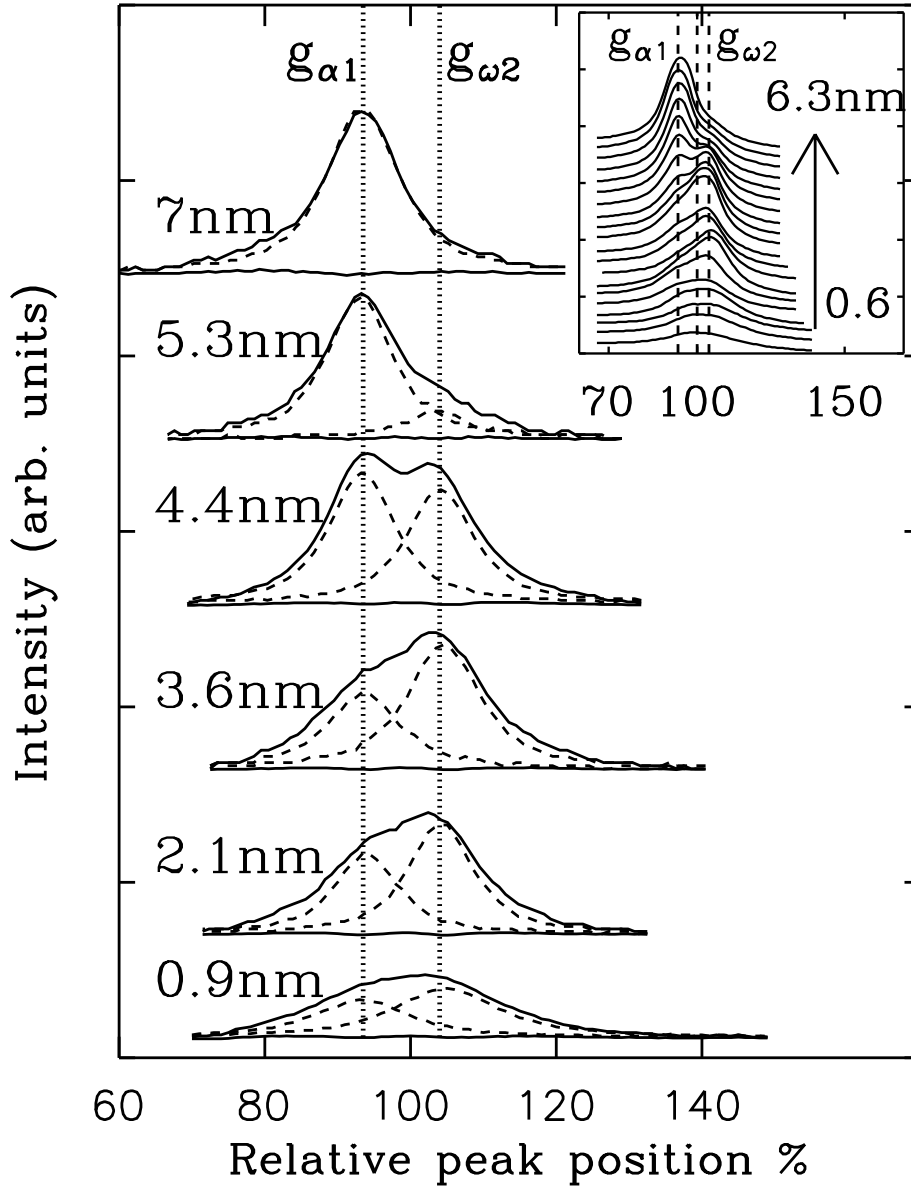


Fig. 4. The intensity profiles of the diffraction streaks recorded along $Al_2O_3[1\bar{1}00]$ axis (figure 2) for titanium films of thickness ranging from 0.9 to 6 nm and their corresponding fits with Voigt functions. The gradual shift in intensity from $g_{\omega 2}$ to $g_{\alpha 1}$ positions is seen on the inserted figure.

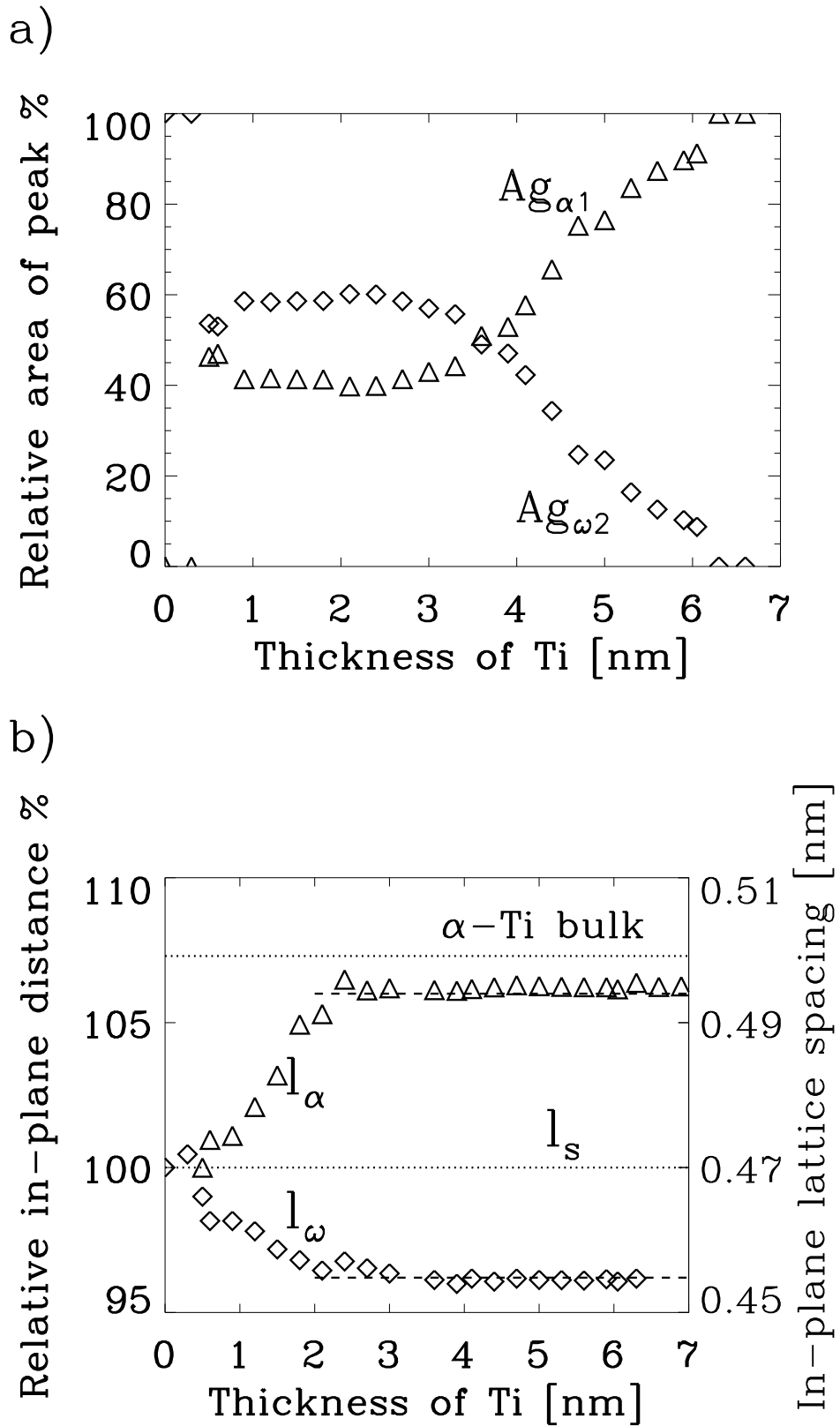


Fig. 5. a) The area of the diffraction peaks of figure 2 given by the fitting procedure. b) The fitted lattice spacings, l_{α} and l_{ω} , relative to the substrate, l_s .

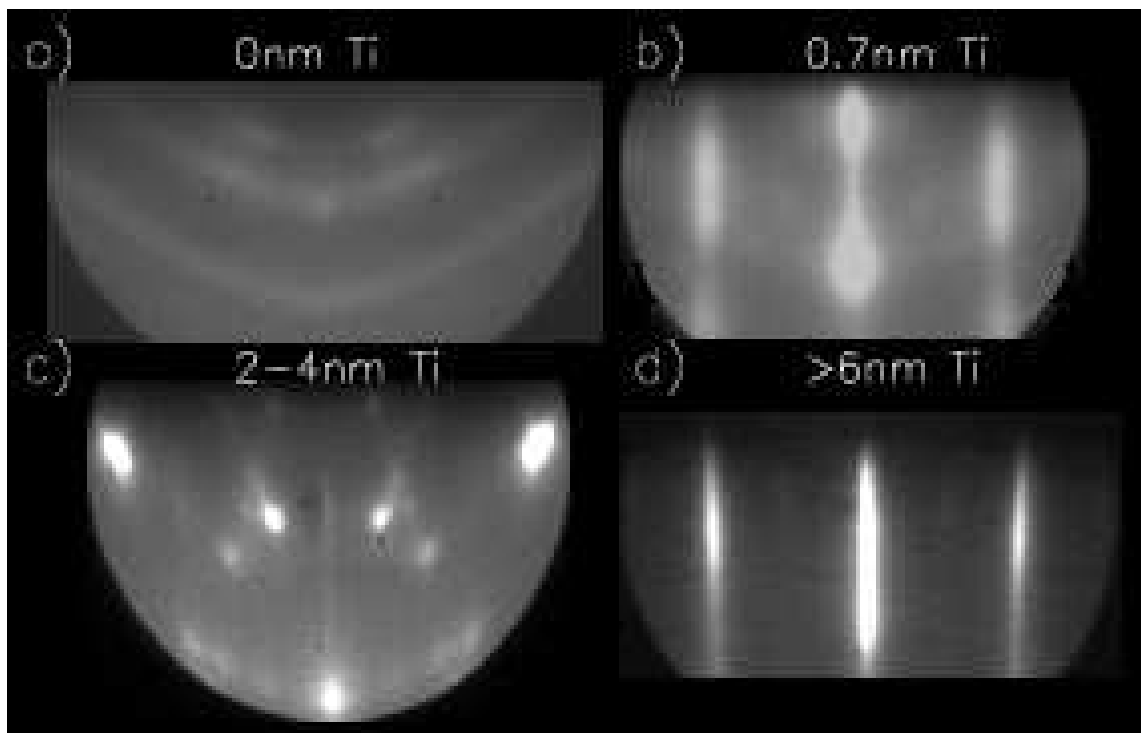


Fig. 6. The diffraction patterns of silver on various Ti/ Al_2O_3 films. Observed with the beam along $Al_2O_3[1\bar{1}00]$.

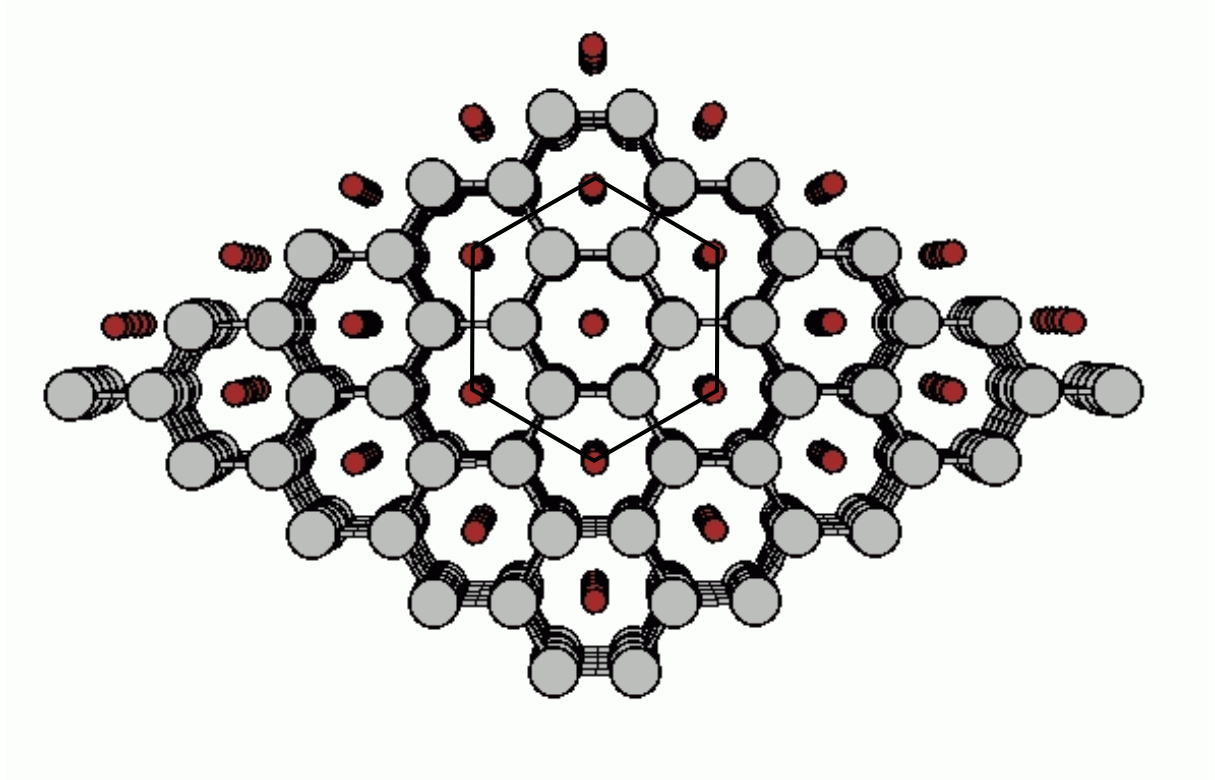


Fig. 7. Schematic top view of the $\omega(0001)$ surface. The structure can be seen as the successive stacking a triangular lattice and a graphite like sheet. The indicated hexagon with an edge of 0.455 nm corresponds to the underlying lattice.

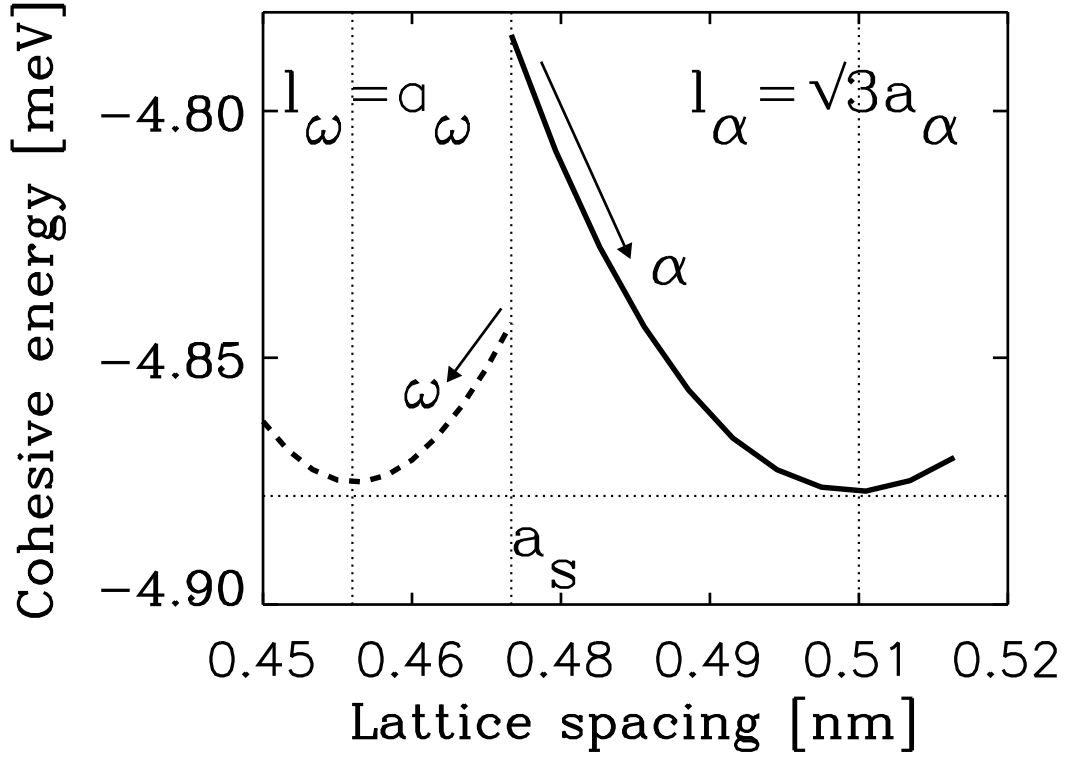


Fig. 8. Calculated total energy per atom of α -Ti (full line) (resp. ω -Ti (dashed line)), for lattice parameters ranging from the one giving the perfect lattice matching with the substrate (*i.e.* $a_\alpha = 0.475/\sqrt{3}$ nm for α -Ti and $a_\omega = 0.459\text{\AA}$ for ω -Ti), up (down) to the equilibrium lattice constant. Note that the total energy of α -Ti is represented as a function of $l_\alpha = \sqrt{3}a_\alpha$ whereas the total energy of ω -Ti is represented as a function of $l_\omega = a_\omega$.