



HAL
open science

Anomalous cooling of the parallel velocity in seeded beams

Alain Miffre, Marion Jacquey, Matthias Büchner, Gérard Tréneç, Jacques Vigué

► **To cite this version:**

Alain Miffre, Marion Jacquey, Matthias Büchner, Gérard Tréneç, Jacques Vigué. Anomalous cooling of the parallel velocity in seeded beams. 2004. hal-00001006v2

HAL Id: hal-00001006

<https://hal.science/hal-00001006v2>

Preprint submitted on 14 Mar 2004 (v2), last revised 16 Jun 2004 (v3)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Anomalous cooling of the parallel velocity in seeded beams

A. Miffre, M. Jacquey, M. Büchner, G. Tréneç and J. Vigué

Laboratoire Collisions Agrégats Réactivité -IRSAMC

Université Paul Sabatier and CNRS UMR 5589 118, Route de Narbonne 31062 Toulouse Cedex, France

e-mail: jacques.vigue@irsamc.ups-tlse.fr

(March 14, 2004)

We have measured the parallel velocity distribution of a lithium supersonic beam produced by seeding lithium in argon. The parallel temperature for lithium is considerably lower than the calculated parallel temperature of the argon carrier gas. We have extended the theory of supersonic cooling to calculate the parallel temperature of the seeded gas, in the limit of high dilution. The theoretical result thus obtained is in good agreement with our observations.

I. INTRODUCTION

We have observed an unexpected cooling effect in an argon supersonic beam seeded with a small amount of lithium: the lithium parallel temperature is smaller than the argon parallel temperature. This effect is surprisingly large, the lithium parallel temperature being 3 times smaller than the calculated argon parallel temperature, whereas the common assumption is that these two parallel temperatures are roughly equal.

Although surprising, this effect can be explained by an extension of the theory giving the terminal parallel temperature in supersonic beams of pure gases. This theory, first developed by Anderson and Fenn in 1965 [1], has been completed by Toennies and Winkelmann [2] in 1977 and Beijerinck and Verster [3] in 1981 (for a brief review, see ref. [2]). In his 1988 review on free jet sources, D.R. Miller [4] discusses the case of expansion of binary mixtures: this case has been studied by various authors but their results, which are most accurate for a dilute heavy species, do not predict a large temperature difference in our case.

The basic phenomenon is due to the exchange of energy between the alkali atoms and the rare gas atoms. Two effects are important: i) the alkali-rare gas interaction has a longer range than the rare gas-rare gas interactions so that the expansion cooling goes on for a longer time for the alkali atoms; ii) the mass ratio plays an important role in the energy exchange between the two species and, as shown below, the parallel temperature of the seeded gas (mass m_2) is considerably lower than the parallel temperature of the carrier gas (mass m_1) when $m_2 \ll m_1$.

In this paper, we first describe our measurement of the parallel temperature of lithium seeded in argon. Then, we recall the theory explaining the final temperature in supersonic beams and we briefly describe its extension to the case of a gas mixture in the limit of a high dilution. We thus get differential equations coupling the parallel

and perpendicular temperatures of the seeded gas to the same quantities of the carrier gas. Finally, we compare our theoretical result to our experiment as well as to another evidence of this anomalous cooling effect.

II. OUR MEASUREMENT

We have measured the parallel temperature of our lithium beam by laser induced fluorescence. The fluorescence of the ${}^7\text{Li}$ isotope (92.5% natural abundance) is excited by a single frequency laser operating at 671 nm, using the ${}^2S_{1/2} - {}^2P_{3/2}$ resonance transition. Two laser beams are used, one near normal incidence measures the transverse velocity distribution, while the other one, with an incidence close to 45° , is sensitive to a combination of the parallel and perpendicular velocity distributions. We use a low laser power density in order to prevent any saturation broadening. Figure 1 presents a typical spectrum of the fluorescence as a function of the laser frequency. Each laser beam excites two resonance lines corresponding to the two hyperfine levels $F = 1$ and $F = 2$ of the ${}^2S_{1/2}$ ground state, with an hyperfine splitting equal to 803.5 MHz, while the very small hyperfine structure of the ${}^2P_{3/2}$ excited state cannot be resolved [5].

The distance between the lines corresponding to the same transition excited by the two laser beams gives access to the mean velocity of the beam, while the linewidths give access to the widths of the velocity distributions. The observed perpendicular velocity distribution depends on the width of the observed part of the lithium beam and our observations correspond to a perpendicular temperature *in the 0.4 – 0.7 K range*. As expected, the parallel velocity distribution is broader than the perpendicular velocity distribution (in our excitation geometry with a 45° incidence, the Doppler width is only 0.71 times the width we would get if the laser and atomic beams were parallel). The measured parallel temperature is $T_{\parallel} = 6.6 \pm 0.6$ K. *Unfortunately, as we cannot measure the parallel temperature of the argon beam, we must deduce its value, 19.0 ± 2.7 K, from the source parameters.*

III. TERMINAL PARALLEL TEMPERATURE IN SUPERSONIC EXPANSIONS

We consider here only the case of monoatomic gases. The starting point is the Boltzmann equation. In the

ccsd-00001006 (version 2) : 14 Mar 2004

frame moving with the hydrodynamic velocity u , the velocity distribution is assumed to be Maxwellian with different parallel and perpendicular temperatures $T_{\parallel 1}$ and $T_{\perp 1}$:

$$f_1(\mathbf{v}) = \left(\frac{m_1}{2\pi k_B T_{\parallel 1}} \right)^{1/2} \times \frac{m_1}{2\pi k_B T_{\perp 1}} \times \exp \left[-\frac{m_1 v_{\parallel}^2}{2k_B T_{\parallel 1}} - \frac{m_1 v_{\perp}^2}{2k_B T_{\perp 1}} \right] \quad (1)$$

m_1 is the atomic mass. We use local cartesian coordinates, the x and y axis being in the perpendicular directions and the z axis being in the parallel direction. The theoretical treatment of Toennies and Winkelmann [2] (noted below TW) and the one of Beijerinck and Verster [3] (noted below BV) are different but equations (14),(15b) and (16) of reference [2] are equivalent to equations (29a,b) and (B6,7) of reference [3]. However, small terms of the order of S_{\parallel}^{-2} , where the parallel speed ratio S_{\parallel} is defined by $S_{\parallel} = \sqrt{mu^2/(2k_B T_{\parallel})}$ are neglected in the BV calculations, which are a good approximation only when the terminal value $S_{\parallel\infty}$ of the parallel speed ratio is large. As we are interested in this case here, we follow the BV calculations because they are simpler. The parallel and perpendicular temperatures are given by:

$$\frac{dT_{\perp 1}}{dz} = -\frac{2T_{\perp 1}}{z} + F \text{ and } \frac{dT_{\parallel 1}}{dz} = -2F \quad (2)$$

where F is due to the collision-induced energy transfer from parallel to perpendicular degrees of freedom:

$$F = \frac{n_1(z)}{2k_B u} \int g \frac{d\sigma_{1,1}(g)}{d\Omega} \Delta E f_1(\mathbf{v}_1) f_1(\mathbf{v}_2) d^3 \mathbf{v}_1 d^3 \mathbf{v}_2 d\Omega \quad (3)$$

\mathbf{v}_1 and \mathbf{v}_2 are the atom velocities before the collision and $g = |\mathbf{v}_1 - \mathbf{v}_2|$ is their relative velocity. The indices 1, 1 recall that the collision considered here involve two atoms of species 1. ΔE is the energy transferred during one collision from the parallel degree of freedom to the perpendicular ones; after averaging over the azimuth of the final relative velocity around the initial relative velocity, ΔE is given by :

$$\Delta E = \frac{m}{8} [g_x^2 + g_y^2 - 2g_z^2] [1 - \cos^2 \chi] \quad (4)$$

where χ is the deflection angle. F is well approximated by a linear function of $(T_{\parallel 1} - T_{\perp 1})$:

$$F \approx \Lambda(z)(T_{\parallel 1} - T_{\perp 1}) \text{ with } \Lambda(z) = 16n_1(z)\Omega_{1,1}^{(2,2)}(T_m)/(15u_{\infty}) \quad (5)$$

Here, $\Omega_{1,1}^{(l,s)}(T_m)$ defined by reference [6] (see equation 7.4-34) is a thermal average at the mean temperature $T_m = (T_{\parallel 1} + 2T_{\perp 1})/3$ of the collision integral $Q_{1,1}^{(l)}$ given by:

$$Q_{1,1}^{(l)} = \int \frac{d\sigma_{1,1}(g)}{d\Omega} (1 - \cos^l \chi) d\Omega \quad (6)$$

When quantum effects can be neglected, if the interaction potential is approximated by a 12-6 Lennard-Jones potential with a $C_6(1, 1)$ long-range coefficient and if the relative kinetic energy is small with respect to the interaction potential well depth ϵ , the $\Omega_{1,1}^{(2,2)}(T)$ integral is given by [3]:

$$\Omega_{1,1}^{(2,2)}(T) = 2.99 (2k_B T/m_1)^{1/2} (C_6(1, 1)/k_B T)^{1/3} \quad (7)$$

Knowing $\Lambda(z)$, which behaves like $z^{-2} T_m^{1/6}$, the differential equations (2) can be integrated numerically. *The parallel temperature tends toward a constant while the perpendicular temperature decreases indefinitely, a situation very far from thermodynamic equilibrium. The perpendicular temperature cools for geometrical reasons discussed by reference [2]. This calculation relates the terminal parallel temperature $T_{\parallel\infty}$ to the density n_0 and temperature T_0 in the source and to the nozzle diameter d . The tradition is to give the terminal value $S_{\parallel\infty}$ of the parallel speed ratio:*

$$S_{\parallel\infty} = A \left[n_0 d (C_6/k_B T_0)^{1/3} \right]^{\alpha} \quad (8)$$

Using SI units, the various values of A and α are $A = 1.412$ and $\alpha = 0.53$ obtained by TW [2], $A = 1.313$ and $\alpha = 0.545$ obtained by BV [3], $A = 1.782$ and $\alpha = 0.495$ obtained also by BV from a fit of experimental $S_{\parallel\infty}$ values for argon expansions. *We will use this semi-empirical form of equation (8) to calculate the argon parallel temperature and we have estimated a 7% error bar on $S_{\parallel\infty}$ from the dispersion of the data set used by BV and by comparison with other measurements [7].*

IV. GENERALIZATION TO THE CASE OF A MIXTURE OF TWO MONOATOMIC GASES

Our calculation is made in the case of a high dilution, i.e. when the seeded gas density n_2 is considerably smaller than the carrier gas density n_1 . We neglect the velocity slip i.e. we assume the same hydrodynamic velocity u for both species. Because $n_2 \ll n_1$, the expansion of the carrier gas is not modified by the presence of the seeded gas and the equations written above remain valid. For the seeded gas, we write similar equations, considering only the collisions with atoms of the carrier gas. *The calculation of the variation of the parallel and perpendicular energies of an atom 2 during a collision with an atom 1 is straightforward but too complex to be detailed here.* In order to linearize the equations, we have remarked that the products $gQ_{1,2}^{(l)}(g) \propto g^{1/3}$ vary slowly with g and we have treated *them as independent of g .* *The values of these products are then chosen to be coherent with the treatment of the pure gas case described in*

part III. We thus get the differential equations relating $T_{\parallel 2}$ and $T_{\perp 2}$:

$$\frac{dT_{\perp 2}}{dz} = -\frac{2T_{\perp 2}}{z} + \Lambda(z)\rho_s \frac{m_1}{M} [T_{\parallel av} - T_{\perp av}] - 4\Lambda(z)\rho_s \rho_o \frac{\mu}{M} [T_{\perp 2} - T_{\perp 1}] \quad (9)$$

$$\frac{dT_{\parallel 2}}{dz} = -2\Lambda(z)\rho_s \frac{m_1}{M} [T_{\parallel av} - T_{\perp av}] - 4\Lambda(z)\rho_s \rho_o \frac{\mu}{M} [T_{\parallel 2} - T_{\parallel 1}] \quad (10)$$

with $T_{\parallel av} = \beta T_{\parallel 1} + \alpha T_{\parallel 2}$, $T_{\perp av} = \beta T_{\perp 1} + \alpha T_{\perp 2}$, $\alpha = m_1/(m_1 + m_2)$ and $\beta = m_2/(m_1 + m_2)$. ρ_s is the ratio of $\Omega_{i,j}^{(2,2)}$ collision integrals:

$$\rho_s = \frac{\Omega_{1,2}^{(2,2)}}{\Omega_{1,1}^{(2,2)}} = \left[\frac{C_6(1,2)}{C_6(1,1)} \right]^{1/3} \times \left[\frac{m_1 + m_2}{2m_2} \right]^{1/2} \quad (11)$$

ρ_o is the ratio of the $l = 1$ and $l = 2$ angle-averaged cross-sections $\rho_o = Q_{1,2}^{(1)}/Q_{1,2}^{(2)}$. Using the low energy deflection function for a 12-6 Lennard-Jones potential (see appendix A of [3]), we have calculated $\rho_o = 1.32$.

Following BV, we have introduced reduced temperatures and a reduced z -coordinate. Then, the numerical integration of the coupled equations gives the temperatures represented in figure 2. At the end of the expansion, the parallel and perpendicular temperatures of the seeded gas are lower than the similar quantities for the carrier gas. *It is difficult to have a very simple explanation of this effect, because of the complexity of equations (9,10). It is nevertheless clear that the parallel temperature of species 2 remains coupled to the perpendicular temperatures when the parallel temperature of species 1 is no more coupled.*

The parallel temperature ratio $T_{\parallel 2}/T_{\parallel 1}$ can be substantially lower than 1. The minimum value for temperature $T_{\parallel 2}$ is reached when the right-hand side of equation (9) vanishes. Neglecting the perpendicular temperatures, we thus get:

$$\frac{T_{\parallel 2, \min}}{T_{\parallel 1\infty}} = \frac{m_2 \rho_o}{m_1 + 2m_2 \rho_o} \quad (12)$$

This limiting value is reached if the number of 1 – 2 collisions is sufficiently larger than the number of 1 – 1 collisions, i.e. if the ratio ρ_s is large enough. We have plotted in figure 3 the terminal value of the ratio $T_{\parallel 2}/T_{\parallel 1}$ for various values of the ratio ρ_s and for different values of the mass ratio m_2/m_1 .

V. COMPARISON WITH EXPERIMENTAL RESULTS

We may now compare our theoretical and experimental results for lithium seeded in argon. In our experiment

with $p_0 = 333$ millibar, $T_0 = 943$ K and a nozzle diameter $d = 200 \mu\text{m}$, the parallel temperature of lithium is $T_{\parallel 2\infty} = 6.6 \pm 0.6$ K. Using equation (8), we deduce the argon parallel temperature from our source conditions, $T_{\parallel 1\infty} = 19.0 \pm 2.7$ K. We thus get a parallel temperature ratio $T_{\parallel 2\infty}/T_{\parallel 1\infty} = 0.35 \pm 0.08$. This value *is in good agreement with* our theoretical result $T_{\parallel 2\infty}/T_{\parallel 1\infty} = 0.31$ obtained with $\rho_s = 2.55$ deduced from the C_6 values of argon-argon and lithium-argon interaction [8].

A similar beam with sodium seeded in argon was built by D. Pritchard and coworkers [9,10] with a source temperature near 1000 K, an argon pressure up to $p_0 = 3$ bar and a nozzle diameter $d = 70 \mu\text{m}$. The measured velocity $u = 1040 \pm 2$ m/s corresponds to $T_0 = 1039$ K and the rms velocity width, deduced from the atomic diffraction pattern, was found equal to $3.7 \pm 0.4\%$, corresponding to $T_{\parallel 2\infty} = 4.1 \pm 0.9$ K. From the calculated argon parallel temperature $T_{\parallel 1} = 7.7 \pm 1.1$ K corresponding to the largest pressure, *we get the parallel temperature ratio $T_{\parallel 2\infty}/T_{\parallel 1\infty} = 0.53 \pm 0.19$, in good agreement with our theoretical result is $T_{\parallel 2\infty}/T_{\parallel 1\infty} = 0.61$, using the C_6 values from reference [8] ($\rho_s = 1.67$).*

VI. CONCLUSION

In this letter, we have described a measurement of the parallel velocity distribution of a supersonic beam of lithium seeded in argon, with a high dilution. The measured lithium parallel temperature 6.6 ± 0.6 K is considerably lower than the calculated parallel temperature of the argon beam near 19.0 ± 2.7 K.

We have extended the theory of the velocity distribution in supersonic beams to a mixture of monoatomic gases, in the high dilution limit. We have established differential equations describing the parallel and perpendicular temperatures of the seeded gas; these equations explain the observed effect as a consequence of several circumstances: role of the mass ratio in energy exchange during collisions of the two species, longer range of the lithium-argon interaction as compared to argon-argon interaction. We have deduced the lowest possible value of the temperature ratio $T_{\parallel 2\infty}/T_{\parallel 1\infty}$. This ratio may be quite low when a light species is seeded in a heavy species with $m_2 \ll m_1$.

We have compared our theory with experimental results corresponding to lithium seeded in argon (our experiment) and sodium seeded in argon (an experiment done by D. Pritchard and coworkers). In both cases the experimental and theoretical values of the ratio of the parallel temperatures are in good agreement, but the error bars on the experimental values are large. A simultaneous measurement of the parallel temperatures of the carrier and seeded gases would be useful to reduce these error bars, thus providing a better test of the approximations done in our theoretical analysis.

VII. ACKNOWLEDGEMENTS

We thank J.P. Toennies for his help and interest in this question. We also thank Ph. Dugourd and M. Broyer for advice concerning the design of the oven of our lithium supersonic beam. We thank Région Midi Pyrénées for financial support.

-
- [1] J. B. Anderson and J. B. Fenn, Phys. Fluids **8**, 780 (1965)
 - [2] J. P. Toennies and K. Winkelmann, J. Chem. Phys. **66**, 3965 (1977)
 - [3] H. C. W. Beijerinck and N. F. Verster, Physica **111C**, 327 (1981)
 - [4] D. R. Miller, in Atomic and molecular beam methods, G. Scoles ed., p. 14-53 (Oxford University Press 1988)
 - [5] C. J. Sansonetti et al., Phys. Rev. **A 52**, 2682 (1995)
 - [6] J. O. Hirschfelder, C. F. Curtiss and R. B. Bird, Molecular theory of liquids, (Wiley, New York, 1954)
 - [7] H. D. Meyer, MPI fr Stromungsforschung, Bericht 5 (Göttingen, 1978)
 - [8] K. T. Tang, J. M. Norbeck and P. R. Certain, J. Chem. Phys. **64**, 3063 (1976)
 - [9] J. Schmiedmayer et al., in Atom interferometry edited by P. R. Berman (Academic Press 1997), p 1
 - [10] C. R. Ekstrom et al., Phys. Rev. A **51**, 3883 (1995)

FIG. 1. Laser induced fluorescence signal as a function of the laser frequency: the dots represent the experimental data and the full curves are the Gaussian fits to each line. The Fabry-Perot used for calibration has a free spectral range equal to 251.4 ± 0.5 MHz. The fluorescence peaks are labeled by the ground state F value ($F = 1$ or 2) and by a letter corresponding to the laser beam: A for the beam near 45° incidence, B for the beam near normal incidence.

FIG. 2. *The reduced temperatures $\tau = \Xi^{-9/11}T/T_0$ are plotted as a function of the reduced z -coordinate $\zeta = 2.48\Xi^{12/11}z/d$ with $\Xi = 0.813n_0d(C_6/k_B T_0)^{1/3}$ (to give an example, for our argon beam seeded with lithium, $\Xi = 32.4$). This plot, which is identical to figure 5 of Beijerinck and Verster [3] for $\tau_{\parallel 1}$ and $\tau_{\perp 1}$, shows the anomalous cooling effect for species 2. We have used $\rho_o = 1.32$ and $\rho_s = 2.55$ corresponding to lithium seeded in argon.*

FIG. 3. The calculated parallel temperature ratio $T_{\parallel 2\infty}/T_{\parallel 1\infty}$ is plotted as a function of the mass ratio m_2/m_1 (spanning the range from hydrogen to rubidium seeded in argon) for several values of the ratio ρ_s , with ρ_o fixed, $\rho_o = 1.32$. From bottom to top, $\rho_s = \infty$, $\rho_s = 2.55$ (*lithium case*), $\rho_s = 1.67$ (*sodium case*). The experimental points corresponding to lithium and sodium seeded in argon are represented by stars.